# FY2045 Quantum Mechanics I

Fall 2023

Henning Goa Hugdal Week 3 General formulation of QM

# Dirac's $\langle bra|ket \rangle$ notation

#### State vector

A quantum mechanical state of a system is described by a state vector

$$|\psi\rangle$$

in a complex, linear vector space  $\mathcal{H}-$  Hilbert space.

### Dual vector

For each vector  $|a\rangle$  we define the dual vector  $\langle a|$  in the dual space  $\mathcal{H}^*$ ,

$$|a\rangle \stackrel{\text{dual}}{\longleftrightarrow} \langle a|,$$

so that we can define the scalar (inner) product of vectors  $|a\rangle$  and  $|b\rangle$ 

$$\langle a||b\rangle \equiv \langle a|b\rangle \in \mathbb{C},$$

with  $\langle a|b\rangle = \langle b|a\rangle^*$ .

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# Interpretation

### Example

Probability amplitude for particle arriving at point *x*:

⟨Particle arrives at x|particle leaves s⟩

or simply

$$\langle x|s\rangle$$
.

Can go through either slit 1 or 2:

$$\langle x|s\rangle = \langle x|1\rangle\langle 1|s\rangle + \langle x|2\rangle\langle 2|s\rangle.$$

Point x

2

Screen

Based on Ch. 3 of Vol. III in the Feynman Lectures.

## Interpretation

The wavefunction is probability amplitude of finding state  $|\psi\rangle$  at point x:

$$\psi(\mathsf{X}) = \langle \mathsf{X} | \psi \rangle.$$

The momentum wavefunction is probability amplitude of finding state  $|\psi\rangle$  with momentum p:

$$\phi(p) = \langle p | \psi \rangle.$$

The probability amplitude of finding state  $|\psi\rangle$  with energy  $E_n$ :

$$\langle \psi_n | \psi \rangle$$
.

## Completeness

*n* linearly independent vectors  $|1\rangle, |2\rangle, |3\rangle, ...$  span  $\mathcal{H}$  if  $\forall |\psi\rangle \in \mathcal{H}$  we have

$$|\psi\rangle = \sum_{k=1}^{n} c_n |k\rangle.$$

Assuming orthonormality  $\langle m|k\rangle=\delta_{mk}$ ,

$$\Rightarrow \ |\psi\rangle = \sum_{k} \langle k|\psi\rangle |k\rangle = \sum_{k} |k\rangle \langle k||\psi\rangle,$$

meaning we have the completeness relation

$$\sum_{k} |k\rangle\langle k| = 1.$$

## Operators

An operator  $\hat{A}$  applied to a vector  $|a\rangle \in \mathcal{H}$  results in a new vector  $|c\rangle \in \mathcal{H}$ ,

$$\hat{A}|a\rangle = |c\rangle.$$

Adjoint or Hermitian conjugate of operator:

$$\langle a|\hat{A}^{\dagger}|b\rangle = \langle b|\hat{A}|a\rangle^* \quad \forall |a\rangle, |b\rangle \in \mathcal{H}.$$

meaning that we have the dual vector

$$\hat{A}|a\rangle \stackrel{\text{dual}}{\longleftrightarrow} \langle a|\hat{A}^{\dagger}.$$

## **Properties**

$$(\hat{A}^{\dagger})^{\dagger} = \hat{A},$$
  

$$(\alpha \hat{A})^{\dagger} = \alpha^* \hat{A}^{\dagger},$$
  

$$(\hat{A}\hat{B})^{\dagger} = \hat{B}^{\dagger} \hat{A}^{\dagger}.$$

For a Hermitian (self-adjoint) matrix

$$\hat{A}^{\dagger}=\hat{A}.$$

# Eigenvectors and eigenvalues

The eigenvectors  $|\alpha\rangle$  and eigenvalues  $\lambda_{\alpha}$  of an operator  $\hat{\bf A}$  are defined by

$$\hat{A}|\alpha\rangle = \lambda_{\alpha}|\alpha\rangle.$$

The set of eigenvectors  $\{|\alpha\rangle\}$  corresponding to physical quantities is assumed to be complete — they form a basis set that span  $\mathcal{H}$ .

The eigenvectors of a Hermitian operator are real.

### **Examples**

Energy:

$$\hat{H}|n\rangle = E_n|n\rangle.$$

Position:

$$\hat{x}|x'\rangle=x'|x'\rangle.$$

Momentum:

$$\hat{p}|p'\rangle = p'|p'\rangle.$$

# Postulates in general formulation

#### Postulate A

Each observable quantity F corresponds to a linear, Hermitian operator  $\hat{F}$  in Hilbert space. The operators for a generalized coordinate  $q_n$  and generalized momentum  $p_n$  fulfill

$$[\hat{q}_n,\hat{p}_n] = \hat{q}_n\hat{p}_n - \hat{p}_n\hat{q}_n = i\hbar.$$

### Postulate C

The expectation value of an observable F, given the state  $|\Psi\rangle$ , is

$$\langle F \rangle = \langle \Psi | \hat{F} | \Psi \rangle.$$

#### Postulate B

Each state of a physical system is represented by a state vector  $|\Psi,t\rangle$  in Hilbert space with length 1,  $\langle\Psi,t|\Psi t\rangle=$  1. The vector fulfills the time-dependent SE

$$i\hbar \frac{\partial}{\partial t} |\Psi, t\rangle = \hat{H} |\Psi, t\rangle.$$

### Postulate D

The measurement of an observable F yields one of the eigenvalues  $f_n$  of the corresponding operator  $\hat{F}$ .

# Position representation

Eigenvectors  $|x'\rangle$  of operator  $\hat{x}$ :

$$\hat{x}|x'\rangle = x'|x'\rangle,$$

with eigenvalue x' taking continuous values.  $|x'\rangle$  is  $\delta$ -function normalized

$$\langle X''|X'\rangle = \delta(X''-X'),$$

with completeness relation

$$\int dx'|x'\rangle\langle x'|=\mathbb{1}.$$

## Position space wavefunction

Given a state vector  $|\psi\rangle$  , the position space wavefunction is given by

$$\psi(\mathsf{X}') = \langle \mathsf{X}' | \psi \rangle,$$

the projection of  $|\psi\rangle$  on the position basis vector  $|\mathbf{x}'\rangle$ .

### **Operators**

For an operator  $\hat{F}$ , which is a function of  $\hat{x}$  and  $\hat{p}$ , we have

$$\langle x''|F(\hat{\rho},\hat{x})|x'\rangle = F\left(\frac{\hbar}{i}\frac{\partial}{\partial x''},x''\right)\delta(x''-x').$$

### Momentum formulation

Eigenvectors  $|p\rangle$  of operator  $\hat{p}$ :

$$\hat{p}|p\rangle = p|p\rangle,$$

with eigenvalue p taking continuous values.  $|p\rangle$  is  $\delta$ -function normalized

$$\langle p'|p\rangle = \delta(p'-p),$$

with completeness relation

$$\int dp|p\rangle\langle p|=1.$$

## Momentum space wavefunction

Given a state vector  $|\psi\rangle$ , the momentum space wavefunction is given by

$$\phi(p) = \langle p | \psi \rangle,$$

the projection of  $|\psi\rangle$  on the momentum basis vector  $|p\rangle$ .

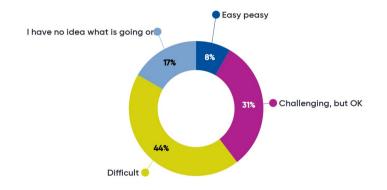
# Relation between $|p\rangle$ and $|x\rangle$

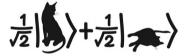
Projecting  $|p\rangle$  on the position basis  $|x\rangle$ , we get the momentum eigenfunctions in the position formulation,

$$\langle x|p\rangle = \frac{1}{\sqrt{2\pi\hbar}}e^{ipx/\hbar}.$$



### How do you find the general formulation?







### What concepts do you find challenging?

48 Responses

