

Computational Physics Homework 1

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Problem 1

a)

The code for differentiating one-dimensional functions using single precision forward-, central-, and extrapolated-difference algorithms is available in “differentiation_1D.h”, attached in the folder.

This code returns the derivative with the least relative error, and save the step sizes h , the derivatives $f'(x)$, the absolute total errors e , as well as the relative errors ϵ of every loop in a file. By choosing some pre-recorded functions (presently unable to enter functions from the commander line), of which the code is in “function_input.h”, and enter from the commander line the position x , the initial step size h_0 , the user defined relative error tolerance ϵ_{tol} , and the maximum of loops `max_loop` of calculation, the calculation will continue until reaching the error tolerance or the

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maximum loop.

In this problem, since we want to see the relationship of relative error ϵ vs. step size h , a medium initial step size $h_0 = 0.5$, an unachivable error tolerance $\epsilon_{\text{tol}} = 1e - 8 = 1 \times 10^{-8}$, and a $\text{max_loop} = 20$ were chosen. But in practice, ϵ may sometimes go to zero accidentally and thus end the calculation, so to get the whole $\log(\epsilon) - \log(h)$ plot, the break procedure was temporarily omitted.

b)

The log-log plots of relative error ϵ vs. step size h of the derivatives of $\cos(x)$ and e^x at $x = 0.1, 10$ in the three methods are shown below in Figure 1 and 2 respectively.

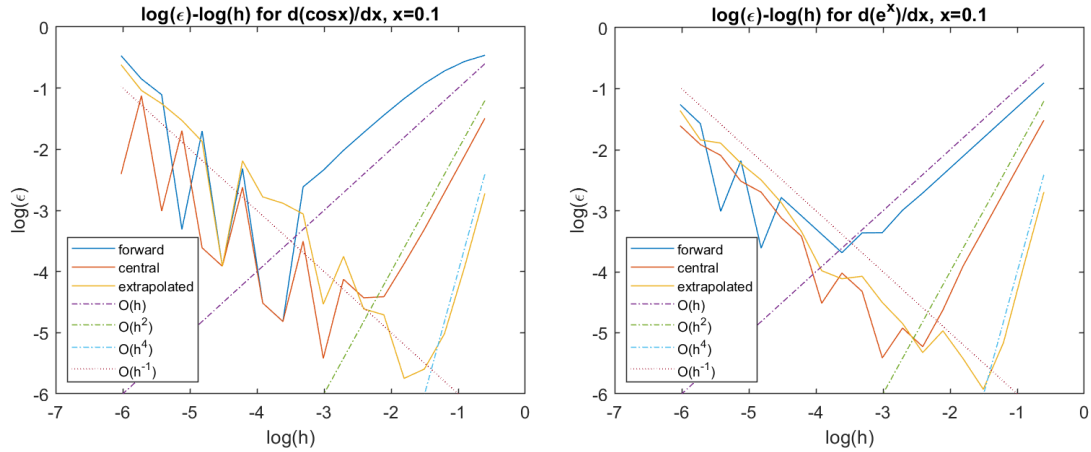


Fig. 1: $\log(\epsilon) - \log(h)$ plots of $\frac{d\cos(x)}{dx}|_{x=0.1}$ and $\frac{de^x}{dx}|_{x=0.1}$. Initial step size $h_0 = 0.5$, calculate for $\text{max_loop} = 20$. $\epsilon_{\text{tol}} = 1e - 8$, but the code was set not to break to see the whole behavior. In these two plots, the error tolerance was not reached, as estimated.

Initial step size is chosen to be $h_0 = 0.5$, calculate for $\text{max_loop} = 20$, relative error tolerance $\epsilon_{\text{tol}} = 1e - 8$, but the code was set not to break to see the whole behavior.

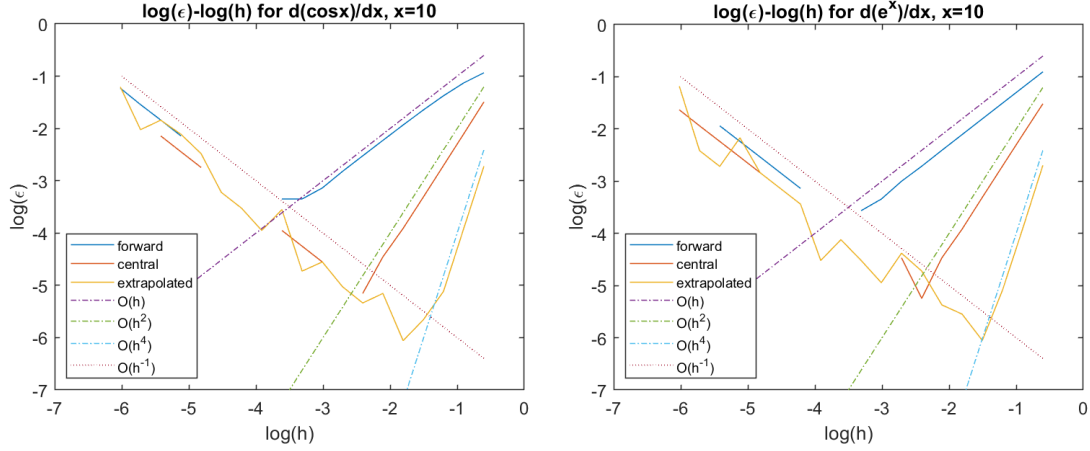


Fig. 2: $\log(\epsilon) - \log(h)$ plots of $\frac{d \cos(x)}{dx}|_{x=10}$ and $\frac{de^x}{dx}|_{x=10}$. Initial step size $h_0 = 0.5$, calculate for $\text{max_loop} = 20$. $\epsilon_{\text{tol}} = 1e-8$, but the code was set not to break to see the whole behavior. However, the tolerance was reached accidentally when some ϵ goes to zero, which correspond to the points missing in the plots.

In Figure 2 of $x = 10$ we can see that, for some step sizes, ϵ accidentally goes to zero, so those points are missing in the plots.

Figure 1 and 2 show that when step size h decreases, before the relative error ϵ goes down to optimal value, $\epsilon \sim h, h^2, h^4$ for forward-, central-, and extrapolated-difference algorithms respectively, which satisfies the theoretical estimates. Below h_{opt} , $\epsilon_r \sim \epsilon_m h^{-1}$ is the leading error, as shown in the plots. Also, it agrees well with the theory that the best relative error for the three methods are $\epsilon_{\text{opt}} \approx 10^{-3.5}, 10^{-5}, 10^{-6}$ respectively, i.e. the number of significant digits obtained matches with the estimates.

c)

As mentioned in b), truncation error ϵ_t manifest itself in the regime $h > h_{\text{opt}}$, and for round-off error ϵ_r is in $h > h_{\text{opt}}$. In the order of the three algorithms above, $h_{\text{opt}} \sim \epsilon_m^{1/2}, \epsilon_m^{1/3}, \epsilon_m^{1/5}$. Since for $\cos(x)$ and e^x at 0.1, 10, $\frac{f}{f''} \approx 1$, so $h_{\text{opt}} \approx 10^{-3.5}, 10^{-2.5}, 10^{-1.5}$, for single precision(machine precision $\epsilon_m \approx 10^{-7}$).

Problem 2

a)

The code for integrating one-dimensional functions using single precision midpoint, trapezoid, and Simpson's rule is available in "integration_1D.h", attached in the folder.

This code returns the integral with the least relative error, and save the number of bins N , the integrals $\int_a^b f(x)dx$, the absolute total errors e , as well as the relative errors ϵ of every loop in a file. By chosing some pre-recorded functions (presently unable to enter functions from the commander line), of which the code is in "function_input.h", and enter from the commander line the lower and the upper bounds a, b , the initial number of bins N_0 , the user defined relative error tolerance ϵ_{tol} , and the maximum of loops max_loop of calculatation, the calculation will continue until reaching the error tolerance or the maximum loop.

In this problem, since we want to see the relationship of relative error ϵ vs. number of bins N , a medium initial number of bins $N_0 = 2$, an unachivable error tolerance $\epsilon_{\text{tol}} = 1e - 8 = 1 \times 10^{-8}$, and a $\text{max_loop} = 20$ were chosen. But in practice, ϵ may

sometimes go to zero accidentally and thus end the calculation, so to get the whole $\log(\epsilon) - \log(h)$ plot, the break procedure was temporarily omitted.

b)

The log-log plots of relative error ϵ vs. number of bins N of $\int_0^1 e^{-t} dt$ in the three methods are shown in Figure 3.

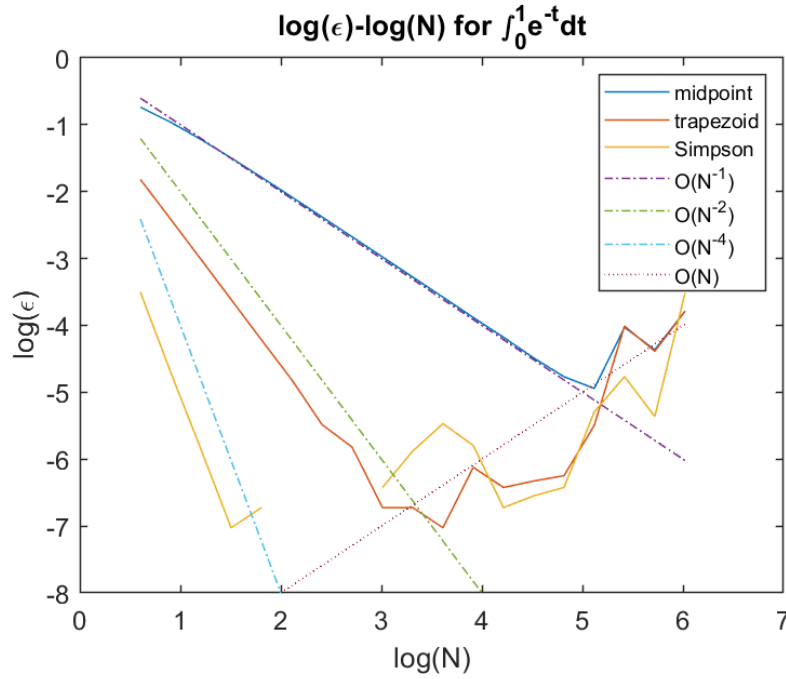


Fig. 3: $\log(\epsilon) - \log(N)$ plots of $\int_0^1 e^{-t} dt$. Initial number of bins $N_0 = 2$, calculate for $\text{max_loop} = 20$. $\epsilon_{\text{tol}} = 1e - 8$, but the code was set not to break to see the whole behavior. However, the tolerance was reached accidentally when some ϵ goes to zero, which correspond to the points missing in the plots.

Initial number of bins is chosen to be $N_0 = 2$, calculate for $\text{max_loop} = 20$, relative error tolerance $\epsilon_{\text{tol}} = 1e - 8$, but the code was set not to break to see the whole

behavior. In Figure 3 we can see that, for some N 's for Simpson's rule, ϵ accidentally goes to zero, so those points are missing in the plots.

c)

Figure 3 shows that when number of bins N decreases, before the relative error ϵ goes down to optimal value, $N < N_{\text{opt}}$, truncation error ϵ_t manifest itself in the regime, so $\epsilon \sim N, N^{-2}, N^{-4}$ for midpoint, trapezoid, and Simpson's rules respectively, which satisfies the theoretical estimates. Above N_{opt} , round-off error is the leading error, as shown in the figures. Similar to differentiation, the optimal number of bins $N_{\text{opt}}(\text{midpoint}) > N_{\text{opt}}(\text{trapezoid}) > N_{\text{opt}}(\text{Simpson's})$, and the best relative error $\epsilon_{\text{opt}}(\text{midpoint}) < \epsilon_{\text{opt}}(\text{trapezoid}) < \epsilon_{\text{opt}}(\text{Simpson's})$, as estimated.

Problem 3

The plot and the log-log plot of the power spectrum $P(k)$ are shown in Figure 4. In the log-log plot, we can clearly see a “baryon wiggle” within $k \in [10^{-1}, 10^{-2}]$.

According to the notes, $\lim_{k \rightarrow 0} P(k) = \lim_{k \rightarrow +\infty} P(k) = 0$. As a good approximation for the whole power spectrum within $k \in [0, +\infty)$, the log-log plots of $P(k)$ for $[0, k_0]$ and $[k_N, +\infty]$ are linear, where N is the number of data points in the file (in this problem is 501), and k_0, k_N are the minimum and maximum wave number in the data. The slopes of the two lines are the same as the derivative of the cubic spline of $\log(P) - \log(k)$ at k_0, k_N .

From the cubic spline result of the $P(k)$ data in the data file provided, we estimate

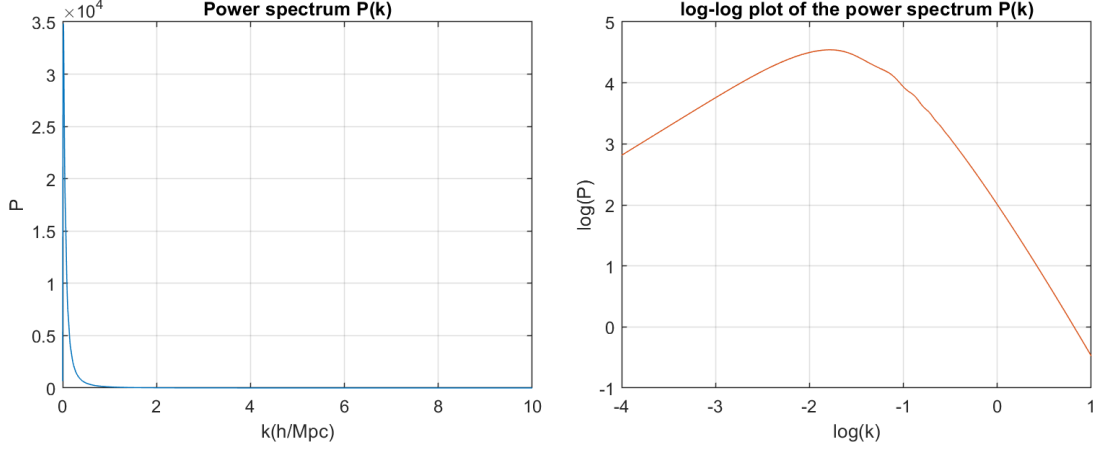


Fig. 4: The plot and the log-log plot of the power spectrum $P(k)$

that

$$P(k)|_{k \in [0, 10^{-4}]} = P(k = 10^{-4}) \cdot \left(\frac{k}{10^{-4}}\right)^{f'_0} = 651.34 \cdot \left(\frac{k}{10^{-4}}\right)^{0.94963} \quad (1)$$

$$P(k)|_{k \in [10, +\infty]} = P(k = 10) \cdot \left(\frac{k}{10}\right)^{f'_N} = 0.33403 \cdot \left(\frac{k}{10}\right)^{-2.6189} \quad (2)$$

which satisfy the boundary condition of the power spectrum. In the data range $[k_0, k_N]$, since the power spectrum is tabulated in logarithmic intervals, assume $x = \log k$, in each interval $x \in [x_i, x_{i+1}] = [\log k_i, \log k_{i+1}]$,

$$P(x)|_{x \in [x_i, x_{i+1}]} = P(x_i) \cdot 10^{a_i(x-x_i)^3 + b_i(x-x_i)^2 + c_i(x-x_i)} \quad (3)$$

here a_i, b_i, c_i are the coefficients of the cubic spline result of the log-log plot of the power spectrum in the i 'th bin $[k_i, k_{i+1}]$

$$\log P(k)|_{k \in [k_i, k_{i+1}]} - \log P(k_i) = a_i(\log k - \log k_i)^3 + b_i(\log k - \log k_i)^2 + c_i(\log k - \log k_i) \quad (4)$$

Now that we have the whole expression of the power spectrum, we can calculate the integral of correlation function in bins.

(1) For $k < k_{\min} = 10^{-4}$, we directly use the given expression

$$\xi_1(r) = \frac{1}{2\pi^2} \int_0^{10^{-4}} dk \cdot k^2 P(k) \frac{\sin(kr)}{kr} \quad (5)$$

using Romberg integration for faster speed. The result of this part

$$\xi_1(r) \approx 8.3547 \times 10^{-12} \quad (6)$$

is quite small relative to the whole integral $\xi(r) > 10^{-4}$ and hardly varies.

(2) For within data range $k \in [k_{\min}, k_{\max}] = [10^{-4}, 10]$, we change the independent variable from k to $x = \log(k)$, new expression

$$\xi_2(r) = \frac{\ln 10}{2\pi^2 r} \int_{10^{-4}}^{10} dx \cdot 10^{2x} P(x) \sin(r \cdot 10^x) = \frac{\ln 10}{2\pi^2 r} \sum_{i=0}^{N-1} \int_{x_i}^{x_{i+1}} dx \cdot 10^{2x} P(x) \sin(r \cdot 10^x) \quad (7)$$

Calculate the integral of each bin using Simpson's rule for stable result.

(3) For $k > k_{\max} = 10$, we still use the direct expression

$$\xi_3(r) = \frac{1}{2\pi^2} \int_{10}^{+\infty} dk \cdot k^2 P(k) \frac{\sin(kr)}{kr} \quad (8)$$

since although for integral over infinite range we can change variable, we will instead get another infinite value and highly oscillatory integral, which is even harder to calculate. Assume $y = 1/k$, that expression is

$$\xi_3(r) = \frac{1}{2\pi^2 r} \int_0^{0.1} dy \cdot \frac{1}{y^3} P\left(\frac{1}{y}\right) \sin\left(\frac{r}{y}\right) \quad (9)$$

The integrated function is highly oscillatory and is not finite when $y \rightarrow 0$. Therefore, we instead use the original expression and set a finite upper limit to get a good approximation of $\xi_3(r)$, since the integral converges. After testing for some $r \in [50, 120]$ in Mathematica, 5000 is chosen to be the upper limit, for fairly good accuracy, fast speed and less error spikes.

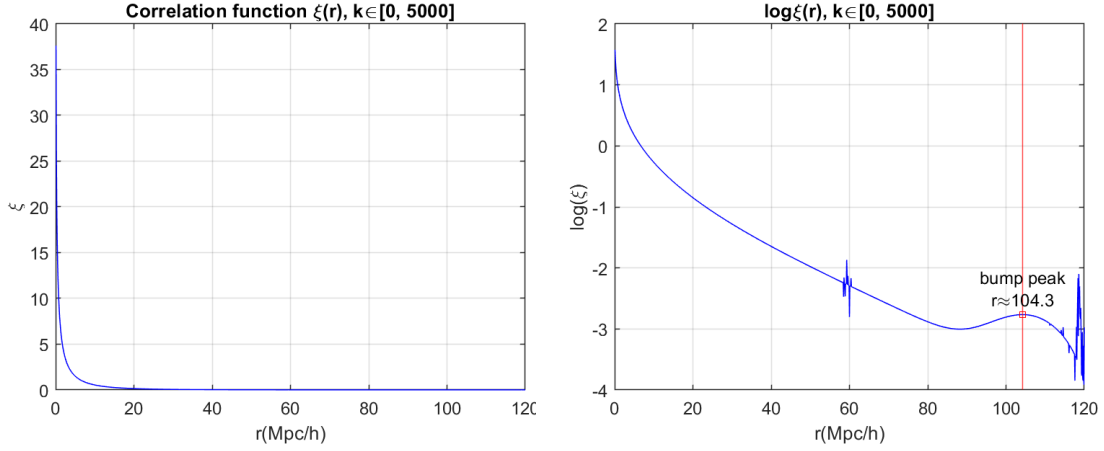


Fig. 5: The plot of the correlation function $\xi(r)$ and $\log \xi(r)$. From the log plot we can see clearer the peak of the bump of $\xi(r)$

Figure 5 is the plot of correlation function $\xi(r)$, integrated over $k \in [0, 500]$. From the plot of $\log \xi(r)$, we can see a clear bump with peak at $r \approx 104.3$. Also from this plot and the output datafile we can see that at around $r = 60, 120$, there are some error spikes in the correlation function, entirely coming from $\xi_3(r)$, maybe because around that scale, the error in each bin of the integral $\xi_3(r)$ tend to be of the same sign and accumulate to a large error, while when away from these scales the errors will cancel thus gives a smooth result.

Figure 6 shows the plot of $r^2 \xi(r)$ over $[0, 120]$ and the required range $[50, 120]$. Although it seems from the output datafile of $r^2 \xi(r)$ that $\xi_3(r)$ is relatively small compared to $\xi_2(r)$, in Figure 7, however, it is obvious that there are some small oscillations, while in Figure 6 there are not. Therefore, $\xi_3(r)$ should not be omitted even it is small.

According to Figure 6, the “baryon acoustic oscillation” (BOA) peak is at around the scale $r = 105.6(\text{Mpc/h})$.

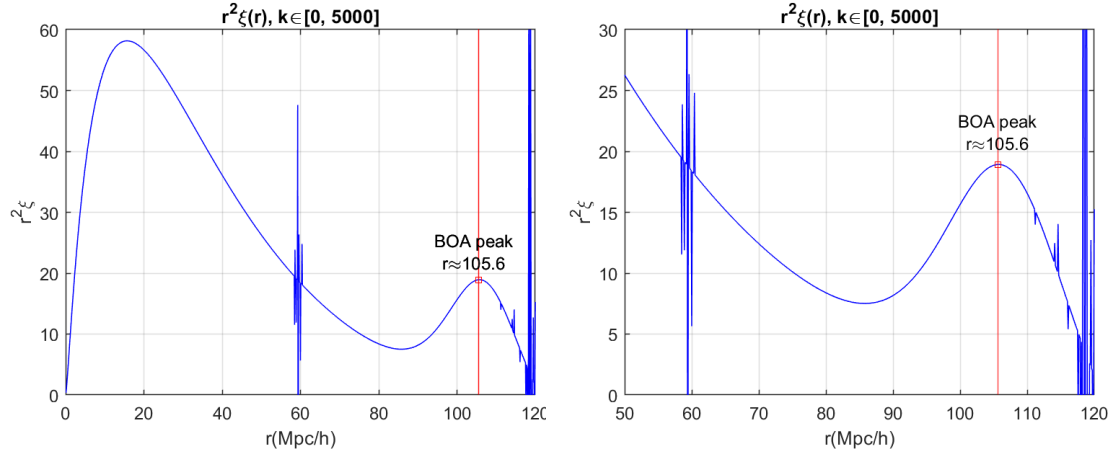


Fig. 6: The plot of $r^2 \xi(r)$, where $\xi(r)$ is the correlation function. The “baryon acoustic oscillation” (BOA) peak is at $r \approx 105.6$. In this plot we can see clearly how much influence the small error spikes in $\xi_3(r)$ will have on the $r^2 \xi(r)$ plot. However, since the upper limit of the integral was chosen to be big enough, so that the plot is still smooth enough in most r ’s except those spikes.

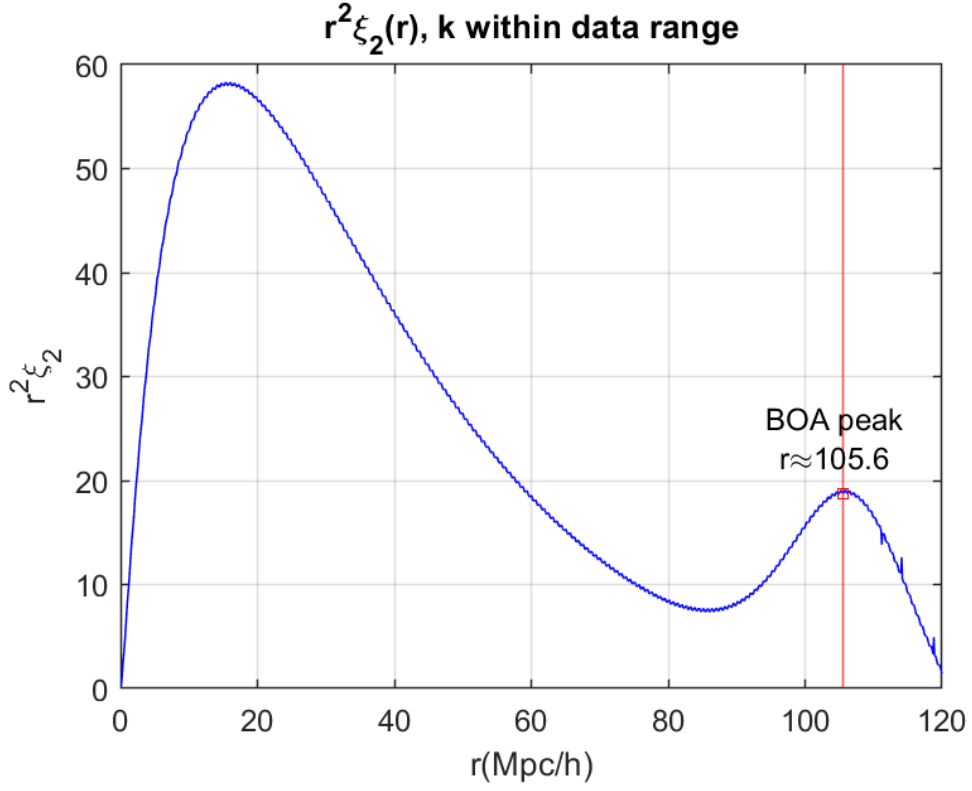


Fig. 7: The plot of $r^2 \xi_2(r)$, integrated only over the data range $[k_0, k_N]$. Small oscillations are clear and cannot be neglected, so the integrals below and over this range should be considered to get a smooth plot of $r^2 \xi(r)$.