

Q: Poisson's equation: $-\Delta u = f$

fundamental solution: $-\Delta u = \delta(x) \cdot x \in \mathbb{R}^d$

green function: $\begin{cases} -\Delta u = \delta(y-x) & x \in U, y \in U \subset \mathbb{R}^d, \text{两者区别在于 } \delta(y-x), \delta(x); \text{边值条件} \\ u|_{\partial U} = 0 \end{cases}$

• ~~fundamental solution~~ 是一个解, green function 是求解的中间变量 $\begin{cases} -\Delta u = 0 \\ u|_{\partial U} = g \end{cases}$

G satisfy green function: $u(y) = \int_U \delta(y-x) \cdot u(x) dx$

$$\begin{aligned} &= \int_U (-\Delta G) u(x) dx \\ &= \int_U (-\Delta u) G dx - \int_{\partial U} \frac{\partial G}{\partial n} u ds + \int_{\partial U} \frac{\partial u}{\partial n} G ds \quad \dots \text{①} \\ &= \int_{\partial U} \left(-\frac{\partial G}{\partial n} \right) g ds + \int_U f G dx. \quad (\text{和 HFB 中 Duhamel 类似}) \end{aligned}$$

$\rightarrow \delta G = \delta(y-x)$, 把 y 当成 const, x 是 G 的变量, g 的变量

At fixed y : $u(y) = \int_{\partial U} \left(-\frac{\partial G}{\partial n} \right) g ds$, 用变量 x, s 过渡到 const y 的直

$$ds = d((\nabla_x u)^2 + x^2)^{\frac{1}{2}}$$

u 单变量 $u(x)$, G 也是: $u(y) = \int_{\partial U} \left(-\frac{\partial G}{\partial n} \right) g(x) \sqrt{(\nabla_x u)^2 + x^2} dx$

即 x 只是 Green function 里的中间变量

• ~~对称~~ satisfy fundamental solution: $-\Delta \bar{\Phi}(x) = \delta(x)$

$$\bar{\Phi}(x) = \frac{1}{d(d-2)} \cdot \frac{1}{|x|^{d-2}}, \quad d \geq 3, \quad \omega_d = \frac{\pi^{d/2}}{\Gamma(\frac{d}{2} + 1)}$$



和 G 都可以求出 $-\Delta u = f$ 的通解,

$$u(x) = \int \bar{\Phi}(x-y) f(y) dy \quad \text{易知}$$

And:

$$\text{①: } \int_{B_r} (\Delta \bar{\Phi}) \cdot 1 dx = \int_{B_r} \bar{\Phi}(0) dx + \int_{\partial B_r} \frac{\partial \bar{\Phi}}{\partial n} \cdot 1 - \int_{\partial B_r} \frac{\partial \bar{\Phi}}{\partial n} \cdot \bar{\Phi} ds = \int_{\partial B_r} \frac{\partial \bar{\Phi}}{\partial n} ds$$

$$\therefore \int_{B_r} (\Delta \bar{\Phi}) \cdot 1 dx = \int_{\partial B_r} \frac{\partial \bar{\Phi}}{\partial n} ds = -1$$

$$\text{②: } \int_{B_r} \bar{\Phi} dx \rightarrow 0, \quad \int_{\partial B_r} \bar{\Phi} dx \rightarrow 0 \text{ as } r \rightarrow 0^+$$

• ~~对称~~ 和 fundamental $\bar{\Phi}$, (Symmetry)

L7 第3-种设置

want: $\bar{\Phi}(x) = \bar{\Phi}(|x|)$ s.t. $\Delta \bar{\Phi} = 0$; 在 ω_d 中证过 $\bar{\Phi}(x) = 0 \Rightarrow \bar{\Phi}(0) = 0$; "0" is rotation

$$\bar{\Phi}(u(x)) = u(r) = u(|x|)$$

• 对称求 fundamental (Symmetry)

还有一种

$\nabla \Phi \cdot \nabla \Phi = 0 \Rightarrow \Phi(Ax) = 0, A \text{ is rotation.}$

$$\nabla^2 U_r(x) = U(r) = U(|x|); \Delta U_r = 0$$

$$\begin{cases} \partial x_i \cdot U_r(x) = U(r) - \frac{\partial r}{\partial x_i}, \text{ 且 } \partial x_i \cdot r = \partial x_i \cdot \sqrt{x_1^2 + x_2^2 + \dots + x_d^2} = \frac{\partial x_i}{\sqrt{2x_i^2}} = \frac{x_i}{r} \\ \partial x_i^2 \cdot U_r(x) = (U''(r) \cdot \frac{\partial r}{\partial x_i}) \cdot \frac{\partial r}{\partial x_i} + U'(r) \cdot \partial x_i \cdot \frac{x_i}{r} \\ = U''(r) \cdot \left(\frac{x_i}{r}\right)^2 + U'(r) \cdot \left(\frac{1}{r} - \frac{x_i}{r^2} \cdot \frac{x_i}{r}\right) \end{cases} \quad \begin{array}{l} \text{注意: } \partial x_i \left(\frac{\partial r}{\partial x_i}\right) = \partial x_i \left(\frac{x_i}{r}\right) \text{ 与 } r \text{ 无关!} \\ * \frac{1}{r} \end{array}$$

$$\text{Summing over } i: \begin{cases} \partial x_i U_r(x) = U'(r) \cdot \partial x_i \cdot r \\ \Delta U_r \partial x_i U_r(x) = U''(r) \frac{\sum x_i^2}{r^2} + U'(r) \left(\frac{d}{r} - \frac{\sum x_i^2}{r^3}\right) + \sum_{i \neq j} \left(U''(r) \frac{x_i x_j}{r^2} - U'(r) \frac{x_i x_j}{r^3}\right) \times 2 \end{cases}$$

$$\text{Since: } \partial x_i U_r(x) = \sum_{j=1}^d \partial x_i U_r(x).$$

$$\partial x^2 U_r(x) = \left(\sum_{i=1}^d \partial x_i U_r(x)\right)' = \sum_{i=1}^d \partial x_i \left(\sum_{j=1}^d \partial x_j U_r(x)\right)$$

$$\text{如果 } i \neq j, \partial x_i \partial x_j U_r(x) = \partial x_i \left(U'(r) \cdot \frac{x_i}{r}\right)$$

$$= U''(r) \cdot \frac{\partial r}{\partial x_i} \cdot \frac{x_i}{r} + U'(r) \left(-\frac{x_i}{r^2}\right) \cdot \frac{\partial r}{\partial x_j}$$

$$= U''(r) \frac{x_i x_j}{r^2} - U'(r) \frac{x_i x_j}{r^3}$$

$$\Rightarrow \Delta U = U''(r) + \frac{d-1}{r} U'(r) = 0$$

$$\text{解 DDE: } (r U''(r) + U'(r)) + (d-1) U'(r) = 0$$

$$\therefore r U' + (d-2) U = \text{const} \quad \text{if } d \geq 2$$

$$\text{then: } U(r) = a \log r + b \quad d \geq 2$$

$$\text{例如: } \int_{0B_1} \frac{\partial U}{\partial n} ds = 1 \Rightarrow a \int_{0B_1} \frac{1}{r^n} dr = 1 \quad ? \quad \therefore c = -\frac{1}{2n}$$

$$\Phi(x) = \begin{cases} -\frac{1}{2} |x| & d=1 \\ -\frac{1}{2n} \log |x| & d=2 \\ \frac{1}{d(d-2)} \frac{1}{2d} |x|^{2-d} & d \geq 3 \end{cases}$$

$d \geq 3$ 的物理意义为在原点的单位正电荷引起的在全空间 R^3 上的电势分布

• $\Phi(x)$ 的物理意义

$$\bar{\Phi}(x) = \begin{cases} -\frac{1}{2\pi} |x| & d=1 \\ -\frac{1}{2\pi} \log|x| & d=2 \\ \frac{1}{d(d-2)\partial_d} \cdot |x|^{2-d} & d \geq 3 \end{cases}$$

$d \geq 3, \bar{\Phi}_d = \frac{\pi^{d/2}}{\Gamma(\frac{d}{2}+1)}$ 表示 unit-ball 的体积.

① $d=3, \bar{\Phi}(x) = \frac{1}{4\pi} \cdot \frac{1}{r}$

$\vec{E} = \frac{q}{4\pi\epsilon_0 r^2} \cdot \frac{\vec{x}}{|\vec{x}|}$

$$\bar{\Phi}(r) = -\frac{q}{4\pi\epsilon_0} \cdot \frac{1}{r}$$

② $d=2$ Coulomb's Law should be $\vec{E} = \frac{q}{2\pi\epsilon_0 r} \cdot \frac{\vec{r}}{|\vec{r}|}$. Since fundamental solution is now $\frac{-1}{2\pi} \log|x|$

基本解在广义函数下满足 $-\Delta \bar{\Phi}(x) = \delta(x)$

here: $\bar{\Phi}(x) = \bar{\Phi}(|x|)$, $\Delta \bar{\Phi}(x) = 0 \forall x \neq 0$ 位势方程的特殊基本条件

不含点源(即)的任意体积分 $\int_V -\Delta \bar{\Phi}(x) dx = -1 = \iiint_V \nabla \cdot \nabla \bar{\Phi} dV$

$$= \iint_S \bar{\Phi}_r ds = 4\pi x^2 \bar{\Phi}_r(x)$$

$$\therefore \bar{\Phi}_r(x) = \frac{1}{4\pi x^2}, \bar{\Phi}(x) = \frac{1}{4\pi x}$$

• 不对称情况 G

$$\begin{cases} -\Delta G(x,y) = \delta(x-y) & x \in U \\ G(\cdot, y)|_{\partial U} = 0 \end{cases}$$

case (1): $U = \mathbb{R}^d - \{(x_1, x_2, \dots, x_d) : x_i > 0\}$ half-space

$G(x,y) = \bar{\Phi}(|x-y|) - \bar{\Phi}(|x-\bar{y}|)$, $\bar{y} = (-y_1, y_2, \dots, y_d)$

且不跨平面 $\frac{+}{-}$, 和物理直觉法

check: ① $-\Delta G(x,y) = -\Delta_x \bar{\Phi}(|x-y|) + \Delta_y \bar{\Phi}(|x-y|)$

$$= -\Delta \bar{\Phi}(|x-y|)$$

$$= \delta(x-y)$$

② $G(x_0, y) = \bar{\Phi}(|x_0-y|) - \bar{\Phi}(|x_0-\bar{y}|) = 0$ where $x_0 = (0, x_2, \dots, x_d)$ 且 $G|_{\partial U} = 0$

问题: $G(x,y) = G(x-y) \checkmark$

$$\begin{cases} -\Delta \bar{\Phi}(x, \bar{y}) = 0 \text{ 否则} \\ \text{因为 } \Delta \bar{\Phi}(x, z) = \delta(x, z) = \begin{cases} 0 & x \neq z \\ \infty & x = z \end{cases} \end{cases}$$

$\therefore |z| \neq 0$ 且 $x_1 \neq 0$ 否则 $x = (x_1, 0, x_3, \dots)$ 是公共边

after symmetry reflection

$\bar{y} \neq x$ forever, since $\bar{y} \notin \text{Domain}$

$$\therefore \Delta \bar{\Phi}(x, \bar{y}) = 0$$

$$\text{Solve } \begin{cases} -\Delta u = 0, & V \\ u = g, & \partial V \end{cases}$$

$$\star: \text{已知 } u(y) = \int_V \delta(x-y) \cdot u(x) dx = \int_V (-\Delta G) \cdot u(x) dx \\ = \int_V (\Delta u) \cdot G dx - \int_{\partial V} \frac{\partial G}{\partial n} u ds + \int_{\partial V} u \frac{\partial G}{\partial n} G ds \\ = - \int_{\partial V} \left(\frac{\partial G}{\partial n} \right) g ds$$

这段恒成立的
case II 和 III 都是

$$\text{here } \int_{\partial V} \left(\frac{\partial G}{\partial n} \right) g ds = \int_{\partial V} \left(-\frac{\partial G}{\partial n} \right) (z_1 \times) g(z) dz \stackrel{G(z)}{=} , z = (z_1, \bar{z}) \\ = \int_{-\infty}^{+\infty} z \partial_{z_1} \Psi(z; x) \Big|_{z_1=0} g(\bar{z}) d\bar{z}$$

? 为什么从 ds 变成 $d\bar{z}$? "哪来, 而怎么找
为什么 ∂V 是 $z_1=0$. 前面也有过这道题"

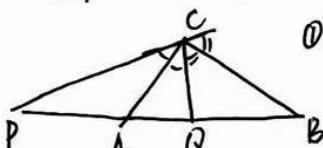
$$\text{Case I: } d=2 \quad \Phi(z; x) = -\frac{1}{4\pi} \log((z_1-x_1)^2 + (z_2-x_2)^2)$$

$$\partial_z \Psi(z; x) \Big|_{z_1=0} = -\frac{1}{2\pi} \cdot \frac{z_1-x_1}{(z_1-x_1)^2 + (z_2-x_2)^2} \Big|_{z_1=0} = \frac{1}{2\pi} \cdot \frac{x_1}{x_1^2 + (z_2-x_2)^2}$$

$$\therefore u(x_1, x_2) = \frac{1}{2} \int_{-\infty}^{+\infty} \frac{x_1}{x_1^2 + (y-x_2)^2} g(y) dy \\ \Rightarrow \text{Poisson Kernel}$$

Case II: $V = B_R(0)$

若点 C 在 circle of Apollonius



① $\frac{CP}{CA} = \frac{AP}{AB} = \frac{PB}{QB}$, $\forall C$, C 所成的圆以 $\frac{AB}{2}$ 为半径, mid(A, B) 为圆心

② $\left| \frac{x-y}{x-q} \right| = \text{const.}, x$ 也形成圆



$$\text{圆的性质} (\quad = \left| \frac{y_0-y}{y_0-q} \right| \\ = \frac{|y_0-y|}{R/|y_1-y|} = \frac{|y_1|}{R} \quad \text{即} \text{②} \text{很清楚}$$

$$\Rightarrow \text{let } G_1(x, y) = \Phi(x-y) - \Phi\left(\frac{|y_1|}{R}(x-\bar{y})\right) \quad \text{since } |x-y| = |x-\bar{y}| \cdot \frac{|y_1|}{R}$$

$$\star \left\{ \frac{\partial \Phi}{\partial x_i}(x-y) = -\frac{1}{d^2} \frac{x_i-y_i}{|x-y|^d}$$

$$\left\{ \frac{\partial \Phi}{\partial x_i}\left(\frac{|y_1|}{R}(x-\bar{y})\right) = -\frac{1}{d^2} \frac{\frac{|y_1|^2}{R^2}(x_i-\bar{y}_i)}{\left|\frac{|y_1|}{R}(x-\bar{y})\right|^d}, \quad \bar{y} = y + \frac{R^2}{|y_1|^2} \cdot \bar{y}_i = y_i - \frac{R^2}{|y_1|^2} y_i$$

$$\text{summing over } i: \frac{\partial G}{\partial n}(x, y) = \sum_{i=1}^d \frac{x_i}{R} \left[\frac{\partial \Phi}{\partial x_i}(x-y) - \frac{\partial \Phi}{\partial x_i}\left(\frac{|y_1|}{R}(x-\bar{y}_i)\right) \right]$$

$$\frac{\partial G}{\partial n} = \frac{1}{d^2} \cdot \frac{1}{|x-y|^d} \times \frac{R^2 - |y_1|^2}{R} \quad \vec{n} = \frac{\vec{R}}{R} \\ \frac{\partial G}{\partial n} = \frac{\partial \Phi}{\partial n}(x-y) - \frac{\partial \Phi}{\partial n}\left(\frac{|y_1|}{R}(x-\bar{y})\right), \quad \frac{\partial \Phi}{\partial n} = \sum_{i=1}^d \frac{\partial \Phi}{\partial x_i} \cdot \frac{\partial x_i}{\partial R}$$

所以应有 $\frac{\partial x_i}{\partial R} = \frac{x_i}{R}$?

擇上面：

$$\frac{\partial \Phi}{\partial n}(x, y) = \frac{\partial \Phi}{\partial n}(x-y) - \frac{\partial \Phi}{\partial n}\left(\frac{|y|}{R}(x-\bar{y})\right)$$

$$= \sum_{i=1}^d \frac{\partial \Phi}{\partial x_i} \cdot \frac{\partial x_i}{\partial n}(x-y) - \sum_{i=1}^d \frac{\partial \Phi}{\partial x_i} \cdot \frac{\partial x_i}{\partial n}\left(\frac{|y|}{R}(x-\bar{y})\right)$$

$$= \sum_{i=1}^d \frac{x_i}{R} \left[\frac{\partial \Phi}{\partial x_i}(x-y) - \frac{\partial \Phi}{\partial x_i}\left(\frac{|y|}{R}(x-\bar{y})\right) \right]$$

$$\frac{\partial x_i}{\partial n}$$

$$\text{if } R = \sqrt{x_1^2 + \dots + x_d^2}, \quad \frac{\partial P}{\partial x_i} = \frac{\partial x_i}{\partial n} = \frac{x_i}{R} \quad \text{是?}$$

\vec{n} : 方向量表示；
也不是这个表达式吗

$$d \geq 3 \quad \Phi(x) = \frac{1}{d(d-2)\cdot 2d} |x|^{2-d}$$

$$\Phi(x, y) = \frac{1}{d(d-2)\cdot 2d} \cdot \sqrt{(x_1-y_1)^2 + (x_2-y_2)^2 + \dots + (x_d-y_d)^2}$$

$$\frac{\partial \Phi}{\partial x_i}(x-y) = \frac{1}{d(d-2)\cdot 2d} - \frac{2-d}{2} \left(\sum (x_i-y_i)^2 \right)^{\frac{d-2}{2}} \cdot 2(x_i-y_i)$$

$$= -\frac{1}{d \cdot 2d} \times \frac{x_i-y_i}{|x-y|^d}$$

$$\begin{aligned} \frac{\partial \Phi}{\partial x_i}\left(\frac{|y|}{R}(x-\bar{y})\right) &= \frac{1}{d(d-2)\cdot 2d} \cdot \frac{2d}{2} \left(\frac{|y|}{R}\right)^{2d} \cdot \sum (x_i-\bar{y}_i)^2 \cdot 2(x_i-\bar{y}_i) \\ &= -\frac{1}{d \cdot 2d} \left(\frac{|y|}{R}\right)^{2d} \times \frac{x_i-y_i}{|x-y|^d} \end{aligned}$$

$$\downarrow \text{代入 } \frac{\partial \Phi}{\partial x_i} \text{ 得, } \frac{\partial \Phi}{\partial n}(x, y) = -\frac{1}{d \cdot 2d} \frac{1}{|x-y|^d} \cdot \frac{R^2 - |y|^2}{R}$$

$$\therefore u(y) = \int_{\partial B} \left(\frac{\partial \Phi}{\partial n} \right) g dS$$

$$= \int_{\partial B} \frac{R^2 - |y|^2}{R} \cdot \frac{1}{d \cdot 2d} \times \frac{1}{|x-y|^d} g(x) dS(x) \quad y \in B_R \quad \text{用 } x \text{ 表示}$$

$$\text{此处 Poisson kernel } k(x, y) = \frac{R^2 - |y|^2}{R \cdot d \cdot 2d} \cdot \frac{1}{|x-y|^d} \quad \text{for } B_R$$

$$\int k(x, y) dx = 1$$