Intensity variation of large solar energetic particle events associated with coronal mass ejections

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[1] We studied the coronal mass ejections (CMEs) and flares associated with large solar energetic particle (SEP) events of solar cycle 23 (1996–2002) in order to determine what property of the solar eruptions might order the SEP intensity. The SEP events were divided into three groups: (1) events in which the primary CME was preceded by one or more wide CMEs from the same solar source, (2) events with no such preceding CMEs, and (3) events in which the primary CME might have interacted with a streamer or with a nearby halo CME. The SEP intensities are distinct for groups 1 and 2 although the CME properties were nearly identical. Group 3 was similar to group 1. The primary findings of this study are as follows: (1) Higher SEP intensity results whenever a CME is preceded by another wide CME from the same source region. (2) The average flare size was also larger for high-intensity SEP events. (3) The intensity of SEP events with preceding CMEs showed a tighter correlation with CME speed. The extent of scatter in the CME speed versus SEP intensity plots was reduced when various subgroups were considered separately. (4) The intensities of energetic electrons were better correlated with flare size than with CME speed. (5) The SEP intensity showed poor correlation with the flare size, except for group 3 events. Since only a third of the events did not have preceding CMEs, we conclude that the majority of SEP producing CMEs propagate through the near-Sun interplanetary medium severely disturbed and distorted by the preceding CMEs. Furthermore, the preceding CMEs are faster and wider on the average, so they may provide seed particles for CME-driven shocks that follow. Therefore we conclude that the differing intensities of SEP events in the two groups may not have resulted due to the inherent properties of the CMEs. The presence of preceding CMEs seems to be the discriminating characteristic of the high- and low-intensity SEP events. 7514 Solar Physics, Astrophysics, and Astronomy: Energetic particles (2114); 7513 Solar Physics, Astrophysics, and Astronomy: Coronal mass ejections; 7519 Solar Physics, Astrophysics, and Astronomy: Flares; 2118 Interplanetary Physics: Energetic particles, solar; KEYWORDS: CME, flare, SEP

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Introduction

[2] Although the close relationship between interplanetary (IP) shocks and the observed time profiles of solar energetic particle (SEP) events was recognized long ago [Obayashi, 1962], the connection between shocks and coronal mass ejections (CMEs) got the attention soon after the discovery of white light CMEs [Stewart et al., 1974]. The importance of CME occurrence for the production of SEPs was first demonstrated by Kahler et al. [1978], leading to the suggestion that SEPs may be accelerated by the shock ahead of CMEs. It is currently thought that the large, gradual, and long-lived SEP events are due to CME-

widely accepted theory that explains all the observed properties of SEPs. For example, the CME speed and SEP intensity are reasonably correlated, yet the scatter is very large: for a given CME speed, the SEP intensity can vary over four orders of magnitude, with no satisfactory explanation in sight [Kahler, 2001; Gopalswamy et al., 2003a]. The presence of preexisting SEPs in the ambient medium and variations of energy spectra among SEP events are thought to be significant factors that can account for 1 to 2 orders of magnitude variation in SEP intensities [Kahler, 2001]. The ambient SEPs are supposed to form a seed population, which are further accelerated by the CMEdriven shocks. The seed particles may originate from impulsive flares [Mason et al., 1999] or from preceding CMEs [Gopalswamy et al., 2002a]. In addition to the presence of seed particles, the physical conditions in the ambient medium may also modify the characteristics of shocks and hence affect the intensity of SEPs. In a preliminary study, Gopalswamy et al. [2003a] found that SEP

driven shocks [Reames, 1999]. However, there is still no

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events of cycle 23 (1997-2001) preceded by wide CMEs from the same source region within a day were more likely to have higher intensities than those without such preceding CMEs. However, they did not evaluate how the correlation between SEP intensity and CME speed was affected by the presence of preceding CMEs. In this paper, we study this correlation for all large SEP events of solar cycle 23 (from 1996 to the end of 2002), identify their solar sources, and study the properties of the associated CMEs. For the subgroup with preceding CMEs, we compare the properties of the SEP-related CMEs with those of preceding CMEs. Since CME-driven shocks are thought to accelerate both electrons (inferred from type II radio bursts) and ions, we extend the study to energetic electrons detected in situ to see if the electron flux is affected by the presence of preceding CMEs. Finally, we examine the correlation between energetic proton and electron intensities with the flare size.

2. Data Selection

2.1. SEP Events

[3] We started with the set of all large proton events (intensity > 10 pfu, with 1 pfu = 1 proton cm⁻² s⁻¹ sr⁻¹) in the >10 MeV channel of the GOES instrument from 1997 to 2002. GOES data are available on line from NOAA Space Environment Center (http://sec.noaa.gov). We excluded six events that occurred during the period of temporary disability (June to October, 1998 and January to February, 1999) of the Solar and Heliospheric Observatory (SOHO) mission because the CME data in this study were obtained by SOHO's Large Angle and Spectrometric Coronagraph (LASCO) [Brueckner et al., 1995]. The C2 and C3 telescopes of LASCO obtain images of the corona from \sim 2 to 32 R_s field of view (FOV) from which the characteristics of CMEs are measured. For 8 events the higher intensity may be due to shock spike or interplanetary (IP) modulation of previous events, so they are not included in the analysis. Table 1 lists the remaining 60 intervals with SEP intensity exceeding 10 pfu. The event number (No.), SEP date, and onset time at GOES are given in columns 1-3, respectively. For each of these SEP events, we were able to identify a unique CME (which we call the primary CME). We have compiled the properties of the primary CMEs such as the time (T2) of first appearance in the LASCO/C2 FOV, speed (V2 km s⁻¹), width (W2 deg), and mass (M in 10^{15} g) in columns 4-7. The masses of the primary CMEs were determined using the standard technique of estimating the CME volume and the number of electrons contained in the CME [Vourlidas et al., 2002]. The solar source location of the CMEs was identified using a number of data sources (images, movies and lists) such as (i) the online Solar Geophysical Data, movies from SOHO's Extreme-ultraviolet imaging telescope (EIT) [Delaboudiniere et al., 1995], movies from the soft X-ray telescope (SXT) on board the Yohkoh spacecraft (until December 2001), microwave movies from the ground-based Nobeyama radioheliograph [Nakajima et al., 1994], and TRACE movies. We also used data collected during the Living With a Star (LWS) Coordinated Data Analysis Workshop (LWS/CDAW) held in July 2002 (http://cdaw.gsfc.nasa.gov/LWS/). The sources identified in this way are listed in Table 1 as a heliographic location (latitude and longitude in column 8). When the

solar source is behind the limb, the position angle can be identified from EIT movies and the solar source is indicated to be beyond the appropriate limb. For example, >NW90 means that the source is behind the northwestern limb of the Sun. The corresponding NOAA active region (AR) number, if available, is given in column 9. From the Solar Geophysical Data, we identified the flare size (X-ray/optical importance in column 10) and onset time of GOES soft X-ray flares (column 11). In columns 12 and 13, we have indicated the onset times of the associated metric and decameterhectometric (DH) type II radio bursts; N denotes the lack of a radio burst. In order to compare the association of the SEP events with the corresponding energetic electron events, we use the in-situ electron measurements taken by the Threedimensional Plasma and Energetic Particle (3DP) instrument [Lin et al., 1995] on the Wind spacecraft. Solid state detectors covering the range between 30 and 500 keV provide the best signal to noise ratio. This energy range is normally grouped into 7 logarithmically binned channels. For each energy channel the intensity of the electron event (in units of electron flux units or efu; 1 efu = 1 electron cm $^{-2}$ $s^{-1} sr^{-1} eV^{-1}$) was determined as the maximal flux observed during the event. Since impulsive electron events very often show velocity dispersion - higher energy electrons arrive before the lower energy ones - the peak fluxes at different energies are generally taken at different times. In this paper, we consider data from only one channel (108 keV), because at this energy the electrons have roughly the same speed as the \geq 10 MeV protons. The proton intensity (I_p) and 108 keV electron flux (I_e) are listed in Table 1.

2.2. Subgroups of SEP Events

[4] In order to assess the influence of the medium, we examined the conditions prevailing the corona into which the primary CMEs were launched. A certain volume of the corona is disturbed if there was a recent CME in that part of the corona. In order to quantify this we searched for other CMEs from the same source region as the primary CME that satisfy the following two criteria: (1) the preceding CMEs should have been launched within 24 hours ahead of the onset time of the primary CMEs, and (2) the preceding CMEs must be wider than average CMEs. Since the average width of non-halo CMEs is ~47°, CMEs having a width \geq 60° are deemed wide. If there were more than one preceding CME satisfying the above criteria, we choose the CME closest in time to the primary CME. Since it takes less than one day to more than 4 days for CMEs to reach Earth [Gopalswamy et al., 2000a, 2001d], choosing 1 day precedence makes sure that the aftermath of the preceding CMEs is not relaxed back to the preeruption conditions when the primary CME starts out. Figure 1 shows an SEP event with a preceding CME. The preceding (CME1) and primary (CME2) CMEs originated from the same solar source (AR 9628) located close to the southwest limb as can be seen in the EIT difference images. In the LASCO/C3 image at 0618 UT, both CME1 and CME2 can be seen. The proton intensity at three energy channels is also shown. The peak intensity in the >10 MeV channel is 2360 pfu. In the CME height-time plots, the onset times of the two CMEs and their difference (ΔT) are marked. The trajectories intersect at a height of \sim 18 R_s. When the primary CME lifts off, the preceding CME is already at a height of \sim 12 R_s.

Table 1. Major Solar Energetic Particle Events of Cycle 23, 1996-2003

No. SEP Data SEP Time CME Time V2, kms W2, deg M5 Source Loc. AR No. Fluer Limp. Fluer Obset Conset Up No. V2, deg M5 Source Loc. AR No. Pluer Limp. Fluer Obset Conset Up No. V2, deg M5 V2, deg V2	Tal	ne 1. Majo	1 Solai Elli	ergene ran	icie Even	is of Cyc	16 2.	5, 1996–200	13						
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 $[^]a Mass$ in unit of 10^{15} g. $^b \!\!>\!\! NW90$ (>SW90) means beyond northwest (southwest) limb.

c? denotes no optical flare information; EPL, eruptive prominence at limb.

^dN, no burst; N?, maybe occulted.

^eFirst instance the burst is seen at or below 14 MHz.

f1 pfu = 1 particle (cm² s sr)⁻¹.
g1 efu = 1 electron (cm² s sr eV)⁻¹.

^hCompound of two events, the second one starting around 0950 UT (see text).

ⁱCME association may be by chance.

^jTime corresponds to previous day.

^kTime corresponds to the next day.

We denote this initial separation between the primary and preceding CMEs as ΔR . The GOES light curves show the flare associated with the primary CME in the 1.0-8.0 Å and 0.5-4.0 A channels, with the heliographic location (S20W84) marked. The flare associated with the preceding CME was not listed in the Solar Geophysical Data, but the eruption was clearly seen in SOHO/EIT images. There were 23 such SEP events in which the primary CMEs were ejected into the aftermath of preceding CMEs, as listed in Table 2. The onset time (T1) of the preceding CME, its speed (V1, km s⁻¹), and width (W1, deg) are given in columns 5-7. The heliographic locations (column 8) and the active region numbers (column 9) are also included. The preceding time interval (ΔT in hours), the initial spatial separation (ΔR in R_s) and the approaching speed (ΔV = $V2-V1 \text{ km s}^{-1}$) are given in columns 10–12. The intensity (I_n in pfu) of the proton event associated with the primary CME is given in column 13. The last column (14) contains information on the kind of energetic activity associated with the preceding CMEs: complex type III bursts (cxIII), simple type III bursts (III or DHIII), metric type II bursts (MII), DH type II bursts (DHII), and minor ($I_p < 10$ pfu) SEP events (minor). It must be pointed out that ΔT was obtained by extrapolating the linear fit to the height-time measurements in the sky plane to 1 R_s, and taking the difference between the two onset times. ΔT is generally not too different from T2-T1. ΔT and ΔV indicate the extent of interaction between the primary and preceding CMEs. For example a small ΔT and a large ΔV correspond to CMEs that interact very close to the Sun. This group of 23 SEP events will be henceforth referred to as "P" events to note that the primary CMEs were preceded by at least one wide CME.

- [5] For a group of 20 events, the opposite situation applies: we cannot find a preceding wide CME for any of these from the same source region as the primary CME. We call these as "NP" events to note that there were "No Preceding" CMEs. Figure 2 shows an example of the NP events (event 17 in Table 1). Three snapshots of the LASCO images are shown superposed on EIT difference images with the source region (S17W09) pointed by arrows. The CME originated from a quiescent filament region. The filament eruption resulted in a prominent EUV and X-ray posteruption arcade with an M-class X-ray flare. The event had a proton intensity of 320 pfu. While there were two large CMEs from other parts of the Sun following the primary event, there was no preceding eruption from the filament region in question.
- [6] The remaining 17 events could not be included in either group due to various reasons described below:
- [7] (1) One event was completely back-sided (event 33 [Gopalswamy, 2002]) so it is difficult to say whether there was a preceding CME or not. Therefore this event was dropped from the statistical analysis.
- [8] (2) For another event (32, 2001 August 09) it was not possible to unambiguously identify the primary CME responsible for the SEP event. There was a candidate CME, a partial halo associated with a filament eruption near the central meridian, as indicated by the posteruption arcade in Figure 3. The CME was gradually accelerating with an average speed of only 479 km s⁻¹. The speed was still under 600 km s⁻¹ by the time the CME left the LASCO field of view at 1932 UT. The CME seems to have lifted off

too early (at 1030 UT, \sim 8.5 h before the SEP onset). The EIT images showed another clear eruption later in the day from close to the disk center (AR 9570 at S17E19) at 1822, but no white light CME was observed. Soft X-ray images from Yohkoh/SXT show that the sources of these two eruptions were separated by a large coronal hole extending from the south polar region into the northern hemisphere (see Figure 3). This coronal hole is likely to deflect the nearby CMEs and prevented easy detection of a possible CME from AR 9570. This coronal hole might have also been responsible for the confinement of the 1030 UT CME to the western hemisphere. Since it is difficult to decide between the two eruptions, we drop this event from further analysis.

[9] (3) For 6 events (03, 19, 20, 48, 49, and 52), there were major preceding eruptions from the source regions indicative CMEs, but the primary CMEs followed so closely that they would have overtaken any preceding CMEs within the occulting disk of LASCO coronagraph. The preceding eruptions were clearly observed in the inner coronal images such as from SOHO/EIT, Yohkoh/SXT and the Nobeyama Radioheliograph [Gopalswamy, 2004a]. Figure 4 shows an example from 2002 March 15 (event 48). Two eruptions can be seen in the EIT difference images at 2200 and 2236 UT corresponding to the two peaks in GOES soft X-ray plot. The first eruption produces a dimming to the northwest of AR 9866, while the second one is more extensive with a dimming to the South of the same AR. The LASCO image at 2306 shows a single CME with a complex front. Such CMEs should belong to group P, but we cannot use them in the statistical analysis since the preceding CMEs could not be measured. The onset time difference between the associated flares can be used, but there is no information on ΔV .

[10] (4) For 7 other events (22, 34, 39, 43, 45, 46, and 57), the primary CME interacted with a nearby streamer. Since streamers are the preeruption stages of CMEs, they may be considered as preceding CMEs of "zero speed". Figure 5 shows an example of the streamer interaction events. This is event 43 in Table 1 that occurred on 2001 December 26 from AR 9742 (N08W54). There was a streamer located above the west limb at PA $\sim 300^{\circ}$ as marked in the 0454 UT LASCO image. The streamer had appeared above the west limb for the first time on 2001 December 18. There were some slow CME events apparently originating from the outer layers of the streamer without affecting its overall structure. In the 0530 UT image on 2002 March 15, the CME can be seen above the west limb with the solar source visible in EIT image as a brightness enhancement. In the image at 0654 UT, the streamer is completely blown off. The primary CME on 2001 December 26 resulted in an SEP event of 779 pfu. certainly an order of magnitude above the threshold we set for high-intensity events. The CME was fast (1446 km s⁻¹) and wide (width $\geq 212^{\circ}$). In a sense these streamer interactions are similar to the close-to-source interactions because in either case we cannot derive the properties of preceding CMEs.

[11] (5) Event 41 should belong to the NP group since there was no preceding wide CME. However, the next event (42), from a nearby region, started during its rise phase, so we could not determine the proton flux. Therefore this event

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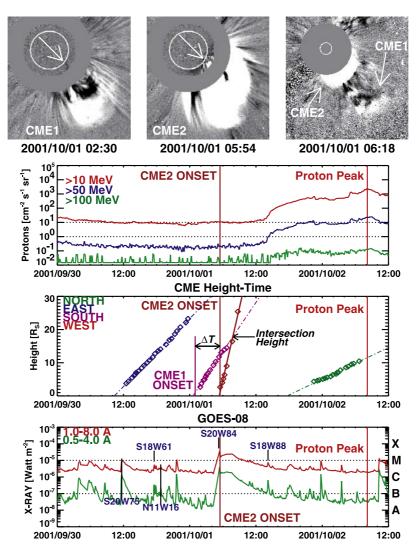


Figure 1. An example of P events that occurred on 2001 October 1 (event 36 in Table 1). (top) LASCO difference images showing the preceding (CME1) and primary (CME2) CMEs. The 0618 UT image shows both the CMEs in a single frame. (top middle) SEP intensity profiles in three energy channels (>10 MeV, >30 MeV, and >100 MeV). The onset time of the primary CME and the time of peak intensity of the protons are marked. (bottom middle) CME activity around the time of the primary CME shown as height-time plots. The plots are color coded to show the direction of the central position angle. The onset time of the preceding CME (CME1) is also marked. The onset time difference (Δ T) between the primary and preceding CMEs is also indicated. The height at which the primary and preceding CMEs intersect is also noted. (bottom) GOES X-ray light curves showing the flare activity associated with the two CMEs. The heliographic locations of the flares are also indicated. The active region from which the CMEs were launched was located at S20W84 at the time of the primary CME.

will be excluded from analysis. Event 42 did not have a preceding CME from the same region, but the interaction of its primary CME with that of event 41 has been documented [Gopalswamy et al., 2003a]. Both were halo CMEs originating from neighboring active regions. For this reason, we treat this event as a special case (halo-halo interaction) in the analysis.

[12] To summarize, there are 57 well-observed large SEP events that can be divided into three groups: (1) 23 P events in which the primary CMEs were preceded by other wide CMEs, (2) 20 NP events in which the primary CMEs were not preceded by other wide CMEs, and (3) 14 O events, in which "Other" types of interaction (below the occulting

disk interactions, streamer interactions and halo-halo interaction) occurred and hence these cannot be included in the NP group. We primarily consider the P and NP groups for studying the influence of CME interaction on SEP intensity. The O events will be used for a consistency check.

3. Analysis and Results

[13] We are interested in the properties of the P and NP events and seek their differences in terms of (1) the intensity of the energetic proton and electron events, (2) the speed, width, mass, source longitude, and kinetic energy of the primary CMEs, (3) the associated flare size, (4) the corre-

Table 2. Properties of Preceding CMEs

			0										
Date	T2	V2	W2	T1	V1	W1	Loc.	AR	ΔT	ΔR	ΔV	I_p	Remarks
97/11/04	0610	785	360	1111 ^a	352	122	S14W33	8100	19.7	36.9	433	72	DHII
97/11/06	1210	1556	360	0420	307	59	S18W63	8100	9.2	15.5	1249	490	III
98/05/02	1406	938	360	0531	542	360	S15W15	8210	9.0	26.3	396	150	III, weak DHII?
98/05/06	0829	1099	190	0002	786	110	S11W65	8210	8.7	36.2	313	210	mII
98/05/09	0335	2331	178	1301 ^a	432	85	S11W90	8210	15.2	35.0	1899	12	cxIII; DHII?
00/06/06	1554	1119	360	1530	929	90	N20E18	9026	0.3	2.5	190	84	mII, DHII
00/07/14	1054	1674	360	2030^{a}	839	62	N22W07	9077	14.7	64.6	835	24000	DHIII
00/11/24	1530	1245	360	0530	994	360	N22W07	9236	10.0	52.5	251	94	mII, DHII, minor
01/03/29	1026	942	360	1927 ^a	258	59	N20W19	9393	15.9	22.3	684	35	DHIII
01/04/02	2206	2505	244	1250	731	155	N19W72	9393	9.0	35.0	1774	1100	mII, DHII, minor
01/04/10	0530	2411	360	1554 ^a	1192	360	S23W09	9415	13.8	86.2	1219	355	mII,DHII, minor
01/04/12	1031	1184	360	2130 ^a	469	74	S19W43	9415	13.8	34.4	715	51	cxIII
01/04/15	1406	1199	167	1754 ^a	830	113	S16W71	9415	20.1	87.1	369	951	mII, cxIII
01/04/26	1230	1006	360	0830	740	91	N17W31	9433	4.3	17.4	266	57	DHIII
01/09/24	1030	2402	360	2154 ^a	633	163	S16E23	9632	13.1	43.7	1769	12900	none
01/10/01	0530	1405	360	0154	478	68	S20W90	9628	4.5	12.1	927	2360	cxIII
01/10/19	1650	901	360	0127	558	254	N15W29	9661	15.9	46.9	343	11	mII, DHII, minor
02/05/22	0350	1494	360	0006	1265	117	S19W56	9948	3.6	22.4	358	820	mII
02/07/15	2130	1307	111	2030	1105	360	N19W01	0030	1.1	7.6	168	234	DHII
02/07/20	2206	2017	360	2042 ^b	1351	83	S13E99 ^c	0039	0.7	6.2	666	28	mII
02/08/22	0206	1013	360	0406^{a}	395	70	S07W62	0069	22.3	55.7	531	36	cxIII
02/08/24	0127	1878	360	1327 ^a	321	103	S02W81	0069	12.3	21.5	1557	317	cxIII, minor?
02/11/09	1331	1793	360	1830 ^a	840	62	S12W29	0180	19.2	84.4	998	404	cxIII
12/11/09	1331	1/93	360	1830"	840	62	S12W29	0180	19.2	84.4	998	404	cxIII

^aTime corresponds to previous day.

lation between the proton intensity and CME speed, and (5) the correlation between energetic electron intensity and CME speed.

3.1. Particle Intensities

[14] Figure 6 shows the distributions of proton and electron intensities of the 57 events in comparison with those of the P and NP events. Energetic electrons were detected only in 45 of the 57 events. For both protons and electrons, the intensity of P events extends to much larger values compared to that of NP events. For example, the median intensity of protons is 210 pfu for the P events, compared to just 29 pfu for the NP events. The results did not change when we considered only the frontside events (the darker histograms in Figure 6). The median value for the entire set of 57 events is 54 pfu, which seems to be an intensity value that demarcates the P and NP events. The intensity difference can also be clearly seen for the electron events. The ratios of various median values are shown in Table 3. The median value of the NP events, I(NP), is half the median value I(ALL) for the 57 events taken together. This result applies for both electrons and protons. On the other hand, the median intensity of the P events, I(P), is four times I(ALL) for protons and twice I(ALL) for electrons. Furthermore, I(P) is \sim 7 times and \sim 4 times I(NP) for protons and electrons, respectively.

[15] On the basis of the distributions shown in Figure 6, SEP events with proton intensity (I_p) ≥ 50 pfu will be referred to as high-intensity events; those with $I_p < 50$ pfu, will be referred to as low-intensity events. From Table 1, we see that the majority (18/23 = 78%) of the P events are of high intensity. On the other hand, only 5 of the 20 (25%) NP events had high intensity. Furthermore, only two of the 20 NP events (10%) had $I_p > 100$ pfu, while 13 of the 23 (57%) P events had $I_p > 100$ pfu. Thus CMEs that propagate into the aftermath of preceding CMEs are

>3 times more likely to be associated with high-intensity SEP events (see also Tables 3 and 4).

3.2. Properties of Primary CMEs

[16] Are there other differences between the CMEs of P and NP events? Figure 7 compares the speed, width, and the source longitudes of the two groups with those of all the 57 events listed in Table 1. First of all, we confirm that the average speed (1468 km s⁻¹) of SEP-producing CMEs is about three times as large as the average speed of the general population [Gopalswamy et al., 2002a]. The average CME speeds of P (1492 km s⁻¹) and NP (1499 km s⁻¹) events are nearly identical. Most of the CMEs in the two groups were halo CMEs (18/23 = 72%)in group P and 15/20 = 75% in group NP). There is no CME in either group with width less than $\sim 120^{\circ}$. Furthermore, the two groups have similar speeds and widths as the entire set of 57 SEP-associated CMEs. As for the longitude distribution, we see that there is a general tendency for the SEP events to occur in the western hemisphere. The tendency is more pronounced for the P events than for the NP events. About half of the NP events originated from beyond the limb, while only three of the group P events were behind the limb. One might argue that behind-the-limb sources mean poor connectivity for detection at Earth and hence the group NP events may be of low intensity. However, as we showed in Figure 6, even if we consider just the frontside events, the intensity difference between P and NP events persisted. In other words, when the frontside events alone are considered, the primary distinction between the two groups is indeed the presence of preceding CMEs. The average speed of the frontside NP CMEs is 1378 km s⁻¹ compared to the 1618 km s⁻¹ for the behind-the-limb NP events. This slight difference (~15%) is not significant because the limb events have little

^bPreceding CME identified in polarized brightness images.

^cEstimated longitude.

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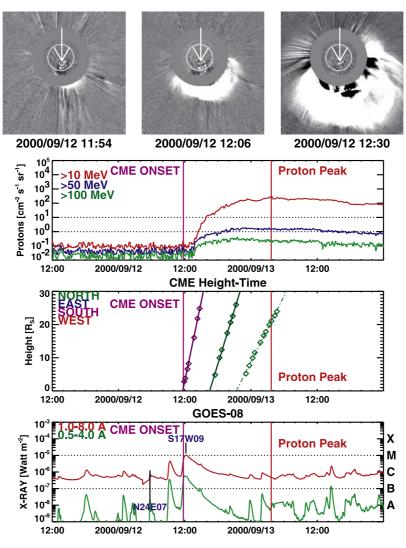


Figure 2. An example of the NP events (event 17 in Table 1). The plots are similar to those of Figure 1, except that there is no preceding CME. The primary CME occurred on 2000 September 12.

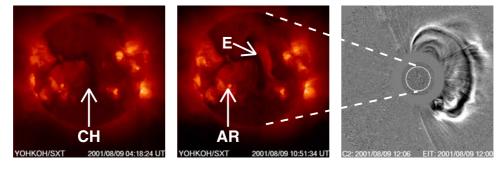


Figure 3. Two Yohkoh/SXT images on 2001 August 09 at 0418 and 1051 UT, showing the coronal hole (CH) and the SXT arcade (E). The LASCO difference image at 1206 UT shows a white light CME associated with the SXT arcade. This CME seems to be too early to be associated with the SEP event of 2001 August 09. Even though the solar source was close to the disk center, the CME moved mostly to the west, probably because of the extended coronal hole immediately to the east of the CME source region. Also marked is active region 9570 (AR), from which an EUV eruption was observed later in the day at 1822 UT by EIT (but not by LASCO).

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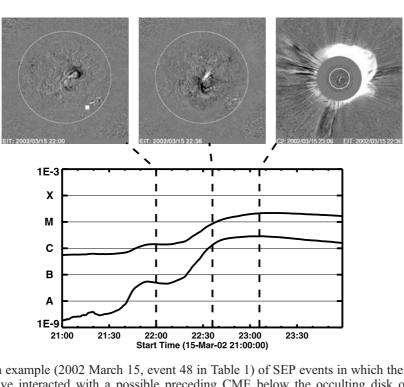


Figure 4. An example (2002 March 15, event 48 in Table 1) of SEP events in which the primary CME is likely to have interacted with a possible preceding CME below the occulting disk of LASCO/C2. Images from SOHO/EIT and LASCO are shown at the top. The times of the images are marked on the GOES X-ray plot. The GOES plot as well as the EIT images clearly show two eruptions from the same region, but LASCO sees only a single CME. The LASCO CME is likely to be a compound CME if both EIT eruptions had associated mass ejections.

projection effects and hence tend to show a higher speed compared to the disk events [Gopalswamy et al., 2000b].

3.3. Mass and Kinetic Energy of CMEs

[17] The CME mass is estimated by determining its volume and the number of electrons contained in it, assuming the CME to be a fully-ionized hydrogen plasma with 10% helium. The average mass of LASCO CMEs is $\sim 1.6 \times 10^{15}$ g, slightly lower than those of pre-SOHO CMEs [see, e.g., *Vourlidas et al.*, 2002]. For the primary CMEs, the average and median masses are much higher than this value (see Figures 8a–8c). For the 50 events for which we were able to make a mass estimate, the masses

were in the range $\sim 5 \times 10^{15}$ g to $\sim 3 \times 10^{16}$ g, with a median value of $\sim 9 \times 10^{15}$ g. The primary CMEs of the P and NP events had nearly the same median mass, suggesting that the CME mass is not a distinguishing characteristic between the two groups. Since there is little difference in the speed and mass between the two groups, we do not expect the kinetic energies to be different either. The median kinetic energy for the 50 SEP events with mass measurements was 8.5×10^{31} erg. For the P and NP groups, the median values were 7.1×10^{31} erg and 1.0×10^{32} erg, respectively. However, the kinetic energy of the SEP-related CMEs is an order of magnitude larger than the average and median kinetic energies of the general popu-

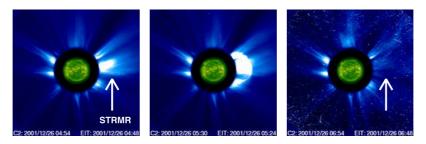


Figure 5. An example (2001 December 26 CME, event 43 in Table 1) of the streamer interaction events. The preexisting streamer is marked "STRMR" in the 0454 UT LASCO image. The CME can be seen in the 0530 LASCO image with the superposed EIT image showing the solar source as a compact brightening near the west limb. The LASCO image at 0654 shows that the streamer is gone with the CME.

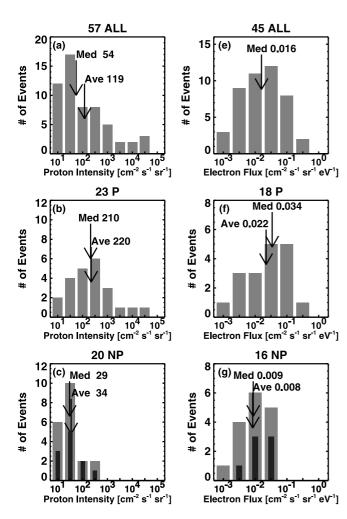


Figure 6. The distribution of intensities for (a–c) proton events and (d–f) electron events. For both protons and electrons, the intensity of P events (Figures 6b and 6e) extends to much larger values than for the NP events (Figures 6c and 6f). The mean ("Ave") and median ("Med") values of the distributions are marked. The darker histograms in Figures 6c and 6f correspond to frontsided events. The mean and median values for the frontsided proton events in Figure 6c are 38 and 34 pfu, respectively; the corresponding values are 0.014 and 0.013 efu for the electron events in Figure 6f.

lation of CMEs [Vourlidas et al., 2002; Gopalswamy, 2004b].

3.4. Flare Size

[18] We now compare the sizes of the flares associated with the P and NP events. We use the peak soft X-ray flux (in units of W $\rm m^{-2}$) during the flare in the GOES 1–8 Å channel as a measure of the flare size. To be observed, the flares have to occur on the front side with no occultation by

Table 3. Ratios of Median Particle Intensities

Ratio	Protons	Electrons
I(P)/I(ALL)	3.89	2.13
I(NP)/I(ALL)	0.54	0.56
I(P)/I(NP)	7.24	3.78

the limb. There were 45 events in all for which the soft X-ray intensity of the flare is available (see Table 1). The distribution of flare sizes is shown in Figures 8g–8i. The flare sizes varied over three orders of magnitude with a median value of 7.1×10^{-5} W m⁻² for the 45 events. However, when the P and NP events are considered separately, there is a clear difference: The NP events have an average peak X-ray flux $(3.9 \times 10^{-5}$ W m⁻²), which is \sim 5 times smaller than that $(2.1 \times 10^{-4}$ W m⁻²) of the P events. In addition to the presence of preceding CMEs, the flare size seems to be another parameter that distinguishes the two groups of events.

3.5. Type II Radio Bursts

[19] Radio bursts of type II are indicative of MHD shocks in the corona (inferred from metric type II bursts) and IP medium (inferred from DH and longer wavelength type II bursts). When we compare the onset times of metric and DH type II bursts with those of CMEs in Table 1, it is clear that they are closely related phenomena. In an earlier report, Gopalswamy [2003] found that 70-80% of the large SEP events were associated with metric type II bursts while 100% of them were associated with DH type II bursts. With the expanded data set of the present study (see Table 1), the earlier results are confirmed: 45 of the 57 SEP events (or 79%) were associated with metric type II bursts while all 60 (or 100%) of them were associated with DH type II bursts. This close relationship between type II bursts and SEP events has been taken to indicate that the same CME-driven shock accelerates electrons responsible for the type II bursts and SEPs detected in situ. It is interesting to note that there is no difference among P, NP and O events in terms of their association with DH type II bursts. The deficit of metric type II bursts is roughly the same for P, NP and O events: 17%, 20%, and 29%, respectively. We conclude that there is no major difference between P and NP events in terms of their association with type II radio bursts.

4. Properties of Preceding CMEs

[20] The speed, width, and longitude distributions of the preceding CMEs (see Table 2) are compared with those of the primary CMEs in Figure 9. The average speed (715 km s⁻¹) of the preceding CMEs is much smaller than that (1492 km s⁻¹) of the primary CMEs. However, the average speed of the preceding CMEs is significantly larger than that of the general population of LASCO CMEs (\sim 488 km s⁻¹ [see *Gopalswamy et al.*, 2003b; *Yashiro et al.*, 2004; *Gopalswamy*, 2004b]). This is expected on the basis of our selection criterion that the preceding CMEs must be wide (\geq 60°) because width and speed are correlated [see, e.g., *Yashiro et al.*, 2004]. What is the implication of the above-average speed of preceding CMEs? According to *Gopalswamy et al.* [2001a], the fast mode speed outside the core of active regions has a broad peak at a heliocentric

Table 4. Intensities of P and NP Events

Group	$I_p \geq 50 \text{ pfu}$	$I_p < 50 \text{ pfu}$	Total
P	18	5	23
NP	5	15	20

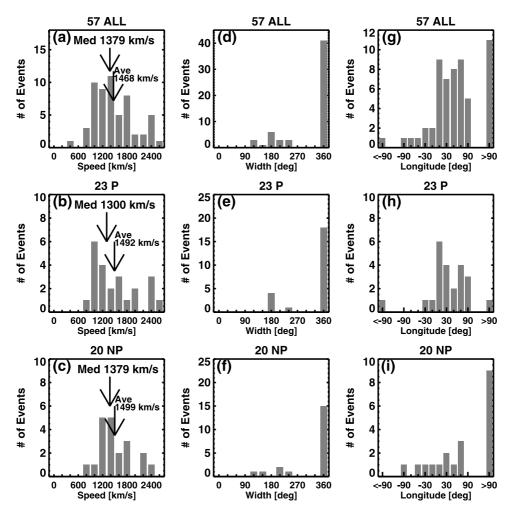


Figure 7. The speed, width, and source longitude distributions of the primary CMEs associated with the P and NP events in comparison with those of all 57 SEP events. The three groups have very similar speeds and widths. The fraction of behind-the-limb CMEs is larger for the NP events.

distance of ~ 3 R_s with a peak value of 550 km s⁻¹. The peak value depends on the actual density and magnetic field values. At coronal levels where metric type II bursts originate, the fast mode speed could be substantially lower than the peak value, so CMEs can easily drive shocks, provided their starting height is low enough. The preceding CMEs, with an average speed well above the peak Alfven speed, are likely to drive shocks near the Sun and hence accelerate some particles, which may serve as seed particles for the shock of the primary CME. In order to show that the preceding CMEs may drive shocks, we examined the radio bursts associated with the preceding CMEs. In Table 2, we have indicated the type of radio phenomena associated with the preceding CMEs (mII - metric type II; DH II - type II bursts in the decameter-hectometric (DH) wavelength regime; cxIII - complex type III bursts generally associated with large CMEs). The majority (13/23 or \sim 57%) of the preceding CMEs had either metric or DH type II bursts, indicative of shocks near the Sun. Four of the 13 CMEs had also minor SEP events ($I_p \le 1$ pfu), because these 4 CMEs had an average speed of 915 km s⁻¹. Figure 9 also shows that three-quarters of the preceding CMEs had widths exceeding 90°. The lower cutoff at 60° is due to the selection criterion used. Thus the above-average width of the preceding CMEs is also consistent with their ability to drive shocks because DH type II bursts are invariably associated with fast and wide CMEs [*Gopalswamy et al.*, 2001c].

4.1. Primary and Preceding CMEs: Spatial and Temporal Relationship

[21] Figure 10 shows the distributions of (a) temporal separation (ΔT), (b) the approaching speed (ΔV), and (c) initial spatial separation (ΔR) between primary and preceding CMEs. ΔT ranges between <1 h and 24 h (the upper limit is due to our selection criterion). Interestingly, the average preceding time was only \sim 11 h. The average approaching speed is relatively high (777 km s⁻¹) suggesting that the primary CMEs would catch up with the preceding CMEs in the near-Sun IP medium. Figure 10c shows that at the lift-off time of the primary CMEs, the preceding CMEs are typically at a heliocentric distance of \sim 35 R_s, close to the outer edge of the LASCO FOV. At the onset time of the primary CMEs, the preceding CMEs were located anywhere from $\Delta R < 2$ to $\sim 90 R_s$, but generally closer to the Sun, than to Earth. The mean and median values of ΔR were 35 and 37 R_s, respectively. Finally, Figure 10d shows the heliocentric distance at which the primary CME would have caught up with the preceding

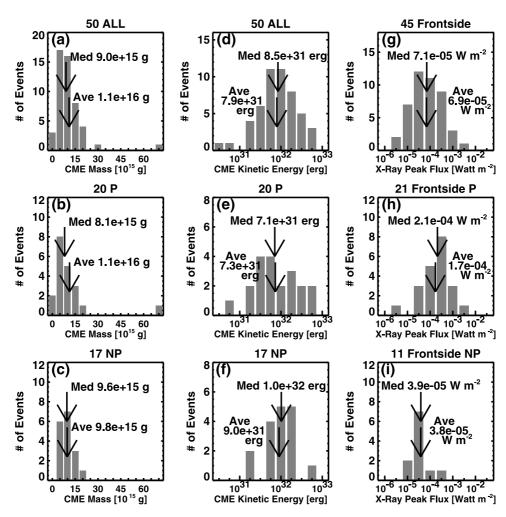


Figure 8. The distributions of CME mass, kinetic energy, and the size of the associated soft X-ray flare for the P and NP events in comparison with those of all 50 events. The mean and median values of the distributions are indicated. The GOES soft X-ray flux (in units of Wm⁻²) in the 1-8 Å channel is used as a measure of the flare size.

CMEs. This distance was obtained by simply extrapolating the height-time plots for primary and preceding CMEs up to the point where the two trajectories would intersect. While this is a reasonable estimate for the limb events, it may not be so for disk events and for those events which show gradual acceleration. Nevertheless, it provides an estimate of the distance at which the CMEs would have completed interaction. The mean value of this "interaction height" is $69 \, R_s$. Gopalswamy et al. [2002a] obtained similar results with less restrictive criteria on the source region and width of the preceding CMEs: the preceding CMEs were required to have a position-angle overlap with the primary CMEs. They obtained a height of intersection to be $\sim 21 \, R_s$, and the temporal separation as $\sim 7 \, h$. They also obtained the average width of the preceding CMEs to be $\sim 50^{\circ}$.

5. Particle Intensity-CME Speed Relationship

[22] The good correlation between SEP intensity and CME speed has been considered as a strong evidence supporting the paradigm that SEPs are accelerated by CME-driven shocks [see, e.g., *Kahler*, 2001, and refer-

ences therein]. The SEP intensity is generally better correlated with CME speed than with the flare size [Gopalswamy et al., 2003a]. However, the scatter is very large, to the extent that for a given CME speed the SEP intensity could vary over 3-4 orders of magnitude. Kahler [2001] suggested that the presence of SEPs in the IP medium might be able to account for one order of magnitude in the scatter. If we accept the paradigm that all large SEP events originate from CME-driven shocks, then the background SEPs in the IP medium might be due to preceding CMEs. In other words, our P events are somewhat similar to Kahler [2001] events that occur when background SEPs are present. The difference, of course, is that we are considering preceding CMEs based on white light observations, while Kahler [2001] considers background SEPs detected in situ.

[23] In order to see the effect of preceding CMEs, we have made scatterplots between the SEP intensity and CME speed for the two groups in Figure 11. We have made two sets of plots, one with all the events and the other with only frontside events ($|longitudes| \le 90^{\circ}$). Clearly, the correlation coefficients (r) for the P (0.43)

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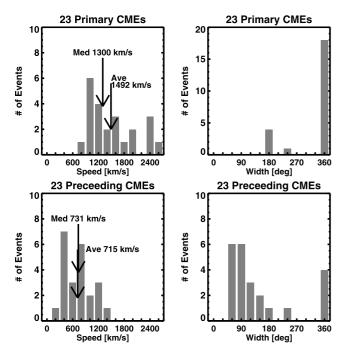


Figure 9. The speed and width distributions of the preceding CMEs as compared to those of primary CMEs. The mean and median values of the distributions are indicated.

and NP (0.58) events are similar to the values obtained before [Kahler, 2001; Gopalswamy et al., 2003a]. When we exclude events from behind the limb, the correlation improves dramatically for the P events (r increases from 0.43 to 0.70) while it slightly decreases for the NP events (from 0.58 to 0.46). The two outliers that caused poorer correlation for the P events are behind-the-limb events. The scattering is also reduced compared to the total plot. We see that the P and NP events occupy distinct regions in the CME speed- I_p plot. When we consider the P and NP events together, the correlation coefficients are 0.36 and 0.52 for all and frontside events, respectively.

[24] The effect of preceding CMEs on the correlation between SEP intensity and CME speed is shown in Figure 12. The largest extent of scatter is seen when the two groups are combined. When front-sided events alone are considered, the extent of scatter did not decrease significantly for the combined data set. When considered separately, the extent of scatter is reduced significantly for the P and NP events. When front-sided events alone are considered, there is further reduction of scatter for the P events, with little change for the NP events.

5.1. Electron Intensity-CME Speed Relationship

[25] The scatterplots between the electron flux and CME speed are shown in Figure 13. The electron flux is hardly correlated with CME speed for both P and NP populations. When we exclude the behind-the-limb events, the correlation improves slightly for the P events, while there is hardly any change for the NP events. From the electron events alone, it is difficult to see the effect of preceding CMEs, but this is consistent with the relationship indicated by proton events. However, the scatterplot

clearly shows the different regions occupied by the P and NP electron events.

5.2. Other Possible Interactions

[26] In the above analysis, we had excluded the O events (streamer interactions, possible interaction below the occulting disk, and halo-halo interaction) because there is little information on the preceding CMEs that could be used for statistical investigation. In order to see how these interactions affect the particle intensity-CME speed correlation, we have made scatterplots for protons (Figure 14) and electrons (Figure 15) including the O events. The correlation coefficient for the O events (r = 0.68) is higher than that for the NP and P groups. Furthermore, when we exclude the single behind-the-limb event in the O group, the correlation coefficient for the non-limb events becomes r = 0.74. The O events clearly merge with the P events and generally move away from the NP events. Only 11 of the 14 O events were associated with energetic electrons and the correlation coefficient between electron intensity and CME speed (r = 0.78) is similar to that of protons (r = 0.74). We conclude that the O-type events clearly belong to the P group. When we combined the O and P groups, the correlation coefficients for protons were 0.45 (all) and 0.59 (front-sided). For electrons, the corresponding correlation coefficients were 0.29 and 0.45.

6. Particle Intensity and Flare Size

[27] The proton intensity is known to be weakly correlated with flare size measured as the peak 1-8 Å X-ray intensity

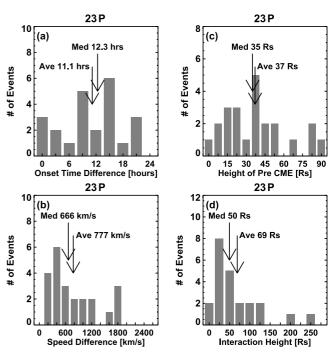


Figure 10. The distribution of (a) temporal separation (ΔT) , (b) approaching speed (ΔV) , (c) initial spatial separation (ΔR) between the primary and preceding CMEs (Table 2), and (d) the height at which the trajectories of the primary and preceding CMEs intersect. The mean and median values of the distributions are indicated.

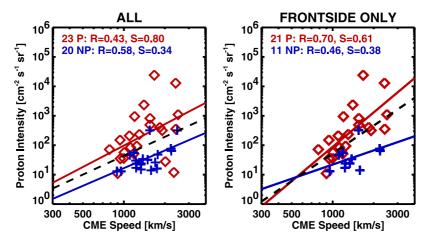


Figure 11. (left) $I_p - V_{cme}$ scatterplot for the P (diamond symbols) and NP (plus symbols) events. (right) The same as in the left but for frontside events (central median distance $\leq 90^{\circ}$). The solid lines are the regression lines to the P and NO data points. The dashed line is the regression line for all the data points. Here and in the following figures, S denotes the standard deviation of the Y variable. R is the correlation coefficient. Note that the P and NP data points occupy different regions of the plot with minimal overlap.

[Gopalswamy et al., 2003a]. We repeated the correlation analysis for the three subgroups (P, NP, and O) and the result is shown in Figure 16. The proton intensity of NP events has no correlation with the flare size at all; for the P events, it has an extremely weak correlation (r=0.12 for all events and 0.22 for front-sided events). The correlation coefficient for the O events is surprisingly high (r=0.71 for all O events and 0.78 for front-sided events). The presence or absence of preceding CMEs seems to make no difference for the correlation between the SEP intensity and flare size. This means the weak correlation found earlier between the SEP intensity and flare size [Gopalswamy et al., 2003a] must be entirely due to the O events.

[28] The correlation between electron intensity and flare size is shown in Figure 17. First of all, we note that the electron flux of NP events is not correlated with the flare size. On the other hand the P events are highly correlated with the flare size (r = 0.74 for all events and 0.86 for front-sided events). The correlation is also very high (r = 0.78) for the O events, consistent with the previous conclusion that the O events behave very similar to the P events.

[29] Table 5 summarizes the various correlations discussed above. We see a clear tendency for the proton intensity of P and O events to have the best correlation with CME speed. On the other hand, electron intensity is best correlated with flare size again for the P and O events.

7. What Causes High Intensity?

[30] The primary result of this paper is that the high-intensity SEP events are predominantly associated with fast and wide CMEs that are preceded by other wide CMEs from the same source region within a day. This means, the primary CMEs must be propagating through the preceding CMEs or their aftermath, rather than through the normal corona/solar wind. During the first few hours of their life time, the shocks driven by the primary CMEs are highly potent in accelerating particles [see, e.g., *Kahler*, 1994]. It is during this phase that the

shocks driven by the primary CMEs propagate through a medium highly disturbed and distorted by the preceding CMEs. Although a detailed analysis of the physical mechanism by which the SEP intensity may be enhanced is beyond the scope of this paper, we would like to discuss a few possibilities. One can roughly divide the effects into three categories: (1) effects related to the modification of the primary shock, (2) effects related to the deviation of the medium from normal solar wind (with Parker-spiral magnetic field), and (3) presence of seed particles in the ambient medium.

[31] It is easy to understand the modification of the primary shock because, shock strength (or Alfvenic Mach number) can change depending on the magnitude and distribution of the plasma density (n) and magnetic field (B) in the disturbed region. The change in the characteristic speed (Alfven speed V_a, since the sound speed is relatively small) due to changes in n and B can be written as: $dV_a/V_a = dB/B - (1/2) dn/n$. Near the Sun, the density enhancement represented by the preceding CME may be important, which lowers Va and makes the shock strong; stronger shocks would result in higher SEP intensity, which is known statistically and from theoretical considerations. This is consistent with the observed DH radio emission enhancement for a duration corresponding to the transit time of the primary shock through the density enhancement of a preceding CME observed by LASCO [Gopalswamy et al., 2001b, 2002b].

[32] A shock propagating through a normal solar wind may have most of its surface threaded by open field lines along which accelerated particles escape. When there is a preceding CME, the shock may be propagating though a mixture of closed and open filed lines. Depending on the separation between the primary and preceding CMEs, some particles escaping the primary shock may be returned for further acceleration by the closed field lines threading the primary shock. Implications of the presence of preceding CMEs as modifiers of the structure of the interplanetary magnetic field (on spatial scales of tenth of an AU) were considered by *Kallenrode* [2002] in a numerical model.

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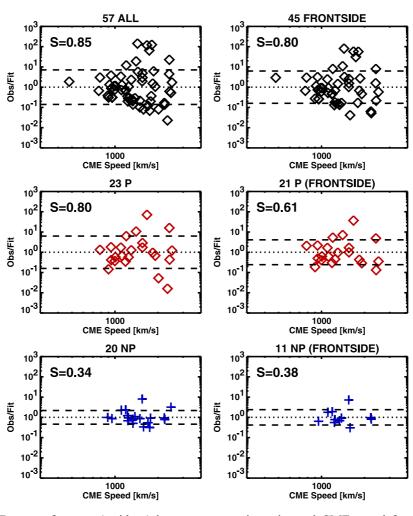


Figure 12. Extents of scatter (residues) between proton intensity and CME speed for the P (middle panels) and NP (bottom panels) in comparison with that for all 57 events (top panels), plotted as the ratio of the observed (Obs) intensity to the intensity given by the regression line (Fit). Plots are made for all events (left panels) and front-sided events (right panels). As before, the right panels correspond to frontside events. The horizontal dotted line represents the regression line. The dashed lines at 1-sigma (S) levels. Note that the scatter is significantly reduced for P and NP events, compared to all events.

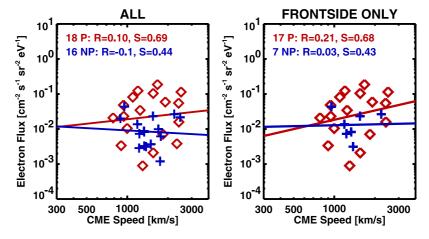


Figure 13. Scatterplots between the 108 keV electron intensity and CME speed for the P (diamond symbols) and NP (plus symbols) events. The solid lines are the best fit lines for the data points in each group.

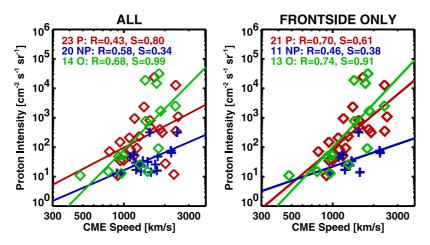


Figure 14. Scatterplots between proton intensity and CME speed for the P, NP, and O groups. The regression line is also shown for each group. The correlation coefficients are also indicated.

[33] Presence of seed particles of both from CMEs [Kahler, 2001] and from preceding flares [Mason et al., 1999] have been considered. These authors base their result on particles detected at 1 AU. We consider the presence of particles in the very early stages in a region directly overlying the primary shocks. It is worth noting that the aftermath of CMEs may contain dense core material which may be cold or hot [see, e.g., Ciaravella et al., 2003]. The cores may also contain nonthermal electrons, which produce moving type IV radio bursts known for more than four decades [see, e.g., Stewart, 1985]. It is likely that these plasmoids with nonthermal electrons may also carry lowenergy ions. When a primary shock encounters a preceding CME with a nonthermal core, it can accelerate these particles further. Presence of type II bursts also indicates the presence of shock-accelerated electrons and the same shocks may accelerate low levels of ions.

[34] The better correlation between electron flux and flare size for the P and O events needs further investigation. It is generally understood that the flare phenomenon observed in H-alpha and X-rays is due to the precipitation of energetic electrons into the chromosphere from the flare site in the corona. On the other hand the electrons detected in situ are

the ones propagating away from the Sun from the acceleration site (which could be the flare site or the CME-driven shock). At present there is no consensus on the source of the energetic electrons detected in situ. While the low-energy electrons responsible for the type II bursts are certainly due to CME driven shocks, the same cannot be said about the higher energy electrons. The present results suggest that the flare size may hold some clue to understand the electron events. The good correlation between flare size and electron flux for the P and O events also suggest that the CME interaction has something to do with the observed electron intensity. The simplest explanation is that the preceding CMEs somehow enable higher intensity to be observed at 1 AU. It is not clear how the CME interaction and flare size work in unison to produce the observed high electron intensity.

8. Discussion and Summary

[35] We studied the large SEP events of solar cycle 23 to understand the correlation between their intensity and the speed of the associated CMEs. While the overall correlation is similar to what was found before [Kahler, 2001;

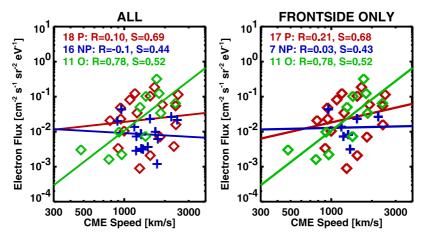


Figure 15. Scatterplots between the 108 keV electron intensity and CME speed for the P, NP, and O groups. The regression line is also shown for each group. The correlation coefficients are also indicated.

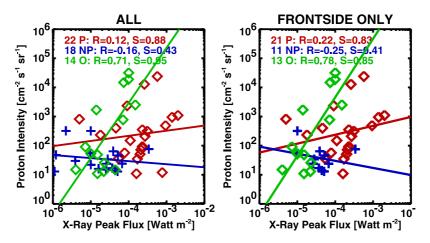


Figure 16. Scatterplots between proton intensity and flare size for the P, NP, and O groups. The regression line is also shown for each group. The correlation coefficients are also indicated.

Gopalswamy et al., 2003a], the correlation improves remarkably for events with preceding CMEs. The result also extends to energetic electrons, although the correlation coefficient is poorer. Excellent correlation was found for the group of events in which the primary CME interacted with a streamer or another eruption very close to the solar surface inferred from non-coronagraphic observations. The energetic electron intensity associated with CMEs preceded by other wide CMEs was better correlated with the flare size than with the CME speed. A combined effect of the preceding CMEs and the flare size seems to have resulted in high electron flux. On the other hand, the proton intensity is not well correlated with flare size except for the O events in which the primary CME interacted with a possible preceding CME close to the Sun.

[36] In addition to the high particle intensity, there may be other implications for the presence of preceding CMEs. Considering the non-NP events, we see that in the majority of the SEP events (37/57 or 65%), the primary CME propagates through the aftermath of preceding CMEs, and not through the normal solar wind. Since the preceding CME is a multithermal plasma (8000 K in the prominence core to \sim 2 MK in the front), the concept of "charge-state" freezing based solely on coronal temperatures may have to

be revised. The observed charge states of SEP ions can be modified due to postacceleration stripping when the ions transit through high-density regions for a sufficiently long time [Reames et al., 1999; Barghouty and Mewaldt, 2000]. In the scenario presented here, the preceding CMEs can act as the high-density region. CMEs near the Sun represent a density enhancement of several times the ambient medium. If we consider a density (n) of $\sim 10^8$ cm⁻³ in the preceding CME, and if the shock transit time (τ) through the density as ~ 0.5 hr (as inferred from radio observations [see Gopalswamy et al., 2001b]), we get $n\tau$ as 1.8×10^{11} cm⁻³s, which is more than adequate for additional stripping [Kovaltsov et al., 2001].

[37] It must be pointed out that we considered only preceding wide CMEs, but there may be other smaller preceding events. We relaxed the width requirement to account for CMEs with width $<60^{\circ}$. There were 14 additional preceding CMEs with widths ranging from 15° to 30°, distributed among the three groups as follows: P (4 events), NP (5 events), and O (5 events). Some primary CMEs in the P group already have multiple preceding CMEs, so allowing for additional preceding CMEs does not change the result. For the O events, the interactions (streamer and under 2 R_s) are spatially closer to the primary

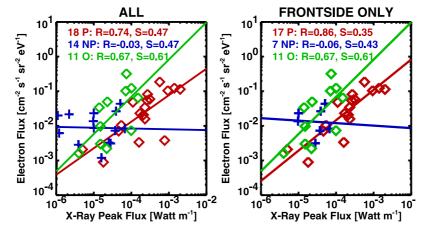


Figure 17. Scatterplots between the 108 keV electron intensity and flare size for the P, NP, and O groups. The regression line is also shown for each group. The correlation coefficients are also indicated.

Table 5. Correlation of Particle Intensities With CME Speed and Flare Size^a

Group	V Versus I _p	V Versus I _e	F Versus I _p	F Versus I _e
ALL	0.45	0.29	0.35	0.53
-Frontside	0.59	0.45	0.36	0.65
P	0.43	0.11	0.12	0.74
-Frontside	0.70	0.22	0.22	0.86
NP	0.58	-0.06	-0.16	-0.03
-Frontside	0.46	0.03	-0.25	-0.06
O	0.68	0.78	0.71	0.67
-Frontside	0.74	0.78	0.78	0.67

^aV, speed; F, flare size.

CME so the interaction with the narrow CME has to occur later. For this reason, we ignore the narrow CMEs preceding the O events. Moving the 5 NP events with preceding narrow CMEs to the P group and repeating the analysis did not change the correlations significantly. The correlation coefficients decreased slightly for the P events and increased slightly for the NP events. The main result that SEP events with preceding CMEs have a higher intensity did not change. Use of wide CMEs is still the best way to study the influence of preceding CMEs because it may be difficult to detect narrow CMEs originating from close to the disk center using coronagraphs.

[38] The requirement that the preceding CME originates from the same active region as that of the primary CME was imposed to make sure that the CMEs interact. Therefore we did not consider interaction of primary CMEs with CMEs from other solar sources. Since CMEs are large-scale structures, the coronal volume affected by fast and wide CMEs is generally very large especially at larger heliocentric distances. Therefore even nearby eruptions can significantly affect particle acceleration and provide seed particles. Shock from one CME can propagate through a CME from a nearby region and may have similar effects [Gopalswamy et al., 2003a].

[39] The main conclusions of this paper are as follows: (1) CMEs preceded by other wide CMEs within a day are more likely to be associated with of high-intensity (>50 pfu) SEP events. (2) When the primary CMEs are not preceded by other wide CMEs, the resulting SEP event is of low intensity ((<50 pfu) and the associated flares are generally weaker. (3) The proton intensity of SEP events with no preceding CMEs had no correlation with flare size; for events with preceding CMEs the correlation was extremely weak, but positive. (4) The flare size seems to play a major role in ordering the energetic electron intensity in those events where CME interaction is important. (5) The electron intensity is generally poorly correlated with the CME speeds, irrespective of the presence of a preceding CME. (6) For SEP events in which the primary CME interacted with a streamer or with an eruption preceding very close in time, both electron and proton intensities were highly correlated with CME speed as well as flare size. This is somewhat puzzling and needs further investigation. (7) The scatter between CME speed and SEP intensity is reduced when the SEP events with and without preceding CMEs are considered separately.

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