Energy profile of lattice vibration at presence of a local drive

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Assuming the normal modes of system to be labeled by (α, q) , with which the equation of motion is written as

$$[\partial_t^2 + K_\alpha(q)]r_\alpha(q,t) = 0 \tag{1}$$

The Green's function in these modes becomes

$$\tilde{G}_{\alpha}(q,\omega) = \frac{1}{-\omega^2 + K_{\alpha}(q)},\tag{2}$$

which satisfies

$$[-\omega^2 + K_{\alpha}(q)]\tilde{G}_{\alpha}(q,\omega) = 1 \tag{3}$$

Next, the ω_0 driving force acting on the center mass at site j=0 is

$$\vec{F}_{ext}(j,\omega) = 2\pi\delta(\omega - \omega_0)\delta_{j,0}(c_x\hat{R}_x + c_y\hat{R}_y)$$
(4)

where (R_x, R_y) is the degree of freedom for the atom at the center of a unit cell. In the eigen-basis, the force can be represented as

$$\vec{F}_{ext}(j,\omega) = 2\pi\delta(\omega - \omega_0) \sum_{\alpha} c_{\alpha} \hat{r}_{\alpha} \int \frac{dq}{2\pi} e^{iqR_j}$$

$$\rightarrow F_{ext,\alpha}(q,\omega) = 2\pi\delta(\omega - \omega_0) c_{\alpha}$$
(5)

where $c_{\alpha} \equiv (c_x \hat{R}_x + c_y \hat{R}_y) \cdot \hat{\alpha}$. Therefore, letting $\vec{r} = r_{\alpha} \hat{\alpha}$,

$$r_{\alpha}(q,\omega) = 2\pi\delta(\omega - \omega_0) \frac{c_{\alpha}}{-\omega_0^2 + K_{\alpha}(q)}$$
(6)

transforming back to real space,

$$r_{\alpha}(j,t) = c_{\alpha} \int \frac{d^2q}{(2\pi)^2} \frac{e^{i(qj-\omega_0 t)}}{-\omega_0^2 + K_{\alpha}(q)}$$
 (7)

Assuming the elastic energy on site j can be written as

$$U(j) = \sum_{\alpha,\beta} r_{\alpha}(j,t) \tilde{K}_{\alpha,\beta} r_{\beta}(j,t)$$
(8)

the time-averaged energy is

$$\langle U(j)\rangle = \frac{1}{2} \sum_{\alpha,\beta} r_{\alpha}(j,0) \tilde{K}_{\alpha,\beta} r_{\beta}(j,0)$$
(9)