# **EECS 16B Midterm 1 Review Session**

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#### **Disclaimer**

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Slides are also posted at @1046 on Piazza.

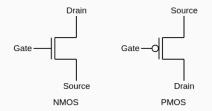
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### **HKN Drop-In Tutoring**

- HKN has office hours Monday, Wednesday, and Friday from 1
   PM 3 PM and 8 PM 10 PM on hkn.mu/ohqueue
- The schedule of tutors can be found at hkn.mu/tutor

# **CMOS Transistors and Logic**

#### **Transistors**

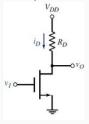


Two varieties of MOSFETs: P-type and N-type Any MOSFET has a characteristic threshold voltage  $V_{th}$  NMOS "turns on" (connects drain to source) when  $V_{GS} > V_{th}$  PMOS turns on when  $V_{GS} < V_{th}$ 

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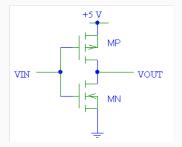
### **NMOS Logic**

We can build an inverter (a circuit that flips a 1 to a 0 and vice versa) with a single N-type MOSFET!



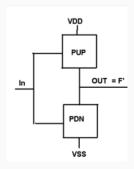
What is  $v_O$  when  $v_I > V_{th}$ ? When  $v_I < V_{th}$ ? Key disadvantage: what is the power dissipated when  $v_I > V_{th}$ ? How can we rectify this?

### **CMOS Logic**



Now, we have the same logical function as the NMOS inverter, but we're using more transistors and much less power is dissipated. Why?

# **CMOS Logic**



This is broadly called Complementary Metal Oxide Semiconductor (CMOS) logic, using a Pull-Up Network of P-type and a Pull-Down Network of N-type MOSFETs.

Now we can build circuits that perform logical functions out of MOSFETs!

#### True or False:

- 1. Power (in the EE16B model) is dissipated in a CMOS circuit only when there is switching
- 2. NMOS devices turn on with a large VGS and off with low VGS
- For a given choice of inputs in a CMOS inverter circuit, we can create a low resistance path from ground to high voltage

#### True or False:

1. Power (in the EE16B model) is dissipated in a CMOS circuit only when there is switching  $\frac{1}{2}$ 

#### True

- 2. NMOS devices turn on with a large VGS and off with low VGS
- 3. For a given choice of inputs in a CMOS inverter circuit, we can create a low resistance path from ground to high voltage

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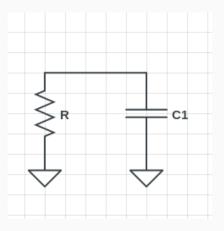
- 2. NMOS devices turn on with a large VGS and off with low VGS **True**
- For a given choice of inputs in a CMOS inverter circuit, we can create a low resistance path from ground to high voltage False

# **RC Circuits**

#### **RC Circuits**

The capacitor and resistor in a NOT circuit form the most basic RC circuit:

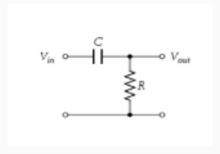
Write down a differential equation describing the circuit below:



#### **RC Circuits**

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Write down a differential equation describing the circuit below:



We have a differential equation describing  $V_{out}$  in terms of  $V_c$ .

$$\frac{dV_c}{dt} = -\frac{1}{RC}V_c$$

How do we actually solve it?

We have a differential equation describing  $V_{out}$  in terms of  $V_c$ .

$$\frac{dV_c}{dt} = -\frac{1}{RC}V_c$$

How do we actually solve it?

Think of the differential operation as a linear operator that scales  $V_c$ , since  $V_c$  is one of its eigenfunctions:

$$\left[\frac{d}{dt}\right]V_c = \lambda V_c$$

Which are the eigenfunctions of differentiation?

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$$Ae^{\lambda t}$$

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$$\frac{dV_c}{dt} = -\frac{1}{RC}V_c$$

The solution to our first order differential equations is therefore:

$$V_c(t) = V_c(0)e^{\frac{-1}{RC}t}$$

# RC Differential Equation: Non-homogenous case

How do you solve a RC circuit with a voltage source?



Applying KCL at the top right node, along with Ohm's law and the capacitor relationship, we get:

$$\frac{dV_c}{dt} = \frac{1}{VC}(V_s - V_c)$$

We can't easily solve this equation, so we change variables to

$$x = V_c - V_s$$

# RC Differential Equation: Non-homogenous case

Now, we have  $\frac{dx}{dt} = -\frac{x}{RC}$ 

We already know how to solve this differential equation, and we get

$$x(t) = x_0 e^{-\frac{t}{RC}}$$

Finally, change back to the original variables by substituting  $V_c$  -  $V_s$  for x

$$V_c(t) = V_c(0)e^{-\frac{t}{RC}} + V_s(1 - e^{-\frac{t}{RC}})$$

# Change of Basis

# **Change of Basis**

In the standard basis, we write vectors as a linear combination of the standard basis vectors.

$$\vec{x} = \begin{bmatrix} x_1 \\ x_2 \end{bmatrix} = x_1 \begin{bmatrix} 1 \\ 0 \end{bmatrix} + x_2 \begin{bmatrix} 0 \\ 1 \end{bmatrix}$$

Likewise, we can write our x vector as a linear combination of some other basis vectors. For example, if our V-basis has basis vectors  $\vec{v_1}$  and  $\vec{v_2}$ :

$$\vec{x} = \tilde{x_1} \cdot \vec{v_1} + \tilde{x_2} \cdot \vec{v_2} = V\vec{\tilde{x}}$$

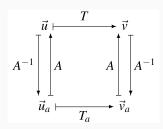
We can go from the V-basis to the standard basis by applying the V-matrix to x, and go from the standard basis to the V-basis by applying  $V^{-1}$ .

# Change of Basis Diagram

- When you are doing change of basis for systems of equations, it is useful to use a diagram mapping transformations between basis
- Given T, to find T<sub>a</sub> you would trace the path from u<sub>a</sub> to v<sub>a</sub>, applying each transformation to the left of the previous:

$$A \Longrightarrow TA \Longrightarrow A^{-1}TA$$
 $T_a = A^{-1}TA$ 

• Step by step, we have  $Au_a = u \implies TAu_a = v \implies A^{-1}TAu_a = v_a$ 



# Diagonalization

# Diagonalization

- The idea is that we want to change into a basis in which the system  $A\vec{x} = \vec{y}$  is represented by a diagonal matrix. So how do we find such basis?
- Remember that for all eigenvalue-eigenvector pairs we have:  $A\vec{v} = \lambda \vec{v}$
- Let's use our eigenvectors as our basis. Doing so we obtain:

$$\vec{x} = V\tilde{\vec{x}}, \vec{y} = V\tilde{\vec{y}}, \text{and} \Lambda \tilde{\vec{x}} = \tilde{\vec{y}}$$

Where the upper-case lambda represents the diagonal matrix with eigenvalues on the diagonals.

• Transforming back to the standard basis, we get:

$$A = V \Lambda V^{-1}$$

### Diagonalization Cont.

- Let's analyze this a bit further, why is it important to be able to do this?
- We see that diagonalizing the matrix makes the system much easier to solve, why?
- Also, we see that there is a "home state" for every system of linearly independent equations, i.e. the space in which the system's components are uncoupled.

2 Minute Break!

# **Differential Equations**

#### **Solving Differential Equations**

Differentiation is linear!

$$\frac{d(c_1x(t)+c_2y(t))}{dt}=c_1\frac{dx(t)}{dt}+c_2\frac{dy(t)}{dt}$$

# Solving systems of differential equations

Write the system in the following form:

$$\frac{dx(t)}{dt} = Ax(t)$$

Find the eigenvalues/eigenvectors of A and transform the system into the eigenbasis. If A has distinct eigenvalues, the general solution is:

$$x(t) = c_1 e^{\lambda_1 t} \vec{v_1} + c_2 e^{\lambda_2 t} \vec{v_2}$$

Then, change back to the standard basis.

Solve the following differential equation:

$$\ddot{x} - \dot{x} - 2x = 0$$

$$x(0) = 2, \dot{x}(0) = 1$$

$$\ddot{x} - \dot{x} - 2x = 0$$
,  $x(0) = 2$ ,  $\dot{x}(0) = 1$ 

Solution:

Write in matrix form and find eigenvectors:

$$\begin{bmatrix} \dot{x} \\ \ddot{x} \end{bmatrix} = \begin{bmatrix} 0 & 1 \\ 2 & 1 \end{bmatrix} \begin{bmatrix} x \\ \dot{x} \end{bmatrix}$$

$$\lambda_1 = 2, v_1 = \begin{bmatrix} 1 \\ 2 \end{bmatrix}, \lambda_2 = -1, v_2 = \begin{bmatrix} -1 \\ 1 \end{bmatrix}$$

Use general form to solve:

$$\begin{bmatrix} x \\ \dot{x} \end{bmatrix} = c_1 e^{2t} v_1 + c_2 e^{-t} v_2$$

$$\ddot{x} - \dot{x} - 2x = 0$$
,  $x(0) = 2$ ,  $\dot{x}(0) = 1$ 

Solution:

Now plug in initial conditions:

$$\begin{bmatrix} 2 \\ 1 \end{bmatrix} = \begin{bmatrix} x(0) \\ \dot{x}(0) \end{bmatrix} = c_1 v_1 + c_2 v_2 = \begin{bmatrix} c_1 - c_2 \\ 2c_1 + c_2 \end{bmatrix}$$

Solving, we get 
$$c_1 = 1, c_2 = -1$$

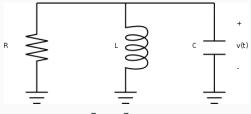
so 
$$x(t) = e^{2t} + e^{-t}$$

$$\ddot{x} - \dot{x} - 2x = 0$$
,  $x(0) = 2$ ,  $\dot{x}(0) = 1$ 

Solution:

Sanity check: compute derivatives and check that the original system is satisfied

$$x(t) = e^{2t} + e^{-t}, \dot{x}(t) = 2e^{2t} - e^{-t}, \ddot{x}(t) = 4e^{2t} + e^{-t}$$
$$\implies \ddot{x} - \dot{x} - 2x = 0, x(0) = 2, \dot{x}(0) = 1$$



Given 
$$x(t) = \begin{bmatrix} v(t) \\ i_L(t) \end{bmatrix}$$

find matrix A such that  $\frac{dx}{dt} = Ax$ 

Direct all currents into ground.

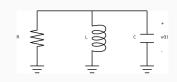
$$i_R + i_L + i_C = 0$$

$$\frac{V}{R} + i_L = -C \frac{dV}{dt}$$

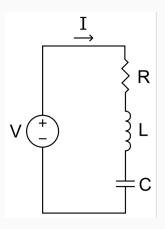
$$-\frac{1}{CR}V + -\frac{1}{C}i_L = \frac{dV}{dt}$$

Also

$$L\frac{di_L}{dt} = V \implies A = \begin{bmatrix} -\frac{1}{RC} & -\frac{1}{C} \\ \frac{1}{L} & 0 \end{bmatrix}$$



At t < 0, the circuit is at steady state. At  $t \geq 0$ , the voltage source is set to 0. Find a differential equation for  $i_L$  for  $t \geq 0$ .

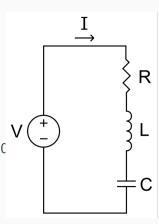


Solution: use KCL/KVL to get

$$i_R = i_L = i_C, V_C + V_R + V_L = 0$$

$$\frac{V_R}{R} = i_L = C \frac{dV_C}{dt}, \frac{1}{C} i_L + R \frac{di_L}{dt} + L \frac{d^2 i_L}{dt} = 0$$

$$\begin{bmatrix} \frac{di_L}{dt} \\ \frac{d^2i_L}{dt} \end{bmatrix} = \begin{bmatrix} 0 & 1 \\ -\frac{1}{LC} & -\frac{R}{L} \end{bmatrix} \begin{bmatrix} i_L \\ \frac{di_L}{dt} \end{bmatrix}$$



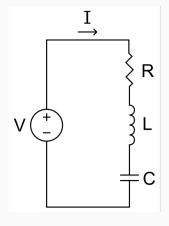
If R = 0, L = C = 1, solve the differential equation:

Initial conditions:  $i_L(0) = 0$ ,

$$\frac{di_L}{dt}(0) = -V_c = -V$$

Differential equation:

$$\begin{bmatrix} \frac{di_L}{dt} \\ \frac{d^2i_L}{dt} \end{bmatrix} = \begin{bmatrix} 0 & 1 \\ -\frac{1}{LC} & -\frac{R}{L} \end{bmatrix} \begin{bmatrix} i_L \\ \frac{di_L}{dt} \end{bmatrix}$$
$$\lambda_1 = i, v_1 = \begin{bmatrix} i \\ -1 \end{bmatrix}$$
$$\lambda_2 = -i, v_2 = \begin{bmatrix} i \\ 1 \end{bmatrix}$$



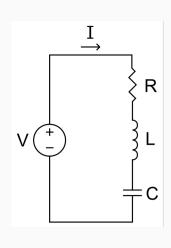
$$\begin{bmatrix} i_L \\ \frac{di_L}{dt} \end{bmatrix} = c_1 e^{it} v_1 + c_2 e^{-it} v_2$$

$$ic_1 + ic_2 = 0, -c_1 + c_2 = -V$$

$$\implies c_1 = \frac{V}{2}, c_2 = -\frac{V}{2}$$

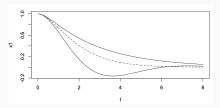
SO

$$i_L = \frac{V}{2}ie^{it} - \frac{V}{2}e^{-it}$$



## What do we do with different types of eigenvalues?

Each interesting "case" of eigenvalues (real, imaginary, and repeated) is complementary to a case of resonance in 2d ODEs, so let's discuss them together:



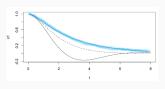
## Overdamping: two real eigenvalues

 When a system has two real eigenvalues, it will have two real eigenvectors, so the solution will be a combination of exponentials in the form

$$ae^{\lambda_1 t} + be^{\lambda_2 t}$$

For negative lambdas, this produces exponential decay.

At large t, greater eigenvalue dominates behavior, so how fast the system approaches to 0 is determined by the larger eigenvalue.



## Underdamping: imaginary or complex eigenvalues

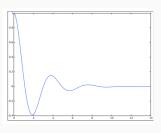
You can still solve the system of differential equations as usual, but your solutions will be complex:

$$(a_1j + b_1)e^{c_1j+d_1} + (a_2j + b_2)e^{c_2j+d_2}$$

 You can rearrange terms and apply Euler's formula to write each term as a real exponential multiplied by a sinusoid

$$e^{d_1}(\alpha_1 sin(...) + \beta_1 cos(...)) + e^{d_2}(\alpha_2 sin(...) + \beta_2 cos(...))$$

If  $d_1$  and  $d_2$  are negative, then you can think of this as a sinusoid where the amplitude decays to 0



.

## Critical Damping: repeated eigenvalue

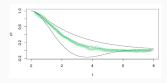
- You only get one linearly independent eigenvector.
- So, arbitrarily choose the second column of your V matrix (making sure it is linearly independent from the eigenvector v<sub>1</sub>)
- Compute  $A_v$  using the change of basis diagram  $A_v = V^{-1}AV$
- This gives you an upper-triangular matrix
- The bottom row of the matrix gives you an equation in one variable (only depends on the second component of  $x_v$ ), which you can solve to get something in the form of  $x_{v,2} = ae^{\lambda t}$

## Critical Damping: repeated eigenvalue

- The top row of the matrix depends on both components of  $x_{v}$ , but you can plug in the value you got for  $x_{v,2}$  and then solve for  $x_{v,1}$
- If you solve this differential equation (which is a diff eq with a non-constant input\*), you'll get an equation in the form:

$$x(t) = c_1 e^{-\lambda t} + c_2 t e^{-\lambda t}$$

The critically damped case gives us the fastest-decaying exponential that doesn't oscillate



\*You can solve this using the formula  $x_D(t) = x_0 e^{\lambda t} + \int_0^t u(\tau) e^{\lambda(t-\tau)} d\tau$ 

2 min break

## **Phasors**

#### **Phasors**

- Phasors express the response of a circuit to a sinusoidal input
- Any real, periodic signal can be expresses as the sum of sinusoids - we can apply this technique very broadly!
- A phasor encodes information about amplitude and phase, but not frequency

$$v_s = A\cos(\omega t + \phi) \rightarrow \tilde{V}_s = \frac{A}{2}e^{j\phi}$$

#### **Impedance**

- Impedance (Z) is a generalized form of resistance expresses a component's response to an alternating current
- Each common passive circuit element (resistor, capacitor, and inductor) has a characteristic impedance

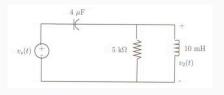
$$Z_R = R$$

$$Z_C = \frac{1}{j\omega C}$$

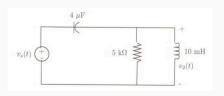
$$Z_L = j\omega L$$

## **Phasor Analysis Procedure**

- 1. Express all time domain signals as cosines
- 2. Convert voltages, currents, and impedances to their phasor equivalents
- 3. Set up phasor equations and solve for unknowns
- 4. Transform back to time domain

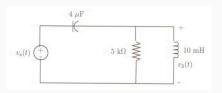


Let 
$$v_s(t) = 5\cos(5000t + \frac{\pi}{4})$$
. Find  $v_2(t)$ .



#### 1. Translate to phasor domain

- $\omega = 5000$
- $v_s(t) = 5\cos(5000t + \frac{\pi}{4}) \rightarrow \tilde{V}_s = \frac{5}{2}e^{j\frac{\pi}{4}}$
- $C = 4\mu F \rightarrow Z_C = \frac{1}{j(4*10^{-6})(5000)} = -j50$
- $R = 5k\Omega \rightarrow Z_R = 5*10^3$
- $L = 10mH \rightarrow Z_L = j(5000)(10*10^{-3}) = j50$

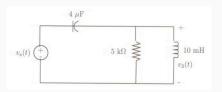


## 2. Solve for $\tilde{V}_2$

• Set up as impedance divider:

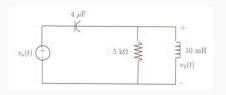
$$\frac{\tilde{V}_{s}}{Z_{C} + (Z_{R} \| Z_{L})} = \frac{\tilde{V}_{2}}{Z_{R} \| Z_{L}}$$

$$\tilde{V}_{2} = \tilde{V}_{s} \frac{Z_{R} \| Z_{L}}{Z_{C} + (Z_{R} \| Z_{L})} = \tilde{V}_{s} \frac{\frac{Z_{R} Z_{L}}{Z_{R} + Z_{L}}}{\frac{Z_{R} Z_{L} + Z_{R} Z_{C} + Z_{L} Z_{C}}{Z_{R} + Z_{L}}}$$



## 2. Solve for $\tilde{V}_2$

$$\tilde{V}_2 = \tilde{V}_s \frac{Z_R Z_L}{Z_R Z_L + Z_R Z_C + Z_L Z_C} = \tilde{V}_s \frac{j2.5 * 10^5}{j2.5 * 10^5 - j2.5 * 10^5 + 2500}$$
$$\tilde{V}_2 = j100 \, \tilde{V}_s = 100 \, \tilde{V}_s e^{j\frac{\pi}{2}}$$



#### 3. Find time-domain solution

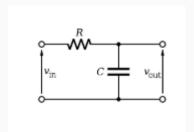
• 
$$\tilde{V}_2 = 100e^{j\frac{\pi}{2}}\tilde{V}_s \rightarrow v_2(t) = 100v_s(t + \frac{\pi}{2}) = 500\cos(5000t + \frac{3\pi}{4})$$

#### **Transfer Function Review**

A *transfer function* of a circuit or system describes the output response to an input excitation as a function of the angular frequency  $\omega$ .

$$H(\omega) = rac{V_{out}(\omega)}{V_{in}(\omega)}$$

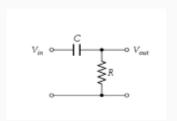
#### **RC Circuits**



$$H(\omega) = \frac{1}{1+j\omega RC}$$

#### Low pass:

$$\lim_{\omega \to \infty} H(\omega) = 0, H(0) = 1$$



$$H(\omega) = \frac{j\omega RC}{1+j\omega RC}$$

## High pass:

$$\lim_{\omega\to\infty}H(\omega)=1, H(0)=0$$

#### **Bode Plot Review**

$$V[dB] = 20\log_{10}(\frac{V}{V_0})$$

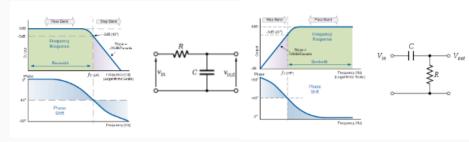
- Bode plots provide a way for us to easily visualize the output response of our system, depending on the input frequency  $\omega$ .
- Because Bode plots are in log scale for  $\omega$ , we are able to take advantage of the properties of logs.
  - $G = XY \implies G[dB] = X[dB] + Y[dB]$
  - $G = \frac{X}{Y} \implies G[dB] = X[dB] Y[dB]$

Hence, we can break our transfer function  $H(\omega)$  into a product of familiar transfer functions (simple poles, quadratic zeros, etc. - "functional forms"), graph them out individually, and then add them together on the graph.

#### **Bode Plot Review**

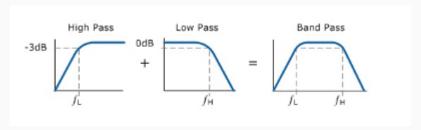
- We want to be able to plot the transfer function as a function of the frequency  $\omega$ .
- However, it's a complex valued function so it's easier to have a plot for the magnitude and one for the phase.
- Because the magnitude plot is in log scale, and the phase plot is in semi-log scale we are able to take advantage of the properties of logs.
- Hence, we can break our transfer function  $H(\omega)$  into a product of familiar transfer functions (simple poles, quadratic zeros, etc. "functional forms"), graph them out individually, and then add them together on the graph.

## RC Circuits, revisited



$$H(\omega) = \frac{1}{1+j\omega RC}$$
  $H(\omega) = \frac{j\omega RC}{1+j\omega RC}$ 

#### **Common Filters and their Bode Plots**

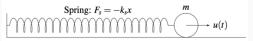


In the high pass filter, the frequencies greater than the corner frequency have a gain of 0dB (so they aren't changed), while frequencies less than the corner frequency have a gain < 0dB (so they are multiplied by something less than 1). In the lowpass filter, the opposite is true.

# **State-Space Representations**

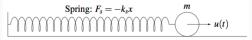
## **State Space Modeling: Example**

Assume we have the following spring system:



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Assume we have the following spring system:



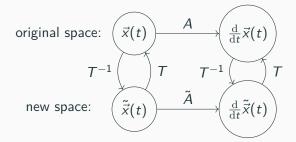
We can model the system as a linear continuous time state space model:

$$\frac{d}{dt}\vec{x}(t) = A\vec{x}(t) + \vec{b}u(t)$$
$$\vec{v}(t) = C\vec{x}(t)$$

in which:

$$\vec{x}(t) = \begin{bmatrix} x(t) \\ v(t) \end{bmatrix}$$
,  $A = \begin{bmatrix} 0 & 1 \\ -\frac{k_s}{m} & 0 \end{bmatrix}$ ,  $\vec{b} = \begin{bmatrix} 0 \\ 1 \end{bmatrix}$ , and  $C = \begin{bmatrix} 0 & 1 \end{bmatrix}$ 

## State Space Modeling:



Always apply operations to right side first

$$\tilde{A} = T^{-1}AT$$

$$\tilde{\vec{x}}(t) = T^{-1}\vec{x}(t)$$

$$\frac{\mathrm{d}}{\mathrm{d}t}\vec{x}(t) = A\vec{x}(t) + \vec{b}u(t)$$

$$\frac{\mathrm{d}}{\mathrm{d}t}\vec{x}(t) = A\vec{x}(t) + \vec{b}u(t)$$

2. Find 
$$\lambda_i$$
 of  $A$ ; let  $\tilde{A} = \begin{bmatrix} \lambda_1 & 0 & \cdots & 0 \\ 0 & \lambda_2 & 0 & \vdots \\ \vdots & 0 & \ddots & 0 \\ 0 & \cdots & 0 & \lambda_n \end{bmatrix}$ 

$$\frac{\mathrm{d}}{\mathrm{d}t}\vec{x}(t) = A\vec{x}(t) + \vec{b}u(t)$$

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3. Find eigenvectors 
$$\vec{v}_i$$
 of  $A$ ; let  $T = \begin{bmatrix} \vec{v_1} & \vec{v_2} & \cdots & \vec{v_n} \end{bmatrix}$ 

$$\frac{\mathrm{d}}{\mathrm{d}t}\vec{x}(t) = A\vec{x}(t) + \vec{b}u(t)$$

2. Find 
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 of  $A$ ; let  $\tilde{A} = \begin{bmatrix} \lambda_1 & 0 & \cdots & 0 \\ 0 & \lambda_2 & 0 & \vdots \\ \vdots & 0 & \ddots & 0 \\ 0 & \cdots & 0 & \lambda_n \end{bmatrix}$ 

- 3. Find eigenvectors  $\vec{v}_i$  of A; let  $T = \begin{bmatrix} \vec{v_1} & \vec{v_2} & \cdots & \vec{v_n} \end{bmatrix}$
- 4. Convert  $\vec{x}(t)$  to  $\tilde{\vec{x}}(t)$  using:  $\tilde{\vec{x}}(t) = T^{-1}\vec{x}(t)$

$$\frac{\mathrm{d}}{\mathrm{d}t}\vec{x}(t) = A\vec{x}(t) + \vec{b}u(t)$$

2. Find 
$$\lambda_i$$
 of  $A$ ; let  $\tilde{A} = \begin{bmatrix} \lambda_1 & 0 & \cdots & 0 \\ 0 & \lambda_2 & 0 & \vdots \\ \vdots & 0 & \ddots & 0 \\ 0 & \cdots & 0 & \lambda_n \end{bmatrix}$ 

- 3. Find eigenvectors  $\vec{v}_i$  of A; let  $T = \begin{bmatrix} \vec{v_1} & \vec{v_2} & \cdots & \vec{v_n} \end{bmatrix}$
- 4. Convert  $\vec{x}(t)$  to  $\tilde{\vec{x}}(t)$  using:  $\tilde{\vec{x}}(t) = T^{-1}\vec{x}(t)$
- 5. Solve  $\frac{d}{dt}\tilde{\vec{x}}(t) = \tilde{A}\tilde{\vec{x}}(t) + T\vec{b}u(t)$

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- 3. Find eigenvectors  $\vec{v}_i$  of A; let  $T = \begin{bmatrix} \vec{v_1} & \vec{v_2} & \cdots & \vec{v_n} \end{bmatrix}$
- 4. Convert  $\vec{x}(t)$  to  $\tilde{\vec{x}}(t)$  using:  $\tilde{\vec{x}}(t) = T^{-1}\vec{x}(t)$
- 5. Solve  $\frac{d}{dt}\tilde{\vec{x}}(t) = \tilde{A}\tilde{\vec{x}}(t) + T\vec{b}u(t)$
- 6. Convert solution back to  $\vec{x}(t)$

## **State Space Modeling:**

Continuous time solution:

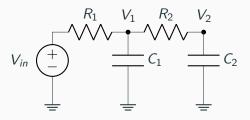
$$\tilde{x}(t) = e^{\lambda t} \tilde{x}(0) + \frac{e^{\lambda t} - 1}{\lambda} u(t) + w(t)$$

Discrete time solution:

$$x_d(i+1) = e^{\lambda \Delta} x_d(i) + \frac{e^{\lambda \Delta} - 1}{\lambda} u(i) + w(i)$$

## **State Space Modeling Example:**

Given the following circuit:



in which  $R_1=2\,\Omega,$   $R_2=\frac{8}{3}\Omega,$   $C_1=1\,C,$   $C_2=\frac{3}{2}C$  solve equations for  $V_1$  and  $V_2$ 

# Stability, Observability, and Controllability

## Stability, Observability, Controllability:

```
given: \vec{x}(i+1) = A\vec{x}(i) + Bu(i) \vec{y}(i) = C\vec{x}(i) in which: \vec{x} is our state, \vec{u} is our input, \vec{y} is what we can observe:
```

## Stability (Discrete time):

Discrete time model: if  $|\lambda_i| < 1$  for all  $\lambda_i$  of A, system is stable intuition: if any  $|\lambda_i| \geqslant 1$ , state vector is increasing each time step will be infinitely magnified over time

## Stability (Continuous time):

Continuous time model:

if the  $\Re\{\lambda_i\}$  < 0 for all  $\lambda_i$  of A are strictly negative, system is stable

intuition: if real part of eigenvalue is positive, state vector is increasing over time and will be infinitely magnified over time

#### Feedback:

```
if system is controllable, we can set: u(t) = K\vec{x}(t) plugging in, we get: \vec{x}(t+1) = (A+BK)\vec{x}(t) we can find the eigenvalues of (A+BK) to check for stability
```

## **Stability Example:**

given the following system:

$$\vec{x}[t+1] = \begin{bmatrix} -5 & 0 \\ 7 & 6 \end{bmatrix} \vec{x}[t] + \begin{bmatrix} 2 \\ -1 \end{bmatrix} u[t]$$

$$\vec{y}[t] = \begin{bmatrix} 1 & 1 \end{bmatrix} \vec{x}[t]$$

## **Stability Check:**

$$\lambda = 6, -5$$

## **Stability Check:**

$$\begin{split} \lambda &= 6, -5 \\ \text{System is unstable} \end{split}$$