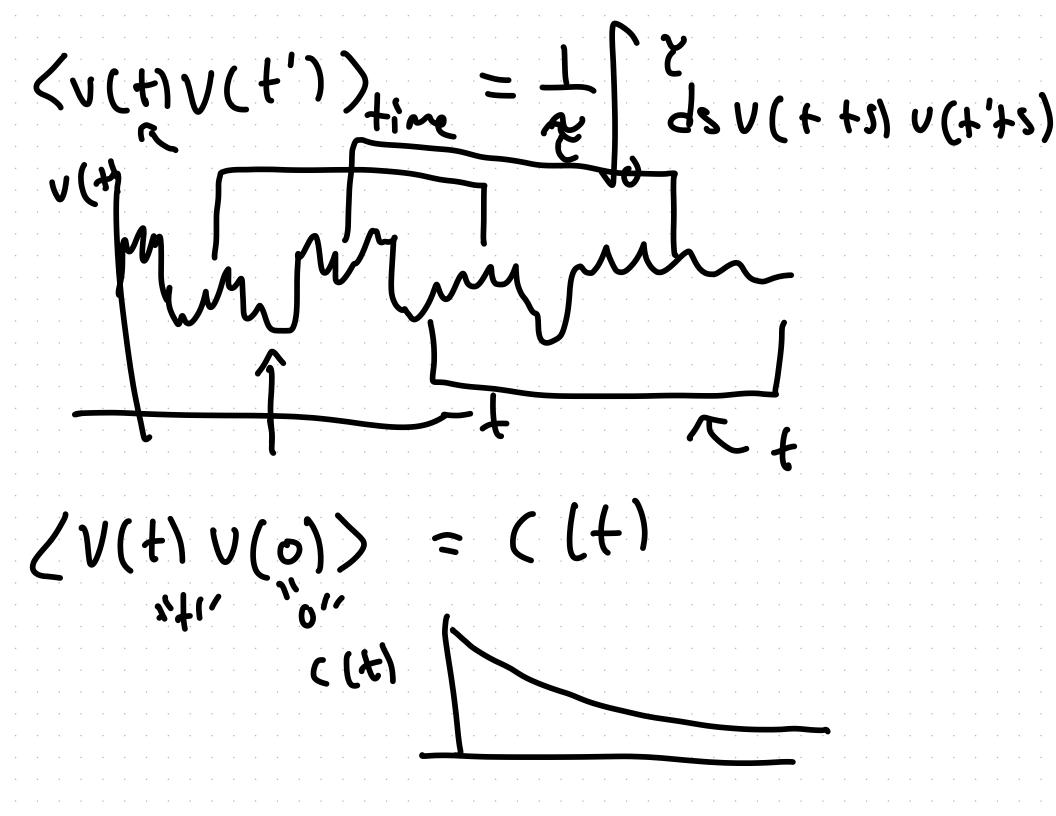
Mon eq -p+2

Langevin equation

$$\frac{dv}{dt} = -\frac{k}{4}v - \nabla N + SF(t)$$
 $V(t) = V(0)e + \frac{1}{4}\int_{0}^{t} dt'e SF(t')$
 $\frac{d}{dt} \langle Sx(t)^{2} \rangle = 2\int_{0}^{t} \langle V(N)V(0) \rangle du = 2D$

MSD $\propto t$ (at long times)



want V(t) $V(\omega)$ want $\int_{0}^{t} \langle V(u)V(\omega) \rangle d\omega$

 $V(t) = V(0) & + \frac{1}{m} \int_{0}^{t} dt' e SF(t')$ accomption that t = 0 is in infinite prest...

in equilibrium

 $\langle v(t) v(t') \rangle = \frac{1}{2} \int_{0}^{2} ds \int_{0}^{2} du, \int_{0}^{2} du_{2}$ $v(t) = \int_{0}^{2} SF(u_{1}) e^{-u_{1}\xi/m} du_{1}$ $\frac{1}{m^2} SF(t-uts) SF(t-uzts)$ replace w/ averge 2BS(+-u,-++uz)

Time average L-> Equilibrium Average $\langle v(t)v(t')\rangle = \frac{c_B^T}{m}e$ $\langle \Delta x(\pi)^2 \rangle = \int_0^{\pi} dt 2 \int_0^{t} \langle u(w)v(o) \rangle dw$

$$\left(\Delta x (2)^{2} \right) = \int_{0}^{2} dt \, 2 \int_{0}^{t} \frac{K_{B}T}{m} e^{-\frac{E}{m}u} \, du$$

$$= \int_{0}^{7} dt \, \frac{2K_{B}T}{m} \int_{\frac{E}{E}} e^{-\frac{E}{m}u} \, du$$

$$= \int_{0}^{7} dt \, \frac{2K_{B}T}{m} \int_{\frac{E}{E}} \left[1 - e^{-\frac{E}{m}u} \right]$$

$$= \frac{2K_{B}T}{E} \left[2 - \frac{m}{E} + \frac{n}{E} e^{-\frac{E}{E}T/m} \right]$$

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at large T MSD(2)2 ZKBT. 2 e before ~ 2 D Einstein relation Einskin Seff D= kst d:ffusion units of D= m² 2/4 energy force/ml

E = 6tt 2a stokes formula E = malius viscosity $D = \frac{K_BT}{6\pi^2 a} \Rightarrow D \cdot 2 = \frac{k_BT}{6\pi a}$

Stokes- Etaskin relation

* Note: violated in 'Glasses'

$$\frac{\langle \Delta x(2)^2 \rangle}{\langle \xi \rangle} = \frac{2 \langle \xi_3 T | 2 - \frac{m}{\xi} + \frac{m}{\xi} e^{-\frac{\xi_2 \gamma_m}{\xi}} \rangle}{\langle \xi_1 \rangle}$$

$$\frac{\langle \Delta x(2)^2 \rangle}{\langle \xi_1 \rangle} = \frac{2 \langle \xi_1 \rangle}{\langle \xi_1 \rangle} \left[\frac{2 \langle \xi_1 \rangle}{\langle \xi_1 \rangle} + \frac{2 \langle \xi_1 \rangle}{\langle \xi_1 \rangle} \right]$$

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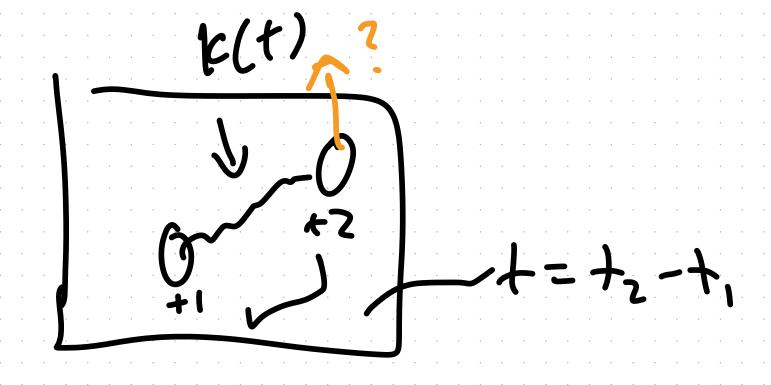
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$$\frac{\langle \Delta x(2) \rangle}{\langle \xi_1 \rangle} = \frac{2 \langle \xi_1 \rangle}{\langle \xi_1 \rangle} + \frac{2 \langle \xi_$$

1: ffusion

Non-Markovian Markovian Means, dynamics only depend on $X = (\overline{q}, \overline{i})$ \mathcal{H} N.M-dependence on history" $-\frac{1}{2}xA \rightarrow -\int_{-\infty}^{\infty} \frac{1}{2} k(t-s)v(s) ds$ $-\frac{1}{2}xA \rightarrow -\int_{-\infty}^{\infty} \frac{1}{2} k(t-s)v(t-s) ds$ $-\frac{1}{2}xA \rightarrow -\int_{-\infty}^{\infty} \frac{1}{2} k(t-s)v(t-s) ds$ $-\frac{1}{2}xA \rightarrow -\int_{-\infty}^{\infty} \frac{1}{2} k(t-s)v(s) ds$



- · Can be non-equilibrium
- · Can be from "coarse graining"

Whole "universe" 5 ysten + bath System Hamiltonian dynamics follows N, U, E bath Integrate out hest in and out N, U, T

K(s)

if memory is exponential

-as

K(s) ~ e

make system merkovian by putling back in degrees of freedom

$$\frac{d\vec{v}}{dt} = -\xi \vec{v} - \nabla u + \delta \vec{F}(t)$$

$$= -\int_{\delta}^{t} ds K(s) \vec{v}(t-s) - \nabla u + \delta \vec{F}(t)$$

Generalized largevis equation

(8FCH8F(+1) = K(+-+1) (508a)eq

is related to the derivative(s)

= matrix

of the "potential"

Harmonic oscilator

 $U(x) = \frac{1}{2}m\omega^2 x^2$

 $u' = \omega^2 x$