RBE549 - Final Exam

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Problem 1

In this problem, we are tasked with solving an iterative optical flow problem. To compute optical flow, we learned an interative method to update u(x, y), v(x, y) at each iteration, according to:

$$\begin{bmatrix} u(x,y) \\ v(x,y) \end{bmatrix}^2 = \begin{bmatrix} \lambda I_x^2 + 4 & \lambda I_x I_y \\ \lambda I_x I_y & \lambda I_y^2 + 4 \end{bmatrix}^{-1} \begin{bmatrix} \sum_{n \in neighbors(x,y)} u^{old}(n) - \lambda I_x I_t \\ \sum_{n \in neighbors(x,y)} v^{old}(n) - \lambda I_y I_t \end{bmatrix}$$
(1)

We are asked to consider a local coordinate frame (x', y') where x' is aligned with the image gradient and y' is perpendicular to the image gradient. Likewise, $(u', v') = (\frac{dx'}{dt}, \frac{dy'}{dt})$ are the image velocities in this frame. In this coordinate frame,

$$I_{x'} = \sqrt{I_x^2 + I_y^2} \text{ and } I_{y'} = 0$$
 (2)

We wish to show that the update equations:

$$\begin{bmatrix} u'(x,y) \\ v'(x,y) \end{bmatrix}^2 = \begin{bmatrix} \lambda I_{x'}^2 + 4 & \lambda I_{x'}I_{y'} \\ \lambda I_{x'}I_{y'} & \lambda I_{y'}^2 + 4 \end{bmatrix}^{-1} \begin{bmatrix} \sum_{n \in neighbors(x,y)} u'^{old}(n) - \lambda I_{x'}I_t \\ \sum_{n \in neighbors(x,y)} v'^{old}(n) - \lambda I_{y'}I_t \end{bmatrix}$$
(3)

...can reduce to

$$u'^{new}(x,y) = \bar{u}'^{old} - \frac{I_{x'}^{2}\bar{u}'^{old} + I_{x'}I_{t}}{I_{x'}^{2} + \frac{4}{\lambda}}$$

$$\tag{4}$$

$$v'^{new}(x,y) = \bar{v}'^{old} \tag{5}$$

. To do this, we first expand our second term to an equivalnet form to make matters easier for us to work with. Specifically, the inverse of a $2x^2$ matrix is:

$$\begin{bmatrix} a & b \\ c & d \end{bmatrix}^{-1} = \frac{1}{ad - bc} \begin{bmatrix} d & -b \\ -c & a \end{bmatrix}$$
 (6)

...and thus...

$$\begin{bmatrix} \lambda I_{x'}^2 + 4 & \lambda I_{x'}I_{y'} \\ \lambda I_{x'}I_{y'} & \lambda I_{y'}^2 + 4 \end{bmatrix}^{-1} = \frac{1}{(\lambda I_{x'}^2 + 4)(\lambda I_{y'}^2 + 4) - \lambda I_{x'}I_{y'}\lambda I_{x'}I_{y'}} \begin{bmatrix} \lambda I_{y'}^2 + 4 & -\lambda I_{x'}I_{y'} \\ -\lambda I_{x'}I_{y'} & \lambda I_{y'}^2 + 4 \end{bmatrix}$$
(7)

...which simplifies to:

$$\frac{1}{4\lambda I_{x'}^2 + 4\lambda I_{y'}^2 + 16} \begin{bmatrix} \lambda I_{y'}^2 + 4 & -\lambda I_{x'} I_{y'} \\ -\lambda I_{x'} I_{y'} & \lambda I_{x'}^2 + 4 \end{bmatrix}$$
(8)

This in turn leads us to an expanded view from earlier:

$$u'^{new}(x,y) = \frac{1}{4(\lambda I_{x'}^2 + \lambda I_{y'}^2 + 4)} \begin{bmatrix} \lambda I_{y'}^2 + 4 & -\lambda I_{x'} I_{y'} \end{bmatrix} \begin{bmatrix} 4\bar{u}'^{old} - \lambda I_{x'} I_t \\ 4\bar{v}'^{old} - \lambda I_{y'} I_t \end{bmatrix}$$
(9)

...We can then expand it into this:

$$u'^{new}(x,y) = \frac{1}{4(\lambda I_{x'}^2 + \lambda I_{y'}^2 + 4)} \left((\lambda I_{y'}^2 + 4)(4\bar{u}'^{old} - \lambda I_{x'}I_t) - (\lambda I_{x'}I_{y'})(4\bar{v}'^{old} + \lambda I_{x'}I_{y'}\lambda I_{y'}I_t) \right)$$
(10)

$$u'^{new}(x,y) = \frac{\lambda I_{y'}^2 4\bar{u}^{old} - \lambda I_{y'}^2 \lambda I_{x'} I_t + 16\bar{u}'^{old} - 4\lambda I_{x'} I_t - (\lambda I_{x'} I_{y'}) 4\bar{v}'^{old} + \lambda^2 I_{y'}^2 I_{x'} I_t}{4(\lambda I_{x'}^2 + \lambda I_{y'}^2 + 4)}$$
(11)

$$u'^{new}(x,y) = \frac{4(\lambda I_{y'}^2 \bar{u}'^{old} + 4\bar{u}'^{old} - \lambda I_{x'} I_t - \lambda I_{x'} I_t y \bar{v}'^{old})}{4(\lambda I_{x'}^2 + \lambda I_{y'}^2 + 4)}$$
(12)

$$u'^{new}(x,y) = \frac{(\lambda I_{y'}^2 \bar{u}'^{old} + 4\bar{u}'^{old} - \lambda I_{x'} I_t - \lambda I_{x'} I_{y'} \bar{v}'^{old})}{(\lambda I_{x'}^2 + \lambda I_{y'}^2 + 4)}$$
(13)

$$u'^{new}(x,y) = \frac{\bar{u}'^{old}(\lambda I_{y'}^2 + 4) - \lambda I_{x'}(I_t + I_{y'}\bar{v}'^{old})}{\lambda (I_{x'}^2 + I_{y'}^2 + \frac{4}{\lambda})}$$
(14)

We established earlier per the problem specifications that $I_{y'} = 0$, so if we plug that in here:

$$u'^{new}(x,y) = \frac{\bar{u}'^{old}(0^2 + 4) - \lambda I_{x'}(I_t + 0\bar{v}'^{old})}{\lambda (I_{x'}^2 + 0^2 + \frac{4}{\lambda})} = 4\bar{u}'^{old} - \frac{\lambda I_{x'}I_t}{\lambda (I_{x'}^2 + \frac{4}{\lambda})}$$
(15)

$$u'^{new}(x,y) = 4\bar{u}'^{old} - \frac{I_{x'}I_t}{I_{x'}^2 + \frac{4}{\lambda}} = \bar{u}'^{old} - \frac{I_{x'}^2\bar{u}'^{old} + I_{x'}I_t}{I_{x'}^2 + \frac{4}{\lambda}}$$
(16)

...which matches what we were seeking to reduce to for $u'^{new}(x,y)$. Doing this for $v'^{new}(x,y)$ we would follow a similar path - we would find that:

$$v'^{new}(x,y) = \bar{v}'^{old} - \frac{I_{y'}}{I_{x'}^2 + I_{y'}^2 + \frac{4}{\lambda}} \left(I_{x'} \bar{u}'^{old} + I_{y'} \bar{v}'^{old} + I_t \right)$$
(17)

...and again, since we were told that $I_{y'} = 0$:

$$v'^{new}(x,y) = \bar{v}'^{old} - \frac{0}{I_{x'}^2 + 0^2 + \frac{4}{\lambda}} (I_{x'}\bar{u}'^{old} + 0\bar{v}'^{old} + I_t) = \bar{v}'^{old} - 0 = \bar{v}'^{old}$$
(18)

...thus proving that:

$$u'^{new}(x,y) = \bar{u}'^{old} - \frac{I_{x'}^2 \bar{u}'^{old} + I_{x'} I_t}{I_x^2 + \frac{4}{\lambda}}$$
(19)

$$v'^{new}(x,y) = \bar{v}'^{old} \tag{20}$$

Problem 2

We can represent an object by its boundary, $(x(s), y(s)), 0 \le s \le S$, where S is the length of the object's boundary and s is distance along that boundary from some arbitrary starting point. Combine x and y into a single complex function z(s) = x(s) + jy(s). The Discrete Fourier Transform (DFT) of z is:

$$Z(k) = \sum_{s=0}^{S-1} e^{-2\pi j \frac{ks}{S}} z(s), 0 \le k \le S - 1$$
(21)

We can use the coefficients Z(k) to represent the object boundary. The limit on s is S-1 because for a closed contour z(S) = z(0). The Inverse Discrete Fourier Transform is:

$$z(s) = \frac{1}{S} \sum_{k=0}^{S-1} e^{+2\pi j \frac{ks}{S}} Z(k), 0 \le s \le S - 1$$
 (22)

\mathbf{A}

Suppose that the object is translated by $(\Delta x, \Delta y)$, that is, $z'(s) = z(s) + \Delta x + j\Delta y$. How is z''s DFT Z'(k) related to Z(k)?

$$Z(k) = \sum_{s=0}^{S-1} e^{-2\pi i j \frac{ks}{S}} z(s), 0 \le k \le S - 1$$
 (23)

Which we can define z'(s) as:

$$z'(s) = z(s) + \Delta x + j\Delta y \tag{24}$$

If we plug this into our original equation, we get...

$$Z'(k) = \sum_{s=0}^{S-1} e^{-2\pi j \frac{ks}{S}} z(s) + \sum_{s=0}^{S-1} e^{-2\pi j \frac{ks}{S}} (\Delta x + j\Delta y)(1)$$
 (25)

Here we see that we have a segment that is equivalent to oru defined Z(k), so we can simplify by expressing:

$$Z(k) + (1)(\Delta x + j\Delta y) \sum_{s=0}^{S-1} e^{-2\pi j \frac{ks}{S}}$$
 (26)

With Δx and $j\Delta y$ isolated, we can use a table of known Fourier Transforms to identify the resulting conversion. Based on our problem's definition ΔX and Δy are both constants, so we can state that:

$$Z'(k) = Z(k) + (\Delta x + j\Delta y)\sigma(\frac{2\pi k}{S})$$
(27)

...where $\frac{1}{S}$ acts as our scaling factor.

\mathbf{B}

In this section, we are asked to suppose that the object is scaled by an integer constant c, that is z'(s) = cz(s). For simplicity, we are to assume that S' = S. How is Z'(k) as the DFT of z' related to Z(k)? Starting with our definition of Z(k), which is our Fourier Transform of z(s) as defined earlier:

$$Z(k) = \sum_{s=0}^{S-1} e^{-2\pi j \frac{ks}{S}} z(s), 0 \le k \le S - 1$$
 (28)

If we then define z'(s) as above, z'(s) = cz(s), we can plug it into our Fourier Transform:

$$Z'(k) = \sum_{s=0}^{S-1} e^{-2\pi j \frac{ks}{S}} cz(s)$$
 (29)

Since c is a constant, we know we can pull it in front of the fourier Transform like so:

$$Z'(k) = c \sum_{s=0}^{S-1} e^{-2\pi j \frac{ks}{S}} z(s)$$
 (30)

...and further simplified:

$$Z'(k) = cZ(k) \tag{31}$$

...which shows that multiplying our time series function by a given scalar simply multiples our frequency domain by the same scalar.

\mathbf{C}

We wish to know what object has $z(s) = \left[x_0 + R\cos(\frac{2\pi s}{s})\right] + j\left[y_0 + R\sin(\frac{2\pi s}{s})\right]$. Below we created a graph drawing the resulting shape, and include the code utilized to generate it. Arbitrary values were chosen for S, r, x_0 , and y_0 for the sake of plotting.

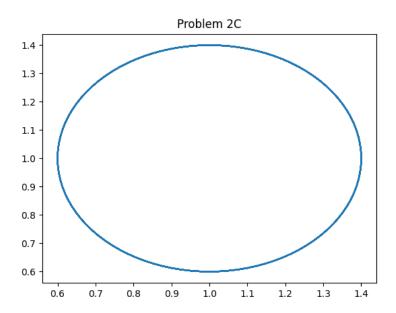


Figure 1: Our resulting shape

```
import numpy as np
from numpy import cos, sin, pi
import matplotlib.pyplot as plt
figure = plt.figure()
plt.title("Problem 2C")
S = 10
r = 4
x0 = 1
y0 = 1
theta = [theta for theta in np.arange(0, S, 0.01)]
        x0 + r * cos(2*pi*theta)/S
        for theta in theta
        y0 + r * sin(2*pi*theta)/S
        for theta in theta
# Plot the results
plt.plot(X, Y)
plt.savefig("./imgs/prob2_c.png")
```

\mathbf{D}

What is Z(k) corresponding to z(s) from Part C? To do this, we begin with our z(s):

$$z(s) = \left[x_0 + R\cos(\frac{2\pi s}{s})\right] + j\left[y_0 + R\sin(\frac{2\pi s}{s})\right]$$
(32)

We can utilize the inverse of Euler's formula; $\cos(x) = \frac{e^{ix} + e^{-ix}}{2}$ and $\sin(t) = \frac{e^{ix} - e^{-ix}}{2x}$. This allows us to expand our starting equation:

$$z(s) = x_0 + R \frac{e^{j\frac{2\pi s}{S}} + e^{-j\frac{2\pi s}{S}}}{2} + jy_0 + R \frac{e^{j\frac{2\pi s}{S}} - e^{-j\frac{2\pi s}{S}}}{2}$$
(33)

$$z(s) = x_0 + jy_0 + \frac{R}{2} \left(e^{j\frac{2\pi s}{S}} + e^{-j\frac{2\pi s}{S}} + e^{j\frac{2\pi s}{S}} - e^{-j\frac{2\pi s}{S}} \right)$$
(34)

$$z(s) = x_0 + jy_0 + \frac{R}{2} \left(2e^{j\frac{2\pi s}{S}} \right)$$
 (35)

$$z(s) = x_0 + jy_0 + Re^{j\frac{2\pi s}{S}}$$
(36)

The Discrete Fourier Transform (DFT) from earlier in our problem we can begin to expand this equation now. Starting with:

$$Z(k) = \sum_{s=0}^{S-1} e^{-2\pi j \frac{ks}{S}} z(s), 0 \le k \le S - 1$$
(37)

...which can lead us to:

$$Z(k) = \sum_{s=0}^{S-1} e^{-2\pi j \frac{ks}{S}} \left(x_0 + jy_0 + Re^{j\frac{2\pi s}{S}} \right)$$
 (38)

...constants can be pulled out, leaving us with:

$$Z(k) = \sum_{s=0}^{S-1} e^{-2\pi j \frac{ks}{S}} (x_0 + jy_0) + R \sum_{s=0}^{S-1} e^{-2\pi j \frac{ks}{S}} e^{j\frac{2\pi s}{S}}$$
(39)

$$Z(k) = (x_0 + jy_0) \sum_{s=0}^{S-1} e^{-2\pi j \frac{ks}{S}} + R \sum_{s=0}^{S-1} e^{-j2\pi(k-1)} S$$
(40)

We know from a Fourier Transform lookup table we can then convert this to:

$$Z(k) = (x_0 + jy_0)\sigma(\frac{2\pi k}{S}) + \sigma(2\pi \frac{k-1}{S})$$
(41)