

VLSI Devices

Lecture 16

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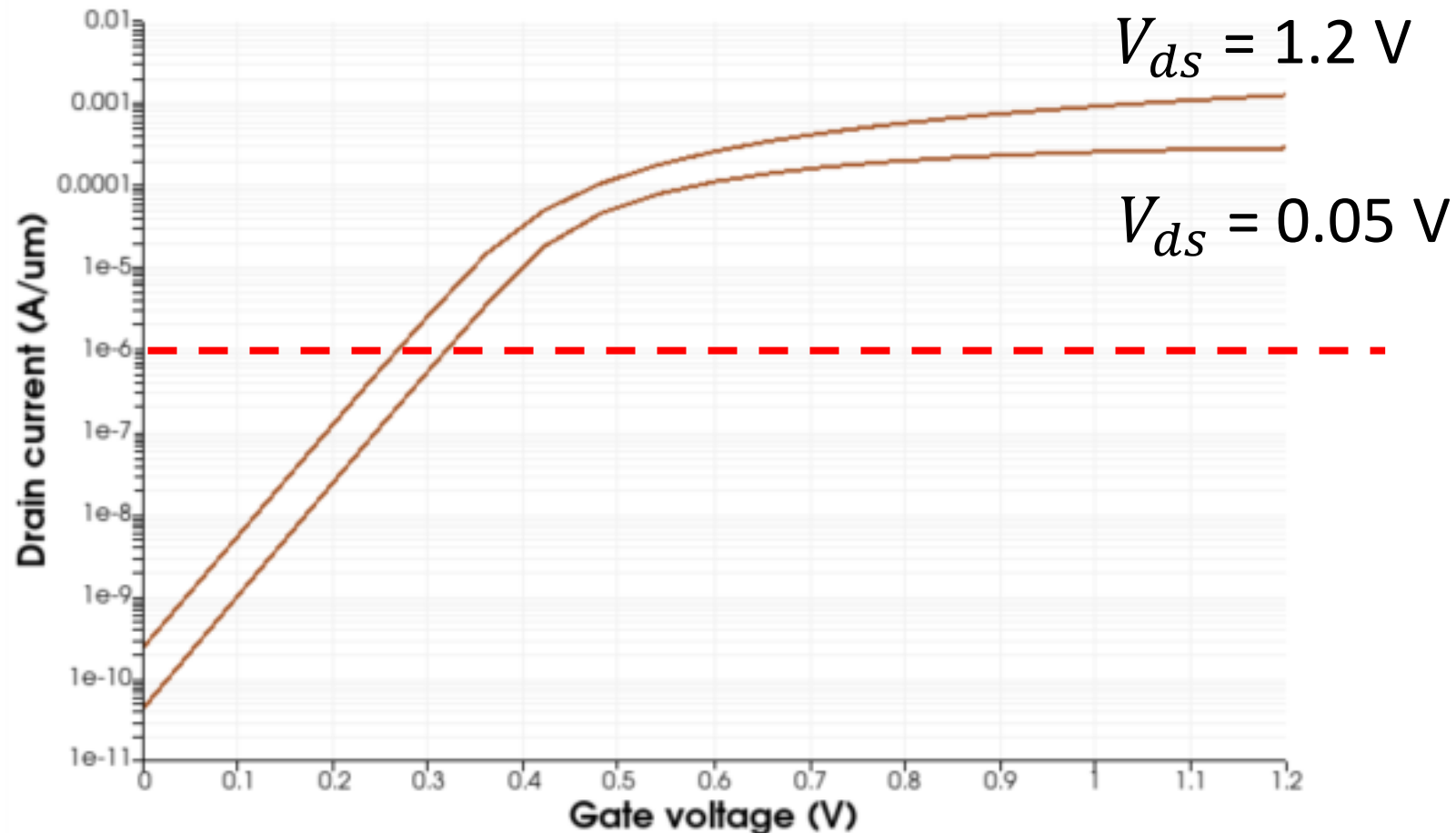
Gwangju Institute of Science and Technology (GIST)

Coverage

- Two YouTube lectures reserved for advanced topics
 - L14: ~~Substrate bias, channel mobility~~
 - L15: 3.2.1
 - L16: 3.2.1 (Continued)
 - L17: Velocity saturation (3.2.2)
 - L18: Channel length modulation and so on (3.2.3, 3.2.4, 3.2.5)
 - L19: MOSFET scaling
 - L20: MOSFET scaling (Continued)
 - L21: Quantum effect (4.2.4)
 - L22: Double-gate MOSFETs (10.3)
 - L23: FinFETs
 - L24: CFETs

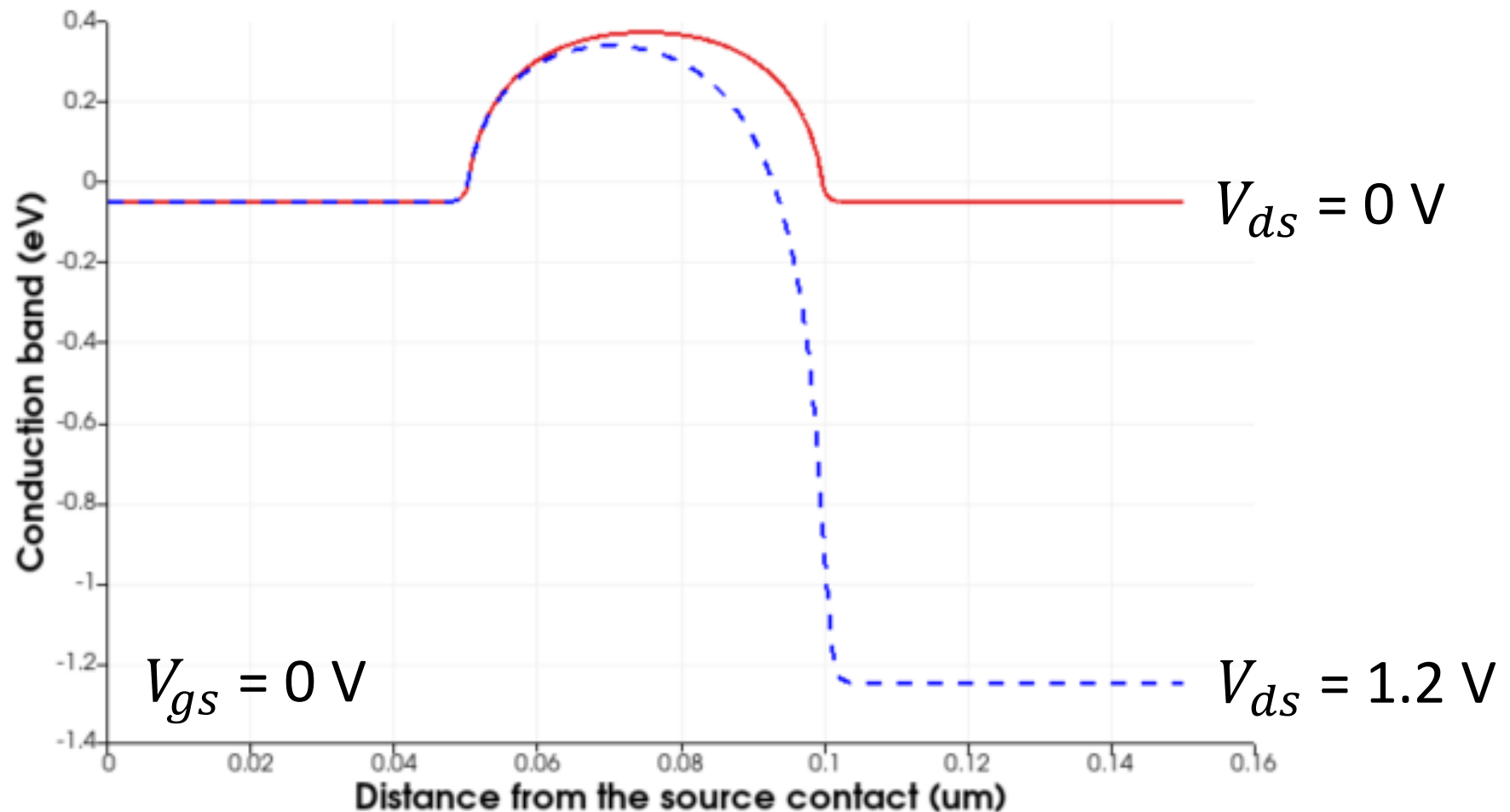
Drain-induced barrier lowering (DIBL)

- Much worse than the long-channel device
 - 45 mV/V @ $I_D = 10^{-6}$ A/ μm



Conduction band, again

- At a high V_{ds} , the energy barrier is further reduced.



Simplified geometry for an analytic solution

- Poisson equation

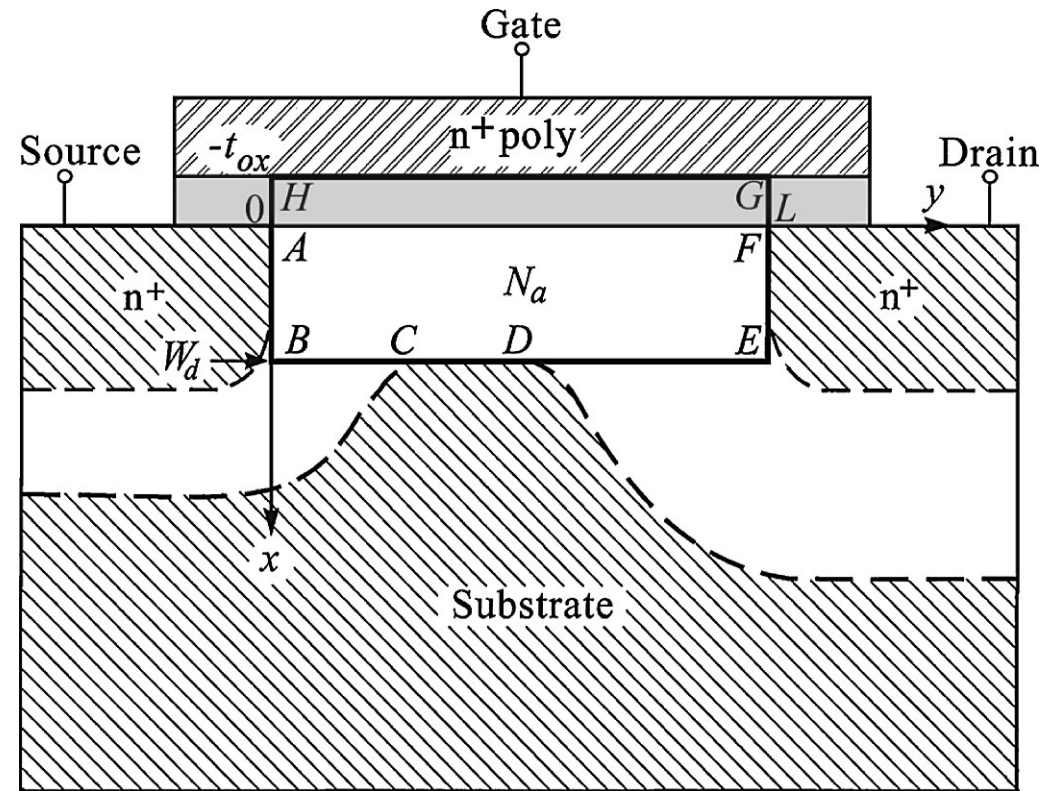
- For AFGH

$$\frac{\partial^2 \phi}{\partial x^2} + \frac{\partial^2 \phi}{\partial y^2} = 0 \quad \text{Taur, Eq. (A9.1)}$$

- For ABEF (Depleted)

$$\frac{\partial^2 \phi}{\partial x^2} + \frac{\partial^2 \phi}{\partial y^2} = \frac{qN_a}{\epsilon_{si}}$$

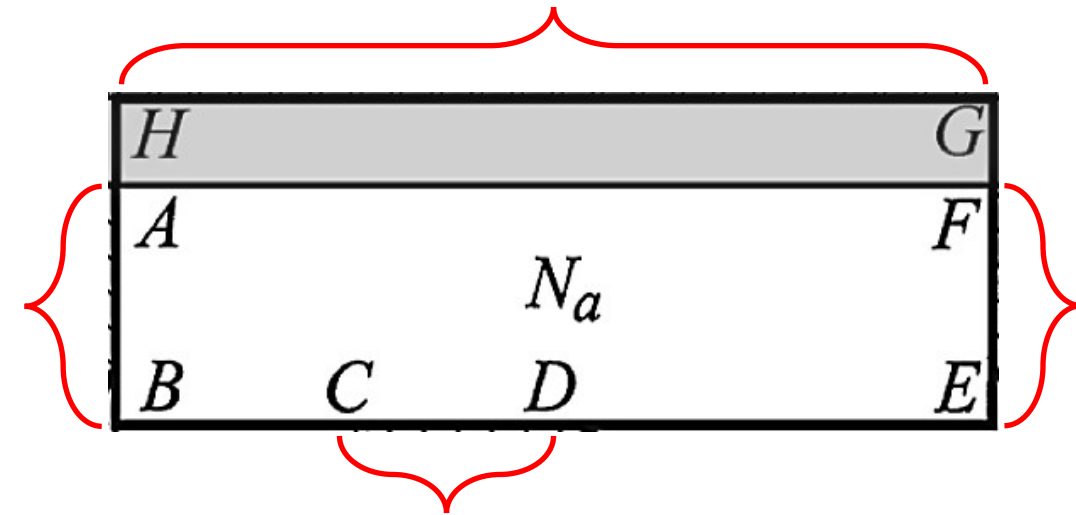
Taur, Eq. (A9.2)



Simplified geometry (Taur, Fig. A9.1)

Boundary conditions

- Potential reference is $\phi(\infty, y) = -\phi_B$.
 - Along GH: $\phi = V_{gs} - V_{fb}$ Taur, Eq. (A9.3)
 - Along AB: $\phi = \phi_{bi} \approx \frac{E_g}{2q} + \phi_B$ Taur, Eq. (A9.4)
 - Along EF: $\phi = \phi_{bi} + V_{ds}$ Taur, Eq. (A9.5)
 - Along CD: $\phi = 0$
Taur, Eq. (A9.6)
- ($3t_{ox}$ for uniform permittivity)



Solution

- Poisson equation with boundary conditions

– Try the following function for the electrostatic potential:

$$\phi(x, y) = v(x, y) + u_L(x, y) + u_R(x, y) + u_B(x, y) \quad \text{Taur, Eq. (A9.9)}$$

Poisson equation
with upper b.c.
(Long-channel)

Laplace equation
to match left b.c.

Laplace equation
to match right
b.c.

Laplace equation
to match bottom
b.c.

Solution, $v(x, y)$

- Actually, it is $v(x)$.

- For the oxide region ($-3t_{ox} \leq x \leq 0$),

$$v(x, y) = \phi_s - \frac{V_{gs} - V_{fb} - \phi_s}{3t_{ox}} x$$

Taur, Eq. (A9.10)

- For the silicon region ($0 \leq x \leq W_d$),

$$v(x, y) = \phi_s \left(1 - \frac{x}{W_d} \right)^2$$

Taur, Eq. (A9.11)

- It is noted that

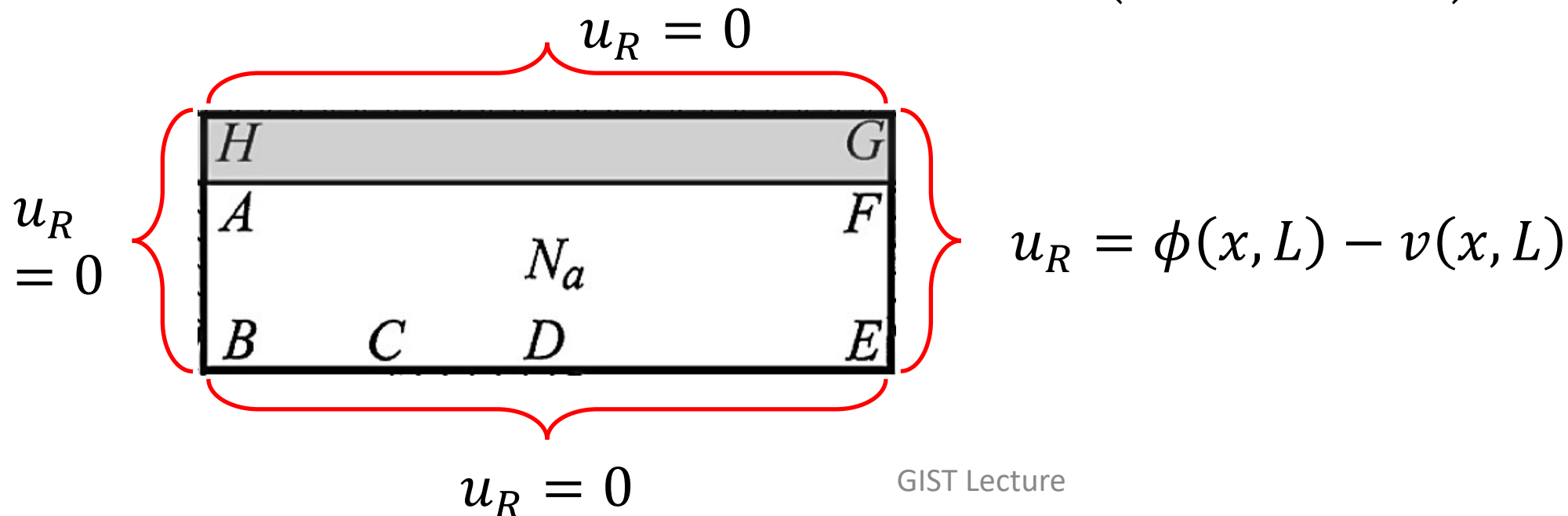
$$\phi_s = \frac{qN_d W_d^2}{2\epsilon_{si}}$$

Taur, Eq. (A9.12)

A mode for $u_R(x, y)$

- For three boundaries, it should vanish.
 - At only one side, it has non-zero values.
 - We can try ($\lambda \equiv W_d + 3t_{ox}$)

$$u_{R,n}(x, y) = \sinh\left(\frac{n\pi y}{\lambda}\right) \sin\left(\frac{n\pi(x + 3t_{ox})}{\lambda}\right)$$



Series expansion of $u_R(x, y)$

- $u_R(x, L) = \phi(x, L) - v(x, L)$ can be expanded with coefficients c_n 's.

– Therefore,

$$u_R(x, y) = \sum_{n=1}^{\infty} c_n \frac{\sinh\left(\frac{n\pi y}{\lambda}\right)}{\sinh\left(\frac{n\pi L}{\lambda}\right)} \sin\left(\frac{n\pi(x + 3_{tox})}{\lambda}\right) \quad \text{Taur, Eq. (A9.15)}$$

- Similar solutions are found for $u_L(x, y)$ and $u_B(x, y)$.

$$u_L(x, y) = \sum_{n=1}^{\infty} b_n \frac{\sinh\left(\frac{n\pi(L - y)}{\lambda}\right)}{\sinh\left(\frac{n\pi L}{\lambda}\right)} \sin\left(\frac{n\pi(x + 3_{tox})}{\lambda}\right) \quad \text{Taur, Eq. (A9.14)}$$

First-order expansion

- Keep only b_1 and c_1 . (u_B is neglected.)

– Then,

$$\begin{aligned}\phi(x, y) &= \phi_s \left(1 - \frac{x}{W_d}\right)^2 \\ &+ \frac{b_1 \sinh\left(\frac{\pi(L - y)}{\lambda}\right) + c_1 \sinh\left(\frac{\pi y}{\lambda}\right)}{\sinh\left(\frac{\pi L}{\lambda}\right)} \sin\left(\frac{\pi(x + 3_{tox})}{\lambda}\right)\end{aligned}$$

Taur, Eq. (A9.22)

- Approximate values for b_1 and c_1 are $\frac{4}{\pi}(\phi_{bi} - a\phi_s)$ and $\frac{4}{\pi}(\phi_{bi} + V_{ds} - a\phi_s)$, respectively. $a \approx 0.4$.

Surface potential

- At $x = 0$,

$$\phi(0, y) = \phi_s + \frac{b_1 \sinh\left(\frac{\pi(L-y)}{\lambda}\right) + c_1 \sinh\left(\frac{\pi y}{\lambda}\right)}{\sinh\left(\frac{\pi L}{\lambda}\right)} \sin\left(\frac{\pi 3t_{ox}}{\lambda}\right)$$

Taur, Eq. (A9.22)

– Let's find the maximum potential.

$$\frac{d}{dy} \phi(0, y) = \frac{\pi}{\lambda} \frac{-b_1 \cosh\left(\frac{\pi(L-y)}{\lambda}\right) + c_1 \cosh\left(\frac{\pi y}{\lambda}\right)}{\sinh\left(\frac{\pi L}{\lambda}\right)} \sin\left(\frac{\pi 3t_{ox}}{\lambda}\right) = 0$$

$$b_1 \cosh\left(\frac{\pi(L-y)}{\lambda}\right) = c_1 \cosh\left(\frac{\pi y}{\lambda}\right)$$

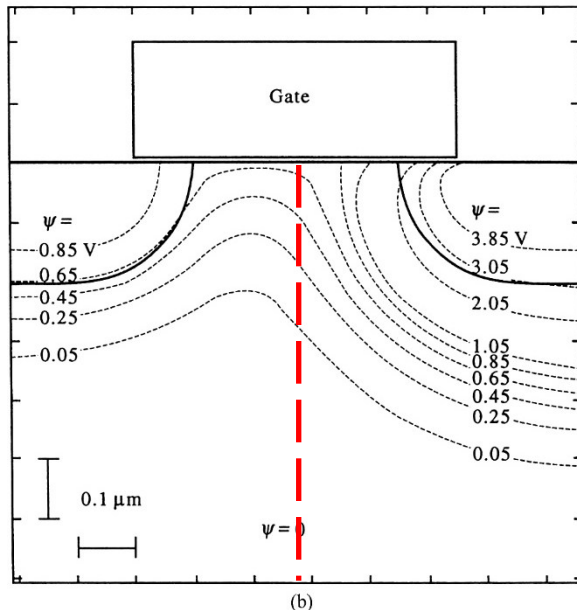
Position for maximum potential, $y = y_c$

- For a positive z , $\cosh z \approx \frac{\exp z}{2}$.

$$\exp \frac{\pi(L^2 - 2y_c)}{\lambda} = \frac{c_1}{b_1}$$

$$y_c = \frac{L}{2} - \frac{\lambda}{2\pi} \ln \frac{c_1}{b_1} \approx \frac{L}{2} - \frac{W_d + 3t_{ox}}{2\pi} \ln \left(1 + \frac{V_{ds}}{\phi_{bi} - a\phi_s} \right)$$

Taur, Eq. (A9.23)



Potential profile (Taur, Fig. 3.20(b))

Maximum potential at $y = y_c$

- Using some approximations,

$$\begin{aligned}\phi(0, y_c) &= \phi_s + 2\sqrt{b_1 c_1} \exp\left(-\frac{\pi L}{2\lambda}\right) \sin\left(\frac{\pi 3t_{ox}}{\lambda}\right) \\ &\approx \phi_s \\ &+ \left(\frac{6\pi t_{ox}}{\lambda}\right) \left(\sqrt{\phi_{bi}(\phi_{bi} + V_{ds})} - \frac{2\phi_{bi} + V_{ds}}{2\sqrt{\phi_{bi}(\phi_{bi} + V_{ds})}} a\phi_s \right) \exp\left(-\frac{\pi L}{2\lambda}\right) \\ &\sim \text{Taur, Eq. (A9.24)}\end{aligned}$$

– Threshold voltage lowering, ΔV_t ,

$$\begin{aligned}\Delta V_t &= \left(\frac{24t_{ox}}{W_{dm}}\right) \left(\sqrt{\phi_{bi}(\phi_{bi} + V_{ds})} - a(2\phi_B) \right) \exp\left(-\frac{\pi L}{2(W_{dm} + 3t_{ox})}\right) \\ &\sim \text{Taur, Eq. (A9.25)}\end{aligned}$$

Typical values in the textbook

- Following Taur, Eq. (3.67),

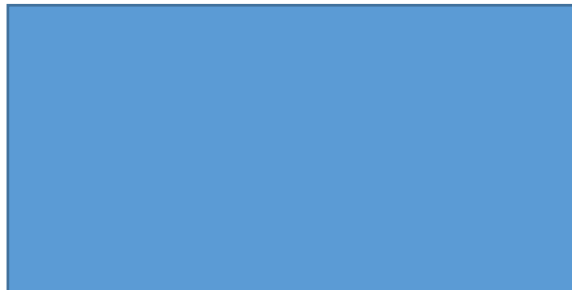
$$\Delta V_t = \left(\frac{24t_{ox}}{W_{dm}} \right) \left(\sqrt{\phi_{bi}(\phi_{bi} + V_{ds})} - a(2\phi_B) \right) \exp \left(-\frac{\pi L}{2(W_{dm} + 3t_{ox})} \right)$$

– Using typical values,

$$0.1 = (2.4)(\sqrt{2} - 0.4) \exp \left(-\frac{\pi L}{2(1.3W_{dm})} \right)$$

– We get $L \approx 2.6W_{dm} \approx 2(W_{dm} + 3t_{ox})$.

– The minimum allowable channel length is $L_{min} \approx 2(W_{dm} + 3t_{ox})$.



Thank you!