

VLSI Devices

Lecture 12

Sung-Min Hong (smhong@gist.ac.kr)

Semiconductor Device Simulation Laboratory

Department of Electrical Engineering and Computer Science

Gwangju Institute of Science and Technology (GIST)

Charge-sheet model

- Simpler model with further approximations

- Consider the previous method to calculate Q_i :

$$Q_i = -q \int_{\delta}^{\phi_s} \frac{(n_i^2 / N_a) \exp\left(\frac{q}{k_B T} (\phi - V)\right)}{E_x(\phi, V)} d\phi$$

- A more simple way? Instead, Q_d is approximated as

$$Q_d = -q N_a W_d = -\sqrt{2\epsilon_{si} q N_a \phi_s} \quad \text{Taur, Eq. (3.15)}$$

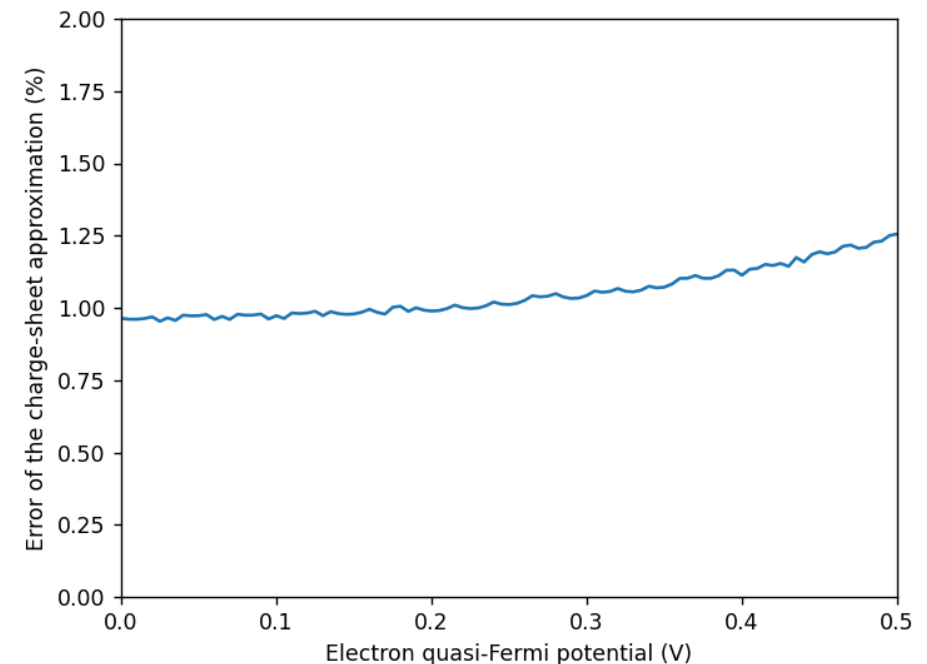
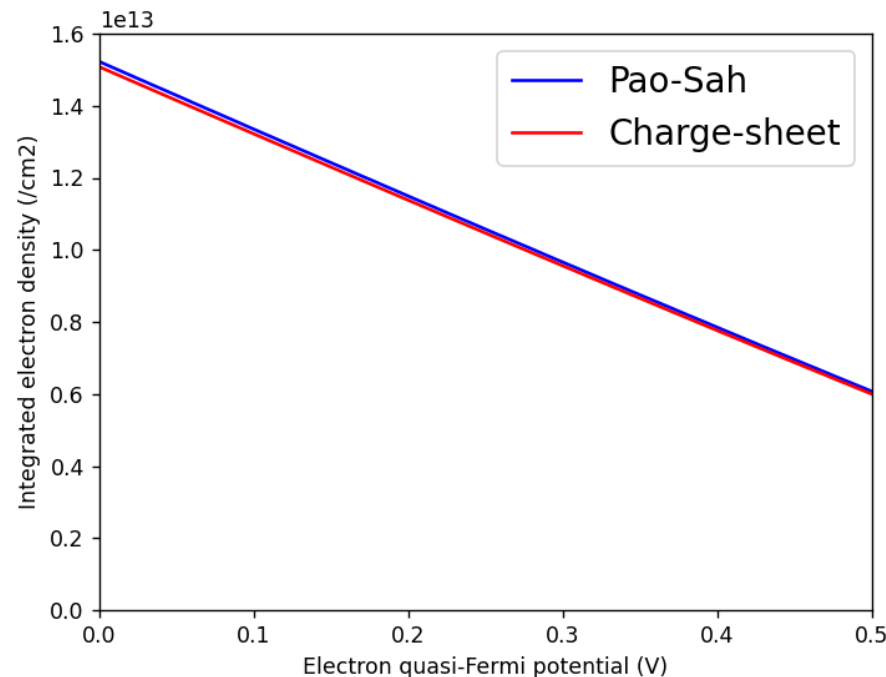
- Then, Q_i can be approximated as

$$Q_i = Q_s - Q_d = -C_{ox}(V_{gs} - V_{fb} - \phi_s) + \sqrt{2\epsilon_{si} q N_a \phi_s} \quad \text{Taur, Eq. (3.16)}$$

(Of course, it is not exact.)

Pao-Sah versus charge-sheet

- Consider a bias condition of $V_{gs} = 1.1$ V and $V_{ds} = 0.5$ V.
 - Calculate ϕ_s within a maximum error of 0.01 mV. Then, apply the charge-sheet approximation.
 - Only about 1 % error




Change of variable

- Now, Q_i can be a function of ϕ_s .
 - Variable change from V to ϕ_s :

$$I_d = \mu_{eff} \frac{W}{L} \int_0^{V_{ds}} [-Q_i(V)] dV = \mu_{eff} \frac{W}{L} \int_{\phi_{s,s}}^{\phi_{s,d}} [-Q_i(\phi_s)] \frac{dV}{d\phi_s} d\phi_s$$

Taur, Eq. (3.17)



Surface potentials at the two ends, $y = 0$ and L . They can be calculated by solving Taur, Eq. (3.14).

$V(\phi_s)$?

- Recall that

$$V_{gs} = V_{fb} + \phi_s + \frac{\sqrt{2\epsilon_{si}k_B T N_a}}{C_{ox}} \left[\frac{q\phi_s}{k_B T} + \frac{n_i^2}{N_a^2} \exp\left(\frac{q}{k_B T}(\phi_s - V)\right) \right]^{1/2}$$

Taur, Eq. (3.14)

- We can rearrange

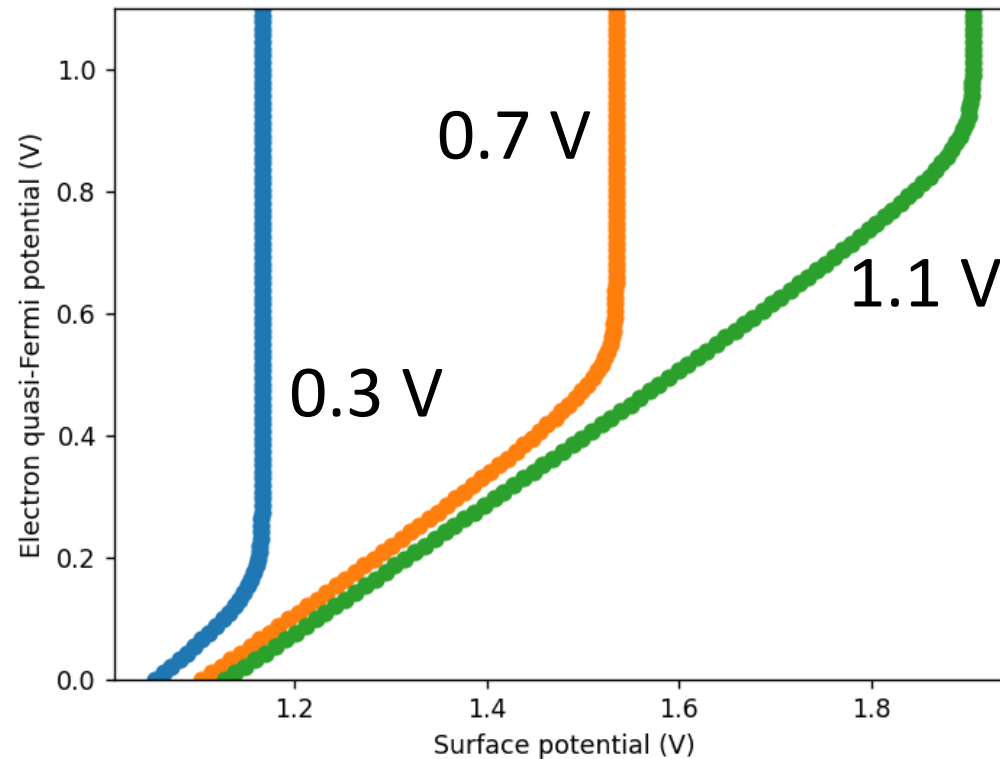
$$\frac{C_{ox}^2 (V_{gs} - V_{fb} - \phi_s)^2}{2\epsilon_{si}k_B T N_a} = \frac{q\phi_s}{k_B T} + \frac{n_i^2}{N_a^2} \exp\left(\frac{q}{k_B T}(\phi_s - V)\right)$$

- Therefore,

$$V = \phi_s - \frac{k_B T}{q} \log \left\{ \frac{N_a^2}{n_i^2} \left[\frac{C_{ox}^2 (V_{gs} - V_{fb} - \phi_s)^2}{2\epsilon_{si}k_B T N_a} - \frac{q\phi_s}{k_B T} \right] \right\}$$

Pao-Sah result

- Consider three gate voltages, $V_{gs} = 0.3 \text{ V}$, 0.7 V , and 1.1 V .
 - Draw the electron quasi-Fermi potential (V) as a function of the surface potential (ϕ_s). See Taur, Fig. 3.3.



$$\frac{dV}{d\phi_s} \text{? (1)}$$

- Recall that

$$V = \phi_s - \frac{k_B T}{q} \log \left\{ \frac{N_a^2}{n_i^2} \left[\frac{C_{ox}^2 (V_{gs} - V_{fb} - \phi_s)^2}{2\epsilon_{si} k_B T N_a} - \frac{q\phi_s}{k_B T} \right] \right\}$$

Taur, Eq. (3.18)

– Therefore,

$$\frac{dV}{d\phi_s} = 1 - \frac{k_B T}{q} \frac{-\frac{C_{ox}^2 (V_{gs} - V_{fb} - \phi_s)}{\epsilon_{si} k_B T N_a} - \frac{q}{k_B T}}{\frac{C_{ox}^2 (V_{gs} - V_{fb} - \phi_s)^2}{2\epsilon_{si} k_B T N_a} - \frac{q\phi_s}{k_B T}}$$

$$\frac{dV}{d\phi_s} \text{? (2)}$$

- Simple rearrange yields

$$\frac{dV}{d\phi_s} = 1 + \frac{2k_B T}{q} \frac{C_{ox}^2 (V_{gs} - V_{fb} - \phi_s) + \epsilon_{si} q N_a}{C_{ox}^2 (V_{gs} - V_{fb} - \phi_s)^2 - 2\epsilon_{si} q N_a \phi_s}$$

Taur, Eq. (3.19)

– It is still very complicated...

Integrand

- When multiplied with $-Q_i(\phi_s)$,

$$\begin{aligned} & (-Q_i(\phi_s)) \frac{dV}{d\phi_s} \\ &= -Q_i(\phi_s) + \frac{2k_B T}{q} \frac{C_{ox}^2(V_{gs} - V_{fb} - \phi_s) + \epsilon_{si} q N_a}{C_{ox}(V_{gs} - V_{fb} - \phi_s) + \sqrt{2\epsilon_{si} q N_a \phi_s}} \end{aligned}$$

- The second term is still very complicated...
- *Is it really important?*

Comparison between two terms

- Let's draw two terms.

- Assume that $N_a = 10^{18} \text{ cm}^{-3}$, $t_{ox} = 1.2 \text{ nm}$, and $V_{fb} = -1.08 \text{ V}$.

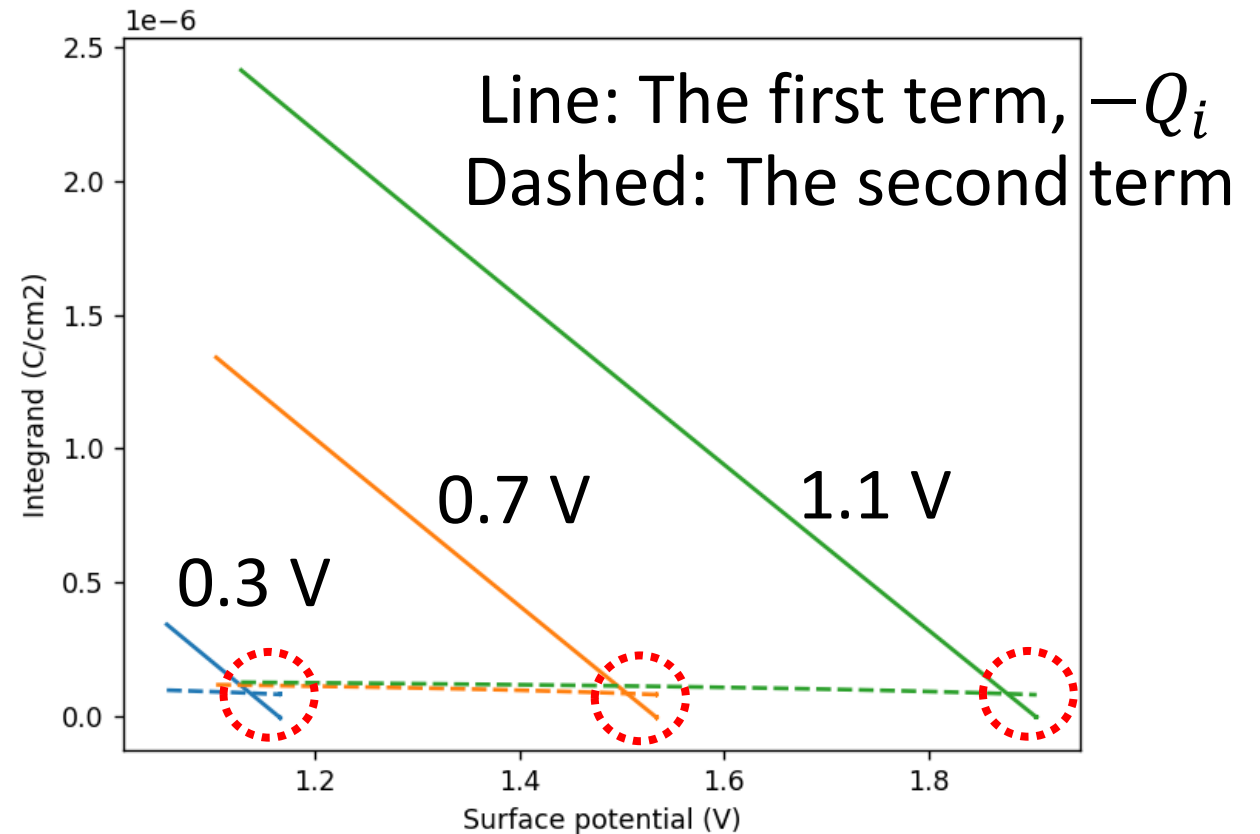
- The second term is small.

- It is meaningful only when $Q_i \approx 0$.

- This is corresponding to

$$C_{ox}(V_{gs} - V_{fb} - \phi_s)$$

$$= \sqrt{2\epsilon_{si}qN_a\phi_s}.$$



Integrand, again

- Within this condition,

$$\begin{aligned} (-Q_i(\phi_s)) \frac{dV}{d\phi_s} &\approx -Q_i(\phi_s) + \frac{k_B T C_{ox} \sqrt{2\epsilon_{si} q N_a \phi_s} + \epsilon_{si} q N_a}{q \sqrt{2\epsilon_{si} q N_a \phi_s}} \\ &= -Q_i(\phi_s) + \frac{k_B T}{q} C_{ox} + \frac{k_B T \sqrt{2\epsilon_{si} q N_a}}{q 2\sqrt{\phi_s}} \end{aligned}$$

Its integration yields $\frac{k_B T}{q} C_{ox} \phi_s$.

Its integration yields $\frac{k_B T}{q} \sqrt{2\epsilon_{si} q N_a \phi_s}$.

Drain current

- Using the previous approximation,
 - We can obtain the following expression:

$$I_d = \mu_{eff} \frac{W}{L} \left\{ C_{ox} \left(V_{gs} - V_{fb} + \frac{k_B T}{q} \right) \phi_s - \frac{1}{2} C_{ox} \phi_s^2 - \frac{2}{3} \sqrt{2 \epsilon_{si} q N_a} \phi_s^{1.5} + \frac{k_B T}{q} \sqrt{2 \epsilon_{si} q N_a} \phi_s \right\} \bigg|_{\phi_{s,s}}^{\phi_{s,d}}$$

Taur, Eq. (3.21)

- Only with $\phi_{s,s}$ and $\phi_{s,d}$, we can evaluate the drain current.

Thank you!