

8. The energy estimate:  $u$  范数<sup>2</sup> +  $\nu u$  范数<sup>2</sup>  $\leq C$  (非齐次项<sup>2</sup> 的积分 + 边界项范数<sup>2</sup>)

$$\begin{cases} \Delta u = \partial_t u - \Delta u = f & \text{in } U \\ u|_{t=0} = \varphi(x) & \text{in } U = (0,1) \times (0,T) \\ u|_{x=0} = u|_{x=1} = 0 & \text{on } \partial U \end{cases}$$

设  $u$  is classical solution, then:  $\sup_{0 \leq t \leq T} \int_0^1 u^2(x,t) dx + 2\nu \int_0^T \int_0^1 u^2(x,t) dx dt \leq C(\int_0^1 \varphi^2(x) dx + \int_0^T \int_0^1 f^2(x,t) dx dt)$   
 $\Rightarrow \sup_{0 \leq t \leq T} \|u(t, \cdot)\|_{L^2(U)}^2 + 2\nu \int_0^T \|\partial u(t, \cdot)\|_{L^2(U)}^2 dt \leq C(\|\varphi\|_{L^2(U)}^2 + \int_0^T \|f(t, \cdot)\|_{L^2(U)}^2 dt)$   
 能量估计 =  $L^2(U)$  范数<sup>2</sup>,  $\int u^2 dx = \|\cdot\|_{L^2(U)}^2$

given:  $\int_0^1 u \cdot (\partial_t u - \Delta u) dx = \int_0^1 u f dx$

LHS:  $\int_0^1 u \cdot \partial_t u dx = \int_0^1 \frac{1}{2} \partial_t (u^2) dx = \frac{1}{2} \frac{d}{dt} \int_0^1 u^2(x) dx$

$\int_0^1 u \cdot (-\Delta u) dx = \int_0^1 |\partial u|^2 dx - \int_0^1 \partial u \cdot \frac{\partial u}{\partial n} \cdot u ds = \int_0^1 |\partial u|^2 dx$ , since  $\partial u \perp u|_{x=0} = u|_{x=1} = 0$

RHS:  $\int_0^1 u f dx \leq \frac{1}{2} \int_0^1 u^2 dx + \frac{1}{2} \int_0^1 f^2 dx$

代入即:  $\frac{1}{2} \frac{d}{dt} \int_0^1 u^2(x) dx + \int_0^1 |\partial u|^2 dx \leq \frac{1}{2} \int_0^1 u^2 dx + \frac{1}{2} \int_0^1 f^2 dx$

$\frac{1}{2} \frac{d}{dt} \int_0^1 u^2(x) dx = \frac{1}{2} \frac{d}{dt} \int_0^1 u^2(x) dx - \frac{1}{2} \int_0^1 u^2(x) dx = \frac{1}{2} \int_0^1 u^2(x) dx - \frac{1}{2} \|\varphi\|_{L^2(U)}^2$   
 $\Rightarrow \|u(t, \cdot)\|_{L^2(U)}^2 + 2\nu \|\partial u\|_{L^2(U)}^2 \leq \|\varphi\|_{L^2(U)}^2 + \|u^2(t, \cdot)\|_{L^2}^2 + \|f^2(t, \cdot)\|_{L^2}^2$

" $\frac{d}{dt}$ " 不是对  $t$  这个变量, 而是所代表的 "first slot" (不然可以随便写)

$\therefore \frac{1}{2} \frac{d}{dt} \int_0^1 u^2(x) dx + \int_0^1 |\partial u|^2 dx \leq \frac{1}{2} \int_0^1 u^2(x) dx + \frac{1}{2} \int_0^1 f^2(x) dx \dots ①$

$\frac{1}{2} \int_0^1 u^2(x) dx \Big|_{t=0}^{t=T} = \frac{1}{2} \int_0^1 u^2(T, x) dx - \frac{1}{2} \int_0^1 u^2(0, x) dx = \frac{1}{2} \|u(T, \cdot)\|_{L^2(U)}^2 - \frac{1}{2} \|\varphi\|_{L^2(U)}^2$

$\Rightarrow (\frac{1}{2} \|u(T, \cdot)\|_{L^2(U)}^2 - \frac{1}{2} \|\varphi\|_{L^2(U)}^2) + \int_0^T \|\partial u(t, \cdot)\|_{L^2(U)}^2 dt \leq \frac{1}{2} \int_0^T \|u(t, \cdot)\|_{L^2(U)}^2 dt + \frac{1}{2} \int_0^T \|f(t, \cdot)\|_{L^2(U)}^2 dt \dots ②$

①  $\Rightarrow$ :

here:  $\begin{cases} G(t) = \frac{1}{2} \int_0^T \|u(t, \cdot)\|_{L^2(U)}^2 dt \\ F(t) = \frac{1}{2} \int_0^T \|f(t, \cdot)\|_{L^2(U)}^2 dt - \int_0^T \|\partial u\|_{L^2(U)}^2 dt \end{cases}$

Lemma: Gronwall's inequality:

$\begin{cases} G'(t) \leq G(t) + F(t) \quad \forall t \\ G(0) = 0 \end{cases}$

$\therefore \int_0^T \|u(t, \cdot)\|_{L^2(U)}^2 dt \leq (e^T - 1) (\int_0^T \|f(t, \cdot)\|_{L^2(U)}^2 dt + \int_0^T \|\partial u\|_{L^2(U)}^2 dt)$  then  $G(t) \leq (e^T - 1) F(t) \quad \forall t$  微分中学过  $F=0$  情况

这么写问题在于没有  $\varphi$ , 我们尝试用 ② 寻找  $G, F(t)$ .

②  $\Rightarrow$  把  $T$  当成变量:  $G(T) = \frac{1}{2} \int_0^T \|u(t, \cdot)\|_{L^2(U)}^2 dt$ ,  $G(0) = 0$

$G'(T) = \frac{1}{2} \|u(T, \cdot)\|_{L^2(U)}^2$

$F(T) = \frac{1}{2} \int_0^T \|f(t, \cdot)\|_{L^2(U)}^2 dt + \frac{1}{2} \|\varphi\|_{L^2(U)}^2$  (4)  $\leq 0$  代入 ② 即可

→  $\frac{1}{2} \frac{d}{dt} \int_{\Omega} u^2(t, x) dx + \int_{\Omega} |Du|^2 dx \leq \frac{1}{2} \int_{\Omega} u^2 dx + \frac{1}{2} \int_{\Omega} f^2 dx$  ; ... HS 和 HS 的简单变形

↓  
对于 first slot 和  $t$  不是  $t$  这个变量, " $t$ " 改成 " $t_0$ " 更合适

$\frac{1}{2} \int_{\Omega} u^2(t_0, x) dx + \int_{\Omega} |Du|^2 dx \leq \frac{1}{2} \int_{\Omega} u^2 dx + \frac{1}{2} \int_{\Omega} f^2 dx$  ; actually holds for  $t_0 \in [0, T]$

$\int_0^{t_0} (\frac{1}{2} \int_{\Omega} u^2(t_0, x) dx) dt + \int_0^{t_0} \int_{\Omega} |Du|^2 dx dt \leq \frac{1}{2} \int_0^{t_0} \int_{\Omega} u^2(t, x) dx dt + \frac{1}{2} \int_0^{t_0} \int_{\Omega} f^2(t, x) dx dt$

here " $t_0$ " is the variable, and " $t$ " only a mediator; 最后开成  $u(t_0, x)$ ;  $t_0$  作用于  $t_0$

$\therefore \int_0^{t_0} (\frac{1}{2} \int_{\Omega} u^2(t_0, x) dx) dt = \frac{1}{2} \int_{\Omega} u^2(t_0, x) dx \Big|_{t=0}^{t=t_0}$

$= \frac{1}{2} \int_{\Omega} u^2(t_0, x) dx - \frac{1}{2} \int_{\Omega} u^2(0, x) dx = \frac{1}{2} \int_{\Omega} u^2(t_0, x) dx - \frac{1}{2} \|\psi(\cdot)\|_{L^2(\Omega)}^2$

$\Rightarrow (\frac{1}{2} \int_{\Omega} u^2(t_0, x) dx - \frac{1}{2} \|\psi(\cdot)\|_{L^2(\Omega)}^2) + \int_0^{t_0} \underbrace{\|Du(t, \cdot)\|_{L^2(\Omega)}^2}_{\geq 0} dt \leq \frac{1}{2} \int_0^{t_0} \int_{\Omega} u^2(t, x) dx dt + \frac{1}{2} \int_0^{t_0} \int_{\Omega} f^2(t, x) dx dt$

$\therefore \frac{1}{2} \int_{\Omega} u^2(t_0, x) dx \leq \frac{1}{2} \int_0^{t_0} \int_{\Omega} u^2(t, x) dx dt + \frac{1}{2} \|\psi(\cdot)\|_{L^2(\Omega)}^2 + \frac{1}{2} \int_0^{t_0} \int_{\Omega} f^2(t, x) dx dt$

Gronwall's inequality:  $\frac{1}{2} \int_0^{t_0} \|u^2(t_0, \cdot)\|_{L^2(\Omega)}^2 dt \leq (e^{t_0}) \frac{1}{2} (\|\psi(\cdot)\|_{L^2(\Omega)}^2 + \int_0^{t_0} \int_{\Omega} f^2(t, x) dx dt)$   $t_0 \in [0, T]$   
let  $C = \max(e^{t_0}) = e^T$   $\square$   $\checkmark$

★ 常用梯度公式: ①  $\int_{\Omega} \nabla \cdot F dx = \int_{\partial \Omega} F \cdot \vec{n} dS$

$\left\{ \begin{aligned} \text{② } \int_{\Omega} \nabla(u \cdot \nabla u) dx &= \int_{\partial \Omega} (u \cdot \nabla u) \cdot \vec{n} dS \\ &= \int_{\partial \Omega} u \cdot (\nabla u \cdot \vec{n}) dS = \int_{\partial \Omega} u \cdot \frac{\partial u}{\partial \vec{n}} dS \end{aligned} \right.$  不怎么用, 略

③  $\nabla(u \cdot \nabla u) = \nabla u \cdot \nabla u + u \cdot \nabla \cdot \nabla u = |\nabla u|^2 + u \cdot \Delta u$   
since  $\nabla \cdot \nabla u = \Delta u$  算子从右到左, 不满足结合律

逆方向也要会推!  
 $\int_{\Omega} \nabla(u \cdot \nabla u) dx = \int_{\Omega} |\nabla u|^2 dx + \int_{\Omega} u \cdot \Delta u dx$  这个常用

$\left\{ \begin{aligned} \text{④ } \int_{\Omega} \nabla \tilde{u} \cdot \nabla v dx &= \int_{\Omega} \nabla v \cdot \nabla u dx \\ &= \int_{\partial \Omega} \tilde{u} \cdot \nabla v \cdot \vec{n} dS = \int_{\partial \Omega} u \cdot \frac{\partial v}{\partial \vec{n}} dS = \int_{\partial \Omega} v \cdot \frac{\partial u}{\partial \vec{n}} dS \end{aligned} \right.$

⑤  $\int_{\Omega} \nabla(u \cdot \nabla v) dx = \int_{\Omega} \nabla u \cdot \nabla v dx + \int_{\Omega} v \cdot \Delta u dx$

其中  $\Delta^2 u = \Delta u$

$\therefore \int_{\Omega} \nabla u \cdot \nabla v dx = \int_{\Omega} \nabla(v \cdot \nabla u) dx - \int_{\Omega} v \cdot \Delta u dx = \int_{\Omega} \nabla(u \cdot \nabla v) dx - \int_{\Omega} u \cdot \Delta v dx$  (以上主要集中在 P3 eigen "不为 0" 项!)



# Th2 Backward HE 问题

$$u_1, u_2 \in C^{1,2}([0, T] \times U) \cap C(\bar{U}) ; u_1, u_2 \text{ solve } \begin{cases} \partial_t u - \Delta u = 0, & \text{in } U \\ u|_{t=T} = \varphi(x), & \text{on } \partial U \quad (\text{usually, HE: } u|_{t=0} = \varphi(x)) \\ u|_{\partial U} = g(t), & t \in [0, T] \end{cases}$$

Th2: 解不稳定,  $\|u(t, \cdot)\| \rightarrow +\infty$  then  $u_1 = u_2$  in  $\bar{U}$ , 即: Backward HE (从末状态  $t=T$  出发) 解不唯一

$$\text{设 } \varphi \text{ solve: } \begin{cases} -\Delta \varphi = \lambda \varphi & U \\ \varphi|_{\partial U} = 0 \end{cases}$$

$$\text{let } u_\lambda(t, x) = e^{(T-t)\lambda} \cdot \varphi(x), \text{ then } u_\lambda \text{ solves: } \begin{cases} \partial_t u = \Delta u, & U \\ u|_{t=T} = \varphi(x), & U \dots \textcircled{1} \end{cases}$$

$$\therefore \|u_\lambda(t, x)\|_{L^2(U)} = e^{(T-t)\lambda} \|\varphi(x)\|_{L^2(U)}$$

given  $t \in (0, T]$ , fix  $t$ ,  $\|u_\lambda\|_{L^2(U)} \rightarrow +\infty$  as  $\lambda \rightarrow +\infty$

$\therefore$  Backward HE 解 unstable;  $u_1 - u_2$  可以极大 "  $\lambda \rightarrow +\infty$  " 在 eigen 章节讲过, 证明不要求

(In HE: 是找不到这样的  $u_\lambda(t, x)$  的, 因为 0 若成立  $\rightarrow$  改成 0,  $\lambda \rightarrow +\infty \|u_\lambda\| \rightarrow 0$ )

Th2: Backward HE 解也唯一

(In HE, unique 和 stable 均用到 Maximum principle; 而 MP 中  $\partial U$  不含  $T$ ,  $\therefore$  反向问题用不了)

$$\text{设: } \begin{cases} w = u_1 - u_2 \\ e(t) = \int_U w^2(t, x) dx = \|w(t, \cdot)\|_{L^2(U)}^2, \quad e(T) = 0 \text{ since } u_1|_{t=T} = u_2|_{t=T} \end{cases}$$

$$\text{then: } \begin{cases} \partial_t w = \Delta w, & U \\ w|_{t=T} = 0, & \partial U \\ w|_{\partial U} = 0, & t \in [0, T] \end{cases}$$

$$e(t) \Rightarrow \begin{cases} \dot{e}(t) = \int_U \frac{d}{dt} w^2 dx = \int_U 2w \cdot \frac{dw}{dt} dx = \int_U 2w \cdot \Delta w dx \dots \textcircled{1} \end{cases}$$

$$\ddot{e}(t) = \int_U 2w_t \cdot \Delta w dx + \int_U 2w \cdot \Delta w_t dx$$

对称公式推导:

$$e''(t) \stackrel{\textcircled{2}}{=} -2 \int_U \frac{d}{dt} | \nabla w |^2 dx$$

$$= -4 \int_U \nabla w \cdot \nabla w_t dx$$

$$= -4 \int_U \nabla(w \cdot \nabla w_t) dx + 4 \int_U w \cdot \nabla \nabla^2 w_t dx \text{ 这样不好}$$

$$= -4 \int_U \nabla w_t \cdot \nabla w dx - 4 \int_U \nabla(w_t \cdot \nabla w) dx + 4 \int_U w_t \cdot \nabla \nabla^2 w dx$$

$$= -4 \int_{\partial U} w_t \cdot \nabla w \cdot \vec{n} ds + 4 \int_U w_t \Delta w dx$$

$$= 4 \int_U | \nabla w |^2 dx$$

$$\therefore |e'(t)|^2 = 4 \left| \int_U w \Delta w dx \right|^2 \leq 4 \int_U w^2 dx \times \int_U | \nabla w |^2 dx = e(t) \cdot e''(t) \quad \forall t \in (0, T) \quad e(t) \text{ convex}$$

①

$e(T) = 0, e'(T) \leq 0 \quad \forall t \in [0, T]$  函数在段  $\log e(t)$  convex why  $e=0$ ? 