# The Dirac equation

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### 1 The Dirac equation

The Dirac Hamiltonian in the presence of an electromagnetic field is

$$H = \alpha \cdot \pi + \beta m + a\phi$$

with the Dirac matrices in the Dirac representation

$$\alpha = \begin{pmatrix} 0 & \sigma \\ \sigma & 0 \end{pmatrix} \qquad \beta = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \qquad H = \begin{pmatrix} m + q\phi & \sigma \cdot \pi \\ \sigma \cdot \pi & -m + q\phi \end{pmatrix}$$

The effective Hamiltonian of the upper-left subspace is

$$\left(m + q\phi + \sigma \cdot \pi \frac{1}{E + m - q\phi} \sigma \cdot \pi\right) \psi = E\psi$$

and let us write  $E = m + \varepsilon$ , such that we obtain

$$\left(q\phi + \sigma \cdot \pi \frac{1}{2m + \varepsilon - q\phi} \sigma \cdot \pi\right) \psi = \varepsilon \psi$$

### 2 The Pauli Hamiltonian

In the case where  $\phi$  vanishes, we have the leading-order correction

$$\frac{1}{2m}\boldsymbol{\sigma} \cdot \boldsymbol{\pi} \, \boldsymbol{\sigma} \cdot \boldsymbol{\pi} = \frac{1}{2m} \boldsymbol{\pi}^2 + \frac{1}{2m} i \boldsymbol{\sigma} \cdot \boldsymbol{\pi} \times \boldsymbol{\pi}$$
$$= \frac{1}{2m} \boldsymbol{\pi}^2 - \frac{q}{2m} \boldsymbol{\sigma} \cdot \nabla \times \mathbf{A} \qquad \boldsymbol{\pi} \equiv \mathbf{p} - q\mathbf{A}$$
$$= \frac{1}{2m} \boldsymbol{\pi}^2 - \frac{q}{2m} \boldsymbol{\sigma} \cdot \mathbf{B}$$

### 3 The spin-orbit coupling

In the case where  $\phi$  does not vanish, we have the leading-order correction

$$q\phi + \frac{1}{2m}\sigma \cdot \pi \sigma \cdot \pi = q\phi + \frac{1}{2m}\pi^2 - \frac{q}{2m}\sigma \cdot \mathbf{B}$$

Now we want the next-to-leading-order correction:

$$\left(q\phi + \frac{1}{2m}\sigma \cdot \pi \sigma \cdot \pi - \sigma \cdot \pi \frac{\varepsilon - q\phi}{4m^2}\sigma \cdot \pi\right)\psi = \varepsilon\psi$$

$$\left(q\phi + \frac{1}{2m}\sigma \cdot \pi \sigma \cdot \pi + \sigma \cdot \pi \frac{q\phi}{4m^2}\sigma \cdot \pi\right)\psi = \varepsilon\left(1 + \frac{1}{4m^2}\sigma \cdot \pi \sigma \cdot \pi\right)\psi$$

such that

$$\left(1 - \frac{1}{8m^2}\sigma \cdot \pi \sigma \cdot \pi\right) \left(q\phi + \frac{1}{2m}\sigma \cdot \pi \sigma \cdot \pi + \sigma \cdot \pi \frac{q\phi}{4m^2}\sigma \cdot \pi\right) \left(1 - \frac{1}{8m^2}\sigma \cdot \pi \sigma \cdot \pi\right) \\
\left(1 + \frac{1}{8m^2}\sigma \cdot \pi \sigma \cdot \pi\right) \psi = \varepsilon \left(1 + \frac{1}{8m^2}\sigma \cdot \pi \sigma \cdot \pi\right) \psi$$

where the effective Hamiltonian is the first line

$$\begin{split} H_{\mathrm{eff}} &= q\phi + \frac{1}{2m}\sigma \cdot \pi \, \sigma \cdot \pi + \sigma \cdot \pi \, \frac{q\phi}{4m^2} \, \sigma \cdot \pi - \frac{1}{8m^2} \, \sigma \cdot \pi \, \sigma \cdot \pi \, q\phi - \frac{1}{8m^2} \, q\phi \, \sigma \cdot \pi \, \sigma \cdot \pi \\ &= q\phi + \frac{1}{2m} \, \sigma \cdot \pi \, \sigma \cdot \pi - \frac{1}{8m^2} \, \sigma \cdot \pi \, \left[ \sigma \cdot \pi, q\phi \right] - \frac{1}{8m^2} \left[ q\phi, \sigma \cdot \pi \right] \sigma \cdot \pi \end{split}$$

and the final result is

$$H_{\text{eff}} = q\phi + \frac{1}{2m}\pi^2 - \frac{q}{2m}\sigma \cdot \mathbf{B} + \frac{1}{8m^2}q\nabla^2\phi + \frac{1}{4m^2}\sigma \cdot q\nabla\phi \times \pi$$

**Remark** We also have the relativistic correction

$$-\frac{1}{8m^3}\sigma \cdot \pi \sigma \cdot \pi \sigma \cdot \pi \sigma \cdot \pi \sim -\frac{1}{8m^3}p^4$$

# 4 Transformation properties

The following is what we want:

$$i\gamma^{\mu} \frac{\partial}{\partial x^{\mu}} \psi(x) = m\psi(x)$$
  
$$\Rightarrow i\gamma^{\mu} \frac{\partial}{\partial x^{\mu}} D(\Lambda) \psi(\Lambda^{-1} x) = mD(\Lambda) \psi(\Lambda^{-1} x)$$

and we write to first order

$$D(\Lambda)\psi(\Lambda^{-1}x) = \left(1 + \frac{i}{2}\omega_{\mu\nu}S^{\mu\nu}\right)\left(1 + \frac{i}{2}\omega_{\mu\nu}L^{\mu\nu}\right)\psi(x)$$

and generally we have

$$U(\Lambda, a) = \left(1 + \frac{i}{2} \omega_{\mu\nu} J^{\mu\nu} - i a_{\rho} P^{\rho}\right)$$