

COSMOLOGY

A NOTE BY CHESTNUT

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Hello, this the notes of when I was studying cosmology, please take it if you needed. The reference books are:

1. *Modern Cosmology*, Scott Dodelson, Fabian Schmidt (Main)
2. *Cosmology*,

Thanks. I also have other notes in English or Chinese with latex. Or may be in the near future with Japneses, ;)

- Theoretical Mechanics (En.) done, in revising
- Quantum Mechanics (En.) done, in revising
- General Relativity (En.) to do, in learning
- Field Theory (En.) to do, to learn
- Electrodynamics (Ch.) doing, in reviewing
- Statistical Mechanics (Ch.) done, in revising
- Some self-solved notes of Japneses M.S. physical exam, want to do, in learning

Welcome to email me if you would like to have or discuss.

2024.03.08

Contents

I	The Expansion of the Universe	1
1	Intro: The concordance model of cosmology	1
1.1	Scale factor	1
1.2	Λ CDM	1
2	The expanding universe	2
2.1	Expanding space	2
2.1.1	The metric	2
2.1.2	The geodesic equation	2
2.2	Distance	2
2.3	Evolution of energy	3
2.4	Cosmic inventory	3
3	The fundamental equations of cosmology	5
3.1	Einstein equations	5
3.2	Boltzmann equation	5
3.3	General theory of small fluctuations	6
3.3.1	Perturbed spacetime	6
3.3.2	The geodesic equation	6
3.3.3	The collisionless Boltzmann equation for photons	6
3.3.4	The collisionless Boltzmann equation for massive particles	6
II	Evolution of structure	7
4	The origin of species	7
4.1	The homogeneous Boltzmann equation revisited	7
4.2	Big Bang nucleosynthesis	7
4.3	Recombination	7
4.4	Dark matter	7
5	Evolution of fluctuations: matter & radiation	8
5.1	The collision Boltzmann equation for photons	8
5.1.1	Distribution function	8
5.1.2	Collision term: Compton Scattering	8
5.1.3	The Boltzmann equation for photon	8
5.2	The Boltzmann equation for cold dark matter	8
5.3	The Boltzmann equation for baryon	9
5.4	The Boltzmann equation for neutrinos	9

6	Evolution of fluctuations: gravity	10
6.1	Scalar vector-tensor decomposition	10
6.2	From gauge to gauge	10
6.3	The Einstein equation for scalar perturbations	10
6.3.1	Ricci tensor	10
6.3.2	Two components of the Einstein equations	10
6.4	Tensor perturbations	10
7	Initial conditions	11
7.1	The horizon problems and a solution	11
7.2	Inflation	11
7.3	Gravitational wave production	11
7.3.1	harmonic oscillator	11
7.3.2	Tensor perturbation	11
7.4	Salar perturbations	11
7.4.1	Salar field perturbations around an unexpected background	11
7.4.2	Super-horizon perturbation	11
7.4.3	Spatially flat slicing	12
7.5	The Einstein-Boltzmann equations at early times	12
III	The Growth of Structure	13
8	Growth of structure: linear theory	13
8.1	Prelude	13
8.2	Large-Scale	13
9	The cosmic microwave background	14
9.1	Overview	14
9.2	Large-scale anisotropies	14
9.3	Acoustic oscillations	14
9.3.1	Tightly-coupled limit of the Boltzmann equations	14
9.3.2	Tightly-coupled solutions	14
9.4	Diffusion damping	14
9.5	Inhomogeneities to anisotropies	14
9.5.1	Free streaming	14
9.5.2	The angular power spectrum	14
9.6	The CMB power spectrum	15
9.7	Cosmological parameters	15
10	The polarized CMB	16
10.1	Acoustic oscillations	16
10.1.1	Tightly-coupled limit of the Boltzmann eq	16
10.1.2	Tightly-coupled solutoins	16
10.2	Diffusion damping	16
10.3	Inhomogeneities to anisotropies	16

11 Probes of structure: tracer	17
11.1 Galaxy clustering	17
11.1.1 Galaxy clustering	17
11.1.2 Redshift-space distortions	17
11.1.3 BAO and Alcock-Paczyński	17
12 Growth of structure: beyond linear theory	18
13 Probes of structure: lensing	19
14 Analysis and inference	20

The Expansion of the Universe

SECTION 1

Intro: The concordance model of cosmology

SUBSECTION 1.1

Scale factor

Introduce scale factor: a for expansion, physical distance = comoving distance $\times a$

Introduce stretching factor: z , observed wavelength = emit wavelength $\times 1/a_{\text{emit}}$

Definition 1 | Redshift:

$$1 + z \equiv \frac{\lambda_{\text{obs}}}{\lambda_{\text{emit}}} = \frac{a_{\text{obs}}}{a_{\text{emit}}} = \frac{1}{a_{\text{emit}}}$$

Universe geometry: Euclidean, open, closed. history of universe: evolution of a with t , determined by energy density.

Definition 2 | Hubble rate: measure how rapidly the scale factor changes.

$$H(t) \equiv \frac{1}{a} \frac{da}{dt}$$

Hubble constant: $H_0 \equiv H(t=0) = \dot{a}$

- Euclidean and matter dominated, $a \approx t^{2/3}$

Definition 3 | Critical density:

$$\rho_{\text{cr}} \equiv \frac{3H_0^2}{8\pi G}$$

Hubble found Distance galaxies are receding from us, i.e. redshifted. comoving motion $\dot{x} = 0$

Law 1 | Hubble-Lemaitre law: distance-redshift relation, the relative velocity is

$$v = \frac{d}{dt}(ax) = \dot{a}x = H_0 d \quad (v \ll c) \quad (1.1)$$

Eq. (1.1) break down

SUBSECTION 1.2

Λ CDM

Early universe hotter and denser, temperature $1\text{MeV}/k_B$, no neutral atoms or bound nuclei. Then light forms (Big Bang Nucleosynthesis, BBN), \rightarrow CMB smooth \rightarrow not completely smooth \rightarrow inhomogeneities (universe structure)

$$H_0 = 100 h \text{ km s}^{-1} \text{ Mpc}^{-1}$$

$$= \frac{h}{0.98 \times 10^{10} \text{ years}}$$

$$= 2.13 \times 10^{-33} \frac{\text{eV}}{h}$$

$$1 \text{ Mpc} = 3.856 \times 10^{-24} \text{ cm}$$

Friedmann equation:

$$H^2(t) = \frac{8\pi G}{3} \left[\rho(t) + \frac{\rho_{\text{cr}} - \rho(t_0)}{a^2(t)} \right]$$

$$\rho_{\text{cr}} = 1.88 h^2 \times 10^{-29} \text{ g cm}^{-3}$$

SECTION 2

The expanding universe

$$\hbar = c = k_B = 1$$

SUBSECTION 2.1

Expanding space

2.1.1 The metric

To describe a curved spacetime, general relativity

Definition 4 Metric: returns the actual physical distance between two infinitesimally close points in spacetime defined in some arbitrary coordinate system.

Proper-time interval

$$ds^2 = \sum_{\mu, \nu=0}^3 g_{\mu\nu} dx^\mu dx^\nu$$

For a Euclidean universe:

$$g_{\mu\nu} = \begin{pmatrix} -1 & 0 & 0 & 0 \\ 0 & a^2(t) & 0 & 0 \\ 0 & 0 & a^2(t) & 0 \\ 0 & 0 & 0 & a^2(t) \end{pmatrix}$$

2.1.2 The geodesic equation

$$\frac{d^2 x^\mu}{d\lambda^2} + \Gamma^\mu_{\alpha\beta} \frac{dx^\alpha}{d\lambda} \frac{dx^\beta}{d\lambda} = 0$$

Christoffel symbol:

$$\Gamma^\mu_{\alpha\beta}$$

Four-dimensional energy-momentum vector: $P^\alpha = (E, \mathbf{P})$

$$P^\alpha = \frac{dx^\alpha}{d\lambda}$$

Basic feature:

- Redshift of light
- Notion of distance
- Evolution of ρ
- Epoch of equality a_{eq}

SUBSECTION 2.2

Distance

Total comoving distance

$$\chi(t) = \int_t^{t_0} \frac{dt'}{a(t')} = \int_{a(t)}^1 \frac{da'}{a'^2 H(a')} = \int_0^z \frac{dz'}{H(z')}$$

Definition 5 Comoving horizon:

$$\eta(t) \equiv \int_0^t \frac{dt'}{a(t')}$$

Angular diameter distance:

$$d_A = \frac{l}{\theta}$$

In Euclidean universe:

t' : Conformal time

l : physical size
 θ : angle subtended

$$\left| \begin{array}{l} \Gamma^0_{00} = 0, \\ \Gamma^0_{0i} = \Gamma^0_{i0} = 0, \\ \Gamma^0_{ij} = \delta_{ij} \dot{a} a, \\ \Gamma^i_{0j} = \Gamma^i_{j0} = \delta_{ij} \frac{\dot{a}}{a} \end{array} \right|$$

$$d_A^{\text{Euc}} = a\chi = \frac{\chi}{1+z}$$

Definition 6 | Luminosity distance in a Euclidean expanding universe:

$$d_L^{\text{Euc}} \equiv \frac{\chi}{a}$$

SUBSECTION 2.3

Evolution of energy

Energy-momentum tensor in isotropic smooth universe:

$$T^\mu{}_\nu = \begin{pmatrix} -\rho & 0 & 0 & 0 \\ 0 & \mathcal{P} & 0 & 0 \\ 0 & 0 & \mathcal{P} & 0 \\ 0 & 0 & 0 & \mathcal{P} \end{pmatrix}$$

\mathcal{P} : pressure

Law 2 | Local energy and momentum conservation,

$$\nabla_\mu T^\mu{}_\nu \equiv 0$$

conservation law in an expanding universe,

$$\frac{\partial \rho}{\partial t} + \frac{\dot{a}}{a} [3\rho + 3\mathcal{P}] = 0$$

Definition 7 | Equation of state parameter

$$w_s \equiv \frac{\mathcal{P}_s}{\rho_s}$$

Evolution of any constituent of s

$$\rho_s \propto \exp \left\{ -3 \int^a \frac{da'}{a} [q + w_s(a')] \right\}$$

$$\stackrel{w_s = \text{const}}{\propto} a^{-3(1+w_s)}$$

SUBSECTION 2.4

Cosmic inventory

Definition 8 | Density parameters

$$\Omega_s \equiv \frac{\rho_s(t_0)}{\rho_{\text{cr}}}$$

$$\rho_s(a) = \Omega_s \rho_{\text{cr}} a^{-3(1+w_s)}$$

c: cold dark matter
b: baryons
 γ : photons
 ν : neutrinos

Photons

Baryons

Dark matter

Neutrinos

Epoch of matter-radiation equality

$$a_{\text{eq}} =$$

$$1 + z_{\text{eq}} = 2.38 \times 10^4 \Omega_{\text{m}} h^2$$

Dark energy

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SECTION 3

The fundamental equations of cosmology

SUBSECTION 3.1

Einstein equations

$$G_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi G T_{\mu\nu} \quad (3.1)$$

Definition 9 Einstein tensor:

$$G_{\mu\nu} \equiv R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R$$

Ricci tensor:

$$R_{\mu\nu} = \Gamma^\alpha_{\mu\nu,\alpha} - \Gamma^\alpha_{\mu\alpha,\nu} + \Gamma^\alpha_{\beta\alpha} \Gamma^\beta_{\mu\nu} - \Gamma^\alpha_{\beta\nu} \Gamma^\beta_{\mu\alpha}$$

Time-time component of the Einstein equations:

$$R_{00} - \frac{1}{2} g_{00} R = 8\pi G T_{00}$$

$$\begin{cases} R_{00} = -3\ddot{a}/a \\ R_{ij} = \delta_{ij}[2\dot{a}^2 + a\ddot{a}] \end{cases}$$

Lead to first Friedmann equations:

Ricci scalar: $R \equiv g_{\mu\nu} R_{\mu\nu}$

$$\left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi G}{3} \rho$$

FLRW metric and assuming a Euclidean universe,

$$\frac{H^2(t)}{H_0^2} = \frac{\rho(t)}{\rho_{\text{cr}}} = \sum_{s=r,m,\nu,DE} \Omega_s[a(t)]^{-3(1+w_s)}$$

To calculate the evolution of the homogeneous universe.

$$\frac{H^2(t)}{H_0^2} = \sum_{s=r,m,\nu,DE} \Omega_s[a(t)]^{-3(1+w_s)} + \Omega_K[a(t)]^{-2}$$

SUBSECTION 3.2

Boltzmann equation

$$\frac{df}{dt} = C[f]$$

Boltzmann equation in an expanding universe

$$\frac{df}{dt} = \frac{\partial f}{\partial t} + \frac{\partial f}{\partial x^i} \cdot \frac{dx^i}{dt} + \frac{\partial f}{\partial p} \cdot \frac{dp}{dt} + \frac{\partial f}{\partial \hat{p}^i} \cdot \frac{d\hat{p}^i}{dt}$$

In homogeneous universe,

$$\frac{df}{dt} = \frac{\partial f}{\partial t} + \frac{p}{E} \frac{\hat{p}^i}{a} \frac{\partial f}{\partial x^i} - H p \frac{\partial f}{\partial p} \quad (3.2)$$

Eq. 3.2 is valid for all particles.

Collision term

SUBSECTION 3.3

General theory of small fluctuations

General theory of small fluctuations

3.3.1 Perturbed spacetime

Perturbed metric:

$$\begin{aligned} g_{00}(\mathbf{x}, t) &= -1 - \Psi(\mathbf{x}, t), \\ g_{0i}(\mathbf{x}, t) &= 0, \\ g_{ij}(\mathbf{x}, t) &= a^2(t)\delta_{ij}[1 + 2\Phi(\mathbf{x}, t)]. \end{aligned}$$

- Ψ : corresponds to Newtonian potential and governs the motion of slow-moving (non-relativistic) bodies;
- Φ : the perturbation to the spatial curvature

$$\begin{aligned} a(t) &\rightarrow \\ a(\mathbf{x}, t) &= q(t)\sqrt{1 + \Phi(\mathbf{x}, t)} \end{aligned}$$

3.3.2 The geodesic equation

Via perturbed metric,

$$P^\mu = \left[E(1 - \Psi), p^i \frac{1 - \Phi}{a} \right]$$

The change in the magnitude of the momentum of a particle as it moves through a perturbed FLRW universe,

$$\frac{dp}{dt} = -[H + \dot{\Phi}]p - \frac{E}{a}\hat{p}^i\Psi_{,i}$$

Energy-momentum tensor in the perturbed universe for a single species with degeneracy factor g :

$$T^0_0(\mathbf{x}, t) = -g$$

3.3.3 The collisionless Boltzmann equation for photons

$m = 0, E = p$

$$\frac{df}{dt} = \frac{\partial f}{\partial t} + \frac{\hat{p}^i}{a} \frac{\partial f}{\partial x^i} - \left[H + \dot{\Phi} + \frac{1}{a}\hat{p}^i\Psi_{,i} \right] p \frac{\partial f}{\partial p} \quad (3.3)$$

3.3.4 The collisionless Boltzmann equation for massive particles

$$\frac{df}{dt} = \frac{\partial f}{\partial t} + \frac{p}{E} \frac{\hat{p}^i}{a} \frac{\partial f}{\partial x^i} - \left[H + \dot{\Phi} + \frac{E}{ap}\hat{p}^i\Psi_{,i} \right] p \frac{\partial f}{\partial p} \quad (3.4)$$

Evolution of structure

SECTION 4

The origin of species

SUBSECTION 4.1

The homogeneous Boltzmann equation revisited

SUBSECTION 4.2

Big Bang nucleosynthesis

SUBSECTION 4.3

Recombination

SUBSECTION 4.4

Dark matter

SECTION 5

Evolution of fluctuations: matter & radiation

SUBSECTION 5.1

The collision Boltzmann equation for photons

Derive how matter, photons and neutrinos behave in a given expanding spacetime with perturbation.

5.1.1 Distribution function

$E = p$, distribution function of Bose-Einstein:

$$f(\mathbf{x}, p, \hat{\mathbf{p}}, t) = \left[\exp \left\{ \frac{p}{T(t)[1 + \Theta(\mathbf{x}, \hat{\mathbf{p}}, t)]} \right\} \right]^{-1}$$

- Θ : perturbation to the distribution function, temperature perturbation

Zeroth-order distribution function is set precisely by the requirement that the collision term vanished.

$$\left. \frac{df}{dt} \right|_{\text{first order}} = -p \frac{\partial f^{(0)}}{\partial p} \left[\dot{\Theta} + \frac{\hat{p}^i}{a} \frac{\partial \Theta}{\partial x^i} + \dot{\Phi} + \frac{\hat{p}^i}{a} \frac{\partial \Psi}{\partial x^i} \right]$$

5.1.2 Collision term: Compton Scattering

Definition 10 Monopole:

$$\Theta_0(\mathbf{x}, t) \equiv \frac{1}{4\pi} \int d\Omega' \Theta(\hat{\mathbf{p}}', \mathbf{x}, t)$$

$$C[f(\mathbf{p})] = -p \frac{\partial f^{(0)}}{\partial p} n_e \sigma_T [\Theta_0 - \Theta(\hat{\mathbf{p}}) + \hat{\mathbf{p}} \cdot \mathbf{u}_b]$$

monopole+dipole \rightarrow photon behave like a fluid.

5.1.3 The Boltzmann equation for photon

$$\dot{\Theta} + \frac{\hat{p}^i}{a} \frac{\partial \Theta}{\partial x^i} + \dot{\Phi} + \frac{\hat{p}^i}{a} \frac{\partial \Psi}{\partial x^i} = n_e \sigma_T [\Theta_0 - \Theta + \hat{\mathbf{p}} \cdot \mathbf{u}_b]$$

$$\eta(t) \equiv \int_0^t \frac{dt'}{a(t')}$$

In terms of the conformal time,

Definition 11 The angle between the wavenumber \mathbf{k} and the photon direction $\hat{\mathbf{p}}$ is denoted as μ :

$$\mu \equiv \frac{\mathbf{k} \cdot \hat{\mathbf{p}}}{k}$$

Definition 12 Optical depth

$$\begin{aligned} \tau(\eta) &\equiv \int_{\eta}^{\eta_0} d\eta' n_e \sigma_T a \\ \tau' &\equiv \frac{d\tau}{d\eta} = -n_e \sigma_T a \end{aligned}$$

In Fourier mode,

$$\Theta' + ik\mu\Theta + \Phi' + ik\mu\Psi = -\tau'[\Theta_0 - \Theta + \mu u_b]$$

SUBSECTION 5.2

The Boltzmann equation for cold dark matter

Collisionless,

Definition 13 | Fluid velocity

$$u_c^i \equiv \frac{1}{n_c} \int \frac{d^3p}{(2\pi)^3} f_c \frac{p \hat{p}^i}{E(p)}$$

continuity equation:

$$\frac{\partial \delta_c}{\partial t} + \frac{1}{a} \frac{\partial u_c^i}{\partial x^i} + 3\dot{\Phi} = 0$$

Euler equation:

$$\frac{\partial u_c^j}{\partial t} + H u_c^j + \frac{1}{a} \frac{\partial \Psi}{\partial x^j} = 0$$

In terms of conformal time η and in Fourier space,

$$\begin{aligned} \delta'_c + i k u_c &= -3\Phi', \\ u'_c + \frac{a'}{a} u_c &= -i k \Psi, \end{aligned}$$

SUBSECTION 5.3

The Boltzmann equation for baryon

Compton scattering.

Definition 14 | Dipole

$$\Theta_1(k, \eta) \equiv i \int_{-1}^1 \frac{d\mu}{2} \Theta(\mu, k, \eta)$$

SUBSECTION 5.4

The Boltzmann equation for newtrinos

Colissionless.

Summary

Equation:

$$\Theta' + i k \mu \Theta + \Phi' + i k \mu \Psi = -\tau' \left[\Theta_0 - \Theta + \mu u_b - \frac{1}{2} \mathcal{P}_2(\mu) \Pi \right]$$

And

$$\delta'_c + i k u_c = -3\Phi', \tag{5.1}$$

$$u'_c + \frac{a'}{a} u_c = -i k \Psi, \tag{5.2}$$

$$\delta'_b + i k u_b = -3\Phi', \tag{5.3}$$

$$u'_b + \frac{a'}{a} u_b = -i k \Psi + \frac{\tau'}{R} [u_b + 3i\Theta_1], \tag{5.4}$$

$$\mathcal{N}' = \tag{5.5}$$

SECTION 6

Evolution of fluctuations: gravity

Non-gravitational to gravitational field

SUBSECTION 6.1

Scalar vector-tensor decomposition

Generally, FLRW spacetime perturbed by a small ammount:

$$\begin{aligned} g_{00}(t, \mathbf{x}) &= -1 + h_{00}(t, \mathbf{x}), \\ g_{0i}(t, \mathbf{x}) &= a(t)h_{0i}(t, \mathbf{x}) = a(t)h_{i0}(t, \mathbf{x}), \\ g_{ij}(t, \mathbf{x}) &= a^2(t)[\delta_{ij} + h_{ij}(t, \mathbf{x})] \end{aligned}$$

Theorem 1 | **Decomposition theorem:** perturbation of each type – scalar, vector, and tensor – evolve independently at linear order.

SUBSECTION 6.2

From gauge to gauge

Consider a scalar field $\phi(x)$, separate into background and perturbation:

$$\phi(x) = \bar{\phi}(t) + \delta\phi(t, \mathbf{x})$$

SUBSECTION 6.3

The Einstein equation for scalar perturbations

Conformal-Newtonian gauge,

6.3.1 Ricci tensor

$$\begin{aligned} \delta R &= -12\Psi\left(H^2 + \frac{\ddot{a}}{a}\right) + \frac{2k^2}{a^2}\Psi + 6\Phi_{,00} \\ &\quad - 6H(\Psi, 0 - 4\Phi_{,0}) + 4\frac{k^2\Phi}{a^2} \end{aligned}$$

6.3.2 Two components of the Einstein equations

$$\delta G = -6H\Phi_{,0} + 6\Phi H^2 - 2\frac{k^2\Phi}{a^2}$$

Evolution equation for Ψ and Φ

$$k^2\Phi + 3\frac{a'}{a}\left(\Phi' - \Psi\frac{a'}{a}\right) = 4\pi G a^2[\rho_c\delta_c + \rho_b\delta_b + 4\rho_\gamma\Theta_0 + 4\rho_\nu\mathcal{N}_0]$$

SUBSECTION 6.4

Tensor perturbations

SECTION 7

Initial conditions

Equations governing perturbation around a smooth background,

SUBSECTION 7.1

The horizon problems and a solution

- $k \gg aH$, leave the horizon
- $k \ll aH$, at the end of inflation for all modes (can observe)

SUBSECTION 7.2

Inflation

Scalar field energy-momentum tensor:

$$T^\alpha{}_\beta = g^{\alpha\nu} \frac{\partial\phi}{\partial x^\nu} \frac{\partial\phi}{\partial x^\beta} - \delta^\alpha{}_\beta \left[\frac{1}{2} g^{\mu\nu} \frac{\partial\phi}{\partial x^\mu} \frac{\partial\phi}{\partial x^\nu} + V(\phi) \right]$$

Assume zero-order: homogeneous; first-order: perturbation;

$$\phi'' + 2aH\phi' + a^2 V_{,\phi} = 0 \quad (7.1)$$

SUBSECTION 7.3

Gravitational wave production

7.3.1 harmonic oscillator

$$\frac{d^2x}{dt^2} + \omega^2 x^2 = 0$$

7.3.2 Tensor perturbation

From 7.1,

$$v'' + \left(k^2 - \frac{a''}{a} \right) v = 0$$

Power spectrum of the primordial tensor perturbation

$$P_h = 16\pi G \frac{|v(k, \eta)|^2}{a^2}$$

SUBSECTION 7.4

Salar perturbations

7.4.1 Salar field perturbations around an unexpected background

$$\delta\phi'' + 2aH\delta\phi' + k^2\delta\phi = 0$$

7.4.2 Super-horizon perturbation

Find the spectrum of scalar perturbations emerging from inflation.

Definition 15 | Curvature perturbation,

7.4.3 Spatially flat slicing

SUBSECTION 7.5

The Einstein-Boltzmann equations at early times

Reduce the Boltzmann equations to when $k\eta \ll 1$,

$$\begin{aligned}\Theta'_0 + \Phi' &= 0 \\ \mathcal{N}'_0 + \Phi' &= 0 \\ \delta'_c &= -3\Phi' \\ \delta'_b &= -3\Phi'\end{aligned}$$

Einstein equation to

$$3\frac{a'}{a}\left(\Phi' - \frac{a'}{a}\Psi\right) = 16\pi G a^2 \rho_r \Theta_{r,0}$$

Summary

Power spectrum of a gauge-invariant quantity

$$P_{\mathcal{R}}(k) = \frac{2\pi G H^2}{\epsilon_{\text{sr}} k^3} \Big|_{aH=k}$$

Or

$$P_{\mathcal{R}}(k) = \frac{2\pi}{k^3} \frac{H^2}{m_{\text{Pl}}^2 \epsilon_{\text{sr}}} \Big|_{aH=k} \equiv 2\pi^2 \mathcal{A}_s k^{-3} \left(\frac{k}{k_{\text{p}}}\right)^{n_s-1}$$

- \mathcal{A}_s : variance of curvature perturbations in a logarithmic wave interval centered around the pivot scale k_{p}
- n_s : scalar spectral index

The Growth of Structure

Goal: a prediction for the linear matter power spectrum.

SECTION 8

Growth of structure: linear theory

SUBSECTION 8.1

Prelude

Three stages of evolution

$$\Phi(\mathbf{k}, a) = \frac{3}{5} \mathcal{R}(\mathbf{k}) \times \{\text{Transfer Function}(k)\} \times \{\text{Growth factor}(a)\}$$

matter density

$$\delta_{\text{m}}(\mathbf{k}, a) = \frac{2}{5} \frac{k^2}{\Omega_{\text{m}} H_0^2} \mathcal{R} \mathbf{k} T(k) D_+(a) \quad (a > a_{\text{late}}, k \ll aH)$$

Linear matter power spectrum at late times:

Important

$$P_{\text{L}}(k, a) = \frac{8\pi^2}{25} \frac{\mathcal{A}_s}{\Omega_{\text{m}}^2} D_+^2(a) T^2(k) \frac{k^{n_s}}{H_0^4 k_{\text{p}}^{n_s-1}} \quad (8.1) \quad \left| \begin{array}{l} \mathcal{A}_s: \\ D_+: \end{array} \right.$$

Closing the Boltzmann hierarchy

SUBSECTION 8.2

Large-Scale

SECTION 9

The cosmic microwave background

SUBSECTION 9.1

Overview

$$\Theta_0'' + k^2 c_s^2 \Theta_0 = F$$

SUBSECTION 9.2

Large-scale anisotropies

SUBSECTION 9.3

Acoustic oscillations

9.3.1 Tightly-coupled limit of the Boltzmann equations

$$\begin{aligned} \Theta_0'' + \frac{a'}{a} \frac{R}{1+R} \Theta_0' + k^2 c_s^2 \Theta_0 &= F(k, \eta), \\ F(k, \eta) &\equiv -\frac{k^2}{3} \Psi - \frac{a'}{a} \frac{R}{1+R} \Phi' - \Phi'' \end{aligned}$$

9.3.2 Tightly-coupled solutions

$$\begin{aligned} \Theta_0(\mathbf{k}, \eta) + \Phi(\mathbf{k}, \eta) &= [\Theta_0(\mathbf{k}, 0) + \Phi(\mathbf{k}, 0)] \cos(kr_s) \\ &\quad + \frac{k}{\sqrt{3}} \int_0^\eta d\tilde{\eta} [\Phi(\mathbf{k}, \tilde{\eta}) - \Psi(\mathbf{k}, \tilde{\eta})] \sin[k(r_s(\eta) - r_s(\tilde{\eta}))] \\ \Theta_1(\mathbf{k}, \eta) &= \frac{1}{\sqrt{3}} [\Theta_0(\mathbf{k}, 0) + \Phi(\mathbf{k}, 0)] \sin(kr_s) \\ &\quad - \frac{k}{3} \int_0^\eta d\tilde{\eta} [\Phi(\mathbf{k}, \tilde{\eta}) - \Psi(\mathbf{k}, \tilde{\eta})] \cos[k(r_s(\eta) - r_s(\tilde{\eta}))] \end{aligned}$$

SUBSECTION 9.4

Diffusion damping

SUBSECTION 9.5

Inhomogeneities to anisotropies

9.5.1 Free streaming

9.5.2 The angular power spectrum

Temperature perturbation

$$T(\mathbf{x}, \hat{\mathbf{p}}, \eta) = T(\eta)[1 + \Theta(\mathbf{x}, \hat{\mathbf{p}}, \eta)]$$

Expand the temperature perturbation in terms of spherical harmonics,

$$\Theta(\mathbf{x}, \hat{\mathbf{p}}, \eta) = \sum_{l=1}^{\infty} \sum_{m=-l}^l a_{lm}(\mathbf{x}, \eta) Y_{lm}(\hat{\mathbf{p}})$$

l, m are conjugate to the real-space unit vector $\hat{\mathbf{p}}$ ¹

¹ as $\mathbf{k} \leftrightarrow \mathbf{x}$

SUBSECTION 9.6

The CMB power spectrum

SUBSECTION 9.7

Cosmological parameters

- Curvature parameter, $\Omega_K \equiv 1 - \Omega_m - \Omega_\Lambda$, often set to zero in the concordance model
- Cosmological constant, parametrized by Ω_Λ
- Normalization of the primordial spectrum, \mathcal{A}_s
- Scalar spectral index, n_s
- Reionization, parametrized by the optical depth τ_{rei} to a redshift after recombination is completed
- Baryon density, $\Omega_b h^2$
- CDM density, $\Omega_c h^2$

SECTION 10

The polarized CMB

SUBSECTION 10.1

Acoustic oscillations

10.1.1 Tightly-coupled limit of the Boltzmann eq

10.1.2 Tightly-coupled solutions

Definition 16 | Sound horizon:

$$r_s(\eta) \equiv \int_0^\eta d\tilde{\eta} c_s(\tilde{\eta})$$

SUBSECTION 10.2

Diffusion damping

SUBSECTION 10.3

Inhomogeneities to anisotropies

SECTION 11

Probes of structure: tracer

SUBSECTION 11.1

Galaxy clustering

11.1.1 Galaxy clustering

$$\chi_{\text{fid}}(z) = \chi(z) + \delta\chi(z)$$

11.1.2 Redshift-space distortions

Redshift-space oversensitivity

$$\delta_{\text{g,RSD}}(\mathbf{k}, t) = [b_1 + f\mu_k^2] \delta_{\text{m}} \mathbf{k}$$

Galaxy power spectrum in redshift space

$$R_{\text{g,RSD}}(k, \mu_k, \bar{z}) = P_{\text{L}}(k, \bar{z}) [b_1 + f\mu_k^2]^2 + P_{\text{N}} \quad (11.1)$$

11.1.3 BAO and Alcock-Paczyński

Observed galaxy power spectrum

$$R_{\text{g,obs}}(\mathbf{k}_{\text{obs}}, \bar{z}) = \left(P_{\text{L}}(k, \bar{z}) [b_1 + f\mu_k^2]^2 \right) \Big|_{\mathbf{k} = ([1+\alpha_{\perp}]k_{\text{obs}}^1, [1+\alpha_{\perp}]k_{\text{obs}}^2, [1+\alpha_{\parallel}]k_{\text{obs}}^3)} + P_{\text{N}} \quad (11.2)$$

SECTION 12

Growth of structure: beyond linear theory

SECTION 13

Probes of structure: lensing

SECTION 14

Analysis and inference
