

DOWNFOLDING RECAST AS A COMPRESSION OF THE ENERGY FUNCTIONAL

Theory

Suppose we start with a quantum system with Hamiltonian H and Hilbert space \mathcal{H} .

Definition 1. $E[\Psi] = \frac{\langle \Psi | H | \Psi \rangle}{\langle \Psi | \Psi \rangle}$

Definition 2. Let \mathcal{H} be a Hilbert space. Then $\mathcal{LE}(H, N)$ is a subset of \mathcal{H} spanned by N vectors given by the lowest energy solutions to $H |\Phi_i\rangle = E_i |\Phi_i\rangle$.

Definition 3. H_{eff} is an operator on the Hilbert space $\mathcal{LE}(\mathcal{H}, N)$. **LKW: Does this need to be linear?**

Definition 4. The effective model $E_{eff}[\Psi] = \frac{\langle \Psi | H_{eff} | \Psi \rangle}{\langle \Psi | \Psi \rangle}$ is a functional from $\mathcal{LE} \rightarrow \mathbb{R}$

Theorem 1. $E[\Psi]$ has a critical point only where Ψ is an eigenstate of H .

Proof. This has been proven already by **LKW: find citation** □

Theorem 2. If $E_{eff}[\Psi] = E[\Psi], \forall |\Psi\rangle \in \mathcal{LE}$, then $H_{eff} |\Phi_i\rangle = E_i |\Phi_i\rangle$ for all eigenstates $|\Phi_i\rangle \in \mathcal{LE}$.

Proof. Suppose that $|\Phi_i\rangle$ is an eigenstate of H . Then

$$\left. \frac{\partial E[\Psi]}{\partial \Psi} \right|_{\Psi=\Phi_i} = 0 \quad (1)$$

LKW: We have to project the derivative into the low-energy subspace. If the functional derivative of E is zero at a point, then certainly the functional derivative of E_{eff} is zero. Then

$$\left. \frac{\partial E_{eff}[\Psi]}{\partial \Psi} \right|_{\Psi=\Phi_i} = 0 \quad (2)$$

for $|\Psi\rangle \in \mathcal{LE}$. Using Theorem 1, $H_{eff} |\Phi_i\rangle = E_i |\Phi_i\rangle$. Similarly, if the derivative is not zero, then the wave function must not be an eigenstate. □

We have thus reduced the problem of finding an effective Hamiltonian H_{eff} that reproduces the low-energy spectrum of H to matching the corresponding energy functionals $E[\Psi]$ and $E_{eff}[\Psi]$. An important result of this is that it is not necessary to diagonalize either of the Hamiltonians; one must only be able to select wave functions from the low-energy space \mathcal{LE} . As we shall see, this can be substantially easier than attaining full eigenstates.

The theory presented above maps coarse-graining into a functional approximation problem. This is still rather intimidating, since even supposing one can generate wave functions in the low-energy space, they are still complicated objects in a very large space. An effective way to accomplish this is through the use of descriptors, $d_i[\Psi]$, which map from $\mathcal{H} \rightarrow \mathbb{R}$. Then we can approximate the energy functional as follows

$$E_{eff}[\Psi] \simeq \sum_i f_i(d_i[\Psi]), \quad (3)$$

where f_i are some parameterized functions. **LKW: This is a compression operation.**

An example

LKW: Simple example: Hubbard \rightarrow Heisenberg?