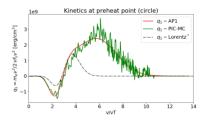
Multi-dimensional High-Order FEM Methods for Electron Kinetics in ALE Hydrodynamics

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SIAM Conference on Computational Science and Engineering (CSE19) February 25 - March 1, 2019 Spokane, Washington, USA

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Motivation - Nonlocal Magneto-Hydrodynamic model (Nonlocal-MHD)

P_N mixed finite element approach to kinetics

4 Conclusions

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Classical MHD

HYDRODYNAMICS
$$local
ightarrow oldsymbol{q}_e = -\kappa_{SH} T_e^{2.5}
abla T_e$$

$$\begin{split} \frac{\mathrm{d}\rho}{\mathrm{d}t} &= -\rho\nabla\cdot\boldsymbol{u}, \\ \rho\,\frac{\mathrm{d}\boldsymbol{u}}{\mathrm{d}t} &= -\nabla(p_i+p_e) + \frac{c}{4\pi}\nabla\times\boldsymbol{B}\times\boldsymbol{B}, \\ \rho\left(\frac{\partial\varepsilon_i}{\partial T_i}\frac{\mathrm{d}T_i}{\mathrm{d}t} + \frac{\partial\varepsilon_i}{\partial\rho}\frac{\mathrm{d}\rho}{\mathrm{d}t}\right) &= -p_i\nabla\cdot\boldsymbol{u} - G(T_i-T_e), \\ \rho\left(\frac{\partial\varepsilon_e}{\partial T_e}\frac{\mathrm{d}T_e}{\mathrm{d}t} + \frac{\partial\varepsilon_e}{\partial\rho}\frac{\mathrm{d}\rho}{\mathrm{d}t}\right) &= -p_e\nabla\cdot\boldsymbol{u} + \nabla\cdot\left(\kappa_{SH}T_e^{2.5}\nabla T_e\right) + \sigma\boldsymbol{E}\cdot\boldsymbol{E} + G(T_i-T_e) + Q_B(\boldsymbol{E}_L), \end{split}$$

MAXWELL EQUATIONS resistive \rightarrow **i** = σ **E**

resistive
$$\rightarrow$$
 j = σ **E**

$$abla imes \mathbf{E} = -rac{1}{c} rac{\mathrm{d} \mathbf{B}}{\mathrm{d} t},$$

$$abla imes \mathbf{B} = rac{4\pi}{c} \sigma \mathbf{E},$$

KINETICS OF ELECTRONS

Landau – Fokker – Planck

$$\frac{\partial f}{\partial t} + \mathbf{v} \cdot \nabla f + (\mathbf{E} + \mathbf{v} \times \mathbf{B}) \cdot \nabla_{\mathbf{v}} f \quad = \quad \Gamma \int \nabla_{\mathbf{v}} \nabla_{\mathbf{v}} (\mathbf{v} - \tilde{\mathbf{v}}) \cdot (f \nabla_{\tilde{\mathbf{v}}} f - f \nabla_{\mathbf{v}} f) \, \mathrm{d}\tilde{\mathbf{v}} + \frac{\nu_{ei}}{2} \frac{\partial^2 f}{\partial \Omega^2}.$$



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Nonlocal-MHD

HYDRODYNAMICS 4D

$$\begin{split} \frac{\mathrm{d}\rho}{\mathrm{d}t} &= -\rho\nabla\cdot\boldsymbol{u}, \\ \rho\frac{\mathrm{d}\boldsymbol{u}}{\mathrm{d}t} &= -\nabla(p_i+p_e)+\boldsymbol{j}(f,\boldsymbol{E})\times\boldsymbol{B}, \\ \rho\left(\frac{\partial\varepsilon_i}{\partial T_i}\frac{\mathrm{d}T_i}{\mathrm{d}t}+\frac{\partial\varepsilon_i}{\partial\rho}\frac{\mathrm{d}\rho}{\mathrm{d}t}\right) &= -p_i\nabla\cdot\boldsymbol{u}-G(T_i-T_e), \\ \rho\left(\frac{\partial\varepsilon_e}{\partial T_e}\frac{\mathrm{d}T_e}{\mathrm{d}t}+\frac{\partial\varepsilon_e}{\partial\rho}\frac{\mathrm{d}\rho}{\mathrm{d}t}\right) &= -p_e\nabla\cdot\boldsymbol{u}-\nabla\cdot\boldsymbol{q}_e(f)+\boldsymbol{j}(f,\boldsymbol{E})\cdot\boldsymbol{E}+G(T_i-T_e)+Q_{\mathrm{IB}}(\boldsymbol{E}_L), \end{split}$$

MAXWELL FOLIATIONS 4D

$$abla imes \mathbf{E} = -rac{1}{c} rac{\mathrm{d} \mathbf{B}}{\mathrm{d} t},$$

$$abla imes \mathbf{B} = rac{4\pi}{c} \mathbf{j}(f, \mathbf{E}),$$

KINETICS OF ELECTRONS 6D

$$\mathbf{v} \cdot \nabla f + (\mathbf{E} + \mathbf{v} \times \mathbf{B}) \cdot \nabla_{\mathbf{v}} f = v \tilde{\nu}_{e} \frac{\partial}{\partial v} (f - f_{MB}(T_{e})) + \nu_{scat} (f_{0} - f).$$

3D velocity space FEM discretizations

AWBS electron kinetic model 7D

$$C_V \frac{\mathrm{d} T_e}{\mathrm{d} t} = -\nabla \cdot \boldsymbol{q}_e(f) + \boldsymbol{j}(f) \cdot \boldsymbol{E} + S_H,$$

$$\boldsymbol{v} \cdot \nabla f + (\boldsymbol{E} + \boldsymbol{v} \times \boldsymbol{B}) \cdot \nabla_{\boldsymbol{v}} f = v \tilde{\nu}_e \frac{\partial}{\partial v} \left(f - f_{MB}(T_e) \right) + \nu_{scat} \left(f_0 - f \right).$$

■ C7 (S_N - discontinuous Galerkin FEM) 3D velocity space in spherical coordinates (v, ϕ, θ)

$$\boldsymbol{n} \cdot \nabla f + \frac{\boldsymbol{E} \cdot \boldsymbol{n}}{v} \frac{\partial f}{\partial v} + \frac{E_{\phi} - v}{v^2} \frac{B_{\theta}}{\partial \phi} \frac{\partial f}{\partial \phi} + \frac{E_{\theta} + v}{v^2 \sin(\phi)} \frac{\partial f}{\partial \theta} = \tilde{\nu_{\theta}} \frac{\partial}{\partial v} \left(f - f_{M} \right) + \frac{\nu_{scat}}{v} \left(f_{0} - f \right).$$

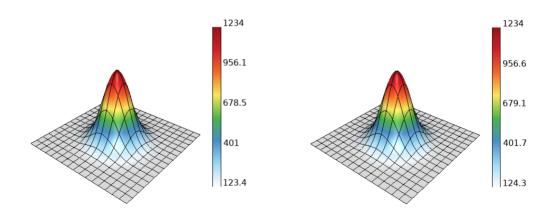
■ AP1 (P_N (VEF $\xi = 1/3$) - continuous Galerkin mixed FEM) *Physicists like it!*

$$\begin{split} \xi \nabla \cdot \mathbf{f_1} + \xi \frac{q_e}{m_e v} \mathbf{E} \cdot \left(\frac{\partial \mathbf{f_1}}{\partial v} + \frac{2}{v} \mathbf{f_1} \right) &= \tilde{v_e} \frac{\partial}{\partial v} \left(\mathbf{f_0} - \mathbf{f_M} \right), \\ \nabla \mathbf{f_0} + \frac{q_e}{m_e v} \mathbf{E} \frac{\partial \mathbf{f_0}}{\partial v} + \frac{q_e \mathbf{B}}{m_e c v} \times \mathbf{f_1} &= \tilde{v_e} \frac{\partial \mathbf{f_1}}{\partial v} - \frac{v_{scat}}{v} \mathbf{f_1}, \end{split}$$

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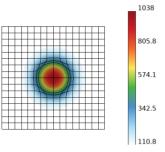
S_N upwind DG, Multigrid Approximate-Ideal-Relaxation (AIR) solver

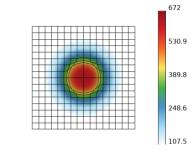
$$\begin{split} \mathbf{M}_{(\tilde{\nu_e} - \frac{E \cdot n_d}{v})} \cdot \frac{\Delta f_d}{\Delta v} &= (\mathbf{n}_d \cdot \mathbf{G} + \mathbf{F}_d) \cdot \left(\tilde{f}_d + \Delta f_d\right) \\ &+ \mathbf{M}_{(\frac{\nu_{SCat}}{v})} \cdot \left(\tilde{f}_d + \Delta f_d\right) \\ &+ \mathbf{S}_{(\nu_{Scat}, E, B)} \cdot \tilde{\mathbf{f}} \\ &+ \mathbf{S}_{(\tilde{\nu_e} \frac{\partial f_M}{\partial v})} \\$$

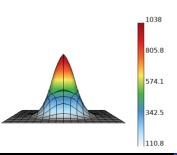


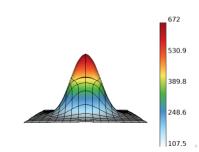
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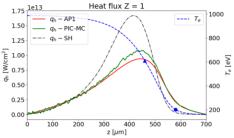


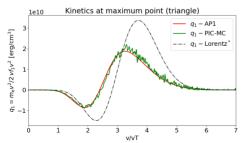


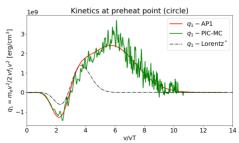
₽ 99€

AP1 formulation, fixed P1 angular discretization PCG(AMG)

$$\begin{array}{lcl} \mathbf{M}_{(\tilde{\nu}e)}^{L_2} \cdot \frac{\mathrm{d}\mathbf{f_0}}{\mathrm{d}v} - \mathbf{V}_{(\frac{\mathcal{E}q_e\mathcal{E}}{m_ev})}^{L_2} \cdot \frac{\mathrm{d}\mathbf{f_1}}{\mathrm{d}v} & = & \mathbf{D}_{(\xi)}^{L_2} \cdot \mathbf{f_1} + \mathbf{M}_{(\frac{\mathcal{E}q_e\mathcal{E}}{m_ev^2})}^{L_2} \cdot \mathbf{f_1} + \mathbf{S}_{(\tilde{\nu}e\frac{\partial \mathbf{f_M}}{\partial v})}^{L_2}, \\ \\ \mathbf{M}_{(\tilde{\nu}e)}^{H_1} \cdot \frac{\mathrm{d}\mathbf{f_1}}{\mathrm{d}v} - \mathbf{V}_{(\frac{q_e\mathcal{E}}{m_ev})}^{H_1} \cdot \frac{\mathrm{d}\mathbf{f_0}}{\mathrm{d}v} & = & \mathbf{G}^{H_1} \cdot \mathbf{f_0} + \mathbf{M}_{(\frac{\nu_{SCalt}}{v})}^{H_1} \cdot \mathbf{f_1} + \mathbf{C}_{(\frac{q_e\mathcal{B}}{m_ev})}^{H_1} \cdot \mathbf{f_1}, \end{array}$$



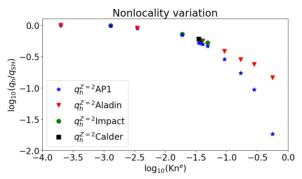




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Electron velocity limit - friction vs. **E** stopping

$$\left(\tilde{\nu_{\text{e}}} - \frac{\textbf{\textit{E}} \cdot \textbf{\textit{n}}}{\textit{\textit{v}}}\right) \frac{\partial \textit{\textit{f}}}{\partial \textit{\textit{v}}} = \textbf{\textit{n}} \cdot \nabla \textit{\textit{f}} + \frac{\nu_{\text{scat}}}{\textit{\textit{v}}} \left(\textit{\textit{f}} - \textit{\textit{f}}_{0}\right) + \frac{\textit{\textit{E}}_{\phi} - \textit{\textit{v}} \; \textit{\textit{B}}_{\theta}}{\textit{\textit{v}}^{2}} \frac{\partial \textit{\textit{f}}}{\partial \phi} + \frac{\textit{\textit{E}}_{\theta} + \textit{\textit{v}} \; \textit{\textit{B}}_{\phi}}{\textit{\textit{v}}^{2}} \frac{\partial \textit{\textit{f}}}{\partial \theta} + \tilde{\nu_{\text{e}}} \frac{\partial \textit{\textit{f}}_{M}}{\partial \textit{\textit{v}}}.$$



E stopping overtakes collisions for Kn > 0.1.

$Kn = \frac{\lambda}{L}$	10 ⁻⁴		10 ⁻²		1
v_{lim}/v_{th}	70.8	22.4	7.3	3.1	1.8

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Adaptive DSA preconditioner for Hydro-Kinetics coupling

AWBS electron kinetic model 7D

$$\begin{split} C_V \frac{\mathrm{d} T_\theta}{\mathrm{d} t} &= -\nabla \cdot \boldsymbol{q}_\theta(f) + \boldsymbol{j}(f) \cdot \boldsymbol{E} + S_H, \\ &= \int_{4\pi} \int_V v \tilde{\nu_\theta} \frac{\partial f}{\partial v} v^4 \mathrm{d} v \mathrm{d} \boldsymbol{n} - \sigma T_\theta^{-0.5} + S_H, \\ \boldsymbol{v} \cdot \nabla f + (\boldsymbol{E} + \boldsymbol{v} \times \boldsymbol{B}) \cdot \nabla_{\boldsymbol{v}} f &= v \tilde{\nu_\theta} \frac{\partial}{\partial v} \left(f - f_{MB}(T_\theta) \right) + \nu_{scat} \left(f_0 - f \right). \end{split}$$

Continuum analysis of local (Kn≪1) transport regime → DIFFUSION

$$C_V \frac{\mathrm{d}T_e}{\mathrm{d}t} = \nabla \cdot \lambda (T_e^{2.5}) \nabla T_e + O(\mathsf{Kn}^2)$$

Fixed-point iteration with preconditioner $\mathbf{E}(\Delta T) = (\mathbf{I} - \mathbf{C}) \cdot \mathbf{M}_{\sigma} (\tilde{T} + \Delta T)^{\alpha} + \mathbf{C} \cdot \mathbf{D}_{\lambda} (\tilde{T} + \Delta T)^{\alpha}$

$$\mathbf{M}_{C_{V}} \cdot \frac{\Delta \mathbf{T_{e}}^{k+1}}{\Delta t} + \mathbf{E}(\Delta \mathbf{T_{e}}^{k+1}) - \mathbf{E}(\Delta \mathbf{T_{e}}^{k}) = \mathbf{K}(f^{k}) - \mathbf{M}_{\sigma}(\tilde{\mathbf{T}_{e}} + \Delta \mathbf{T_{e}}^{k})^{\alpha}.$$

Finally, we get an unconditionally stable backward Euler (SDIRK) fast iterating scheme

$$\mathbf{M}_{C_V} \cdot \frac{\Delta \mathbf{T_e}^{k+1}}{\Delta t} + (\mathbf{I} - \mathbf{C}) \cdot \mathbf{M}_{\sigma} (\tilde{\mathbf{T}} + \Delta \mathbf{T_e}^{k+1})^{\alpha} + \mathbf{C} \cdot \mathbf{D}_{\lambda} (\tilde{\mathbf{T}_e} + \Delta \mathbf{T_e}^{k+1})^{\alpha} = \mathbf{C} \cdot \mathbf{D}_{\lambda} (\tilde{\mathbf{T}_e} + \Delta \mathbf{T_e}^{k})^{\alpha} + \mathbf{K} (t^k) - \mathbf{C} \cdot \mathbf{M}_{\sigma} (\tilde{\mathbf{T}_e} + \Delta \mathbf{T_e}^{k})^{\alpha},$$

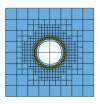
where the adaptive coefficient handles diffusion $\mathbf{C} \xrightarrow{\mathsf{Kn} \ll 1} 1$ and nonlocal transport $\mathbf{C} \xrightarrow{\mathsf{Kn} > 1} 0$ and $\mathbf{C} \in (1,0)$ in between.

T. Haut et al, SIAM, submitted (2018) / arXiv:1810.11082.



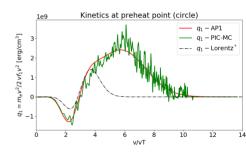
Conclusions

- 7D microscopic world of electrons in hydro simulations.
- S_N high-order DG finite element approach.
- P₁ high-order mixed finite element approach.
- **E** field dominated stopping (P₁ fails).
- Adaptive DSA preconditioner for Hydro-Kinetics coupling (deep learning on **C**).
- Algebraic Multigrid solver pAIR scales $log(P)^{1.22}$.





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MS85 - Developments in Algebraic Multigrid for Nonsymmetric and Hyperbolic Problems - Ben S. Southworth https://github.com/CEED/Laghos/tree/master/amr https://mfem.org



Nonlocal Ohm's law

We found that the quality of the electric field evaluation is essential for a correct plasma modeling

Nonlocal current in plasma

$$\mathbf{j}(f, \mathbf{E}) = \mathbf{e} \int v \mathbf{f_1} v^2 \, \mathrm{d}v = \mathbf{e} \int v \frac{\nu_{ei}^2 \mathbf{E}^* + \omega_B \, \omega_B \cdot \mathbf{E}^* + \nu_{ei} \, \omega_B \times \mathbf{E}^*}{\nu_{ei} (\omega_B^2 + \nu_{ei}^2)} v^2 \, \mathrm{d}v, \tag{1}$$

where $m{E}^* = v rac{v_e}{2} rac{\partial \mathbf{f_1}}{\partial v} - v
abla \cdot \mathbf{f_2} - rac{q_e}{m_e} m{E} rac{\partial f_0}{\partial v}$ is an effective electric field in plasma. A comparison to the *Generalized Ohm's law* shows a correct local behavior of (1), especially that $abla \times
abla \cdot \mathbf{f_2} \sim
abla \cdot \mathbf{f_2} \sim \nabla n_e \times \nabla n_e$

Generalized Ohm's law vs. nonlocal Ohm's law

$$\mathbf{E} = (\mathbf{R}_{T} - \nabla \rho_{e}) + \frac{\mathbf{j}}{\sigma} + \mathbf{j} \times \mathbf{B},$$

$$\mathbf{E} v \frac{q_{e}}{m_{e}} \frac{\partial f_{0}}{\partial v} = v^{2} \nabla \cdot \mathbf{f}_{2} + v^{2} \frac{\nu_{e}}{2} \frac{\partial f_{1}}{\partial v} - v \nu_{ei} f_{1} + v f_{1} \times \frac{q_{e}}{m_{e} c} \mathbf{B}.$$

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