

Computational Microelectronics

Lecture 6 Diffusion

Sung-Min Hong (smhong@gist.ac.kr)

Semiconductor Device Simulation Laboratory
School of Electrical Engineering and Computer Science
Gwangju Institute of Science and Technology

Diffusion – Electric Field

Drift

- Up to now, we have considered only the diffusion term.

- The flux was written as ($\# \text{ cm}^{-2} \text{ sec}^{-1}$)

$$\mathbf{F}_C = -D\nabla C$$

- If we consider a charged dopant ion, the drift should be also considered.

$$\mathbf{F}_C = -D\nabla C + C\mathbf{v}$$

- The drift velocity of a positively charged dopant ion is given as $\mathbf{v} = \mu\mathbf{E}$. (μ is the mobility.) For a negatively charged one, $\mathbf{v} = -\mu\mathbf{E}$.

Alternative form

- Manipulation for combining two terms

- The flux is now given as

$$\mathbf{F}_C = -D\nabla C \pm \mu C \mathbf{E}$$

- Also, with the electrostatic potential, $\phi(\mathbf{r})$, it can be written as

$$\mathbf{E} = -\nabla\phi$$

- Moreover, we adopt the Einstein relation, $D = V_T\mu$. (V_T is the thermal voltage.)

- Then,

$$\mathbf{F}_C = -D \left(\nabla C \pm \frac{1}{V_T} C \nabla \phi \right)$$

- It looks good, but how can we calculate ϕ ?

Approximation for ϕ

- At equilibrium, (← How can we assume the equilibrium?)

- We can express the electron density, n , as

$$n = n_{int} \exp \frac{\phi}{V_T}$$

- In other words, $\phi = V_T \log \frac{n}{n_{int}}$. Then,

$$\nabla \phi = V_T \nabla \left(\log \frac{n}{n_{int}} \right)$$

- By using the above expression, we have

$$\mathbf{F}_C = -D \left[\nabla C \pm C \nabla \left(\log \frac{n}{n_{int}} \right) \right]$$

Approximation for ϕ

- Finally,

$$\mathbf{F}_C = -DC \left[\nabla \log C \pm \nabla \left(\log \frac{n}{n_{int}} \right) \right]$$

- Well, instead of ϕ , now we have n .
- In a rigorous sense, we must calculate n . (However, it costs.)
- Under the charge neutrality at equilibrium,

$$\begin{aligned} N_D^+ + p &= N_A^- + n \\ np &= n_{int}^2 \end{aligned}$$

- Then, the electron density is obtained as

$$n = \frac{(N_D^+ - N_A^-) + \sqrt{(N_D^+ - N_A^-)^2 + 4n_{int}^2}}{2}$$

Discretization

- In 1D,

$$F_C = -DC \left[\frac{\partial}{\partial x} \log C \pm \frac{\partial}{\partial x} \left(\log \frac{n}{n_{int}} \right) \right]$$

- How can we discretize the flux?
- We adopt the “logarithmic flux” approximation.

$$F_{C,i+0.5} = -D \sqrt{C(x_i)C(x_{i+1})} \frac{1}{x_{i+1} - x_i} \left[\begin{array}{l} \log C(x_{i+1}) \pm \log \frac{n(x_{i+1})}{n_{int}} \\ -\log C(x_i) \mp \log \frac{n(x_i)}{n_{int}} \end{array} \right]$$

- It is nonlinear.

HW#6

- Due: AM08:00, September 18
- Problem#1

– The intrinsic carrier concentration in silicon is given by

$$n_{int} = 3.1 \times 10^{16} T^{1.5} \exp\left(-\frac{0.603 \text{ eV}}{k_B T}\right) \text{ cm}^{-3}$$

Draw the electron density, n , as a function of temperature, T . The temperature varies from 300 K to 1200 K.

Assume the charge neutrality at equilibrium. Consider three values of $N_D^+ - N_A^-$: 10^{20} cm^{-3} , 10^{18} cm^{-3} , and 10^{16} cm^{-3}

Nonlinearity

Discretized diffusion equation

- It is much more difficult.

- Although it can be generally written as

$$\frac{C(x_i, t_k) - C(x_i, t_{k-1})}{\Delta t} \Delta x = F_{C,i+0.5} - F_{C,i-0.5}$$

- Now, the flux terms are nonlinear,

$$F_{C,i+0.5} = -D \sqrt{C(x_i, t_k) C(x_{i+1}, t_k)} \frac{1}{\Delta x} \left[\log C(x_{i+1}, t_k) \pm \log \frac{n(x_{i+1}, t_k)}{n_{int}} \right. \\ \left. - \log C(x_i, t_k) \mp \log \frac{n(x_i, t_k)}{n_{int}} \right]$$

- Note that even the electron density is nonlinear.

Is it $Ax = b$?

- Try to construct a matrix, A , for the flux term.
 - What is the entry for $C(x_i, t_k)$?
- Recall the previous case.
 - In our previous simple problem, the discretized form was

$$\frac{C(x_i, t_k) - C(x_i, t_{k-1})}{\Delta t} \Delta x$$
$$= -D \frac{1}{\Delta x} [C(x_{i+1}, t_k) - C(x_i, t_k)] + D \frac{1}{\Delta x} [C(x_i, t_k) - C(x_{i-1}, t_k)]$$

Nonlinearity is the key.

- At present, we cannot solve the problem, because it is nonlinear.
 - Let us learn how to solve a nonlinear problem!
- An example, calculation of n under the charge neutrality.

$$\begin{aligned}N_D^+ + p &= N_A^- + n \\ np &= n_{int}^2\end{aligned}$$

- By eliminating p , we find an equation of

$$n^2 - (N_D^+ - N_A^-)n - n_{int}^2 = 0$$

- Of course, we know the solution. However, instead of using the known formula, just calculate n with a numerical method.

Newton method

- First, we assume an initial solution, n_0 .
 - Of course, there is no guarantee that n_0 is the solution. Therefore,
$$n_0^2 - (N_D^+ - N_A^-)n_0 - n_{int}^2 = r_0 \neq 0$$
 - However, we assert that an improved solution, n_1 , is the real solution. Our assertion can be written as
$$n_1^2 - (N_D^+ - N_A^-)n_1 - n_{int}^2 = 0$$
 - We can take the difference between two equations:
$$(n_1 + n_0)(n_1 - n_0) - (N_D^+ - N_A^-)(n_1 - n_0) = -r_0$$
 - We CANNOT solve, because it is still a nonlinear equation of n_1 .
 - Instead, we can solve the following (approximate) equation:
$$2n_0(n_1 - n_0) - (N_D^+ - N_A^-)(n_1 - n_0) = -r_0$$

Iteration

- Unfortunately, the solution, $n_1 - n_0$, is not exact.
 - Anyway, from n_0 and $n_1 - n_0$, we can calculate (inaccurate) n_1 .
 - Then, from n_1 , we again assert that n_2 is the real solution.
 - Again, n_2 will not be perfect.
 - Even though it is not perfect, it may approach to the solution.
 - Repeat this procedure until the error is sufficiently reduced.

Thank you!