Auto-correlations of spins in East Model

Gang Huang

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East Model is an example of kinetically constrained Ising model.

Our model consists kinetically constrained Ising spins, $n_i=0$ or 1 on a 1-D lattice, with PBC.

No static interactions between spins. (Eisinger, Jaeckle 1993; Faggionato, Martinelli, Roberto, and Toninelli 2012)

A given spin may only flip if the neighboring spin to the right is up, with an acceptance ratio A,

$$A = \begin{cases} e^{-\beta}, \text{ for } 01 \to 11\\ 1, \text{ for } 11 \to 01. \end{cases}$$
 (1)

East model shows glassy dynamics. (Pitts and Andersen, 2001) Define auto correlation function

$$C(t) = \frac{\langle \delta n_i(t) \delta n_i(0) \rangle}{\langle \delta n_i(0)^2 \rangle}.$$
 (2)

 $\langle \cdots \rangle$: Average for the equilibrium state, n_i : occupation number of site i.

Relation between β and c

Assumption (A simplifying case):

- The spins are statistically independent at equilibrium;
- All sites have the same distribution function for the occupation number, ie, $\forall i, t = \infty, P(n_i = 0) = 1 e^{-\beta}, P(n_i = 1) = e^{-\beta}.$

Relation between β and c

Denote $c = \langle n_i \rangle_{eq}$ (Eisinger and Jaeckle, 1991),

$$c = \frac{e^{-\beta}}{1 + e^{-\beta}} \Rightarrow \beta = \ln \frac{c}{1 - c} \text{ for } c < 1/2.$$
 (3)

(Wu Jianlan, 2004, JPC)

$$c = 0.10 \rightleftharpoons \beta = 2.20$$

 $c = 0.20 \rightleftharpoons \beta = 1.39$
 $c = 0.30 \rightleftharpoons \beta = 0.84$
 $c = 0.48 \rightleftharpoons \beta = 0.05$ (4)

C(t) for 1-dimensional East model

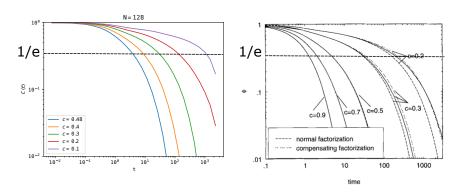


Figure: C(t) from MC simulations and from Eisinger and Jaeckle (1993).

Exact C(t) for 1-dimensional East model

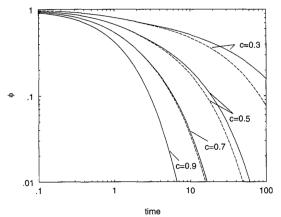


Figure: Exact C(t) (solid lines) from numerical results for finite chains (**Eisinger and Jaeckle**, 1991).

β -dependence of C(t)

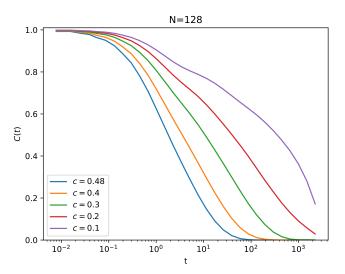


Figure: t-dependence and β -dependence of C(t)



β -dependence of relaxation time τ

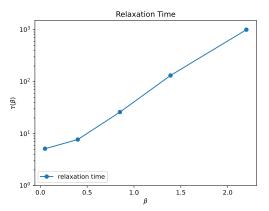


Figure: β -dependence of C(t). N = 128.

c-dependence of relaxation time au

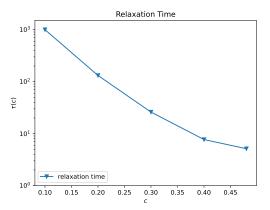


Figure: *c*-dependence of C(t). N = 128.

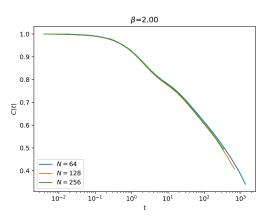


Figure: c-dependence of C(t).

For C(t):

τ(β) ~

•
$$\phi(t) = \frac{\langle \delta n_i(t) \delta n_i(0) \rangle}{c(1-c)}$$
 (Eisinger and Jaeckle, 1991, where $\sigma_i = \pm 1$).

•
$$C(t) = \frac{\langle \delta n_i(t) \delta n_i(0) \rangle}{\langle \delta n_i(0)^2 \rangle}$$
. What is the difference between $C(t)$ and $\phi(t)$? $n_i = \sigma_i = 0$ or 1.

Theoretical object: to derive an equation of C(t) (kinetic theory of C(t))

In general,

$$\frac{dC(t)}{dt} + \omega C(t) + \omega^{-1} \int_0^t d\tau M^{irr}(t - \tau) \frac{dC(\tau)}{d\tau} = 0,$$
 (5)

where $\omega=k(0),\ k(t)=-\frac{dC(t)}{dt},\ M^{\rm irr}$ is irreducible memory function. In MCT, assume that

$$M^{irr}(t) = \sum_{n} a_n(c) [C(t)]^n. \tag{6}$$

Some approximations:

- $M_{\rm K}^{\rm irr}(t) = c(1-c)[C(t)]^2$ (Kawasaki1995)
- $M_{\rm FI}^{\rm irr}(t) = c(1-c)C(t)$ (Eisinger and Jaeckle 1993)
- diagrammatic representations of the irreducible memory function (Pitts and Andersen, 2001).

More about C(t): Detailed balance

Let the state
$$\alpha=(s_1,s_2,\cdots,s_N)$$
. In equilibrium, $\frac{\partial P_{\alpha}(t)}{\partial t}=0 \Rightarrow P_{\alpha}(t)W_{\alpha \to \beta}=P_{\beta}(t)W_{\beta \to \alpha}$, where

- $P_{\alpha}(t)$ is the probability of the chain being in state α .
- From detailed balance, $W_{a \to b} = e^{-\beta}$; $W_{b \to a} = 1 \Rightarrow P_b(t)/P_a(t) = e^{-\beta}$ (see Figure).
 - $\boxed{0}$ $\boxed{1}$ $\boxed{\cdot\cdot}$ $\boxed{}$ $\boxed{}$ state a
 - ☐ ☐ ☐ ☐ ☐ ☐ ☐ ☐ state b

Figure: Two states of the 1-D finite East model with N sites.

More about C(t): Detailed balance condition (In general)

We now restrict attention to systems for which there is a stationary distribution function $\rho_0(\Gamma)$ s.t. the transition probabilities obey the DBC (Pitts and Andersen, 2001):

$$W(\Gamma', \Gamma)\rho_0(\Gamma) = W(\Gamma, \Gamma')\rho_0(\Gamma'). \tag{7}$$