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# Lecture4: Diode

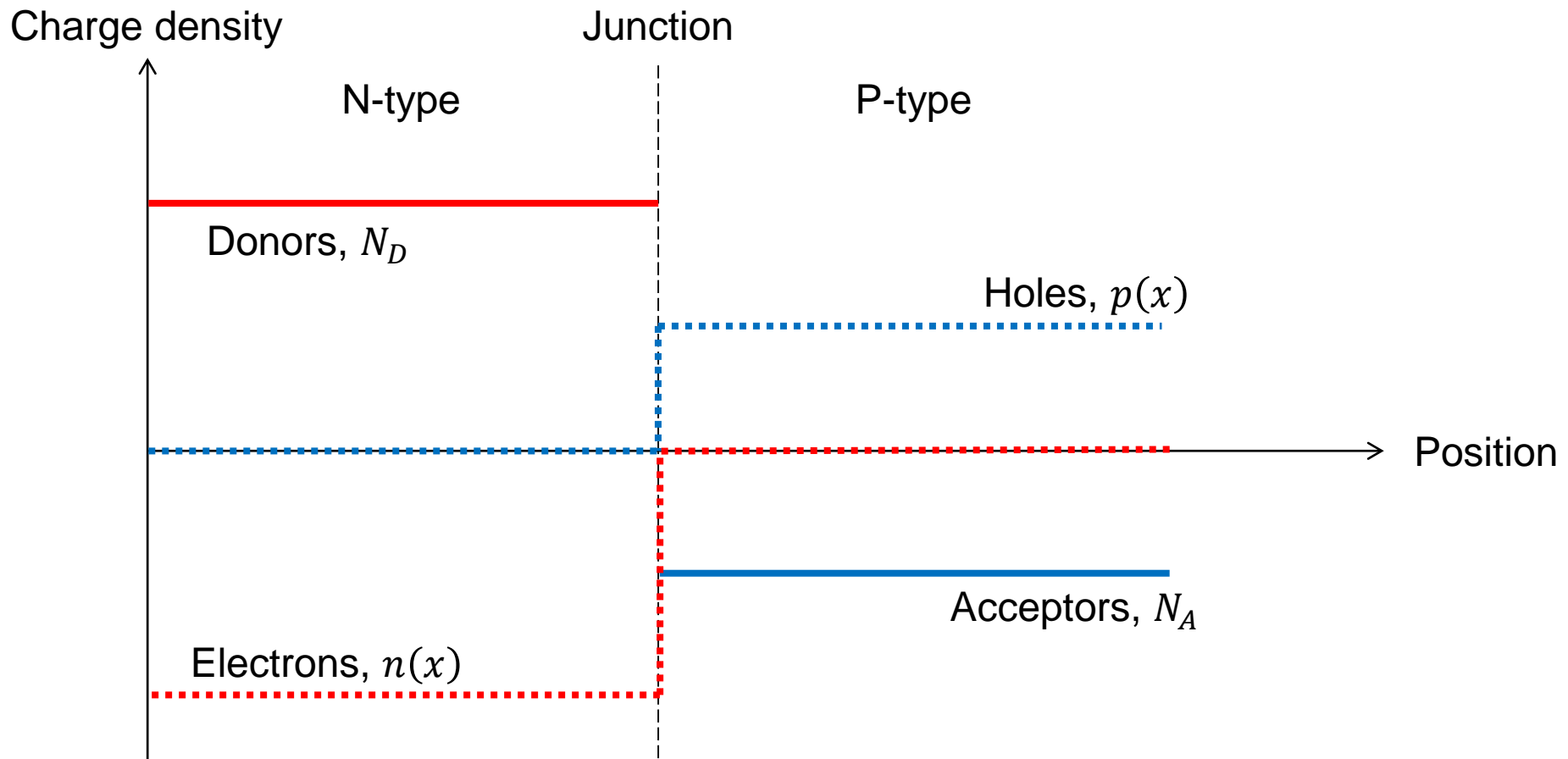
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# This is NOT a solution.

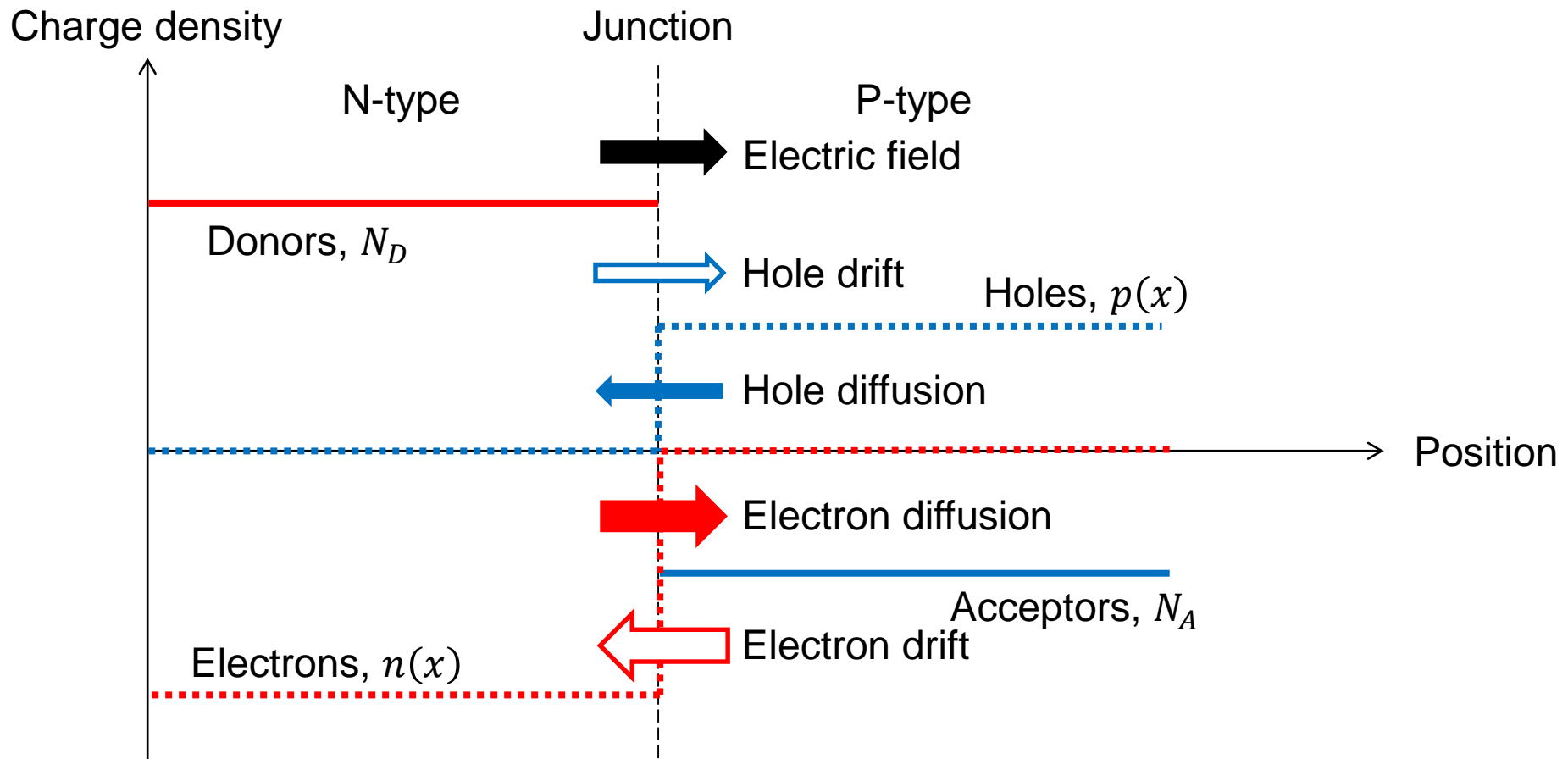
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- $n = N_D$  in the N-type region and  $p = N_A$  in the P-type
  - We need an electric field near the junction.



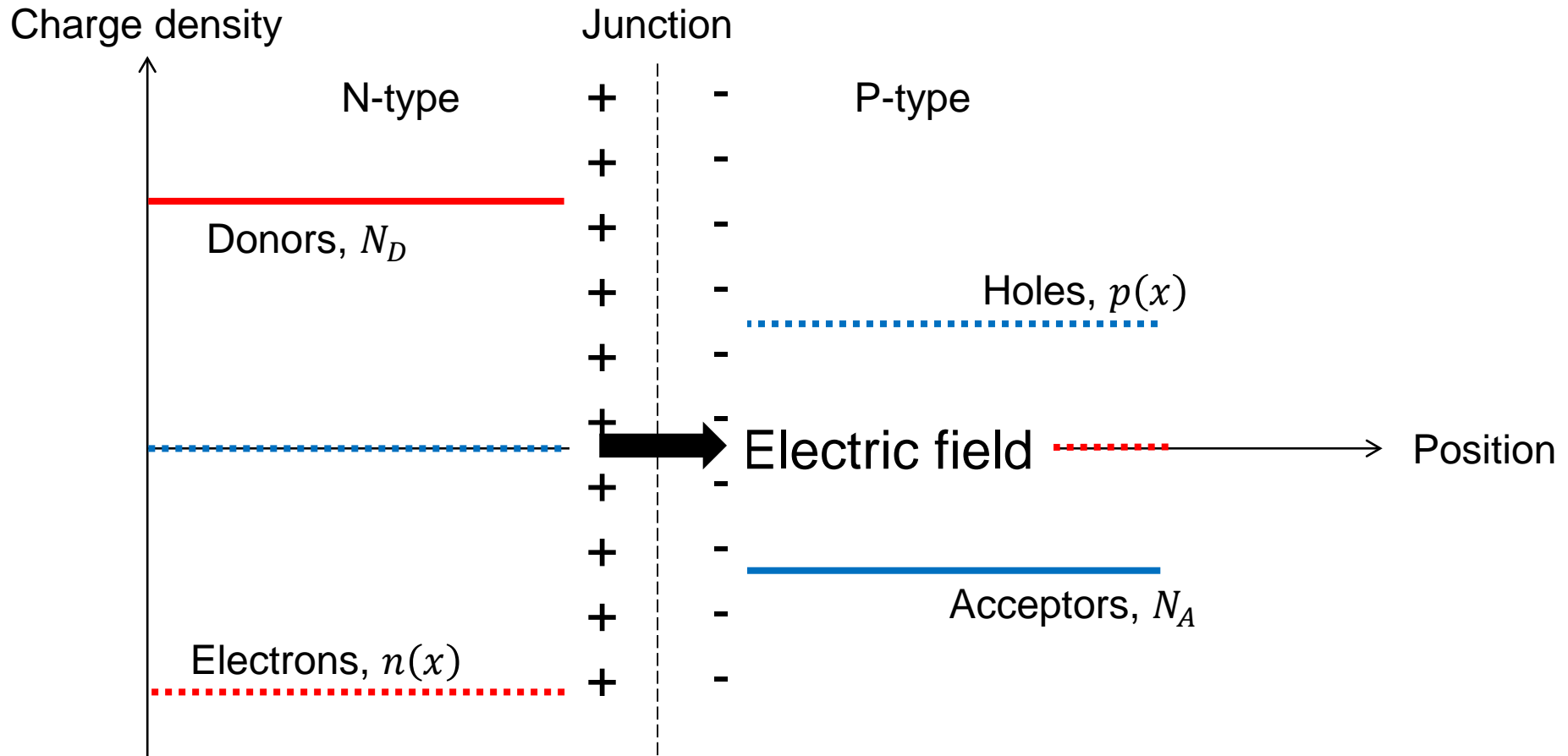
# Direction and location of E

- Electric field from the N-type to P-type
  - Only near the junction, we have the electric field.



# Localized E near junction

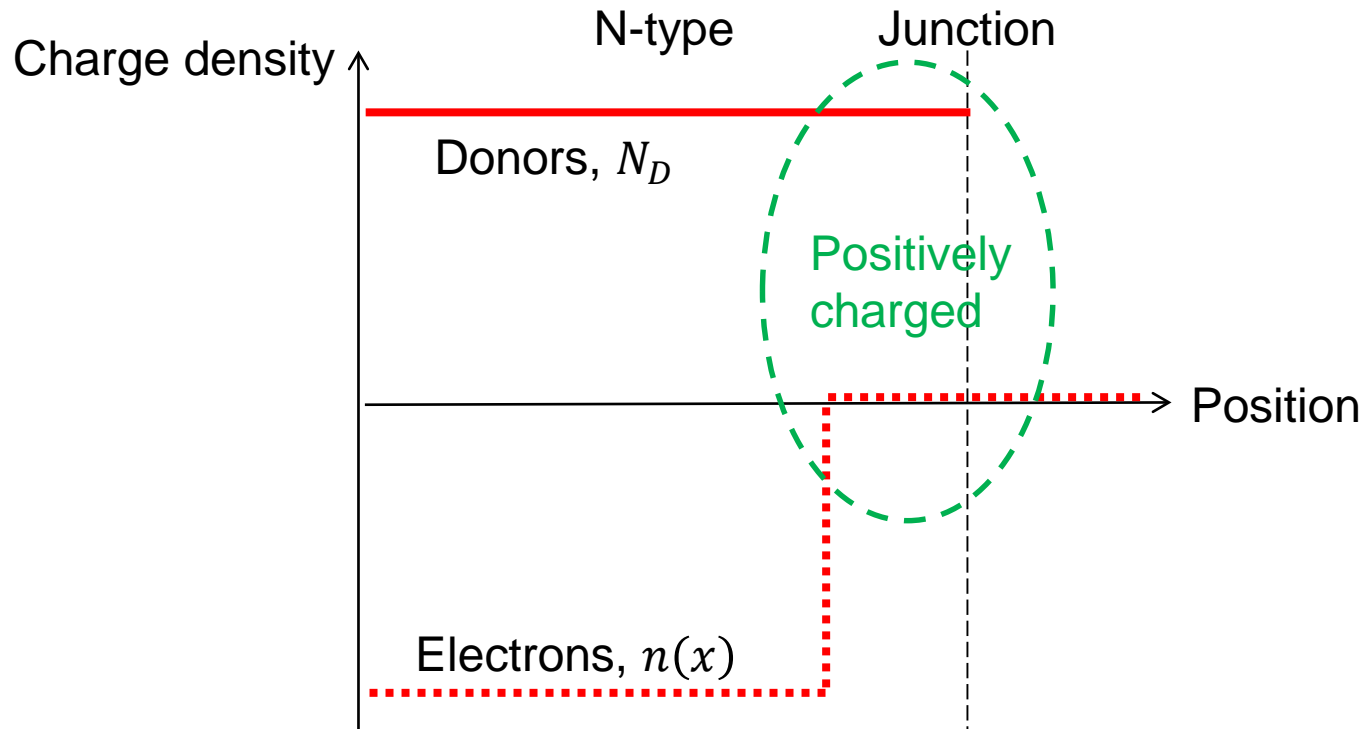
- (+) charges in the N-type region and (-) charges in the P-type region



# Space charge

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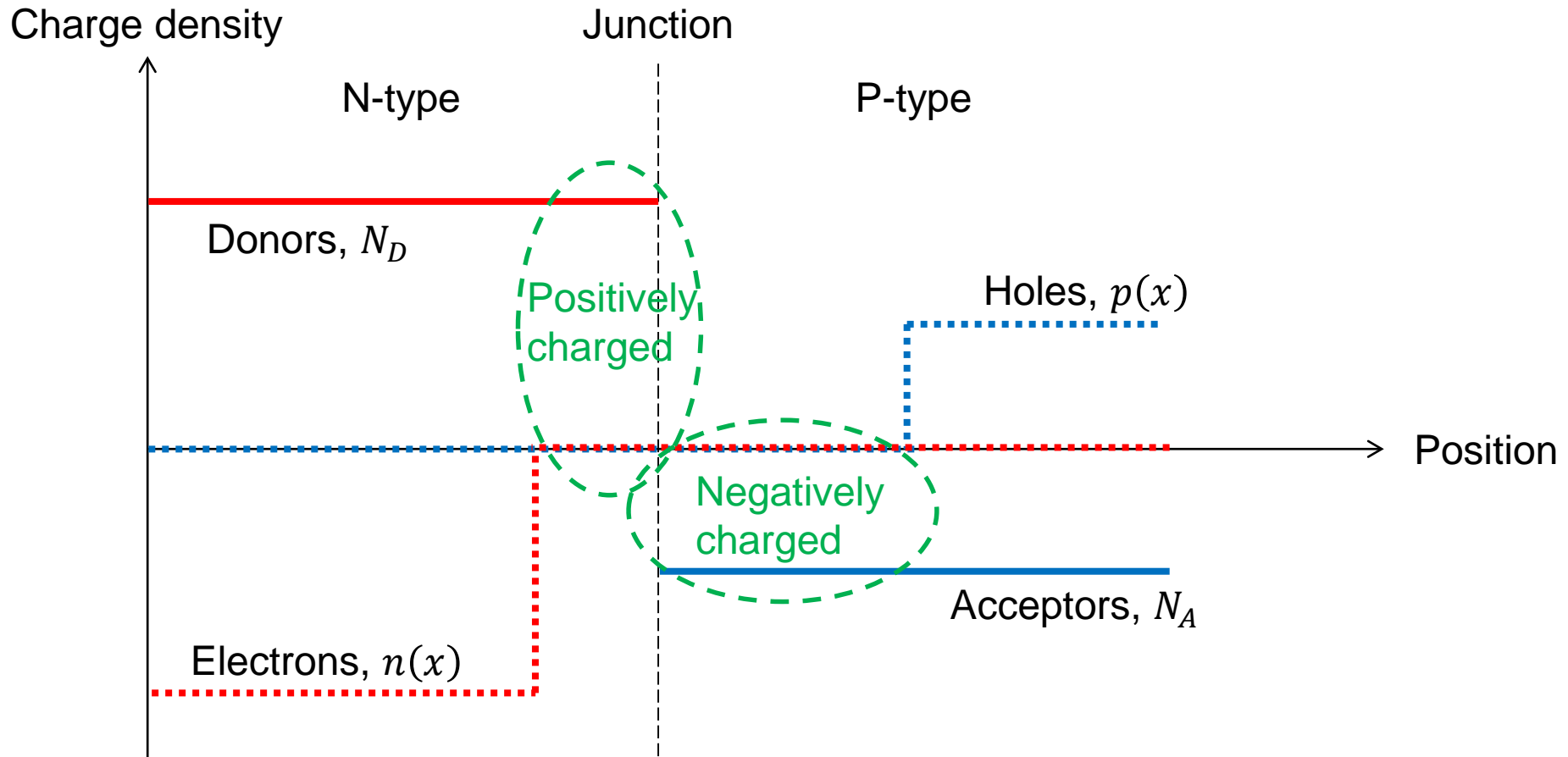
- How can we have (+) charges in the N-type region?
  - Immobile donor atoms? (Positively charged)
  - Mobile electrons! (Negatively charged)



# This is a solution.

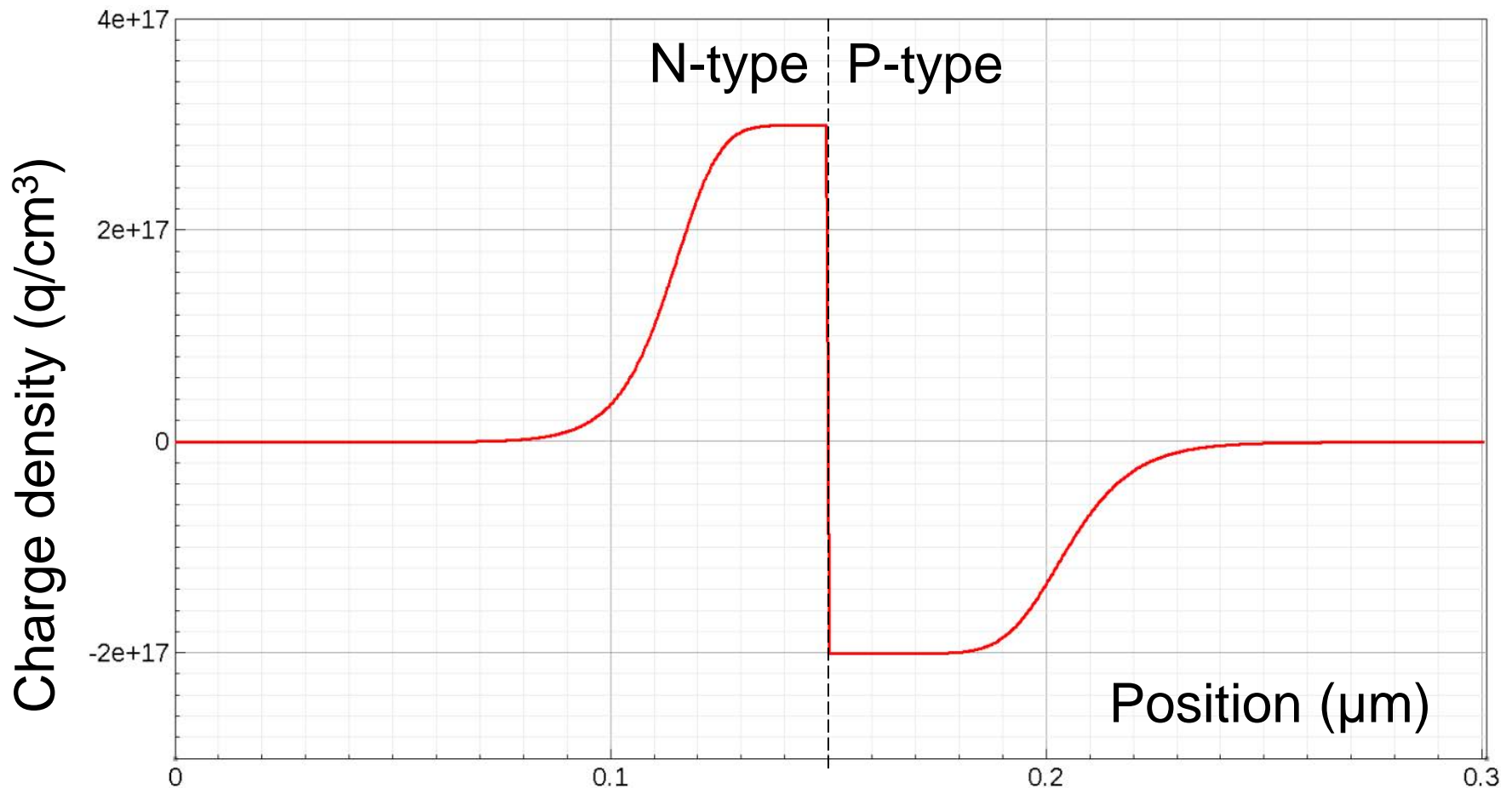
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- Depletion region
  - Calculation of the depletion width is of interest.



# Numerical example

- An example of  $N_D = 3 \times 10^{17} \text{ cm}^{-3}$  and  $N_A = 2 \times 10^{17} \text{ cm}^{-3}$ 
  - Depletion approximation is not perfect, but quite good.

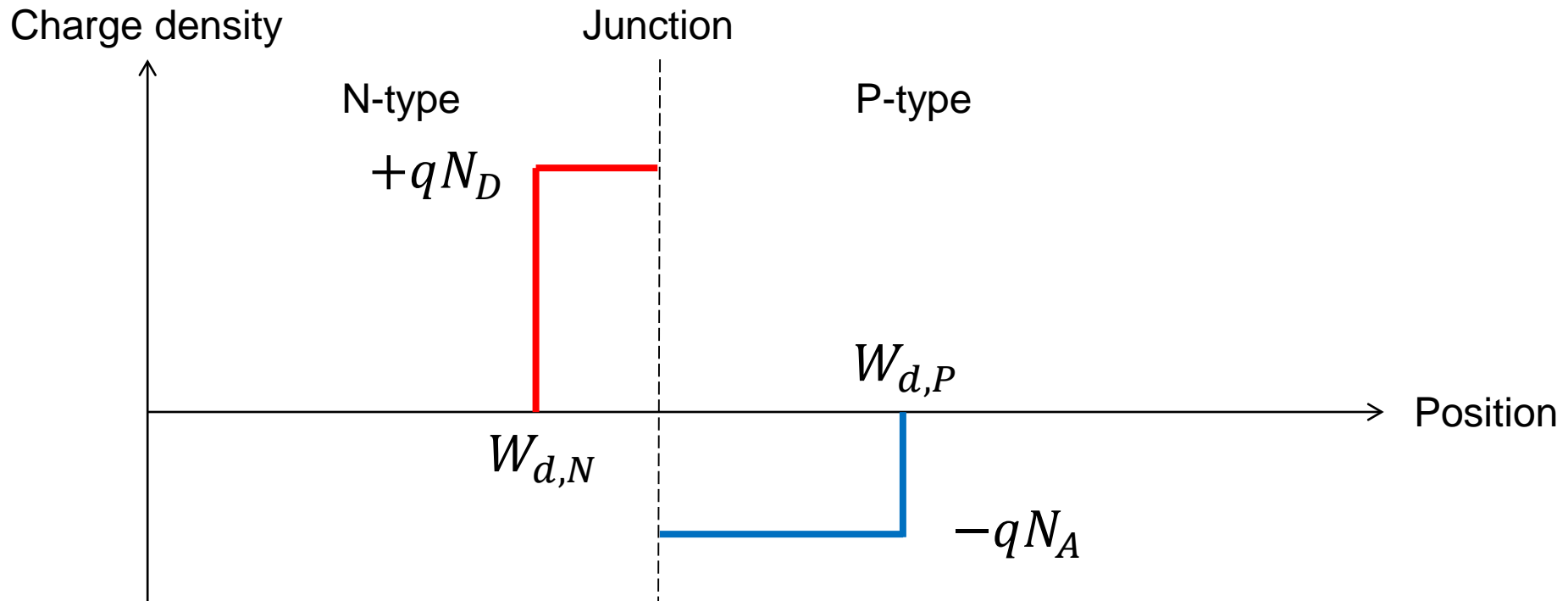


# Depletion width

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- Total depletion width,  $W_d$ 
  - Sum of  $W_{d,N}$  (N-type) and  $W_{d,P}$  (P-type)
  - For the charge balance,

$$N_D W_{d,N} = N_A W_{d,P}$$





# Electric field

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- Solve  $\frac{dE_x}{dx} = \frac{\rho}{\epsilon_{si}}$  where  $\rho = +qN_D$  or  $-qN_A$ .

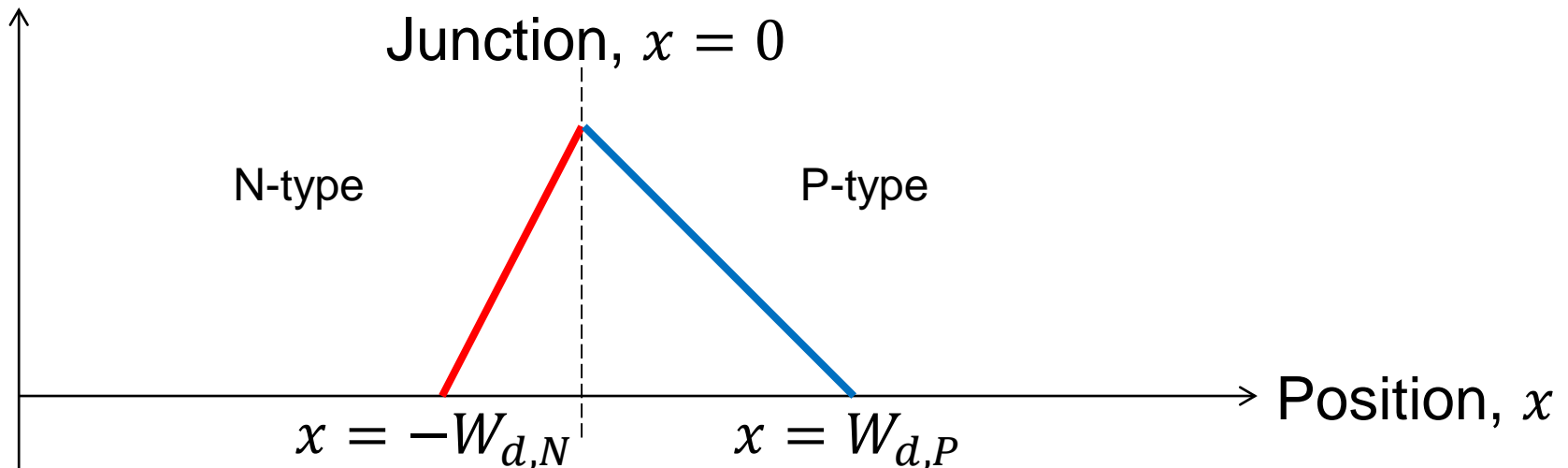
- For the N-type region,  $(-W_{d,N} \leq x \leq 0)$

$$E_x(x) = \frac{qN_D}{\epsilon_{si}} (x + W_{d,N})$$

- For the P-type region,  $(0 \leq x \leq W_{d,P})$

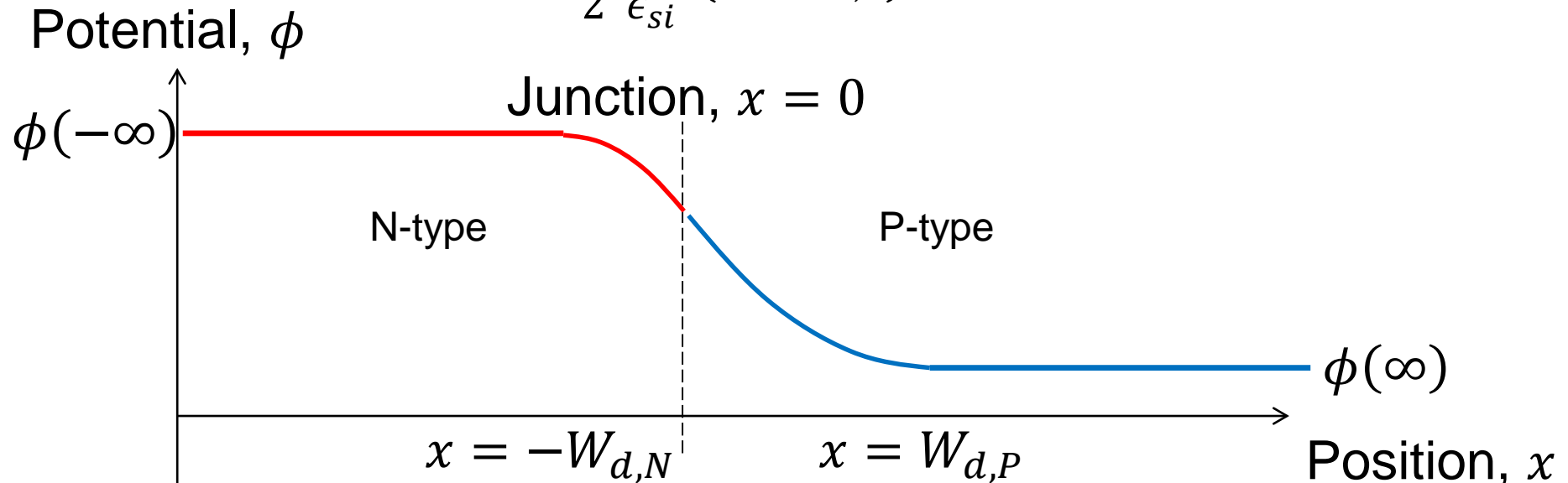
$$E_x(x) = -\frac{qN_A}{\epsilon_{si}} (x - W_{d,P})$$

Electric field,  $E_x$



# Electrostatic potential

- Solve  $-\frac{d\phi}{dx} = E_x$ .
  - For the N-type region,  $(-W_{d,N} \leq x \leq 0)$ 
$$\phi(x) = -\frac{1}{2} \frac{qN_D}{\epsilon_{si}} (x + W_{d,N})^2 + \phi(-\infty)$$
  - For the P-type region,  $(0 \leq x \leq W_{d,P})$ 
$$\phi(x) = \frac{1}{2} \frac{qN_A}{\epsilon_{si}} (x - W_{d,P})^2 + \phi(\infty)$$



# Depletion width

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- Let us define that  $V_{bi} \equiv \phi(-\infty) - \phi(\infty)$ .
  - It is called the built-in potential.
  - At  $x = 0$ , two curves for the electrostatic potential must meet together.
  - Then, the depletion width,  $W_d = W_{d,N} + W_{d,P}$ , can be written as

$$W_d = \sqrt{\frac{2\epsilon_{si}}{q} \left( \frac{1}{N_A} + \frac{1}{N_D} \right) V_{bi}}$$

# Built-in potential

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- Remaining task is to identify  $V_{bi}$ .

- At equilibrium, we have the following relation:

$$n = n_i \exp\left(\frac{\phi}{V_T}\right), \quad p = n_i \exp\left(-\frac{\phi}{V_T}\right)$$
$$V_T = \frac{k_B T}{q}$$

- In the N-type region, far away from the junction,

$$n(-\infty) = N_D = n_i \exp\left(\frac{\phi(-\infty)}{V_T}\right)$$

- In the P-type region, far away from the junction,

$$p(\infty) = N_A = n_i \exp\left(-\frac{\phi(\infty)}{V_T}\right)$$

- The built-in potential becomes

$$V_{bi} = V_T \log \frac{N_D N_A}{n_i^2}$$

# Hu, Example 4-1

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- Assume that  $N_D = 10^{20} \text{ cm}^{-3}$  and  $N_A = 10^{17} \text{ cm}^{-3}$ .

- Remember that

$$V_{bi} = V_T \log \frac{N_D N_A}{n_i^2}$$

- In this example, the built-in potential becomes about 1.012 V.

- Then, the depletion width becomes

$$W_d = \sqrt{\frac{2\epsilon}{q} \left( \frac{1}{N_A} + \frac{1}{N_D} \right) V_{bi}}$$

- It is about 115 nm. Moreover,  $W_{d,P} \gg W_{d,N}$ .

# Breathtaking balance

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- Diffusion is suppressed by a repelling electric field (the built-in field)
  - In this case, the current density vanishes.
- Even stronger electric field?
- Weaker electric field?