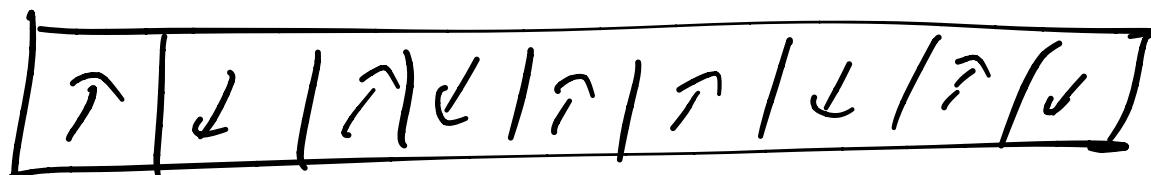


Helix-Coil Model

Previously discussed Ising model,
a good model for phase transitions
and many other processes



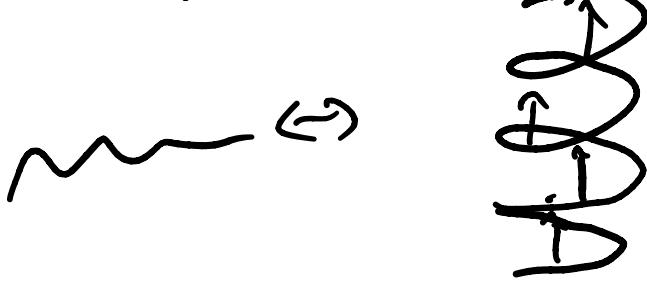
Spins $s_i = \pm \frac{1}{2}$

$$E = \sum_{i=1}^N -J s_i s_{i+1} - h s_i$$

neighbor coupling field, energy

What is H-C model?

α -Helix is a common SS-element in proteins



consider each residue as M or C conformer

doesn't make much sense for each residue to be independent, but consider for now

Lets say $E_{coil} = 0$

$$E_H = -E \text{ (favorable)} \quad , \quad \downarrow \text{sum over states}$$

for one residue $g = \sum_{n=0}^{\infty} e^{-\beta E_n}$ (2 state model)

$$= 1 + e^{+\beta E}$$

$$P_H = \frac{e^{+\beta E}}{g} \quad P_C = \frac{1}{g}$$

$$K = P_H/P_C = e^{+\beta E}$$

(think of ΔE as ΔG
see book for more details)

$$g_{site} = 1 + k$$

$$\text{and } P_H = \frac{k}{1+k}$$

Connection to Ising model

$$\Delta E = h \propto$$

$$K = \frac{e^{+\beta h_L}}{e^{-\beta h_L}} = e^{\beta L} \quad \uparrow \text{Ising} \downarrow \text{Ising}$$

Now from partition function perspective

$$q = 1 + k$$

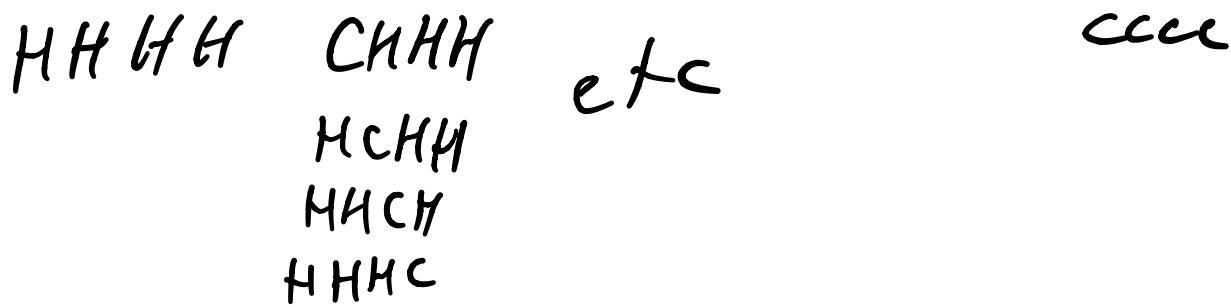
N residues: $Q = q^N = (1+k)^N$

Consider short peptide as example

e.g. 4 sides

$$Q = (1+k)^4 = 1 + 4k + 6k^2 + 4k^3 + k^4$$

coeffs $\binom{N}{n}$ number of HC strings



$$= \sum_{n \geq 0} e^{-\beta E_n} \quad \text{if counted all states}$$

$$= \sum_{\substack{\varepsilon \\ \infty \\ \varepsilon}} w(\varepsilon) e^{-\beta \varepsilon} \quad \varepsilon = -\gamma_H \cdot \epsilon$$

$\gamma_H = 0 \rightarrow N$ list by energy

How do we get $\langle f_H \rangle$ from partition func

$$\text{See that } f_H = \frac{n_H}{N} \left(= \frac{E_{\text{total}}}{NE} \right)$$

$$\text{if } Q = \sum_{n_H=0}^N \binom{N}{n_H} e^{+\beta E n_H} = \sum_{n_H=0}^N \binom{N}{n_H} k^{n_H}$$

$$f_H = \left\langle \frac{n_H}{N} \right\rangle = \sum_{n_H=0}^N \frac{n_H}{N} P(n_H) = \frac{\sum_{n_H=0}^N \binom{N}{n_H} \left(\frac{n_H}{N}\right) k^{n_H}}{Q}$$

$$\frac{\partial \ln Q}{\partial k} = \sum_{n_H=0}^N \binom{N}{n_H} n_H k^{n_H-1} = \frac{N}{k} f_H$$

$$f_H = \frac{k}{N} \frac{\partial \ln Q}{\partial k} \quad \text{but } Q = (1+k)^N$$

$$= \frac{k}{1+k} \text{ as before for one site}$$

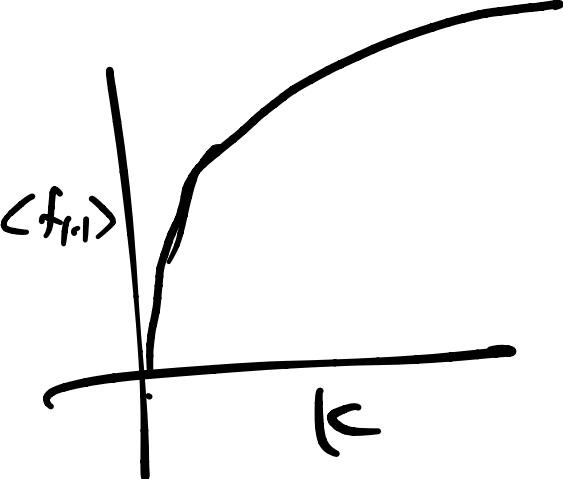
How does this connect to energy descriptions

$$\langle E_{\text{tot}} \rangle = -\langle n_H \rangle E$$

$$\text{and } \langle E \rangle = -\frac{\partial \ln Q}{\partial \beta} = -\frac{\partial \ln Q}{\partial K} \cdot \frac{\partial K}{\partial \beta}$$

$$\text{if } K = e^{+\beta E}$$

$$\frac{\partial K}{\partial \beta} = +E e^{+\beta E} = +E K$$



$$\text{so } \langle E \rangle_{\text{tot}} = -E K \frac{\partial \ln Q}{\partial K}$$

$$\text{so } f_H = \langle n_H / N \rangle = -\frac{\langle E_{\text{tot}} \rangle}{N E} = \frac{K}{N} \frac{\partial \ln Q}{\partial K}$$



Heteropolymers:

proteins are chains of AA with different identities - 20 natural AAs

Each site could be different and at this independent level, different k's

So $k_A, k_B \dots$

$$q_i = (1 + k_i)$$

and $Q_N = \prod_{i=1}^N q_i$

e.g. $(1 + k_A)^{N_A} (1 + k_B)^{N_B}$ for
two types

Interacting model

Should really depend on neighbours

right now $C H C H H C$

$$\beta F = (1)(k)(1)(k)(k)(1) = k^3$$

But can add coupling γ for neighbouring H

$$\beta F = (1)(k)(1)(k\gamma)(k)C = k^3\gamma$$

adding $\tilde{\tau} > 1$ makes stretches of helices

CF Ising model $\tilde{\tau}$ is like J

but only 0's & 1's

first prob-exact analogy

Ziffer model can solve in 1d exact partition function - first physical approx

Suppose $\tilde{\tau}$ very large so eg a chain w/ 3 helices much more likely to have CCCHHH than HCCCHC
eg

Contiguous

$$\# = N - n_H + 1$$

for $n_H = 1 \rightarrow N$
(0, $\# = 1$)

$$\text{So } Q = 1 + \sum_{n_H=1}^N (N-n_H+1) K^{n_H} z^{n_H-1}$$

$$= 1 + \frac{1}{z} \sum_{n_H=1}^N (N-n_H+1) (kz)^{n_H}$$

Skip

$$= 1 + \frac{1}{z} \left[(N+1) \sum (kz)^{n_H} - \sum n_H (kz)^{n_H} \right]$$

$$-k \sum z^{n_H} \frac{d}{dk} k^{n_H}$$

"

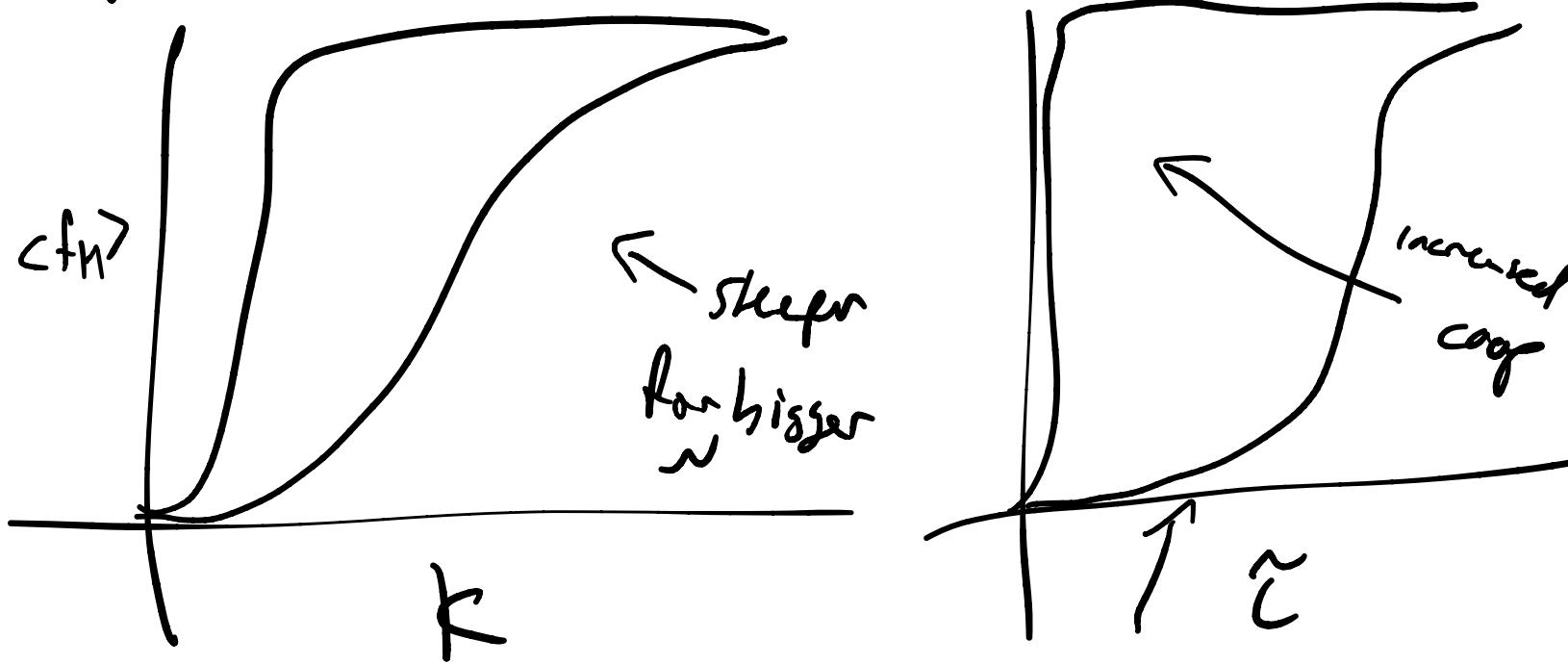
$$-k \frac{\partial}{\partial k} \left(\sum (z^{n_H} k^{n_H}) \right)$$

$$\sum_{n=1}^N x^n = \frac{(x^N - 1)x}{x - 1} \quad \leftarrow$$

then can get f_H as $\frac{K}{N} \frac{\partial \ln Q}{\partial k}$ still

complicated formula!

physically



$\tilde{c} = 1$ curve is wrong
we're in high \tilde{c} limit
for solution

How do we solve exactly?

Method called transfer matrices

partition function has 2^n terms

$$Q = \sum_{x_1=0,1} \sum_{x_2=0,1} \dots \sum_{x_N=0,1} \frac{N}{T} k^{x_1} c^{x_1 x_{i+1}}$$

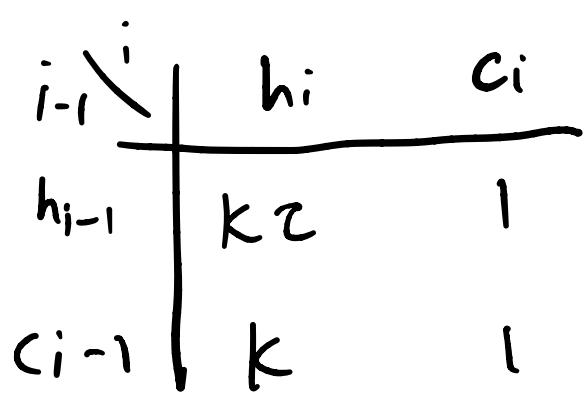
2^n terms

$$x_i = \begin{cases} 0 & \text{helix} \\ 1 & \text{coil} \end{cases}$$

matrix mult $\begin{pmatrix} a & b \\ c & d \end{pmatrix} \begin{pmatrix} w & x \\ y & z \end{pmatrix}$

$$= \begin{pmatrix} aw + bx & ax + bz \\ cw + dy & cx + dz \end{pmatrix} \quad \text{Creates sums}$$

transformation matrix represents possibilities for neighbors



$$\omega = \begin{pmatrix} k^z & 1 \\ k & 1 \end{pmatrix}$$

value at pos i

Every time we multiply ω
 we get sums of terms representing
 Q , if first pos is a helix or coil

These end up in bottom row, easiest
 to see w/ example

$$\omega = \begin{pmatrix} k^2 & 1 \\ k & 1 \end{pmatrix}$$

\curvearrowleft first is helix \curvearrowright first is coil

$$Q = q = 1 + k$$

Chain of \mathcal{L} :

CC	CH	HC	HH
$1 + k$		$k + 2k^2$	

$$\omega^2 = \begin{pmatrix} k^2 & 1 \\ k & 1 \end{pmatrix} \begin{pmatrix} k^2 & 1 \\ k & 1 \end{pmatrix}$$

$$= \begin{pmatrix} - & - \\ k^2 + k & k + 1 \end{pmatrix}$$

And this continues. To get bottom row sum

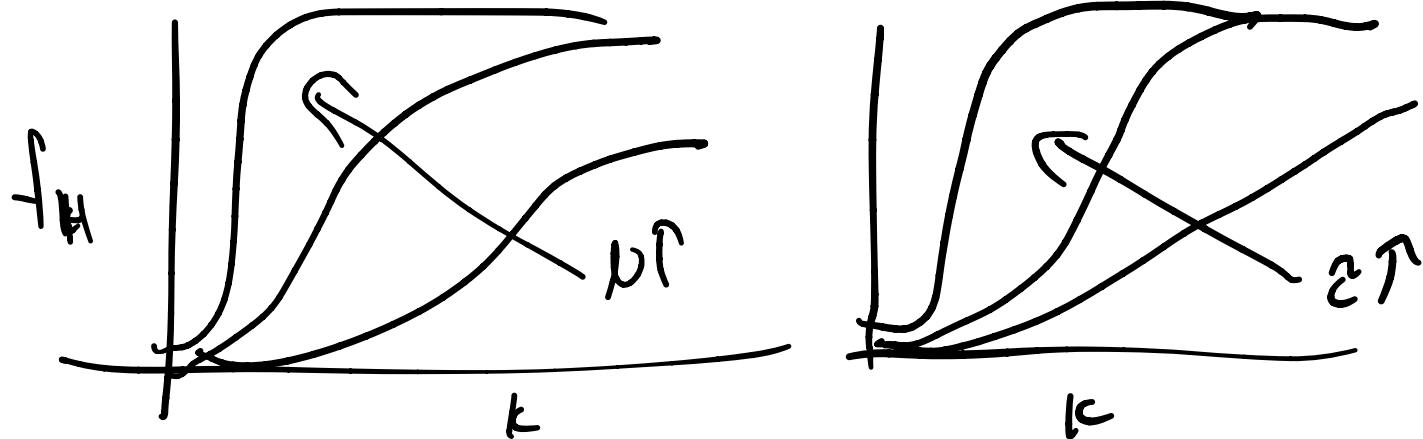
$$(0 \ 1) \begin{pmatrix} a & b \\ c & d \end{pmatrix} = (0+c \quad 0+d) = (c \ d)$$

$$(c \ d) \begin{pmatrix} 1 \\ 1 \end{pmatrix} = c + d$$

$$Q = (0 \ 1) \omega^N(1)$$

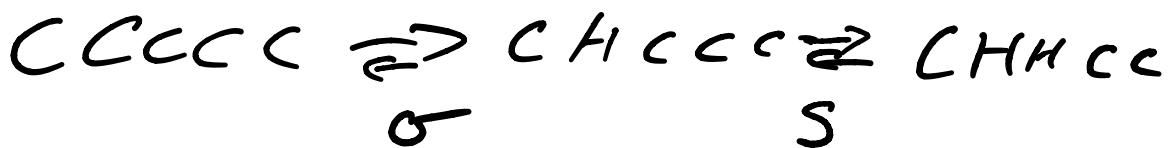
can generate on computer

and get $f_H = \frac{K}{N} \frac{\partial \ln Q}{\partial K}$



Other versions / extns

e.g. Zimm-Bragg



σ is cost of lone helix

$$0 < \sigma < S$$

S is bonus for adjacent H

$$\omega = \begin{pmatrix} k\tau & 1 \\ k & 1 \end{pmatrix}$$

$$\omega - I\lambda = \begin{pmatrix} k\tau - \lambda & 1 \\ k & 1 - \lambda \end{pmatrix}$$

$$\begin{aligned} 11 &= (k\tau - \lambda)(1 - \lambda) - k \\ &\geq \lambda^2 - \lambda(k\tau + 1) + k(\tau - 1) \end{aligned}$$

$$\lambda_2 = \frac{(k\tau + 1) \pm \sqrt{(k\tau + 1)^2 - 4k(\tau - 1)}}{2}$$