## Mean first passage time of active particles

## The SDE

The SDEs governing the dynamics of the particle are

$$\frac{d}{dt}\vec{r}(t) = \vec{\xi}(t) + \nu_0(r)\vec{p}(t) \tag{1}$$

$$\frac{d}{dt}\vec{p}(t) = \vec{\eta}(t) \times \vec{p}(t), \tag{2}$$

with  $\langle \vec{\xi}(t) \vec{\xi}(t') \rangle = 2D_t \delta(t - t') \mathbb{1}$ . Averaging out the orientational degrees of freedom results in a course grained equation for the position:

$$\frac{d}{dt}\vec{r}(t) = \vec{\xi}(t) + \vec{\chi}(t). \tag{3}$$

The noise due to the course graining of the equation is approximately Gaussian with exponential time correlations,  $\langle \vec{\chi}(t) \vec{\chi}(t') \rangle = \frac{1}{3} \nu_0(r)^2 e^{-2D_r|t-t'|} \mathbb{1} = \frac{1}{6D_r \tau} \nu_0(r)^2 e^{-|t-t'|/\tau} \mathbb{1}$  with  $\tau = \frac{1}{2D_r}$ .

## The Fokker Planck Equation

Equation 3 has colored noise, but an approximate FPE can be obtained using Fox's method:

$$\partial_t P(\vec{r}, t) = D_t \nabla^2 P(\vec{r}, t) + \frac{1}{6D_r} \vec{\nabla} \cdot \left[ \nu_0(r) \vec{\nabla} \left( \nu_0(r) P(\vec{r}, t) \right) \right] \tag{4}$$

$$=-\frac{1}{6D_r}\vec{\nabla}\cdot\left[\left(\vec{\nabla}v_0(r)\right)v_0(r)P(\vec{r},t)\right] + \nabla^2\left[\left(D_t + \frac{v_0(r)^2}{6D_r}\right)P(\vec{r},t)\right]. \tag{5}$$

Terms of order  $\frac{1}{D_z^2}$  and higher are neglected.

The adjoint of the FP operator is

$$L^{\dagger} = \frac{1}{6D_r} \nu_0(r) \nu_0'(r) \hat{r} \cdot \vec{\nabla} + \left( D_t + \frac{\nu_0(r)^2}{6D_r} \right) \nabla^2, \tag{6}$$

with  $\vec{\nabla} v_0(r) = \hat{r} v_0'(r)$ .

## The mean first passage time

The mean first passage time  $(\tau(r))$  is governed by the adjoint of the FP operator:

$$\frac{1}{6D_r}v_0(r)v_0'(r)\hat{r}\cdot\vec{\nabla}\tau(r) + \left(D_t + \frac{v_0(r)^2}{6D_r}\right)\nabla^2\tau(r) = -1,\tag{7}$$

with boundary conditions

$$\tau(a) = 0$$
 and  $\hat{r} \cdot \vec{\nabla} \tau(r) \Big|_{r=R} = \partial_r \tau(r) \Big|_{r=R} = 0,$  (8)

where a is the radius of the inner, absorbing shell, and R is the radius of the outer, reflecting shell. Using the gradient and divergence in spherical coordinates eq. 7 becomes

$$-1 = \frac{1}{6D_r} v_0(r) v_0'(r) \frac{\partial}{\partial r} \tau(r) + \left( D_t + \frac{v_0(r)^2}{6D_r} \right) r^{-2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial}{\partial r} \tau(r) \right)$$
(9)

$$= \left[ \frac{1}{6D_r} \nu_0(r) \nu_0'(r) + \frac{2}{r} \left( D_t + \frac{\nu_0(r)^2}{6D_r} \right) \right] \frac{\partial}{\partial r} \tau(r) + \left( D_t + \frac{\nu_0(r)^2}{6D_r} \right) \frac{\partial^2}{\partial r^2} \tau(r). \tag{10}$$

The derivatives with respect to  $\theta$  and  $\phi$  are zero. Equation 9 can be written as

$$f(r)\frac{\partial}{\partial r}\tau(r) + \frac{\partial^2}{\partial r^2}\tau(r) = -g(r), \tag{11}$$

with  $f(r) = \frac{v_0(r)v_0'(r)}{6D_rD_r + v_0(r)^2} + 2r^{-1}$  and  $g(r) = \frac{6D_r}{6D_rD_t + v_0(r)^2}$ . Equation 11 can be written as a full derivative.

$$\frac{\partial}{\partial r} \left( exp \left[ \int_{h}^{r} dx f(x) \right] \frac{\partial}{\partial r} \tau(r) \right) = -g(r) exp \left[ \int_{h}^{r} dx f(x) \right]. \tag{12}$$

Integrating both sides from r to R and using the boundary condition  $\partial_r \tau(r)\Big|_{r=R} = 0$  gives

$$\frac{\partial}{\partial r}\tau(r) = \int_{r}^{R} dy \ g(y) \ exp\left[\int_{r}^{y} dx f(x)\right],\tag{13}$$

and integrating from a to r gives the equation for  $\tau(r)$ :

$$\tau(r) = \int_{a}^{r} dz \int_{z}^{R} dy \ g(y) \ exp\left[\int_{z}^{y} dx f(x)\right]. \tag{14}$$

Using the definition of f(r) the exponent can be calculated.

$$\int_{z}^{y} dx f(x) = \int_{z}^{y} dr \frac{\nu_0(r)\nu_0'(r)}{6D_r D_r + \nu_0(r)^2} + 2r^{-1}$$
(15)

$$= \frac{1}{2} \int_{\nu_0(z)}^{\nu_0(y)} d\left[\nu_0(r)^2\right] \frac{1}{6D_r D_r + \nu_0(r)^2} + 2 \int_z^y dr \frac{1}{r}$$
 (16)

$$= \frac{1}{2} ln \left( \frac{6D_r D_r + \nu_0(y)^2}{6D_r D_r + \nu_0(z)^2} \right) + 2 ln \left( \frac{y}{z} \right)$$
 (17)

$$exp\left[\int_{z}^{y} dx f(x)\right] = \frac{y^{2}}{z^{2}} \sqrt{\frac{6D_{r}D_{r} + \nu_{0}(y)^{2}}{6D_{r}D_{r} + \nu_{0}(z)^{2}}}$$
(18)

Plugging this in equation 14 gives

$$\tau(r) = \int_{a}^{r} dz \int_{z}^{R} dy \frac{6D_{r}}{6D_{r}D_{t} + \nu_{0}(y)^{2}} \frac{y^{2}}{z^{2}} \sqrt{\frac{6D_{r}D_{r} + \nu_{0}(y)^{2}}{6D_{r}D_{r} + \nu_{0}(z)^{2}}}$$
(19)

$$= \int_{a}^{r} dz \, z^{-2} \left( D_{r} + \frac{\nu_{0}(z)^{2}}{6D_{R}} \right)^{-\frac{1}{2}} \int_{z}^{R} dy \, y^{2} \left( D_{r} + \frac{\nu_{0}(y)^{2}}{6D_{R}} \right)^{-\frac{1}{2}}$$
 (20)