## Finite difference method 1

In the finite difference method, we discetize the exact solution of the Helmholtz problem u(x) on a grid R defined on a domain  $\Omega \subseteq \mathbb{R}^d$ , d=1,2,3. Suppose we have a one dimensional problem with  $\Omega=(0,1)$ . We want to approximate u in a fixed number of points, say N. We create a uniform grid

$$X_h = \{x_i \mid x_i = hi, \ i = 0, 1, \dots, N\}$$
 (1)

where we take our gridsize h = 1/N. We denote  $u_i = u(x_i)$ . Now, we want to discretize our equation for u on  $\Omega$ 

 $\frac{\mathrm{d}^2 u}{\mathrm{d}x^2} + fu = g, \quad \text{on } \Omega.$ (2)

We can approximate the second derivate of u in a point  $x_i$  by using the Taylor expansion of u at  $x_i$ . The value of u at points  $x_{i+1}$  and  $x_{i-1}$  are given by

$$u_{i+1} = u_i + h \frac{\mathrm{d}u}{\mathrm{d}x} + \frac{h^2}{2} \frac{\mathrm{d}^2 u}{\mathrm{d}x^2} + \frac{h^3}{6} \frac{\mathrm{d}^3 u}{\mathrm{d}x^3} + \frac{h^4}{24} \frac{\mathrm{d}^4 u}{\mathrm{d}x^4} + \mathcal{O}(h^5)$$

$$u_{i-1} = u_i - h \frac{\mathrm{d}u}{\mathrm{d}x} + \frac{h^2}{2} \frac{\mathrm{d}^2 u}{\mathrm{d}x^2} - \frac{h^3}{6} \frac{\mathrm{d}^3 u}{\mathrm{d}x^3} + \frac{h^4}{24} \frac{\mathrm{d}^4 u}{\mathrm{d}x^4} + \mathcal{O}(h^5).$$
(4)

$$u_{i-1} = u_i - h \frac{\mathrm{d}u}{\mathrm{d}x} + \frac{h^2}{2} \frac{\mathrm{d}^2 u}{\mathrm{d}x^2} - \frac{h^3}{6} \frac{\mathrm{d}^3 u}{\mathrm{d}x^3} + \frac{h^4}{24} \frac{\mathrm{d}^4 u}{\mathrm{d}x^4} + \mathcal{O}(h^5). \tag{4}$$

By adding these equations, we obtain an estimate for  $d^2u/dx^2$  in the point  $x_i$ ,

$$\frac{\mathrm{d}^2 u}{\mathrm{d}x^2}(x_i) \approx \frac{1}{h^2} \left[ u_{i+1} - 2u_i + u_{i-1} \right]. \tag{5}$$

This method provides an approximation of order  $\mathcal{O}(h^2)$ , since the local error is

$$-\frac{h^2}{12}\frac{\mathrm{d}^4 u}{\mathrm{d}x^4}(x_i).$$

When we refine our grid, and let  $h \to 0$ , the error in our approximation to the exact solution u goes to 0 as rapidly as  $h^2$ . Of course, the computational effort necessary will increase as we make our grid smaller.

## 1.1 Five-point formula

We can easily expand the finite difference approach to two dimensions, by approximating  $\nabla^2 u =$  $d^2u/dx^2 + d^2u/dy^2$  as

$$\nabla^2 u(x,y) \approx \frac{1}{h^2} \left[ u(x+h,y) + u(x-h,y) + u(x,y+h) + u(x,y-h) - 4u(x,y) \right]. \tag{6}$$

This equation is known as the five-point formula, since it incorporates the value of u at five different points in the plane. An arrangement of points used for the approximation of a differential equation is called a stencil. The stencil used in the five-point method is shown in Figure 1 and is named the five-point stencil.

The five-point formula can be modified to account for an irregular grid. However, this method is often only order  $\mathcal{O}(h)$  accurate, where h is a bound on the stepsize. This method is commonly used at the boundary of the domain. For small h, one can shift the boundary points such that the regular five-point method can be used. The errors introduced by this transformation have been shown to be no greater than those introduced by the use of an irregular grid.

## Linear System

In the two dimensional case, we transform the Helmholtz problem to

$$\frac{1}{h^2}\left[u_{i+1,j} + u_{i-1,j} + u_{i,j+1} + u_{i,j-1} - 4u_{ij}\right] + f_{ij}u_{ij} = g_{ij},\tag{7}$$

alternatively written as

$$-u_{i+1,j} - u_{i-1,j} - u_{i,j+1} - u_{i,j-1} + (4 - h^2 f_{ij})u_{ij} = -h^2 g_{ij},$$
(8)

where

$$u_{ij} = u(ih, jh),$$
  $f_{ij} = f(ih, jh)$  and  $g_{ij} = g(ih, jh).$ 

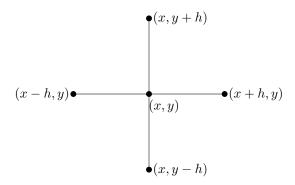


Figure 1: Five-point stencil.