

Special Topics on Basic EECS I

VLSI Devices

Lecture 21

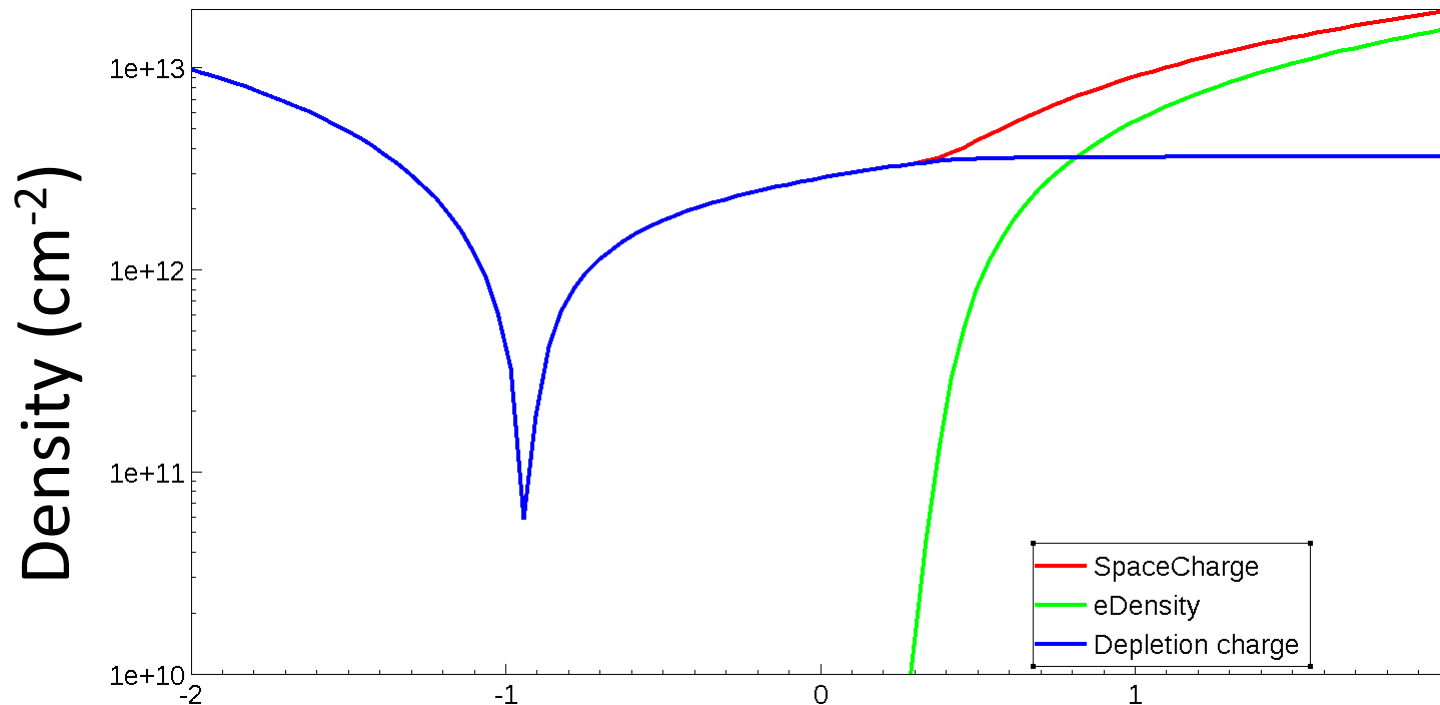
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Maximum depletion width

- Therefore, maximum depletion width becomes

$$W_d = \sqrt{\frac{4\epsilon_{si}\phi_B}{qN_a}} = \sqrt{\frac{4\epsilon_{si}k_B T \ln(N_a/n_i)}{q^2 N_a}} \quad \text{Taur, Eq. (2.190)}$$



← Depletion charge does not increase.

Gate voltage (V)

Beyond threshold voltage

It's not perfectly fixed.

- The surface potential is almost fixed. (Surface potential pinning)
 - Small additional change in ϕ_s induces an exponential increase of the electron density.

- Remember that $n = n_i \exp\left(\frac{q\phi}{k_B T}\right)$.

- When $\phi_s = 2\phi_B$,

$$n(0) = n_i \exp\left(\frac{q\phi_B}{k_B T}\right) = p(\infty)$$

- Additional potential ($\Delta\phi$) yields

$$n(0) = p(\infty) \exp\left(\frac{q\Delta\phi}{k_B T}\right)$$

It's a high density.

General relation beyond depletion approx. (1)

- With the depletion approximation, we obtained

$$Q_d = -qN_aW_d = -\sqrt{2\epsilon_{si}qN_a\phi_s}$$

- We can do much better!

- A generation relation for $Q_s = Q_d + Q_i$

Inversion
charge



- The Poisson equation

$$\frac{d^2\phi}{dx^2} = -\frac{q}{\epsilon_{si}}[p(x) - n(x) - N_a]$$

Taur, Eq. (2.175)

General relation beyond depletion approx. (2)

- Following Taur's notation,

For a while, $\phi(\infty) = -\phi_B$ is used as the reference value.
Therefore,

$$n(x) = n_i \exp\left(\frac{q\phi(x)}{k_B T}\right) \Rightarrow n(x) = n(\infty) \exp\left(\frac{q\phi(x)}{k_B T}\right) \quad \text{Taur, Eq. (2.178)}$$

$$p(x) = n_i \exp\left(-\frac{q\phi(x)}{k_B T}\right) \Rightarrow p(x) = p(\infty) \exp\left(-\frac{q\phi(x)}{k_B T}\right) \quad \text{Taur, Eq. (2.177)}$$

– The Poisson equation

$$\frac{d^2 \phi}{dx^2} = -\frac{q}{\epsilon_{si}} \left[N_a \left(\exp\left(-\frac{q\phi}{k_B T}\right) - 1 \right) - \frac{n_i^2}{N_a} \left(\exp\left(\frac{q\phi}{k_B T}\right) - 1 \right) \right] \quad \text{Taur, Eq. (2.179)}$$

General relation beyond depletion approx. (3)

- Multiplying $\frac{d\phi}{dx} dx$,

– The Poisson equation

$$\frac{d\phi}{dx} d\left(\frac{d\phi}{dx}\right)$$

$$= -\frac{q}{\epsilon_{si}} \left[N_a \left(\exp\left(-\frac{q\phi}{k_B T}\right) - 1 \right) - \frac{n_i^2}{N_a} \left(\exp\left(\frac{q\phi}{k_B T}\right) - 1 \right) \right] d\phi$$

– Integrate the above equation.

$$\int_0^{-E_x(x)} \frac{d\phi}{dx} d\left(\frac{d\phi}{dx}\right)$$

Taur, Eq. (2.180)

$$= -\frac{q}{\epsilon_{si}} \int_0^{\phi(x)} \left[N_a \left(\exp\left(-\frac{q\phi}{k_B T}\right) - 1 \right) - \frac{n_i^2}{N_a} \left(\exp\left(\frac{q\phi}{k_B T}\right) - 1 \right) \right] d\phi$$

General relation beyond depletion approx. (4)

- (Square of) Electric field

– From $\frac{1}{2} E_x^2(x) = -\frac{q}{\epsilon_{si}} \left[-N_a \frac{k_B T}{q} \exp\left(-\frac{q\phi}{k_B T}\right) - N_a \phi + N_a \frac{k_B T}{q} - \frac{n_i^2}{N_a} \frac{k_B T}{q} \exp\left(\frac{q\phi}{k_B T}\right) + \frac{n_i^2}{N_a} \phi + \frac{n_i^2}{N_a} \frac{k_B T}{q} \right]$, we get

$$\begin{aligned} E_x^2(x) &= \frac{2k_B T N_a}{\epsilon_{si}} \left[\left(\exp\left(-\frac{q\phi}{k_B T}\right) + \frac{q\phi}{k_B T} - 1 \right) \right. \\ &\quad \left. + \frac{n_i^2}{N_a^2} \left(\exp\left(\frac{q\phi}{k_B T}\right) - \frac{q\phi}{k_B T} - 1 \right) \right] \end{aligned}$$

Taur, Eq. (2.181)

General relation beyond depletion approx. (5)

- At $x = 0$, we have $\phi(0) = \phi_s$.
 - Then,

$$\begin{aligned} E_s^2 &= \frac{2k_B T N_a}{\epsilon_{si}} \left[\left(\exp \left(-\frac{q\phi_s}{k_B T} \right) + \frac{q\phi_s}{k_B T} - 1 \right) \right. \\ &\quad \left. + \frac{n_i^2}{N_a^2} \left(\exp \left(\frac{q\phi_s}{k_B T} \right) - \frac{q\phi_s}{k_B T} - 1 \right) \right] \end{aligned}$$

General relation beyond depletion approx. (6)

- At $x = 0$, we have $\phi(0) = \phi_s$.

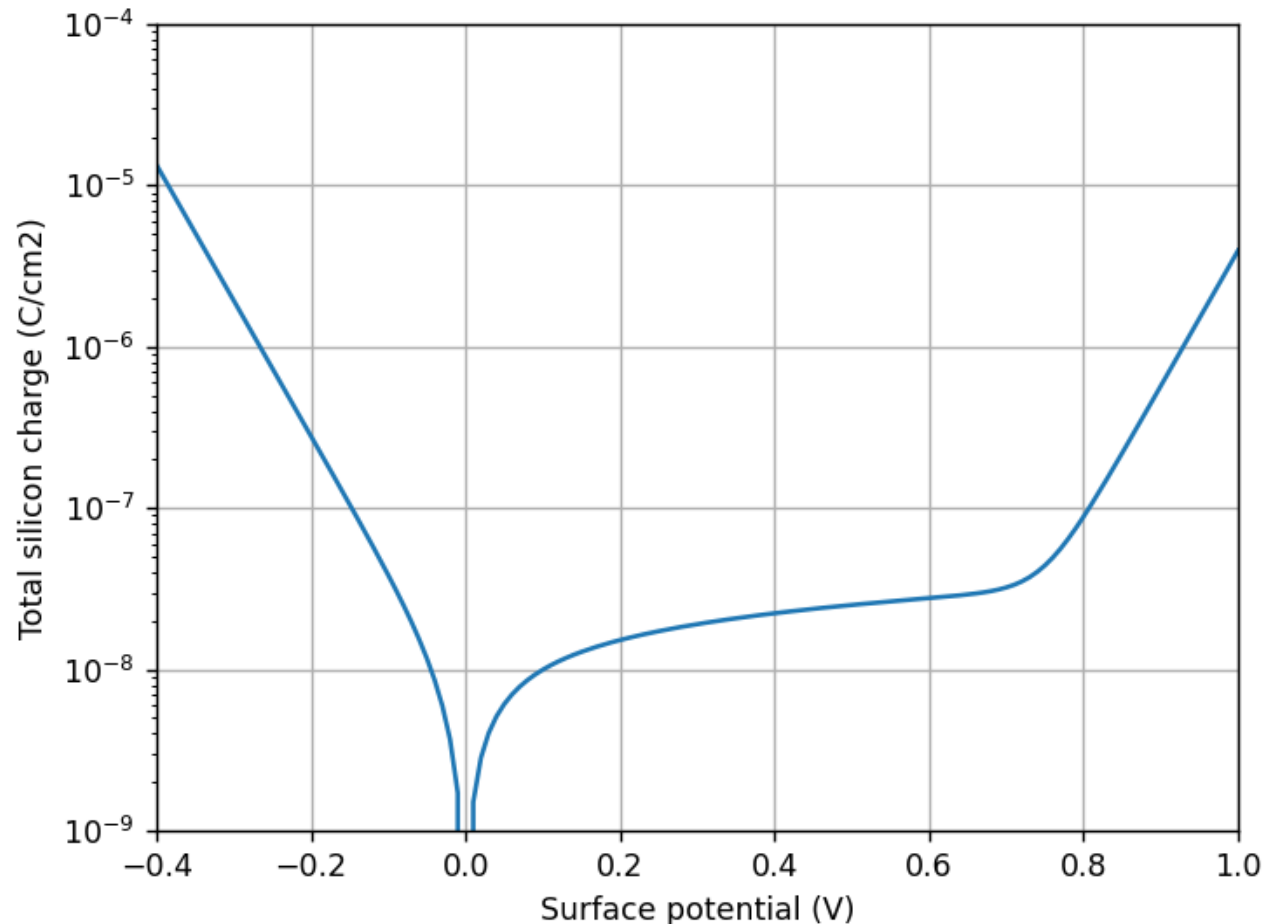
– From $Q_s = -\epsilon_{si}E_s$,

$$Q_s = \pm \sqrt{2\epsilon_{si}k_B T N_a} \left[\left(\exp\left(-\frac{q\phi_s}{k_B T}\right) + \frac{q\phi_s}{k_B T} - 1 \right) + \frac{n_i^2}{N_a^2} \left(\exp\left(\frac{q\phi_s}{k_B T}\right) - \frac{q\phi_s}{k_B T} - 1 \right) \right]^{1/2}$$

Taur, Eq. (2.182)

Homework#4

- Draw $|Q_s|$ as a function of ϕ_s .
 - Assume that N_a is $4 \times 10^{15} \text{ cm}^{-3}$. ϕ_s varies from -0.4 V to 1.0 V.

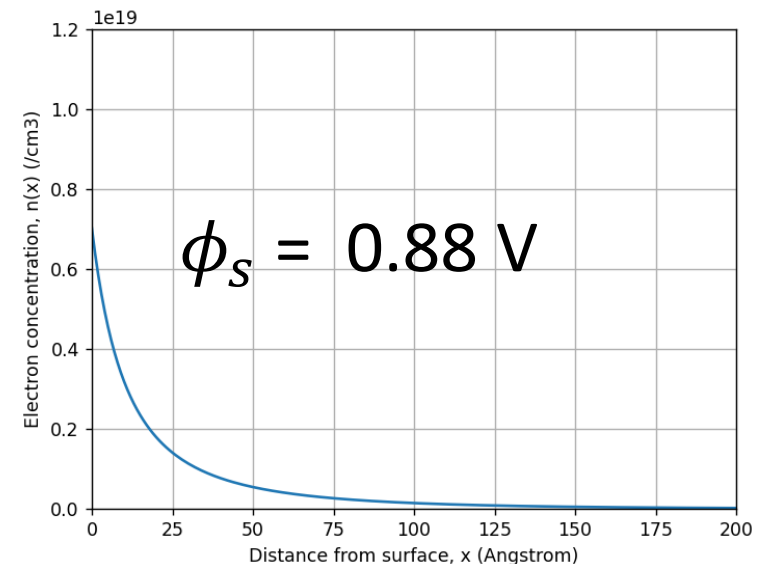
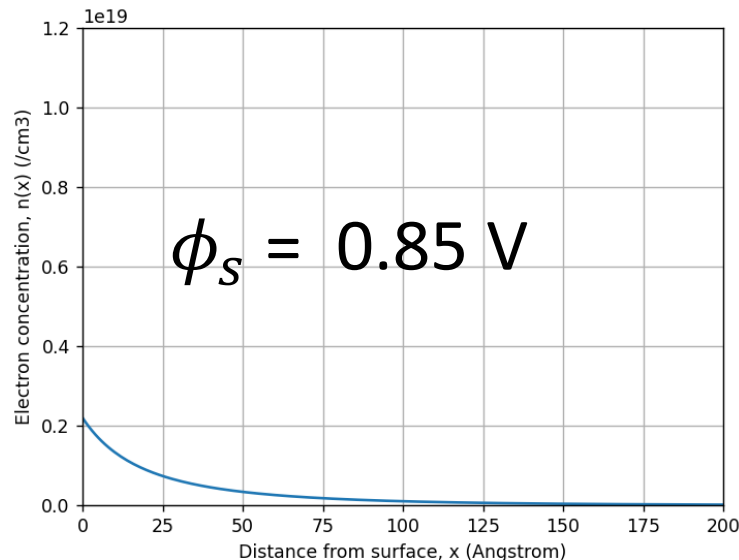


Strong inversion

- Beyond strong inversion,

$$\frac{d\phi}{dx} \approx - \sqrt{\frac{2k_B T N_a}{\epsilon_{si}} \left(\frac{q\phi}{k_B T} + \frac{n_i^2}{N_a^2} \exp\left(\frac{q\phi}{k_B T}\right) \right)} \quad \text{Taur, Eq. (2.191)}$$

- The electrons are distributed extremely close to the surface with an inversion-layer width less than 50 Å.



MOS equation

- Up to now, $Q_s(\phi_s)$ is found. We can control only V_g .

Total silicon
charge per
unit area

- Relation between V_g and ϕ_s

$$V_g - V_{fb} = V_{ox} + \phi_s = -\frac{Q_s}{C_{ox}} + \phi_s \quad \text{Taur, Eq. (2.195)}$$

$\frac{\epsilon_{ox}}{t_{ox}}$, oxide capacitance per unit area

- In general, $Q_s(\phi_s)$ is known. We can solve the above equation.

Taur, Eq. (2.182)

Thank you!