

暗物质II

毕效军

中国科学院高能物理研究所

“2017 年理论物理前沿暑期讲习班——暗物质、
中微子与粒子物理前沿，

2017/7/25

Outline

- 暗物质profile确定
- Subhalo计算
- Axion简介
- 温暗物质sterile neutrino
- Nonthermal DM
- SIDM
- Fuzzy dm

看什么信号?
Gamma,
e+, pbar; 什么
实验探测?

看什么地方? 暗物质
信号, 背景强度, 天
体环境等

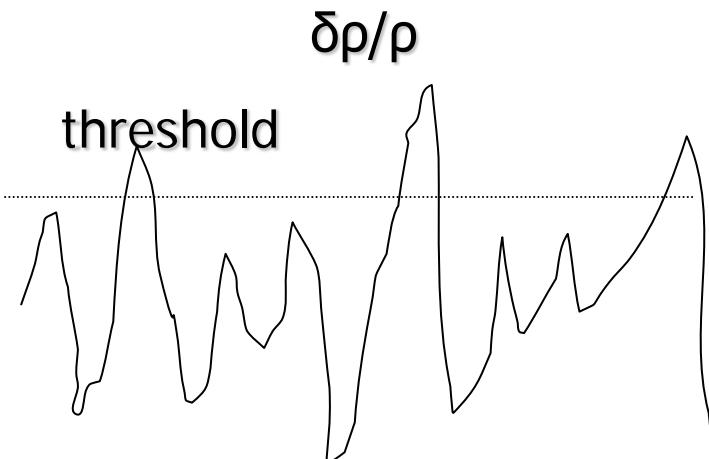
$$\frac{dN}{dE} = \frac{\langle \sigma v \rangle}{2m_\chi^2} \sum_f B_f \frac{dN^f}{dE} \frac{\int \rho^2 dV}{4\pi d^2}$$

粒子物理模型; 相
互强度, 末态?

暗物质在宇宙空间的分布

How does structure form?

- For CDM, structure form hierarchically bottom up, i.e., smaller object forms earlier (at smaller scale, the fluctuation is larger), they merges to form bigger and bigger structures.
- The smaller object have larger concentration parameter, since earlier epoch has greater density.



Structure formation – cold dark matter

Numerical Simulation VS Observation

→ Cold Dark Matter

Distribution of the dark matter

- N-body simulation is extensively adopted to study the evolution of structure with non-linear gravitational clustering from the initial conditions.
- The reliability of an N-body simulation is measured by its mass and force resolution.
- Simulations suggest a *universal* dark matter profile, same for *all masses, epochs and initial power spectra*.

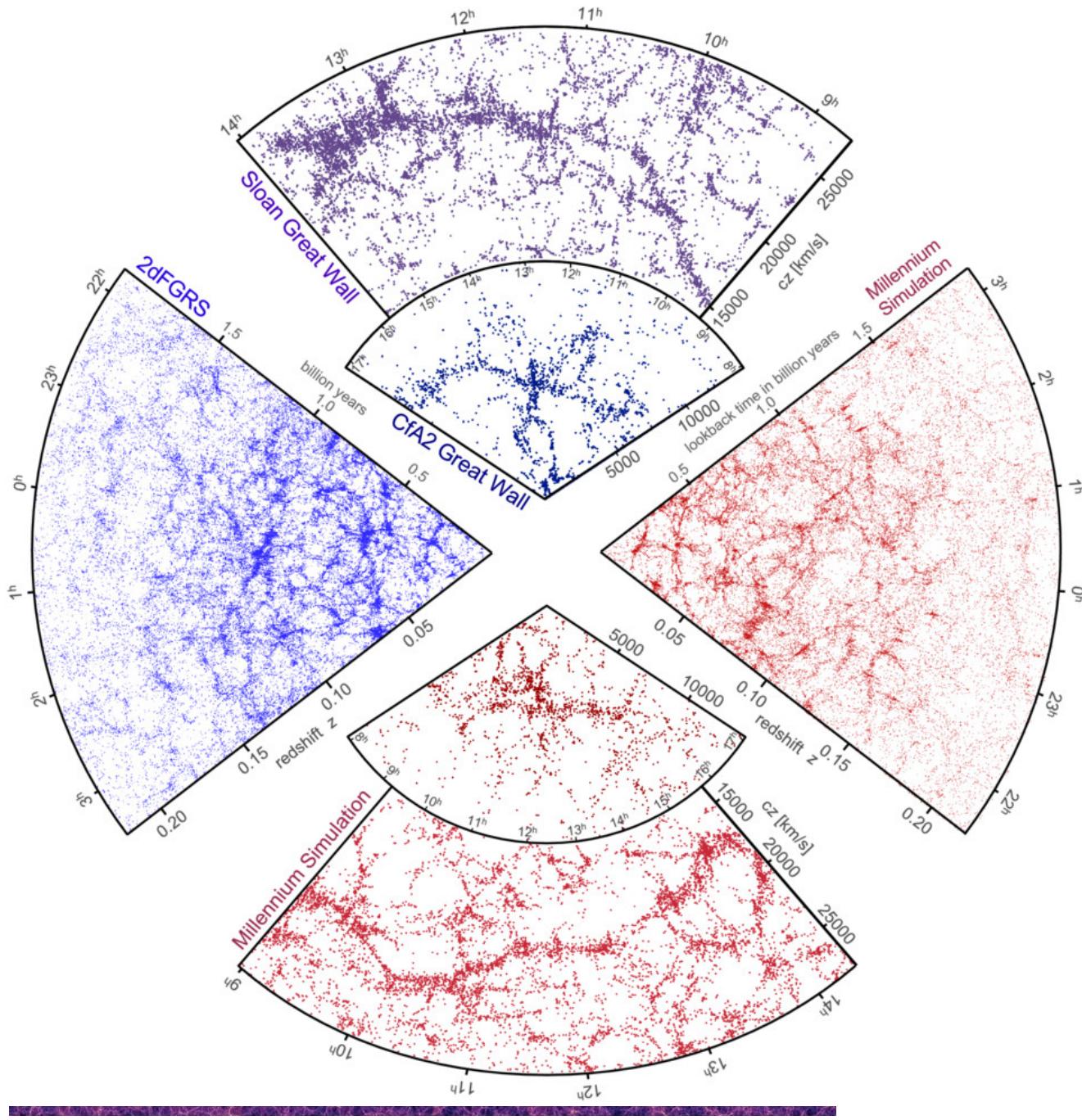
$R = 6.0$ Mpc

$z = 10.155$

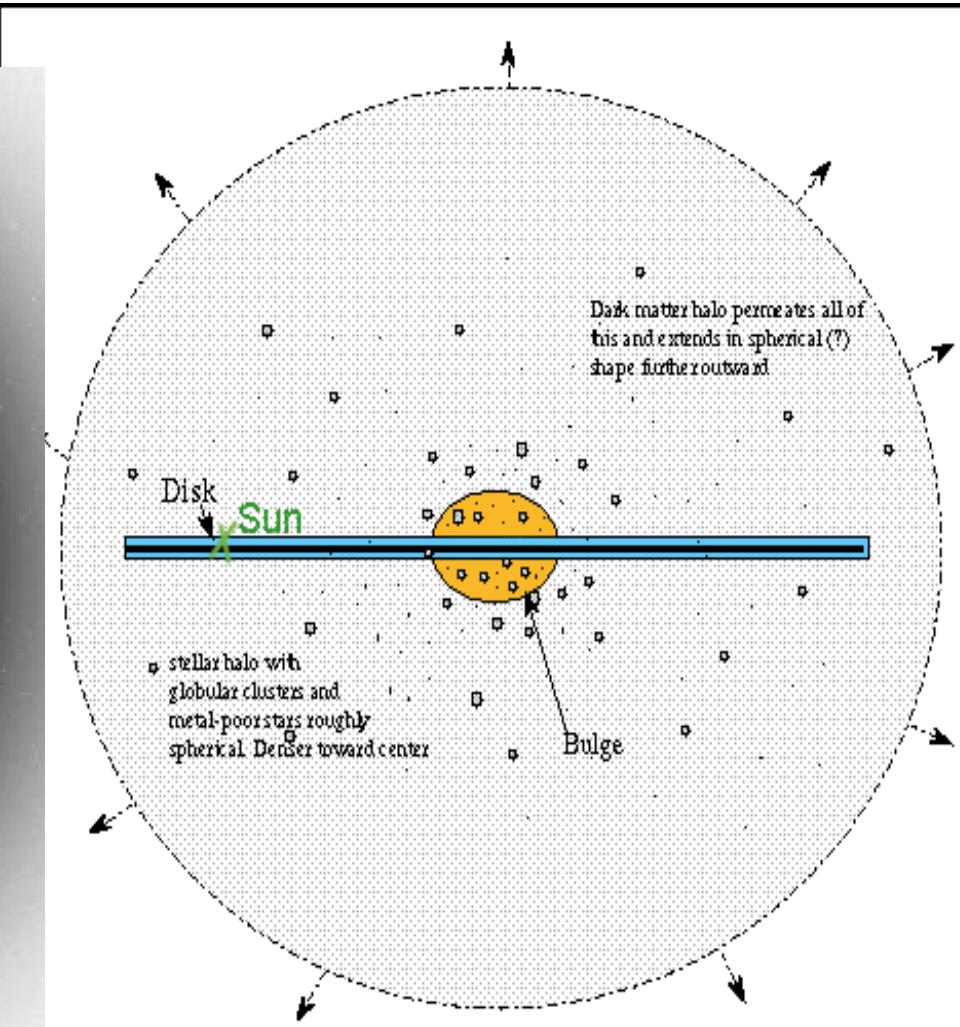
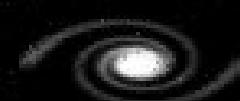


$a = 0.090$

diemand 2003

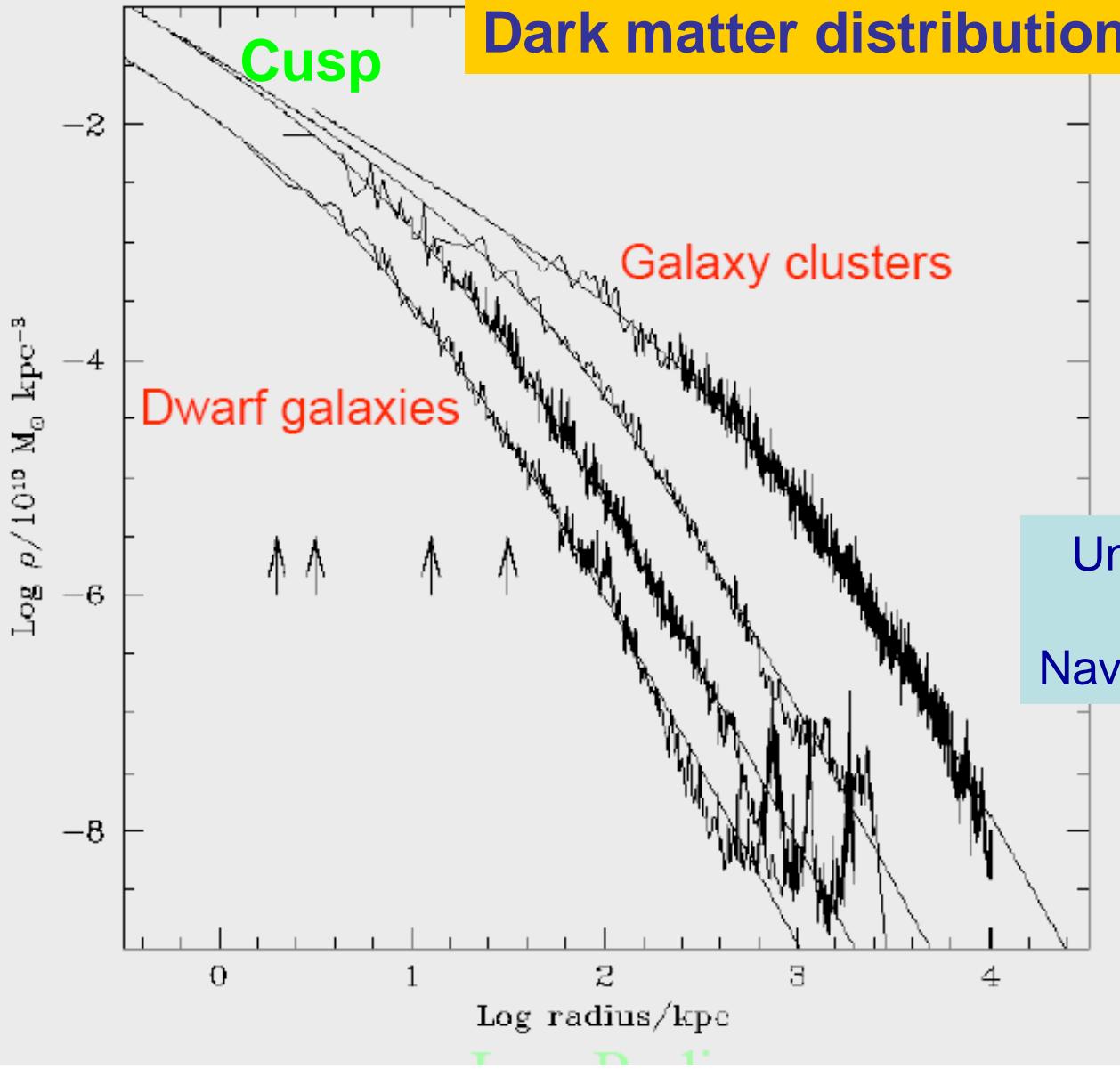


Galaxy configuration



Dark matter halo

Dark matter distribution—Density profile



Universal Density Profile
NFW profile
Navarro, Frenk, White 1997

$$\frac{\rho(r)}{\rho_{crit}} = \frac{\delta_c}{(r/r_s)(1+r/r_s)^2}$$

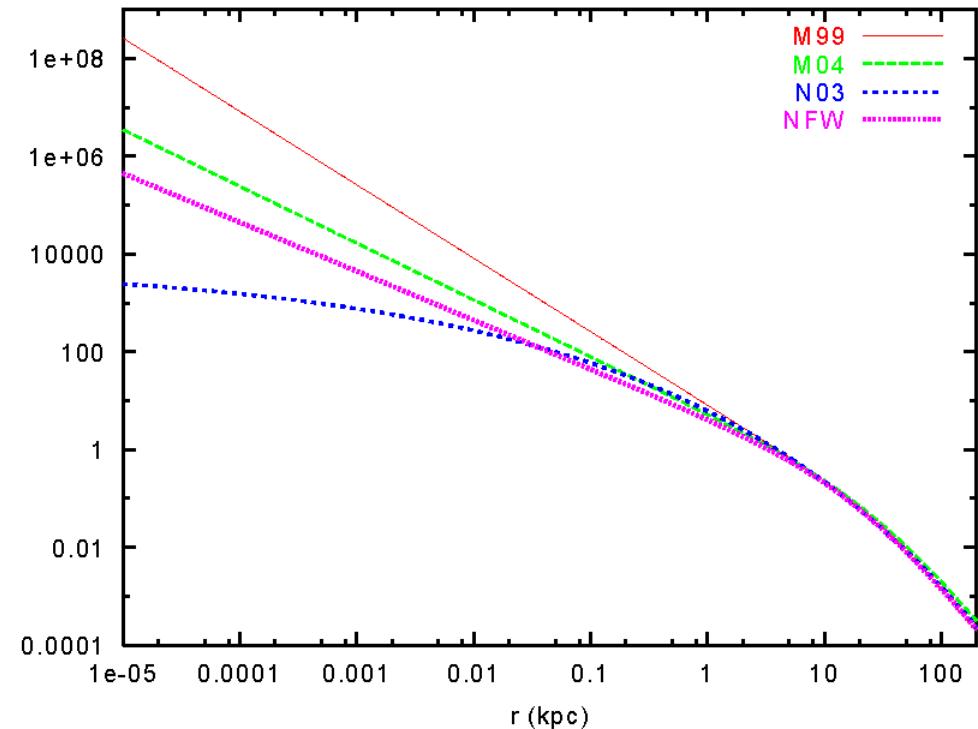
Profiles of dark matter

- Two generally adopted DM profiles are the NFW and Einasto profiles
- They have same density at large radius, while different slope as $r \rightarrow 0$

NFW:
$$\rho_\chi(r) = \frac{\rho_s}{(r/r_s) \left[1 + \left(r/r_s \right)^2 \right]^2}$$
$$\rho_{NFW} \xrightarrow{r \rightarrow 0} r^{-1}$$

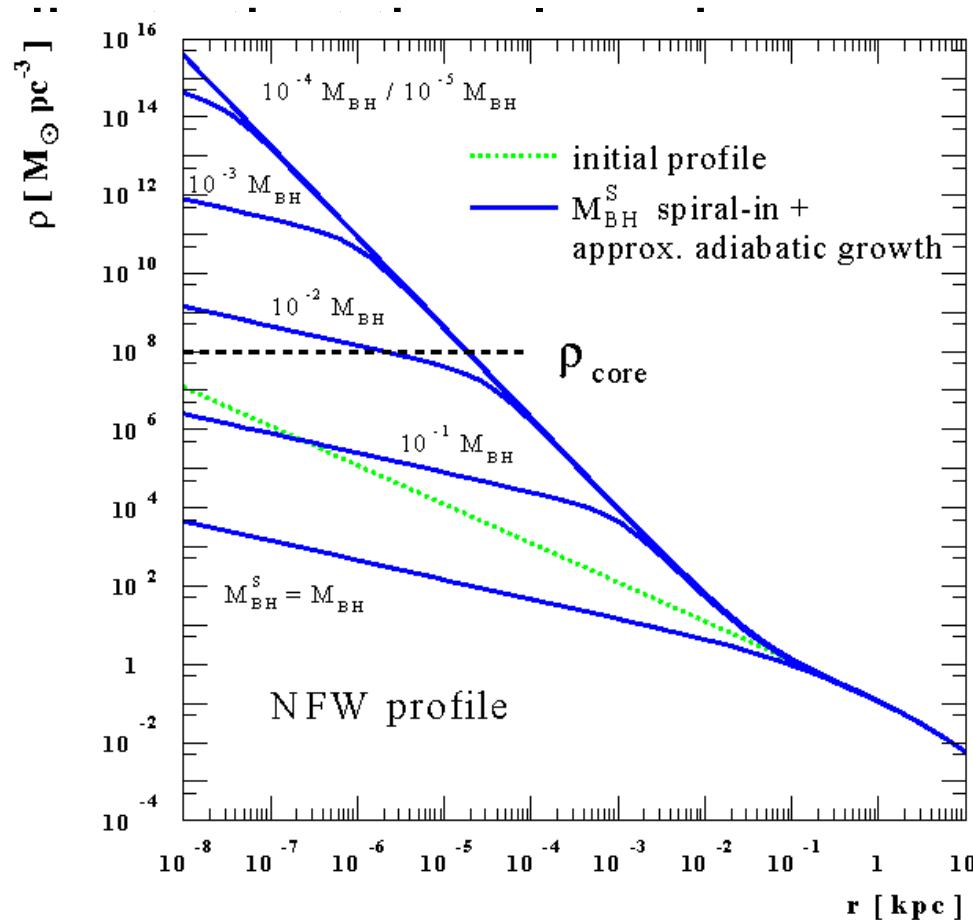
Einasto: $\rho(r) = \rho_s e^{-\frac{2}{\alpha} \left[\left(\frac{r}{r_s} \right)^\alpha - 1 \right]}$

Isothermal: $\rho(r) = \frac{\rho_s}{1 + (r/r_s)^2}$



Some recent developments

- The slope at the inner most radius is under debate. The most recent simulations including baryon matter seem in between the NFW and 0311231, Reed, 03125 be not universal, depe [Reed].
- The central super mas the central cusp heavil initial mass and adiaba

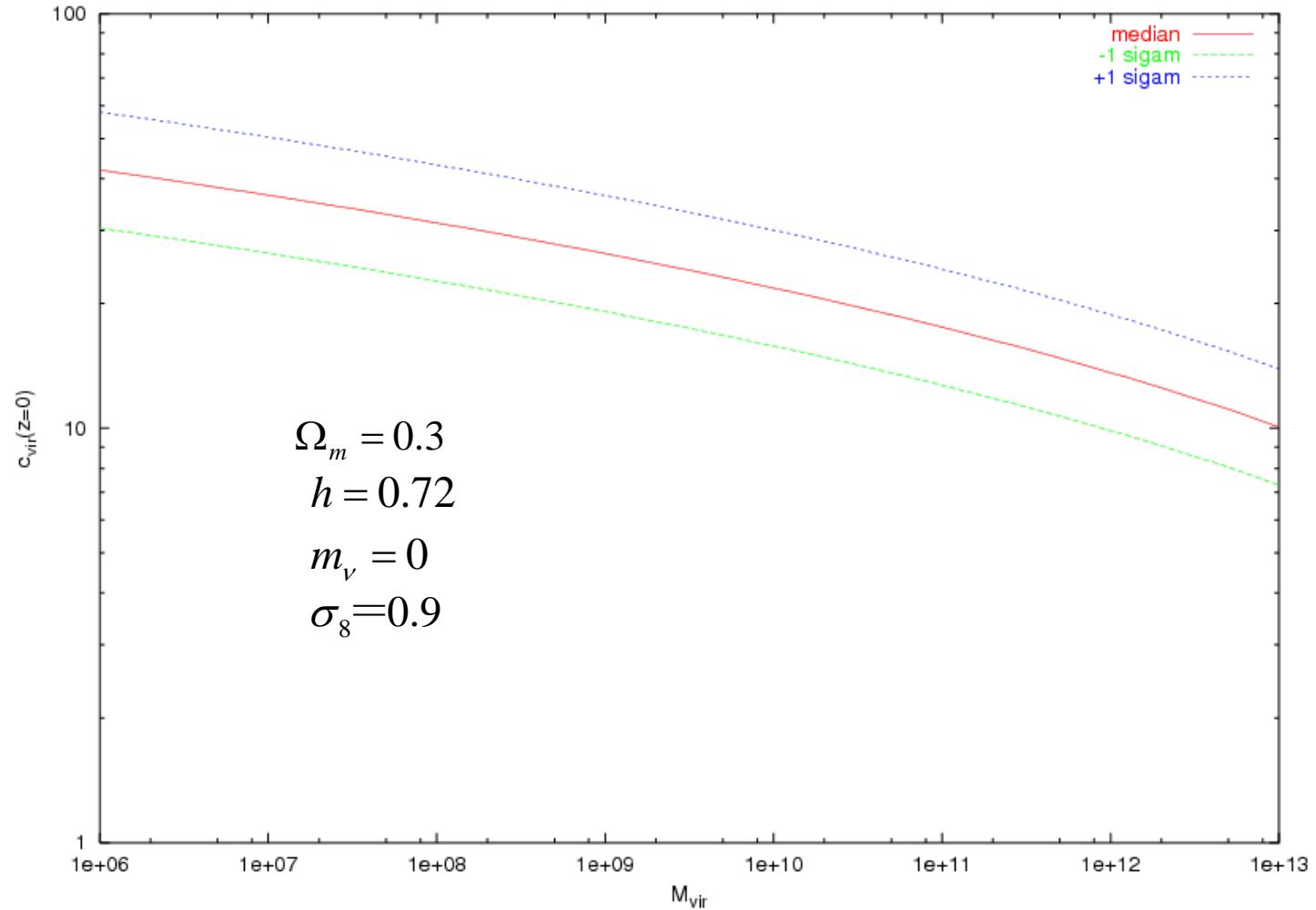


How to determine the dark matter density profile

- The profile is specified by the two scale parameters ρ_s , r_s . They are determined by the virial mass of the structure and the concentration parameter by $c=r_{\text{vir}}/r_s$.
- The concentration parameter represents how the dark matter centrally concentrated. The larger concentration the more centrally concentrated.
- The concentration parameter is determined by the evolution of the structure.

Concentration parameter via the virial mass – fitted to N-body simulation

From the figure, the concentration parameter decreases with the virial mass. It is determined once the cosmological parameters are specified.



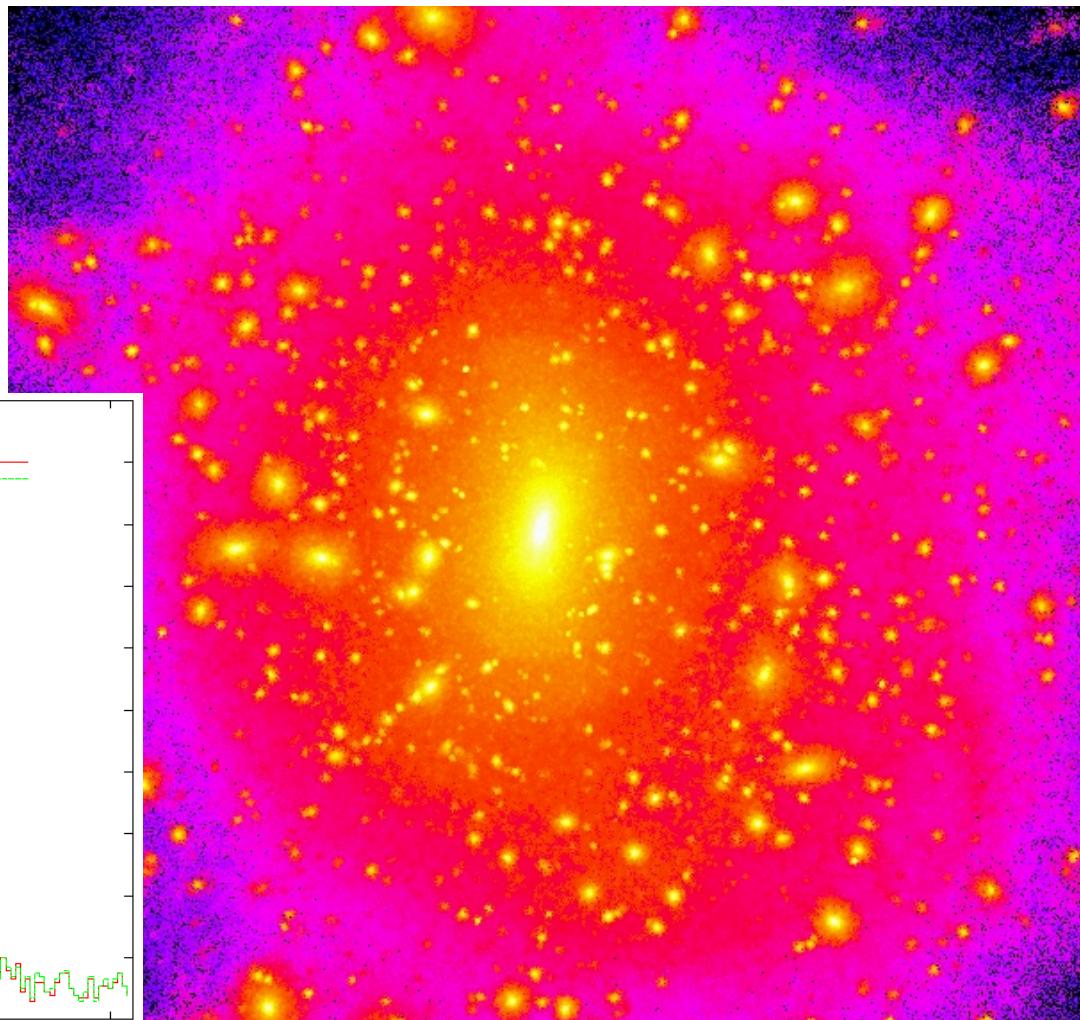
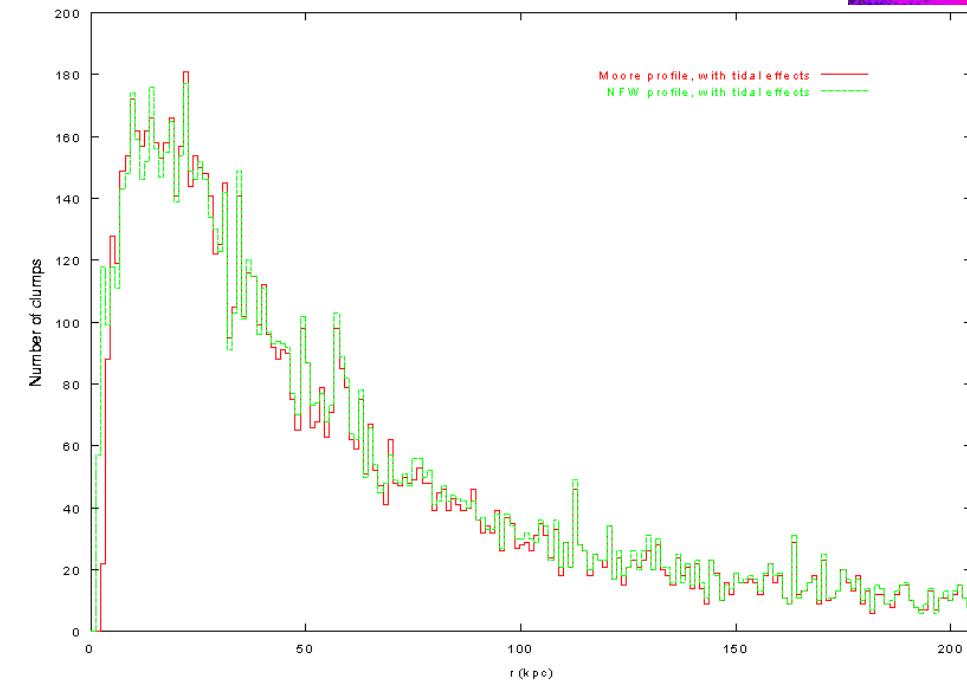
Determine the profile parameters

- If a halo viral mass is M_{vir} at redshift z , we have and
$$r_{\text{vir}} = \left(\frac{3M_{\text{vir}}}{4\pi\Delta_{\text{vir}}(z)\rho_\chi(z)} \right)^{1/3}$$
$$\rho_\chi(z) = \rho_\chi(1+z)^3$$
$$\Delta_{\text{vir}}(z) = (18\pi^2 + 82y - 39y^2)/(1+y), \quad y = \Omega_m(z) - 1 \quad \text{and} \quad \Omega_m(z) = \frac{\Omega_m(1+z)^3}{\Omega_m(1+z)^3 + \Omega_\Lambda}$$
- Here Δ represents the average density in a viralized halo to the matter density of the universe.
- The concentration parameter is $c_{\text{vir}}(z) = c_{\text{vir}}(z=0)/(1+z)$
- Then we have $r_s = \frac{r_{\text{vir}}}{c_{\text{vir}}(2-\gamma)}$ for a generalized NFW
$$\rho(r) = \frac{\rho_s}{(r/r_s)^\gamma(1+r/r_s)^{3-\gamma}}$$
- We have now slightly different models to determine the concentration parameter J. S. Bullock et al 2001, A. V. Maccio et al. 2008

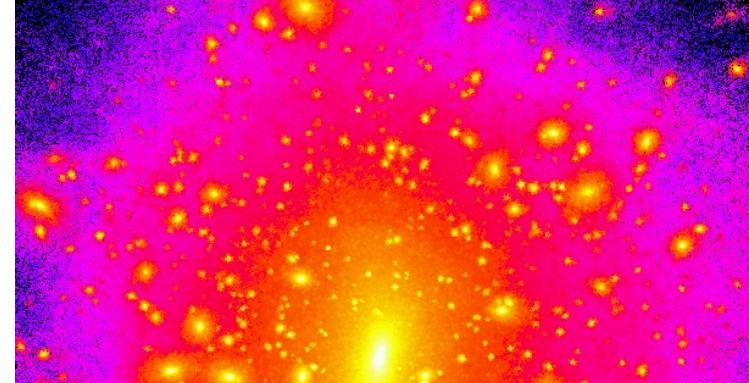
Substructure

Moore et al

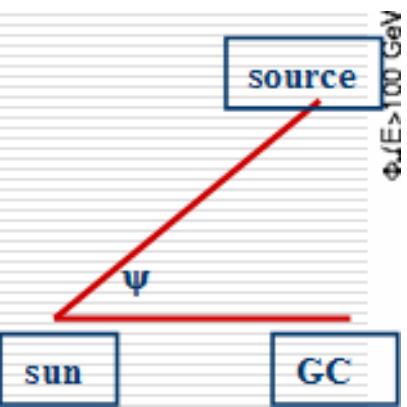
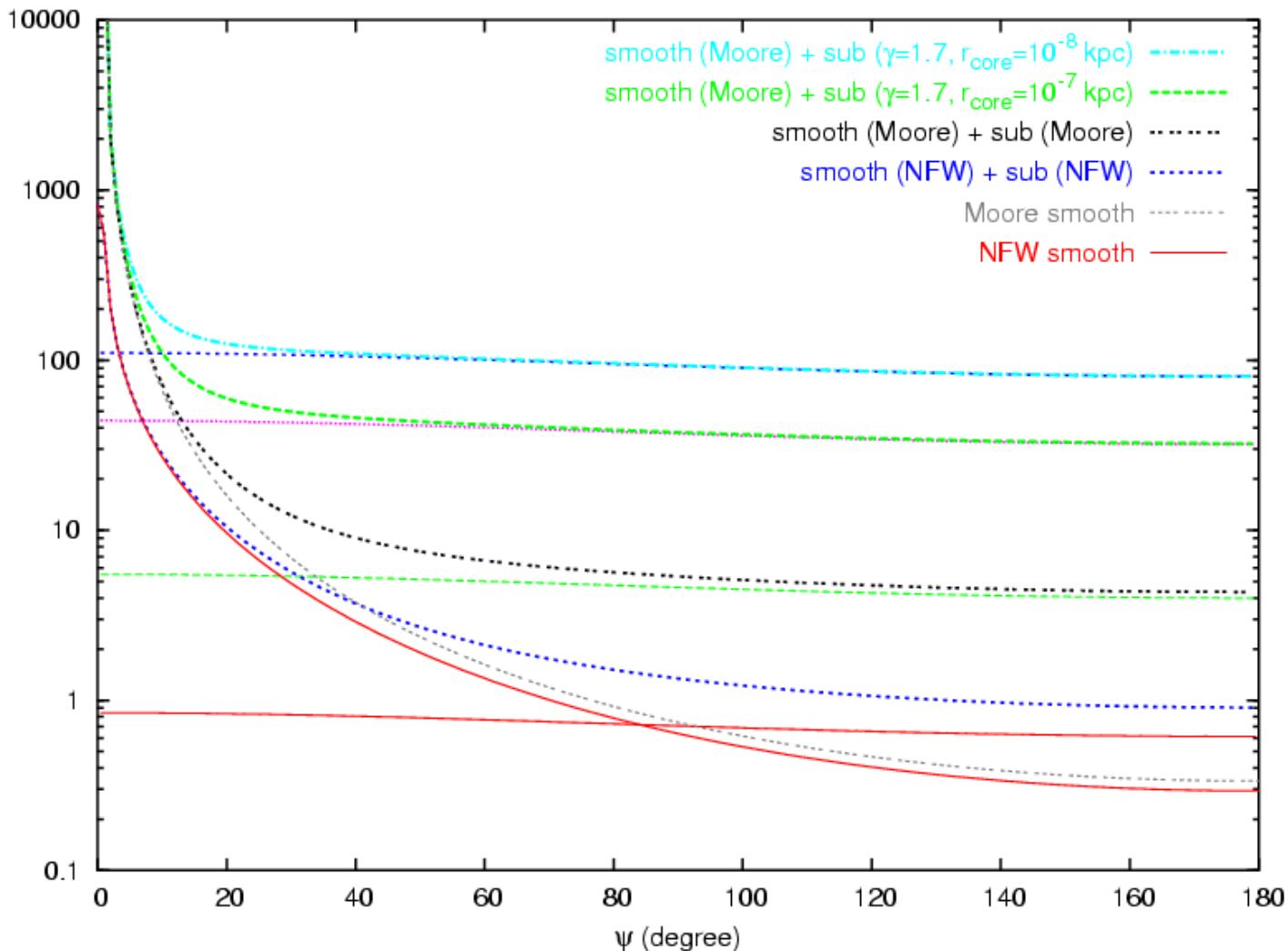
- Will the smaller objects formed earlier survive today in the dark halo?
- A wealth of subhalos exist due to high resolution simulations.
- The number density is as:



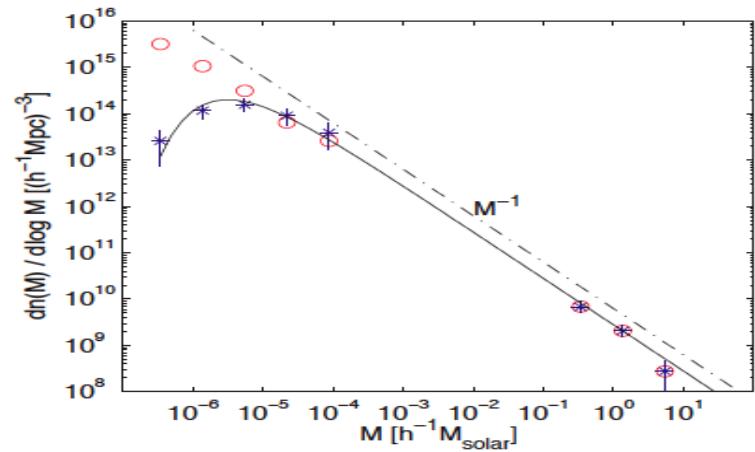
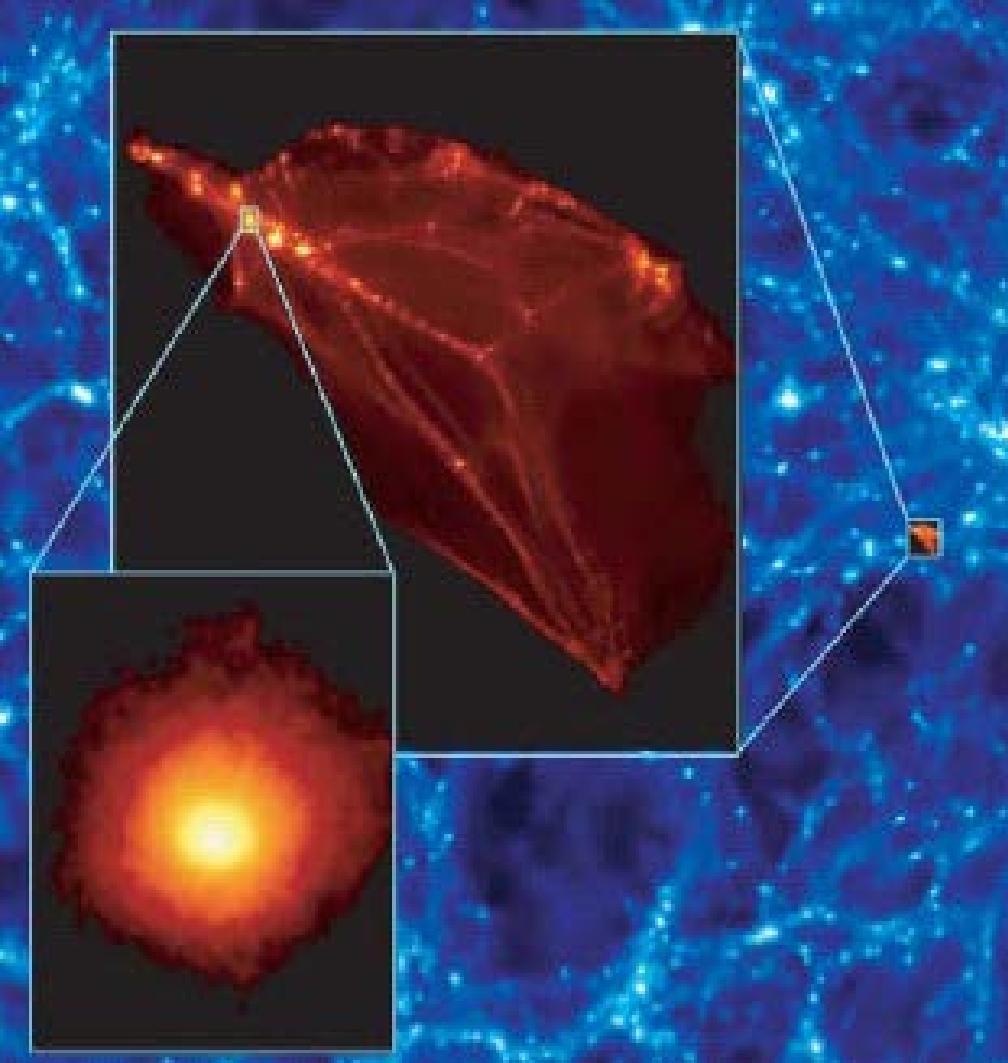
With and without subhalos



$$\Phi^{cosmo} = \int_{l.o.s} \rho^2(r) dl$$



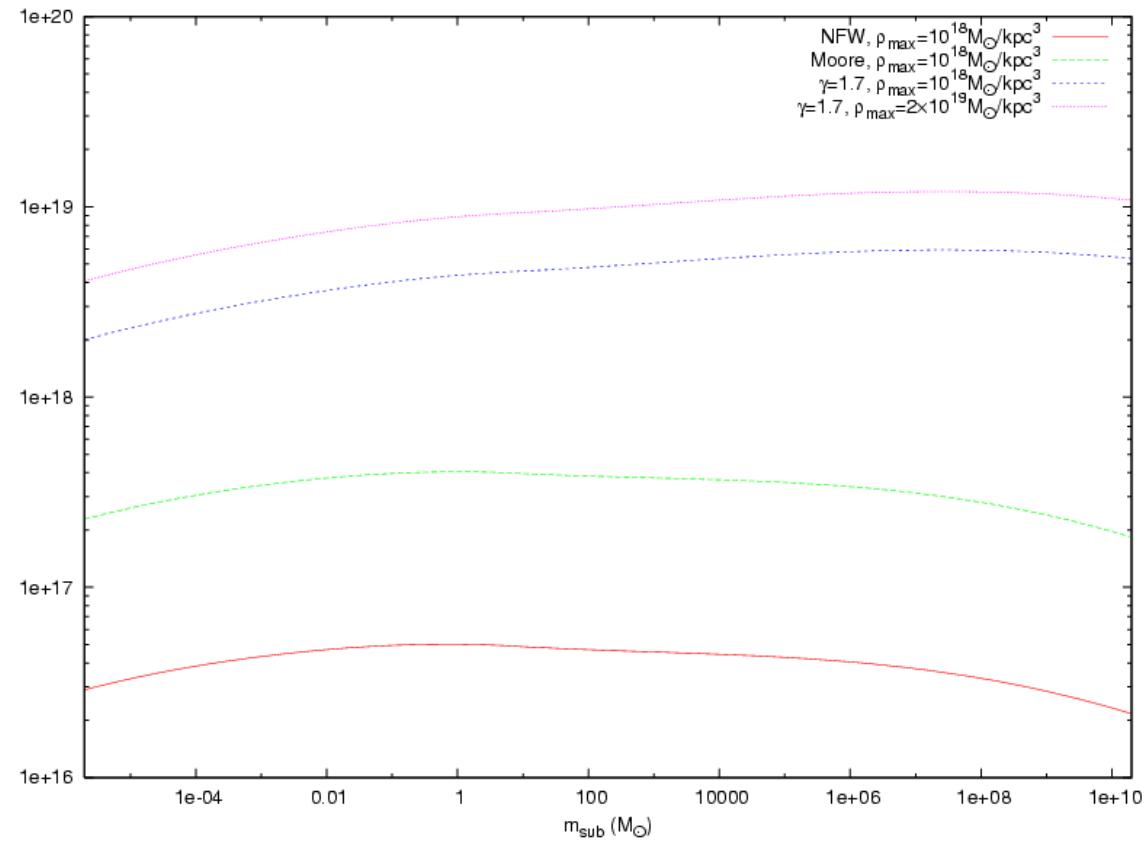
The first generation object



Diemand, Moore & Stadel, 2005:

- Depending on the nature of the dark matter: for neutralino-like dark matter, the first structures are mini-halos of $10^{-6} M_{\odot}$.
- There would be zillions of them surviving and making up a sizeable fraction of the dark matter halo.
- The dark matter detection schemes may be quite different!

Contributions from different mass ranges of subhalos



$$\frac{dN}{dm} \propto m^{-1.9}$$

$$\frac{d\mathcal{P}_V(r)}{4\pi r^2 dr} = K_V \times \left[1 + \left(\frac{r}{r_H} \right)^2 \right]^{-1}$$

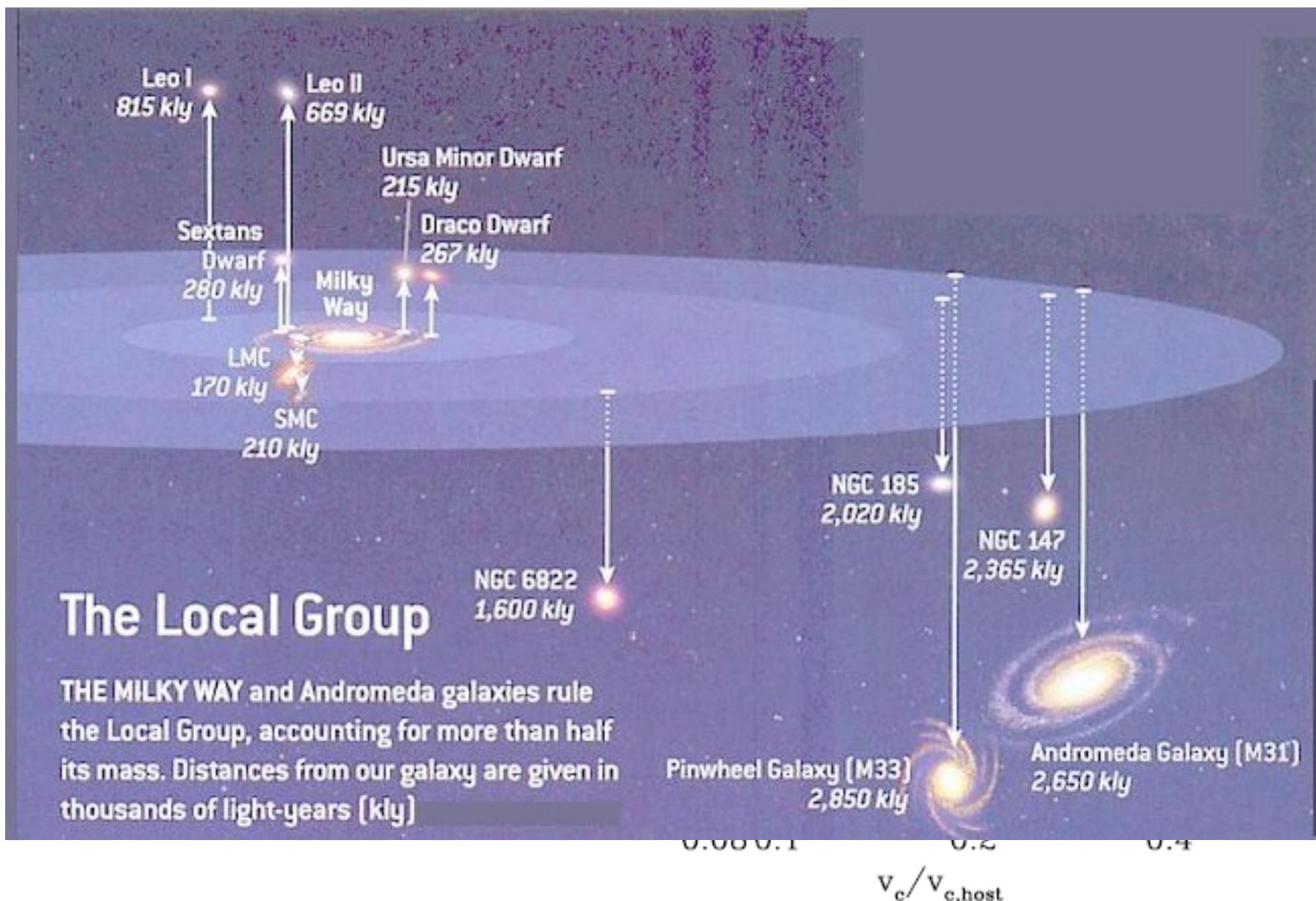
□ Inconsistencies in Λ CDM paradigm

Scale < 1Mpc

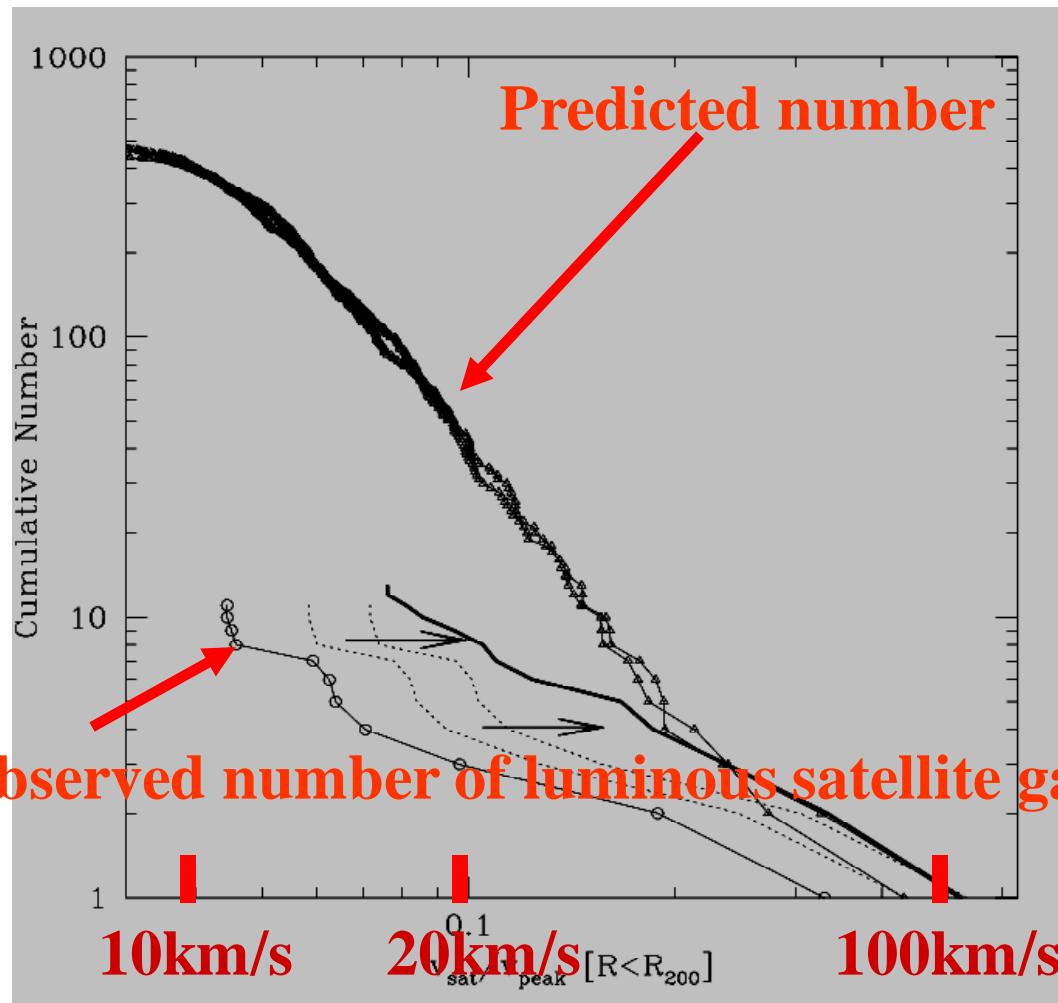
- Substructure number problem
 - DM haloes contain a huge number of subhalos
 - the observed number of dwarf galaxies in the voids appears to be far smaller than expected from CDM
- Cusp problem
 - DM halos have a density profile with sharp slope
 - the density profile inferred from galaxy rotation curves are significantly shallower

Difficulties of CDM: number of satellites

Jing (2001)



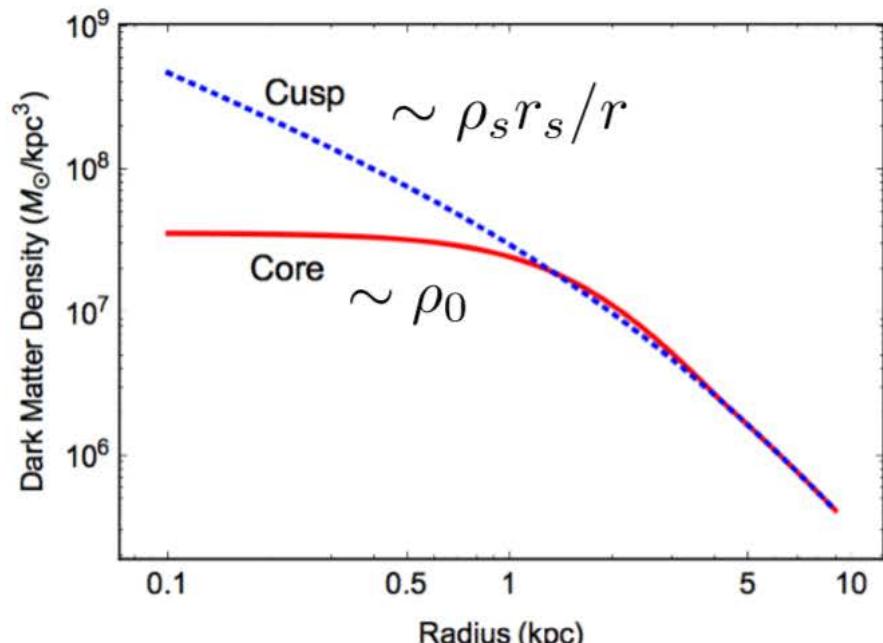
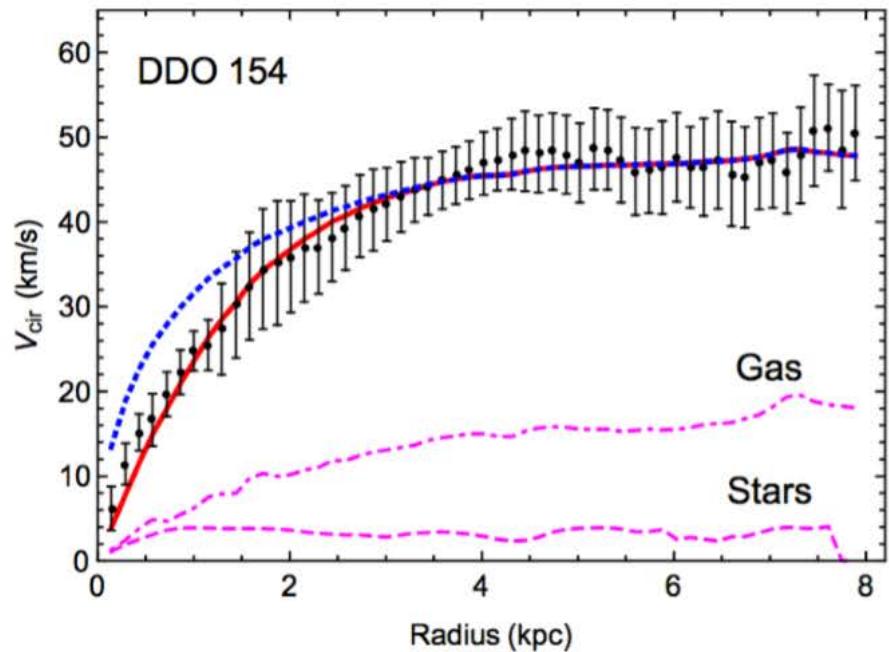
- Satellite galaxies are seen in Milky Way, e.g. Sagittarius, MCs



- The predicted number of substructures exceeds the luminous satellite galaxies: dark substructures?

Core vs. Cusp Problem

- DM-dominated systems (dwarfs, LSBs)



$$\frac{\rho_s}{r/r_s(1+r/r_s)^2}$$

Navarro, Frenk, White (1996)

universal density profile, NFW profile
 ρ_s and r_s are strongly correlated

Problems of CDM on sub-galactic scale

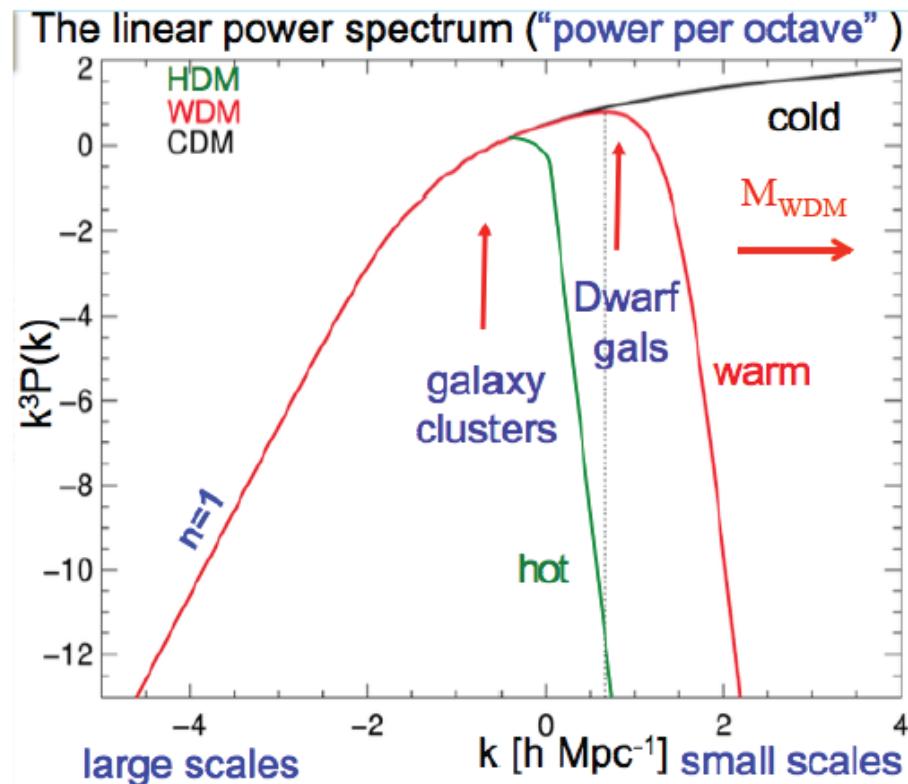
- The simulation over-predicts the number of subhalos. The number of observed dwarf satellites is an order of magnitude smaller than predicted.
- Solutions are proposed, including inflation potential suppressing the small scale fluctuation, star formation suppressed in small subhalos, and new form of dark matter candidates, such as **self-interacting dark matter, non-thermal production, warm dark matter, fuzzy dark matter** and so on.
- Milli-lensing by halo substructure seems favor the CDM scenario. It is still controversial.

How to define “cold”

- Definition of cold, warm or hot depends on the effect of their “free-stream” motion on the formation of objects
 - Cold dark matter that has effectively zero thermal velocity
 - Hot dark matter (eV neutrinos) that washes out fluctuations on cluster scale (10 Mpc/h);
 - Warm dark matter (sterile neutrinos) that washes out fluctuations on galaxy scale (1 Mpc/h);
- $$\lambda_{\text{FS}}^{\text{co}} = \int_0^t \frac{v(t') dt'}{a(t')}$$

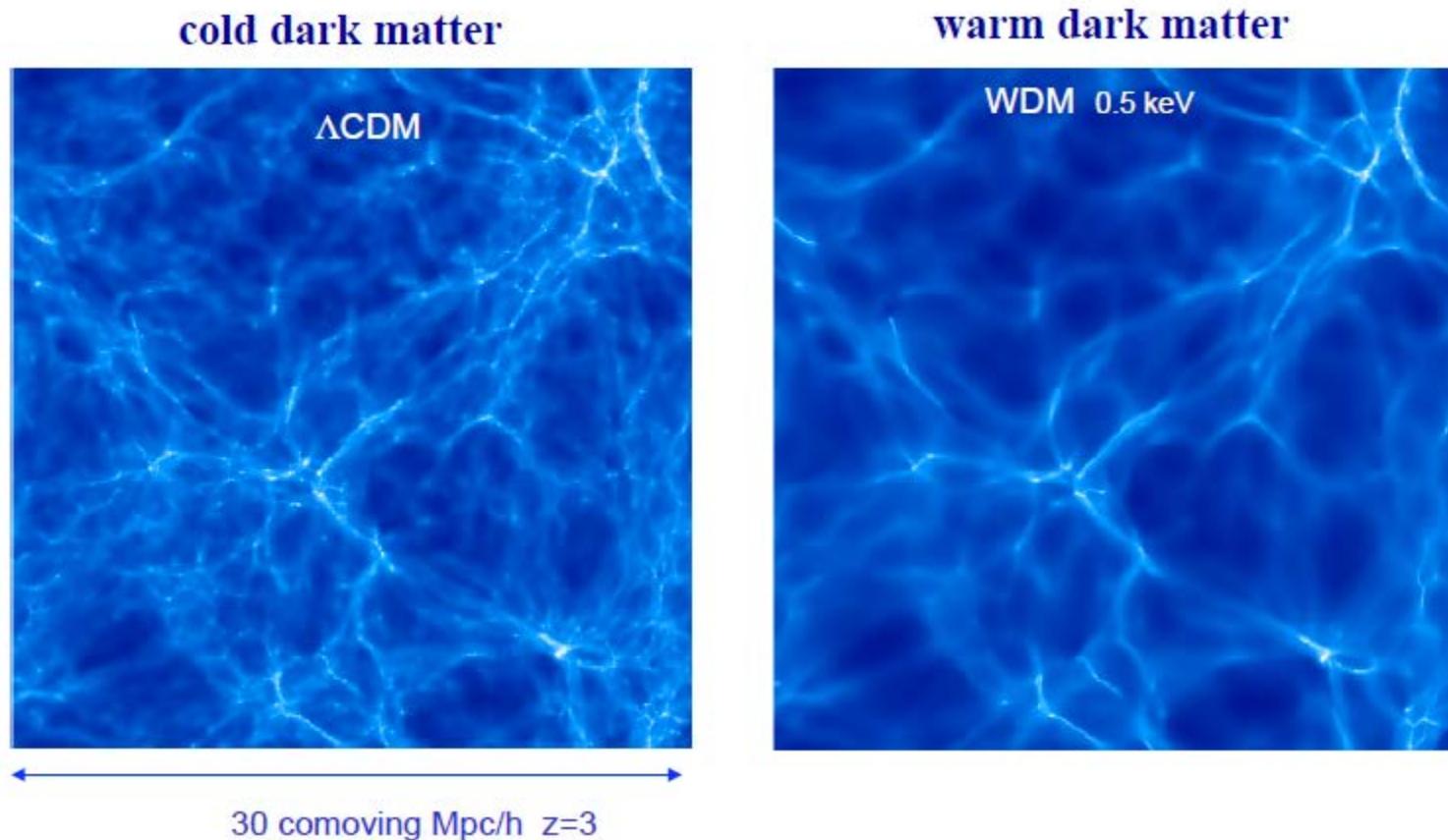
What does warm mean?

Cut-off in the matter power spectrum on astrophysically interesting scales due to free-streaming?



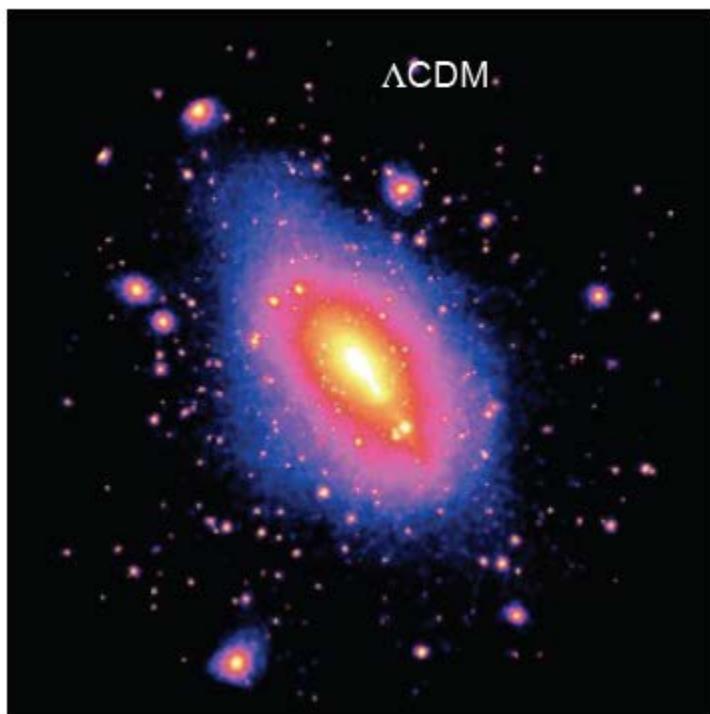
courtesy of
Carlos Frenk

Free-streaming erases structure

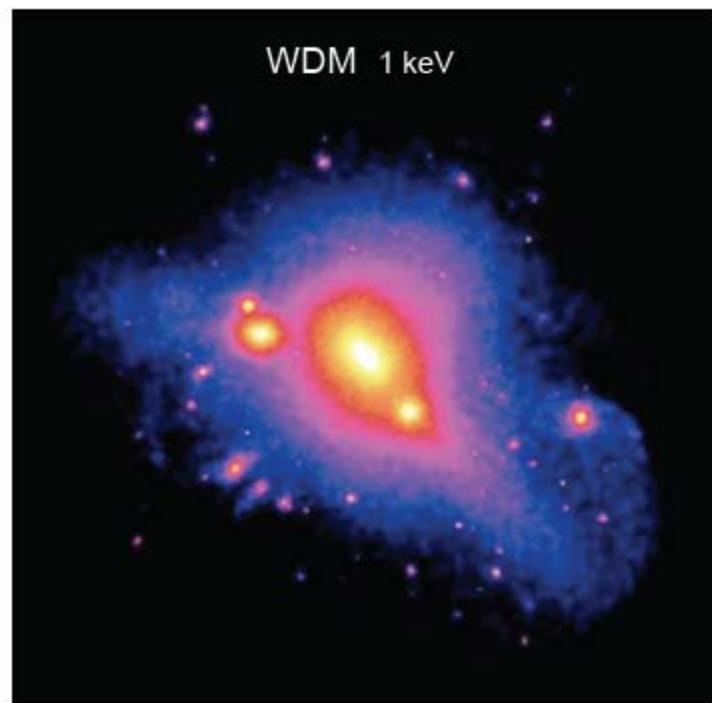


Less subhalos

cold dark matter



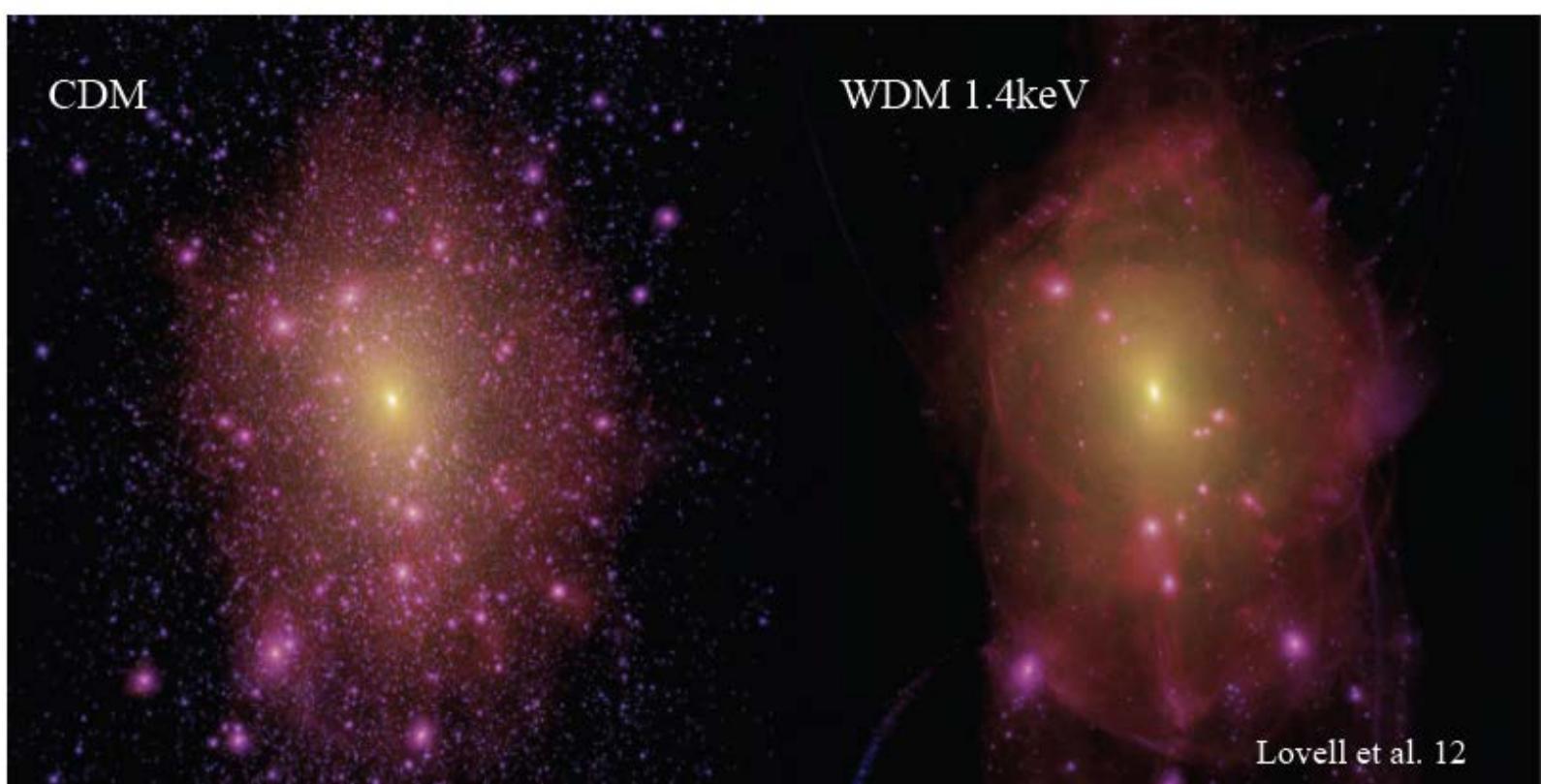
warm dark matter



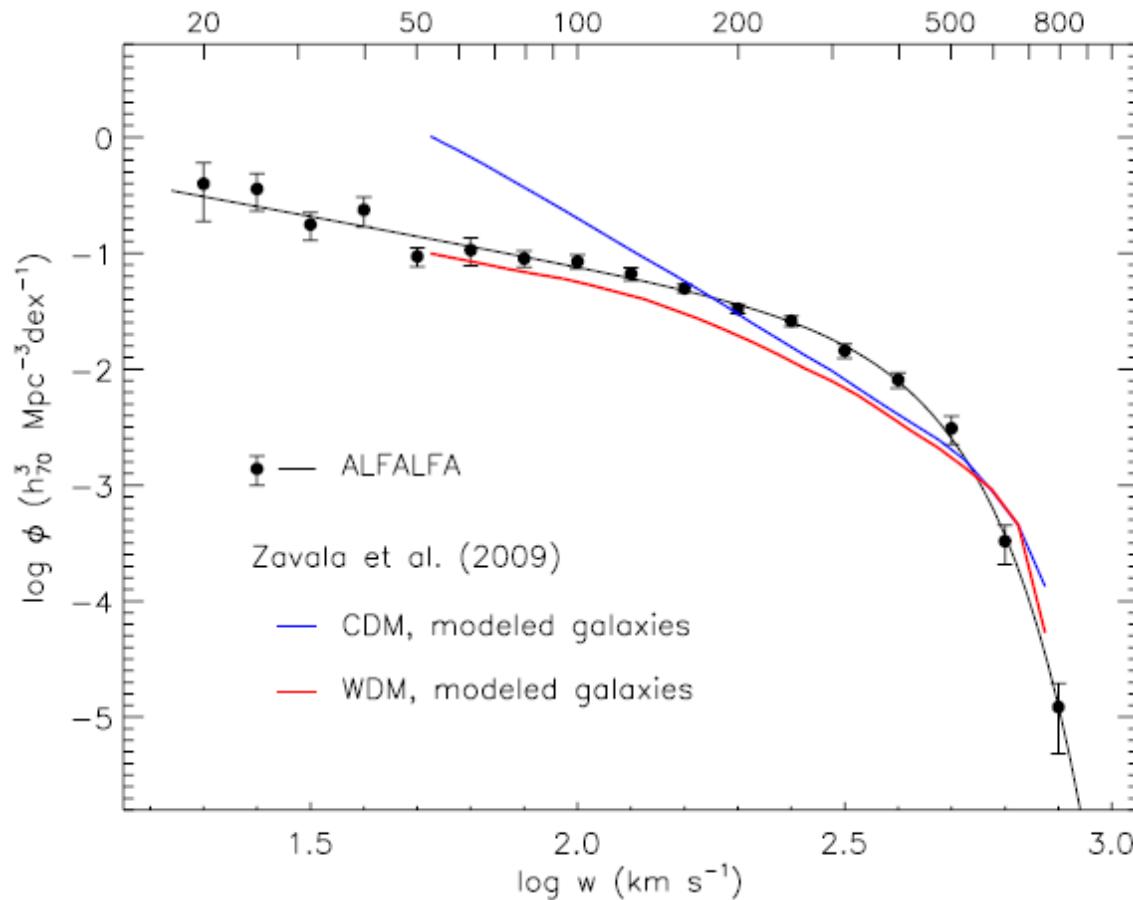
Maccio & Fontanot 2009

More smaller structures

cold dark matter

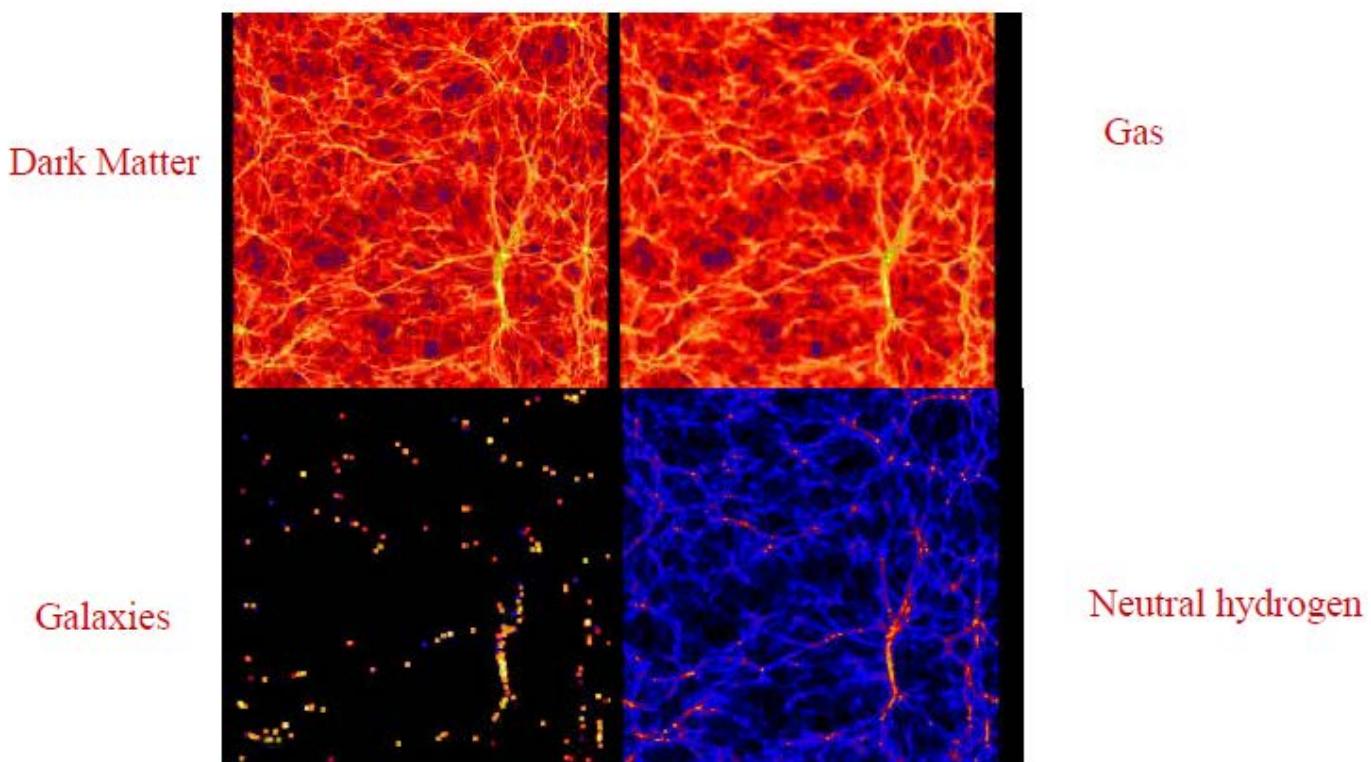


How cold is dark matter: velocity width function of galaxies (ALFALFA survey)



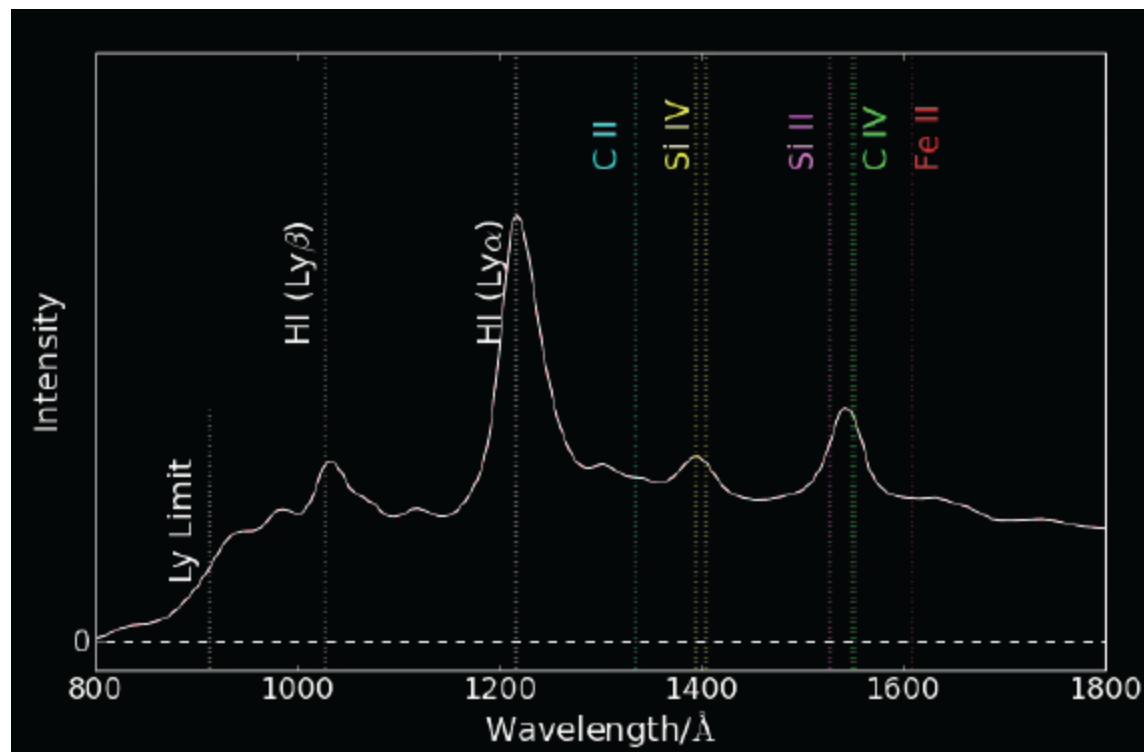
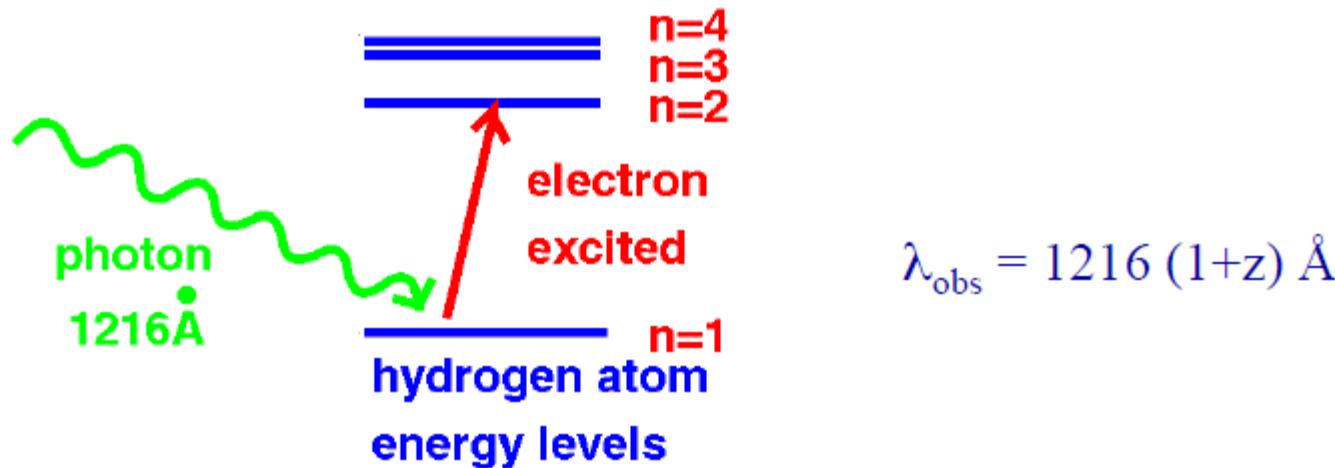
Papastergis et al. (2011)

Observation of the structure



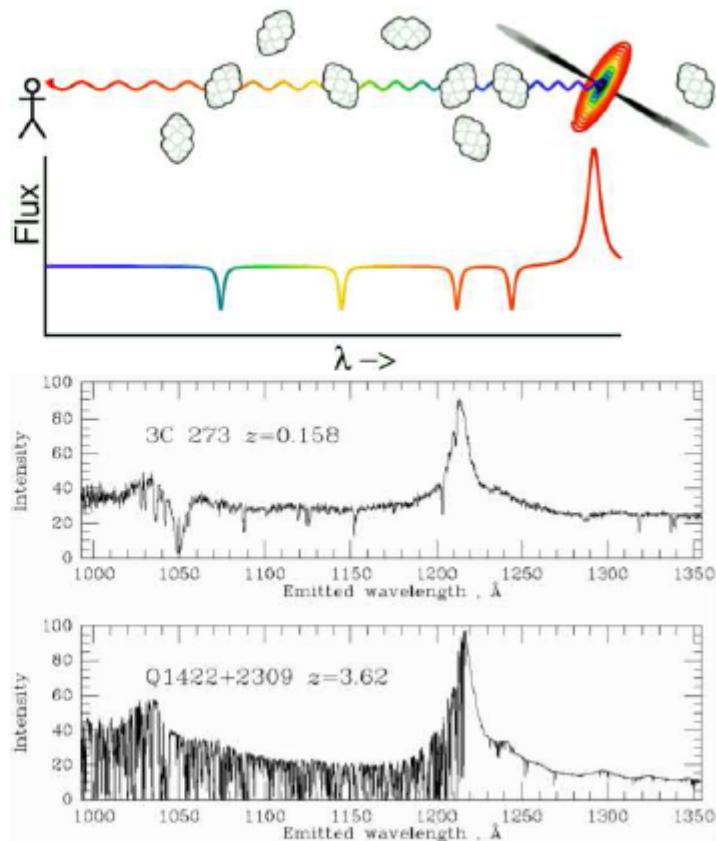
Neutral hydrogen is an excellent tracer of the matter distribution.

Ly α absorption by neutral hydrogen

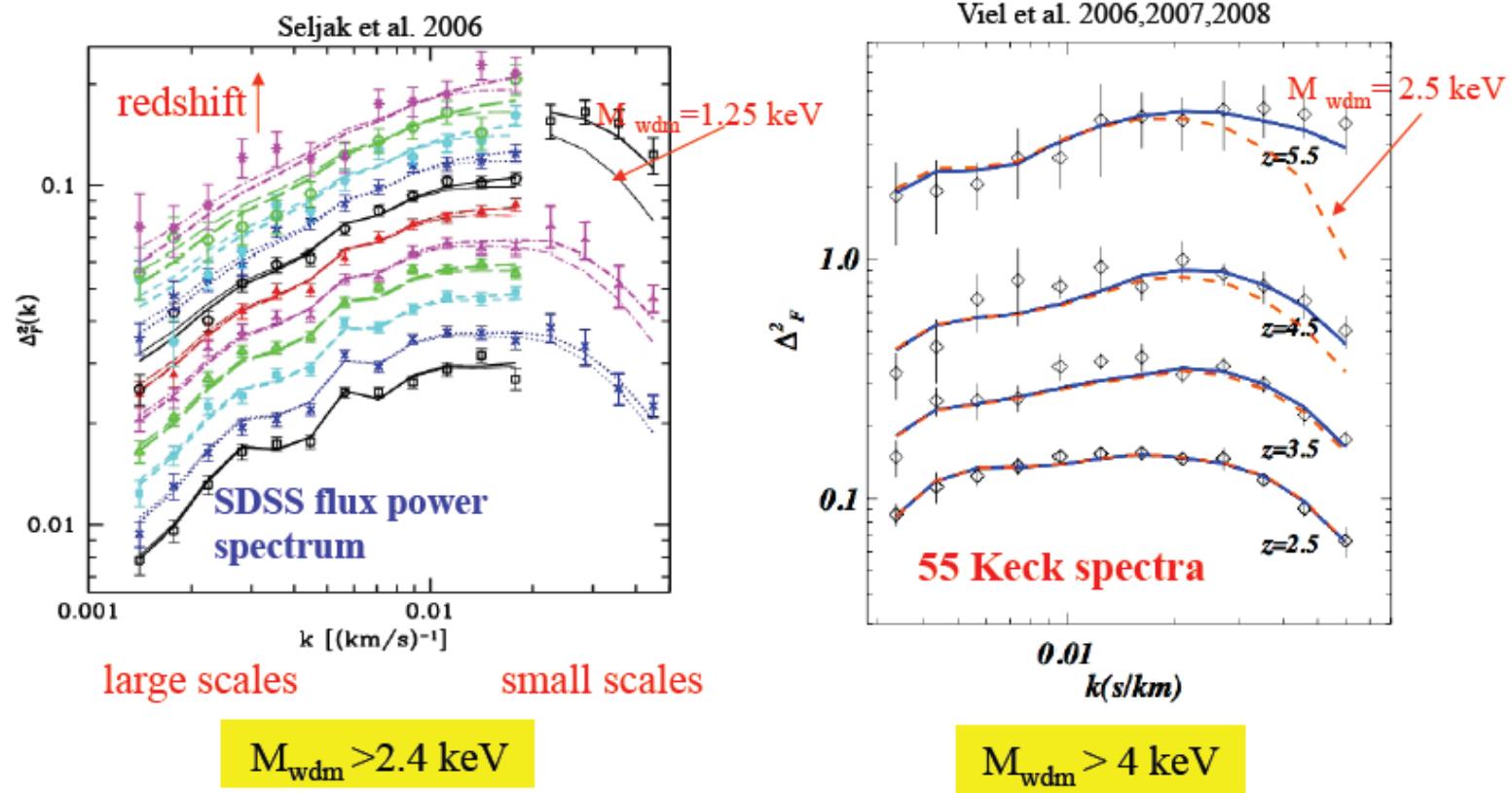


- To probe the DM properties at small scales one can use Lyman- α forest data:
- Red-shifted absorption Lyman- α line in the spectra of distant QSOs
- Neutral hydrogen traces DM distribution at red-shifts $z \sim 2 - 4$.
- Allows to measure one-dimensional non-linear power spectrum:

$$P_{1D} = \int_k^\infty P_{3D}(k) \frac{k dk}{2\pi}$$



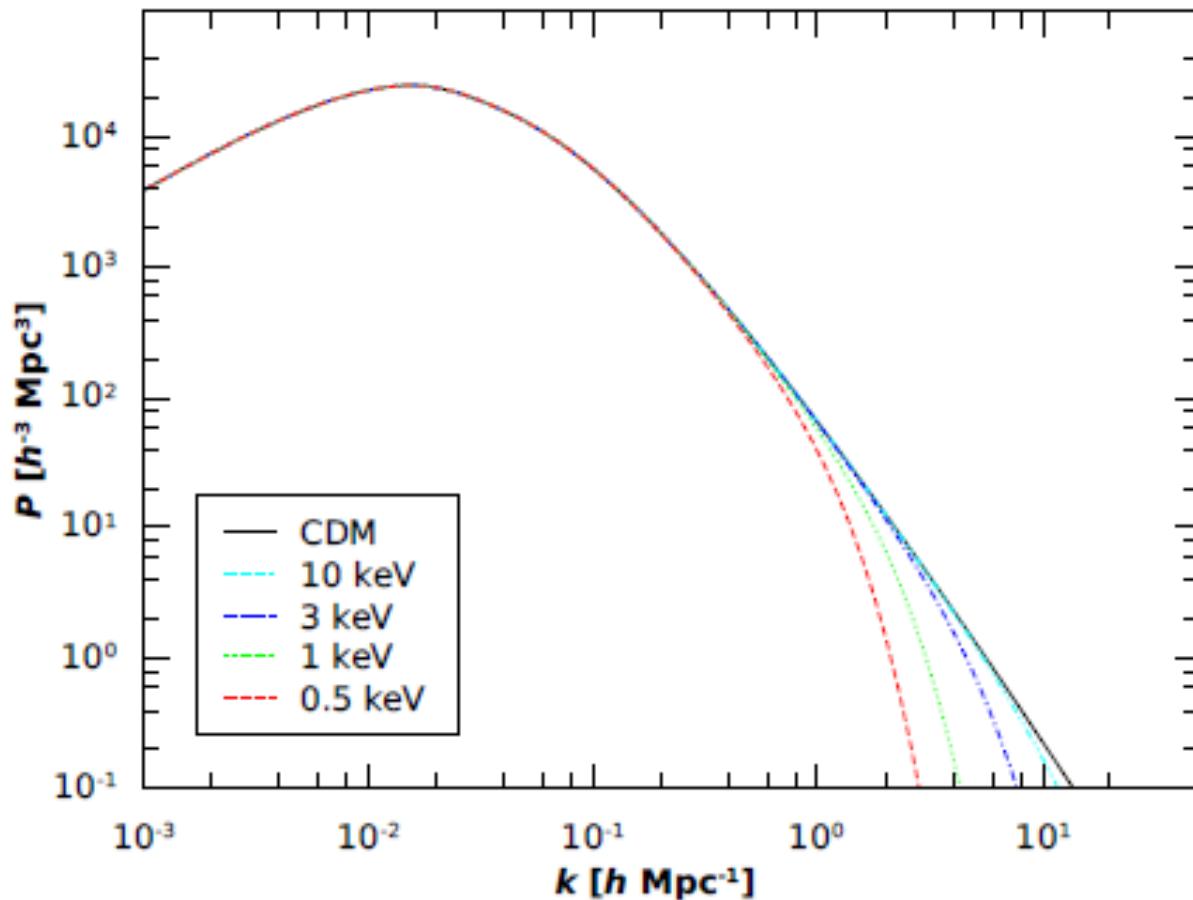
Observational results



CDM or WDM? A summary

- 1) Does CDM predictions contradict observations?
 - CDM simulations are **pure DM**. Pure N-body is not enough
 - Astronomers observe **luminous** matter.
 - Baryonic feedback can be essential
 - Example: not all DM halos can acquire baryons
- 2) Any WDM simulations (N-body or hydrodynamical) should
 - properly include primordial velocities of the particles
 - use correct power spectrum of initial density perturbations.
- 3) WDM is ruled out by Lyman- α ?
 - **No**
- 4) DM with keV mass still allowed?
 - **Yes**

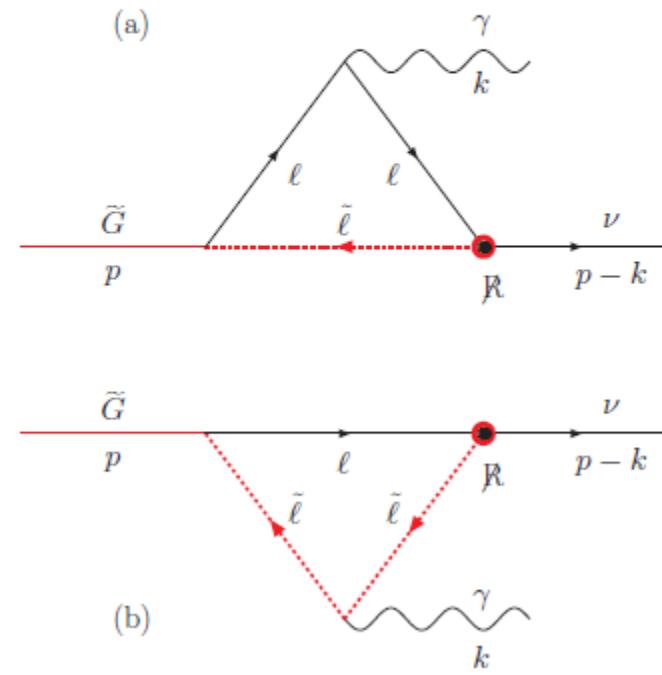
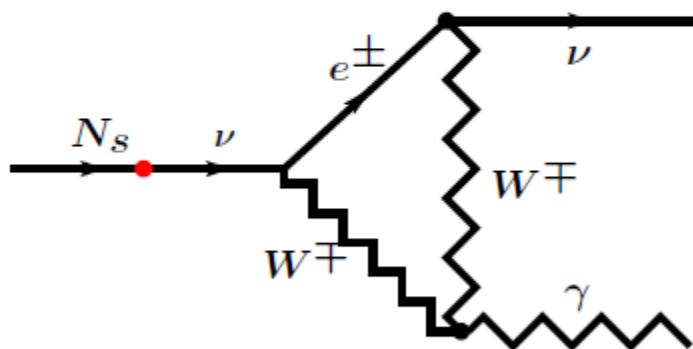
☐ Sterile Neutrino as WDM



the linear matter power spectrum

Decaying DM

DM with **radiative signatures**: $\text{DM} \rightarrow \gamma + \nu, \gamma + \gamma, e^+ + e^- \dots$

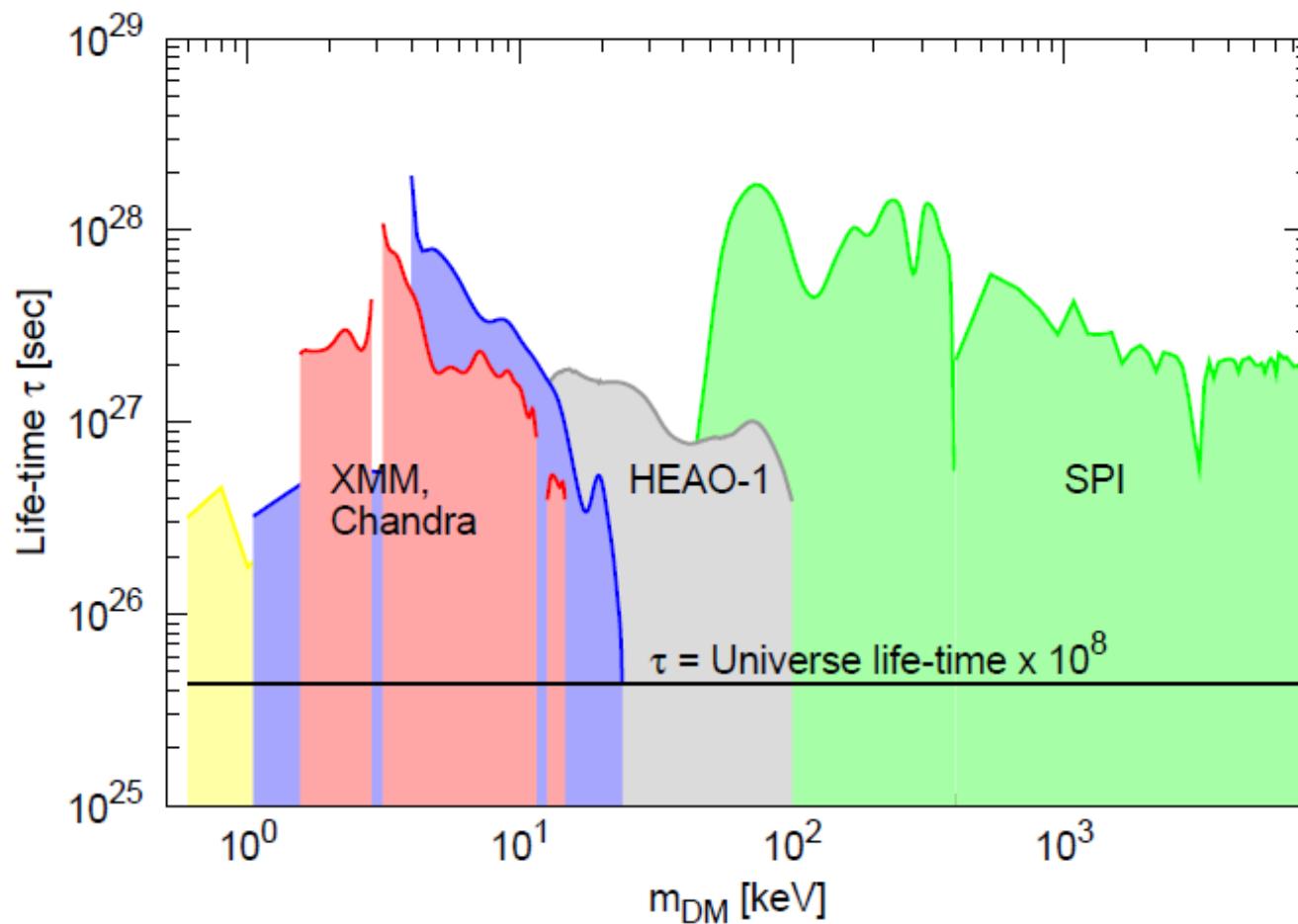


Properties of decaying DM

- Decaying DM should interact **superweakly** $\sim \theta \cdot G_F$ and $\theta \ll 1$
- Radiative decay channel : $\text{DM} \rightarrow \gamma + \nu$
- Photon energy $E_\gamma = \frac{m_{\text{DM}}}{2}$
- Life-time $\tau = 1/\Gamma \gg$ life-time of the Universe
- Flux from DM decay:

$$F_{\text{DM}} = \frac{E_\gamma}{m_{\text{DM}}} \frac{\Gamma \mathcal{M}_{\text{DM}}^{\text{fov}}}{4\pi D_L^2} \approx \frac{\Gamma \Omega_{\text{fov}}}{8\pi} \int_{\text{line of sight}} \rho_{\text{DM}}(r) dr \quad (z \ll 1, \quad \Omega_{\text{fov}} \ll 1)$$

Restriction on the life time of sterile neutrino



非热产生暗物质宇宙模型

WIMPs两种产生机制：

1. 热平衡退耦
2. 非热产生

- Decay of topological defects such as cosmic string
- Decay of an unstable heavy particle
- produced in the reheating process in a scenario of inflation at low energy scale

...

Lin, Huang, Zhang, Brandenberger (2001);
Bi, Brandenberger, Gondolo, Li, Yuan, Zhang (2009)

NTDM momentum distribution

$$f(p) = \frac{A}{\sqrt{2\pi}\sigma} \exp\left(-\frac{(p - p_c)^2}{2\sigma^2}\right)$$

The current velocity of the NTDM particle,

$$v_0 = \frac{T_0}{T_d} \frac{M}{2m}$$

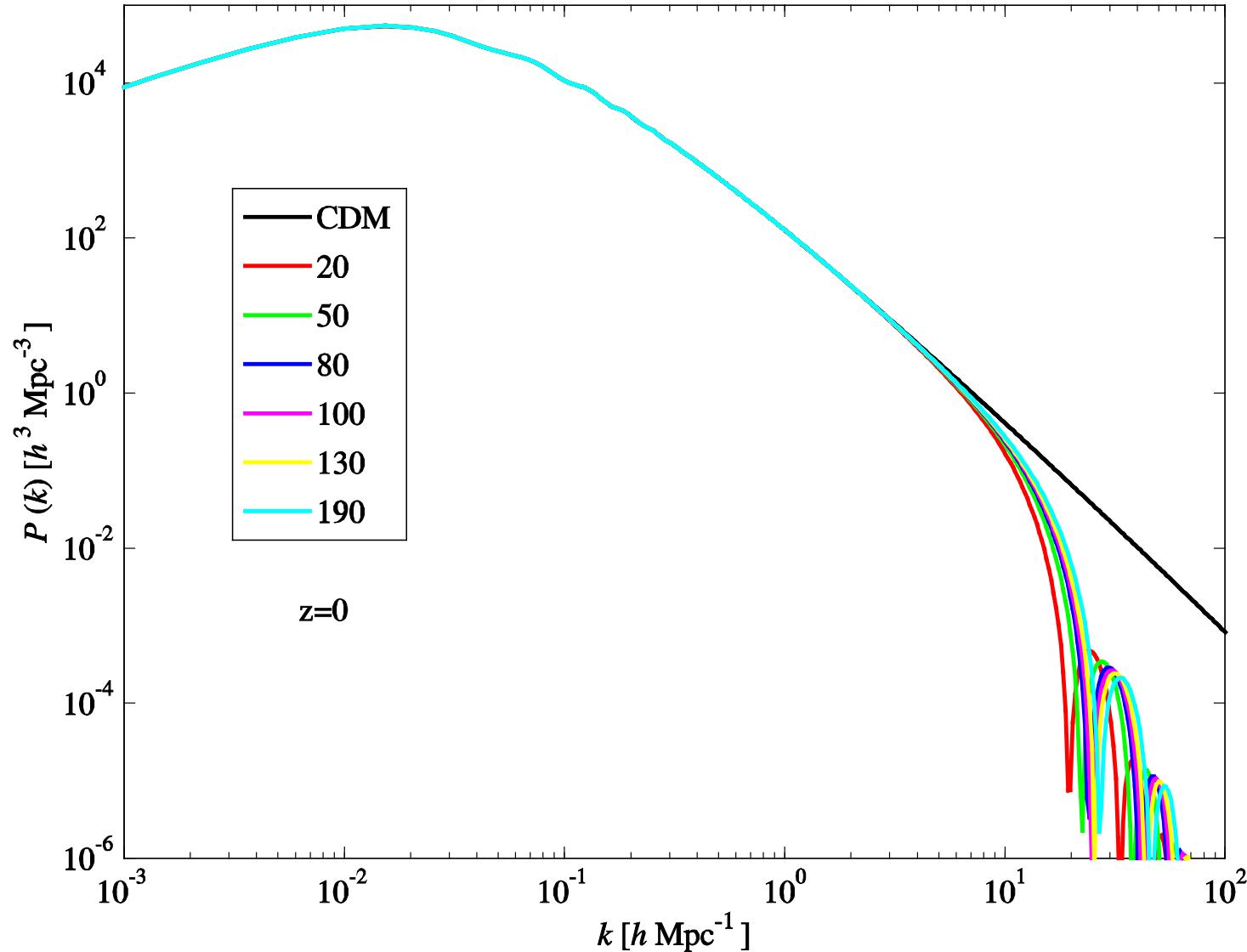
NTDM comoving free-streaming length

$$\begin{aligned} \lambda_f &= \int_{t_i}^{t_{EQ}} \frac{v(t')}{a(t')} dt' \simeq \int_0^{t_{EQ}} \frac{v(t')}{a(t')} dt' \\ &\simeq 2v_0 t_{EQ} (1 + z_{EQ})^2 \ln \left(\sqrt{1 + \frac{1}{v_0^2 (1 + z_{EQ})^2}} + \frac{1}{v_0 (1 + z_{EQ})} \right) \end{aligned}$$

Free streaming length

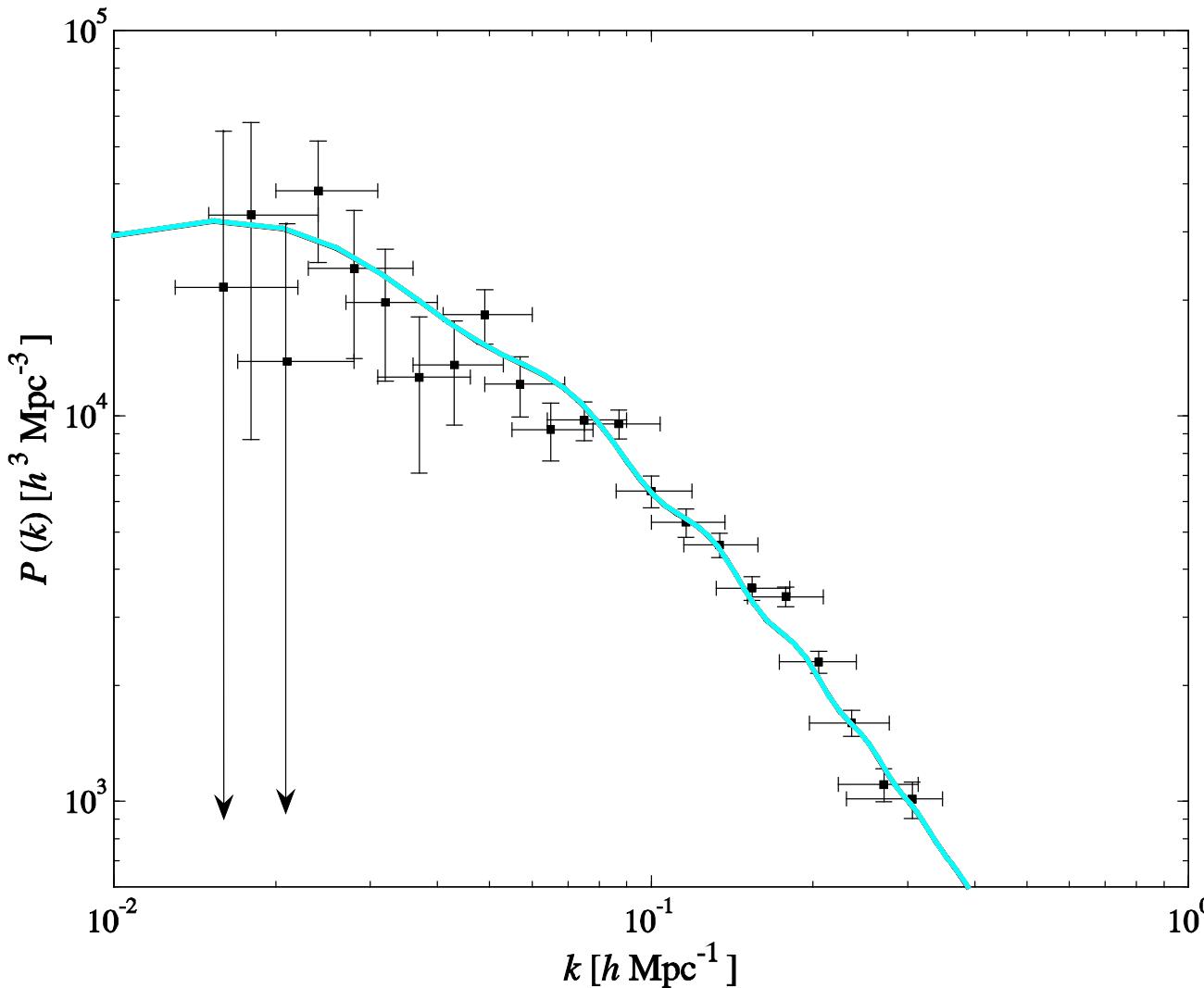
$m(\text{GeV})$	$M(\text{GeV})$	v_0	$\lambda_f(\text{Mpc})$
20	5.81×10^7	3.96×10^{-7}	1.93
50	1.25×10^8	1.36×10^{-7}	0.76
80	1.85×10^8	7.86×10^{-8}	0.47
100	2.22×10^8	6.06×10^{-8}	0.37
130	2.77×10^8	4.46×10^{-8}	0.28
190	3.79×10^8	2.87×10^{-8}	0.19

■ PERTURBATION EVOLUTION



不同的 m_{NTDM} 对应的 $P(k)$

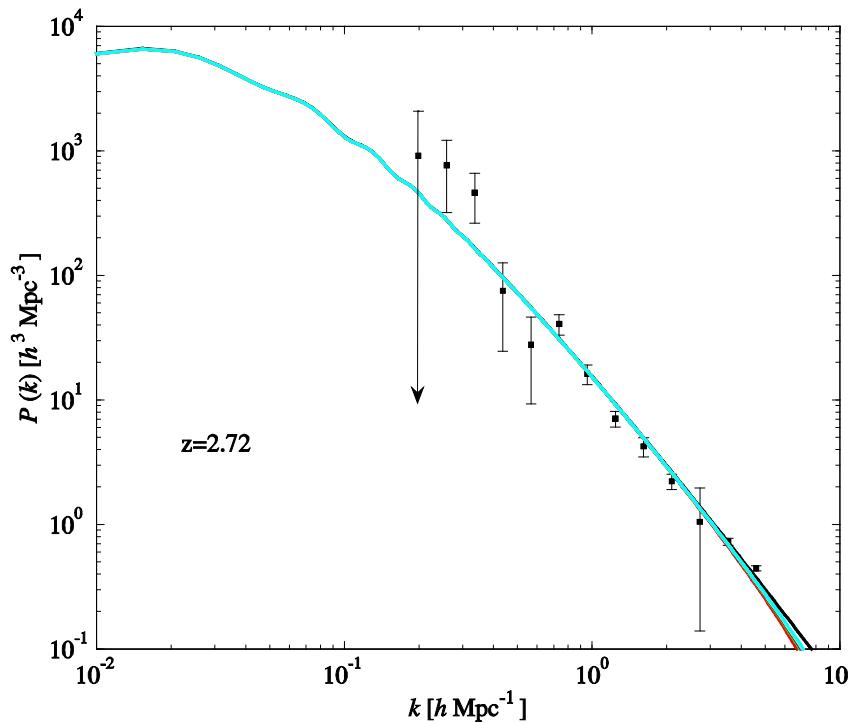
Constraining NTDM with SDSS galaxies



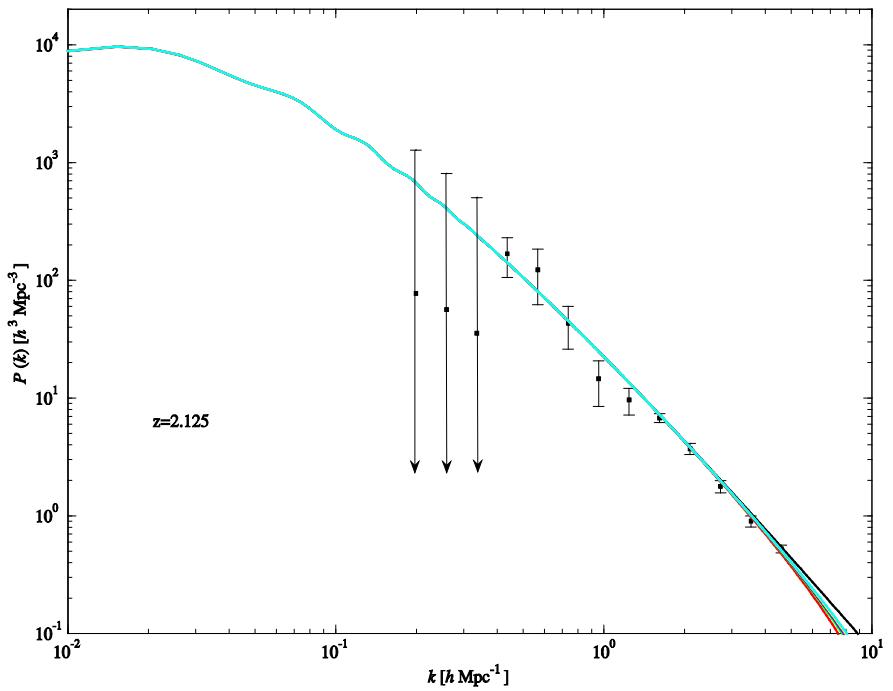
$k < 0.2 h^{-1} \text{Mpc}$

$m_\chi > 20 \text{GeV}$

Constraining NTDM with LYMAN- α Forest



the sample with median redshift
 $z = 2.72$ (Croft et al., 2002)



the LUQAS sample with median
redshift $z = 2.125$ (Kim et al. 2004)

■ N-BODY Simulations for NTDM

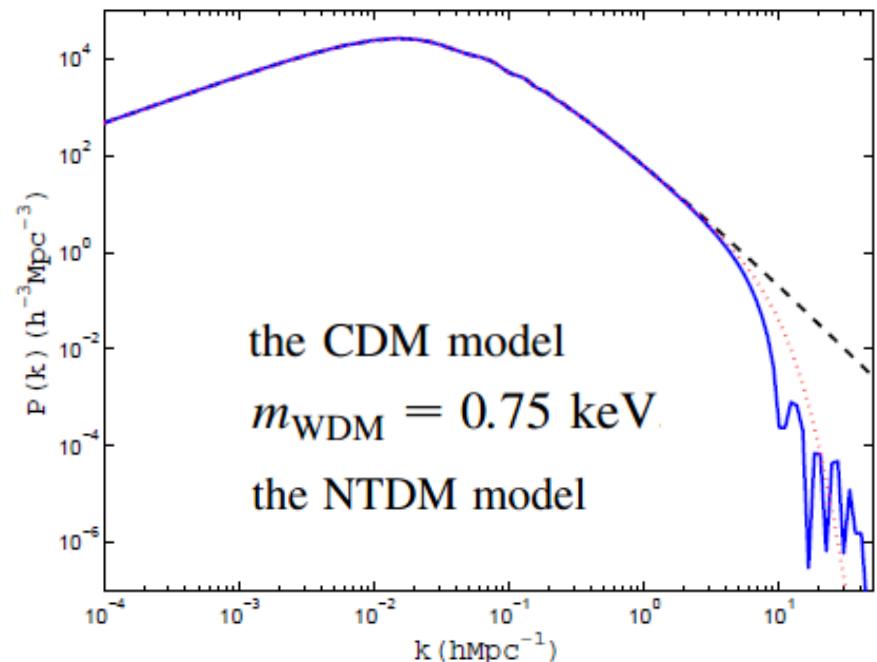
$\Omega_m = 0.28$ $\Omega_\Lambda = 0.72$, $\Omega_b = 0.046$, $h = 0.7$, $n_s = 0.97$, $\sigma_8 = 0.82$

$v_0 = 0.66 \times 10^{-7}$

► Linear power spectra

► N-Body Simulation

$L_{\text{box}} = 64 \text{ h}^{-1}\text{Mpc}$ $N = 256^3$

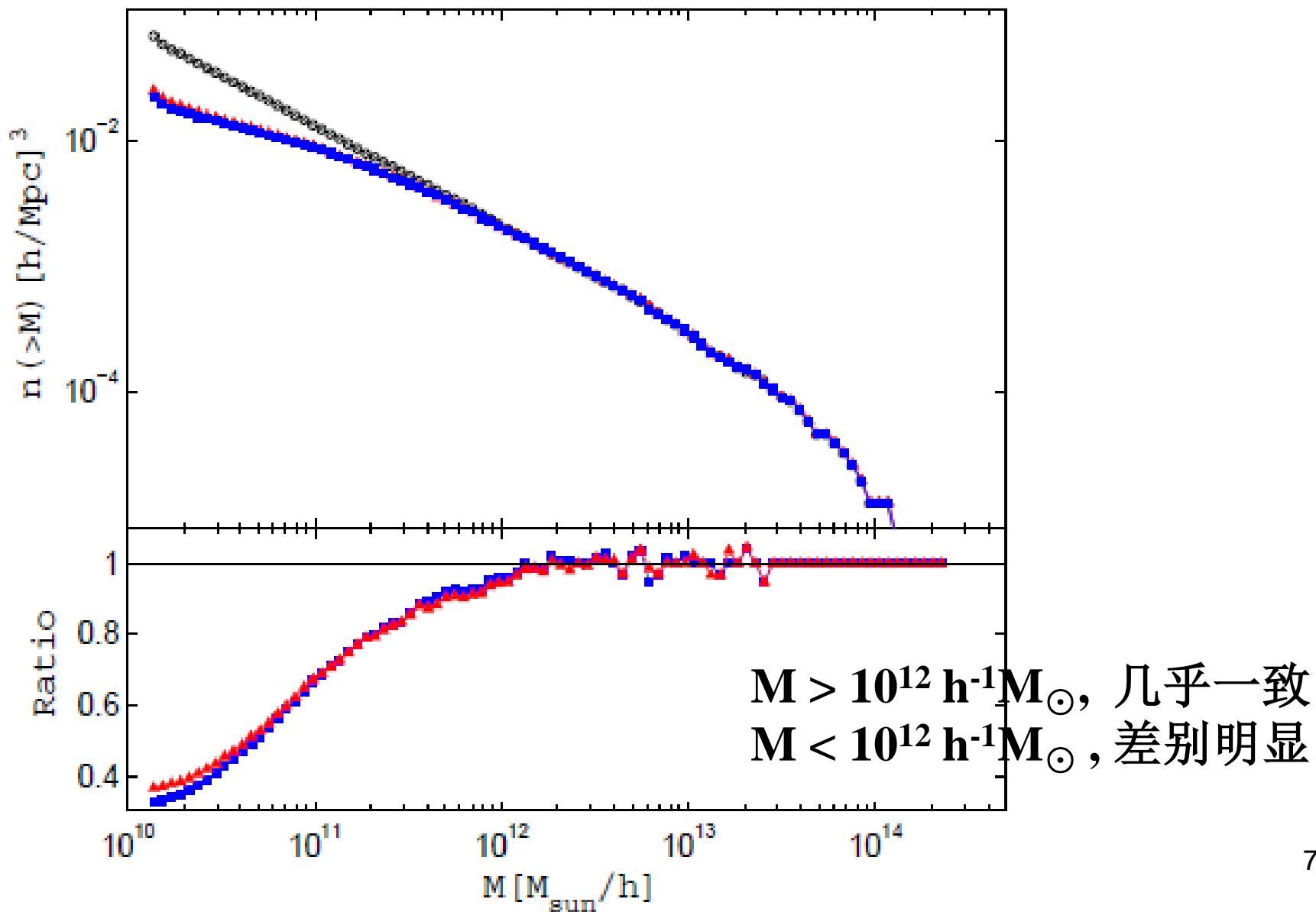


Generate ICs for N-body simulation

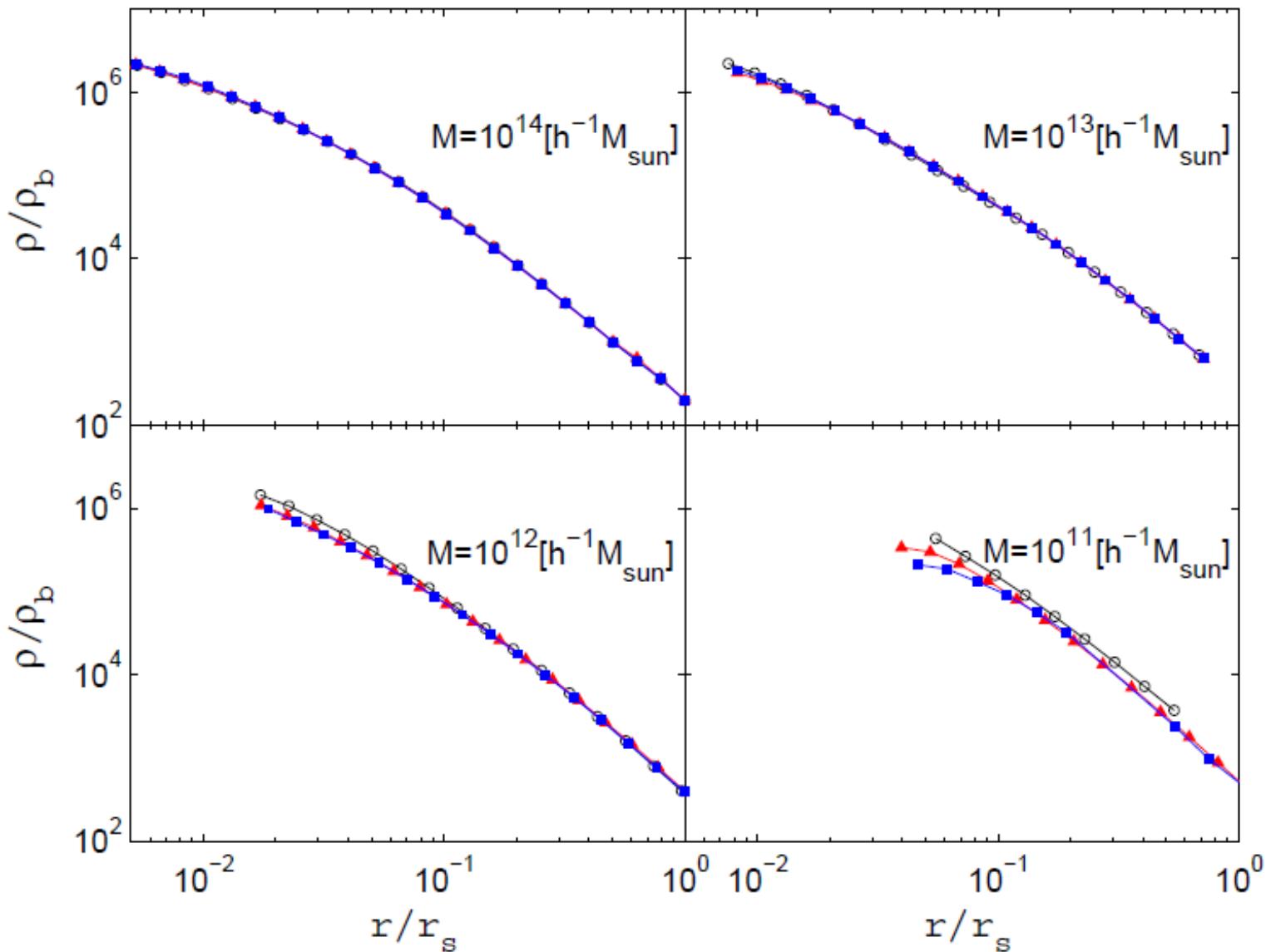
Output particle distribution and velocity at z=49

GADGET-2

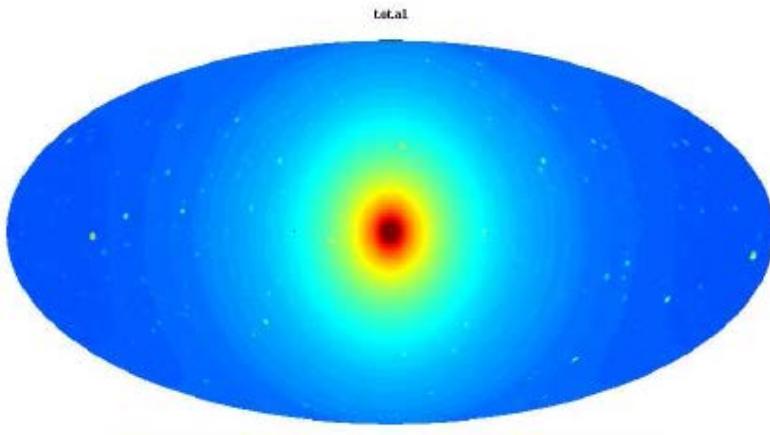
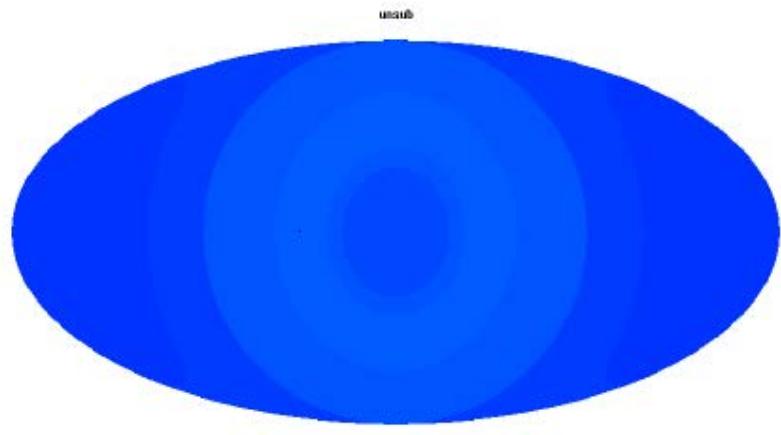
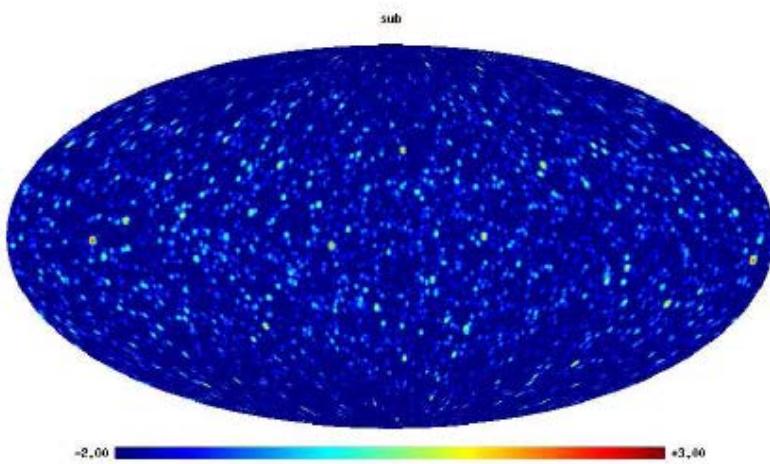
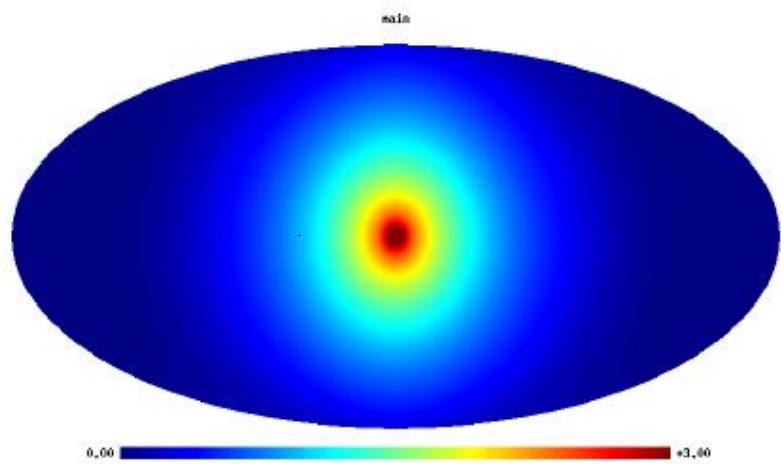
Number Density of Halos



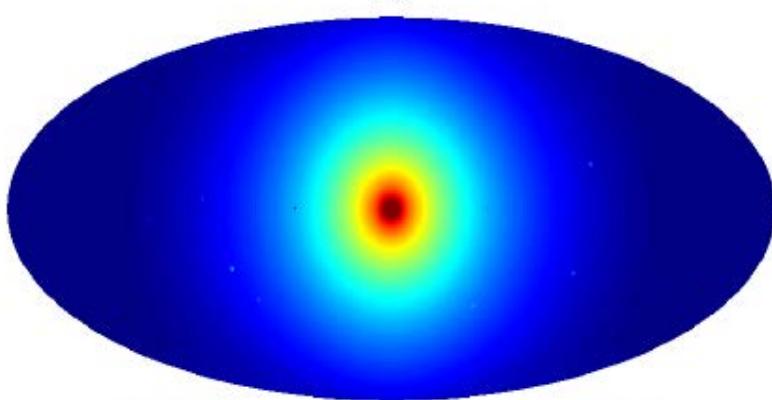
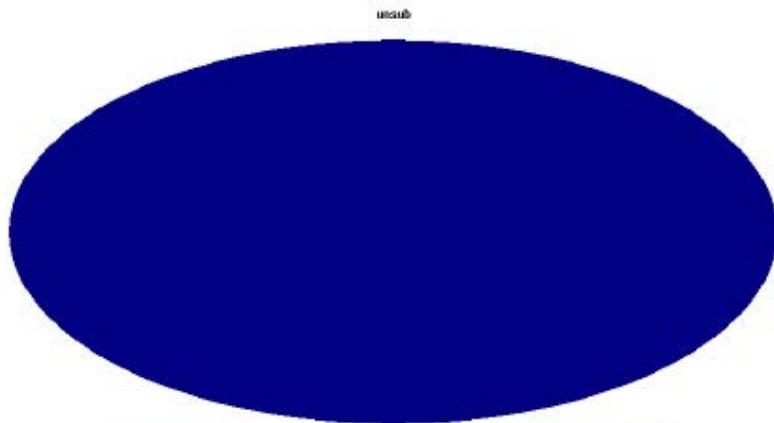
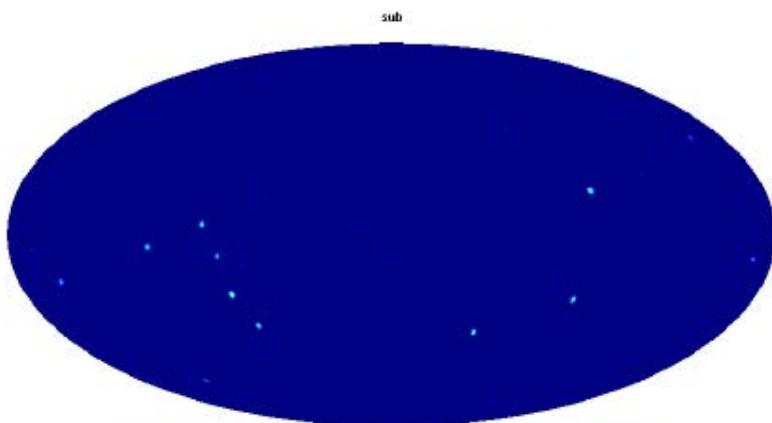
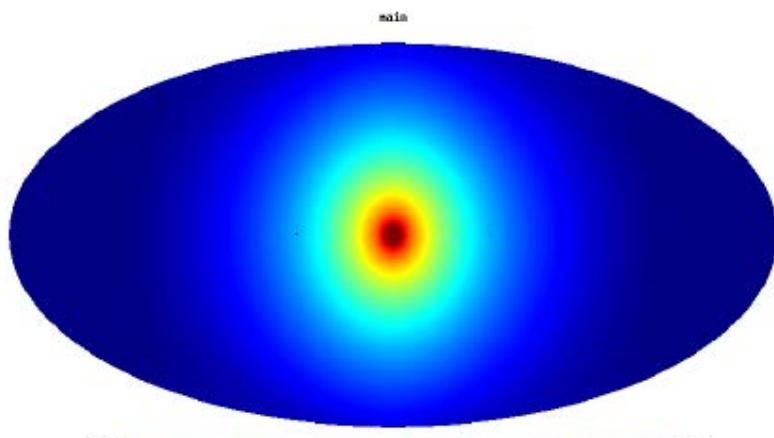
Density profiles



Spatial skymaps: CDM



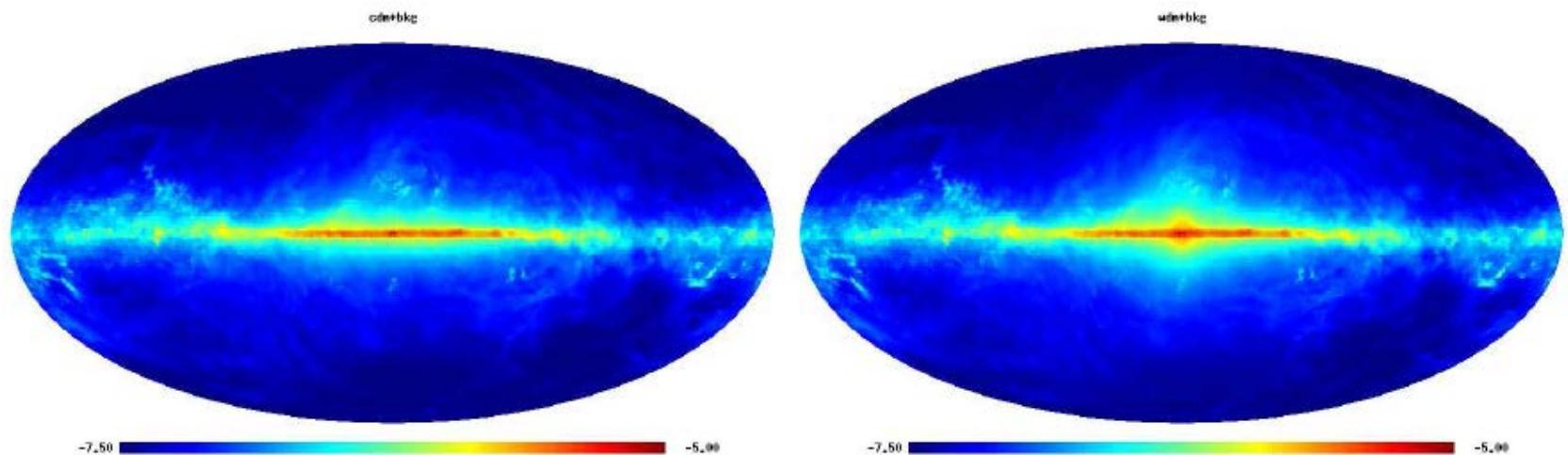
Spatial skymaps: WDM



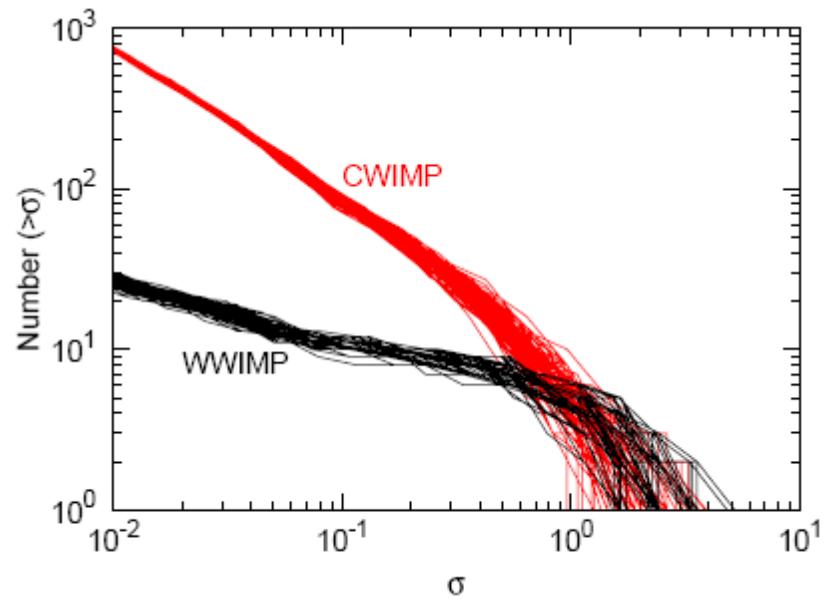
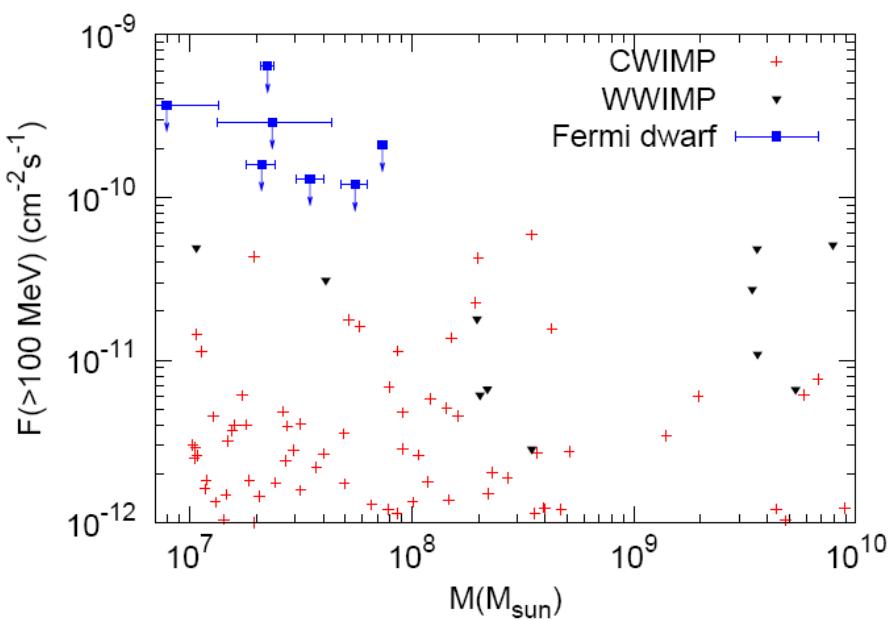
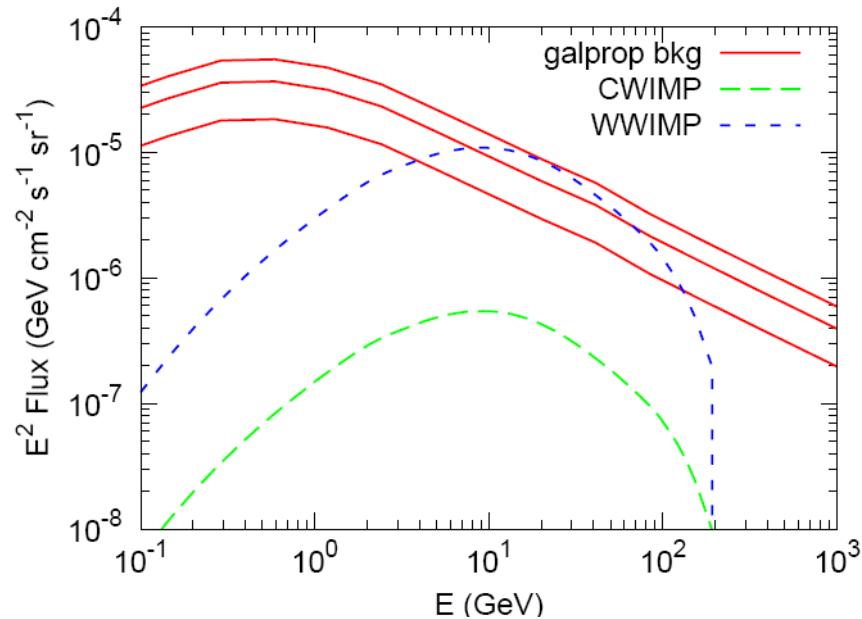
Two supersymmetric benchmark models

	m_0	m_{H_u}	m_{H_d}	$m_{1/2}$	A_0	$\tan\beta$	$\text{sign}(\mu)$	$m_{\tilde{\chi}_1^0}$	$\langle\sigma v\rangle$
Warm WIMP	1200	1300	788	500	-1000	40	+	211	2.70×10^{-25}
Cold WIMP	1200	1300	824	500	-1000	40	+	211	1.38×10^{-26}

Total skymaps with diffuse background ($E > 10$ GeV)

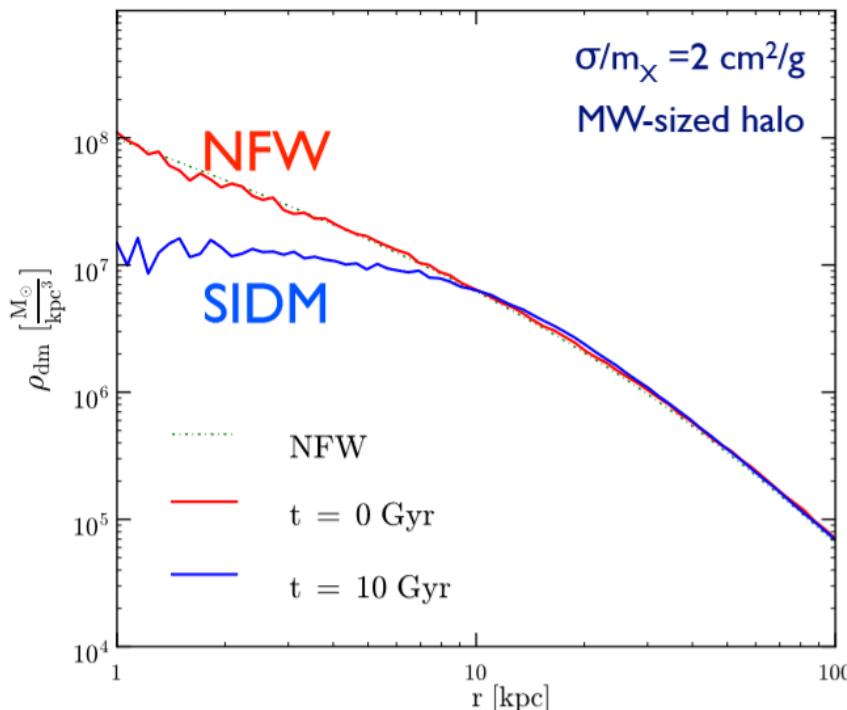


Detectability comparison



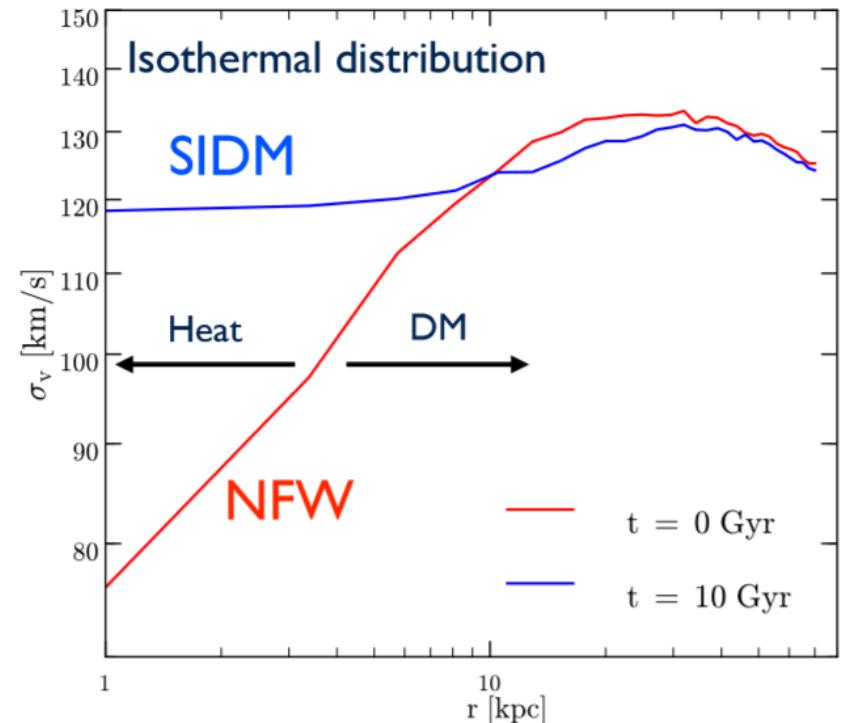
Self-Interacting Dark Matter

- Self-interactions thermalize the inner halo



$$\sigma/m_X \sim 1 \text{ cm}^2/\text{g}$$

$$\Gamma \simeq n\sigma v = (\rho/m_X)\sigma v \sim H_0$$

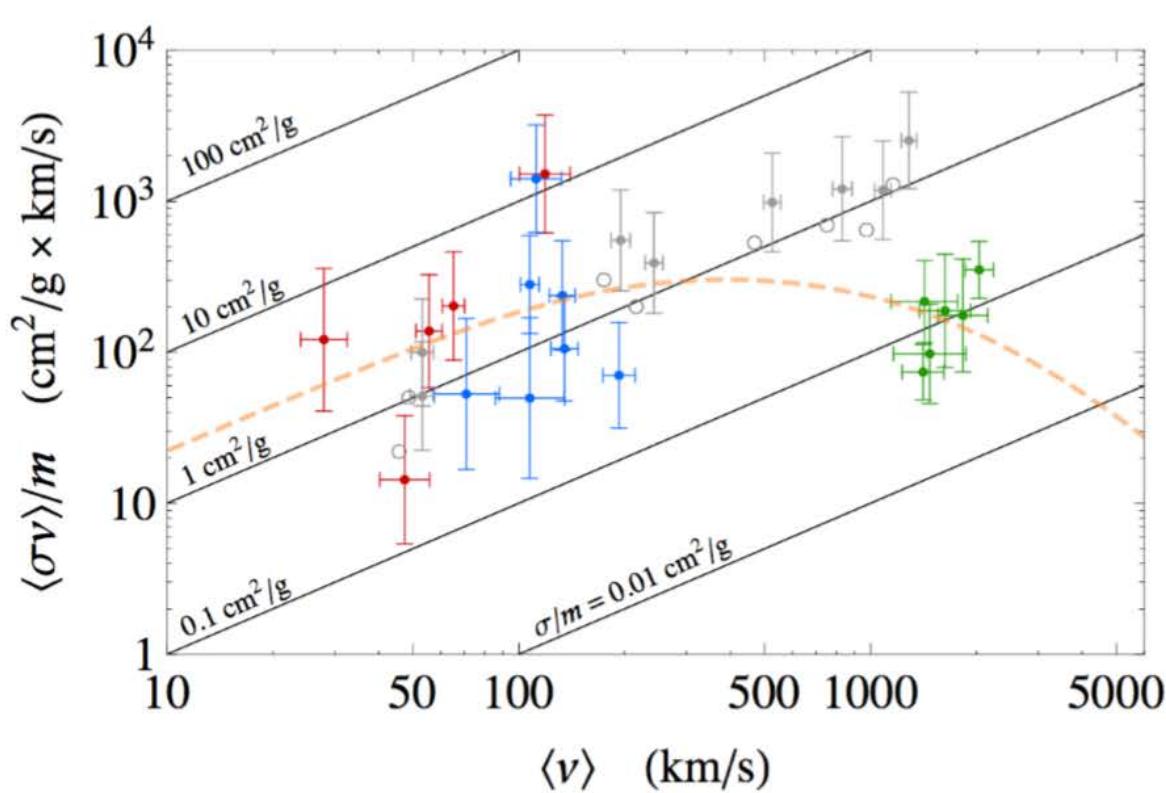


with Huo+(UCR SIDM code)

Tulin, HBY (2017) for a review

SIDM from Dwarfs to Clusters

- Consider 5 THINGS dwarfs (red), 7 LSBs (blue), 6 galaxy clusters (green)
- 8 simulated halos with $\sigma/m=1 \text{ cm}^2/\text{g}$ (gray) for calibration



Galaxies: $\sim 2\text{-}3 \text{ cm}^2/\text{g}$

Clusters: $\sim 0.1 \text{ cm}^2/\text{g}$

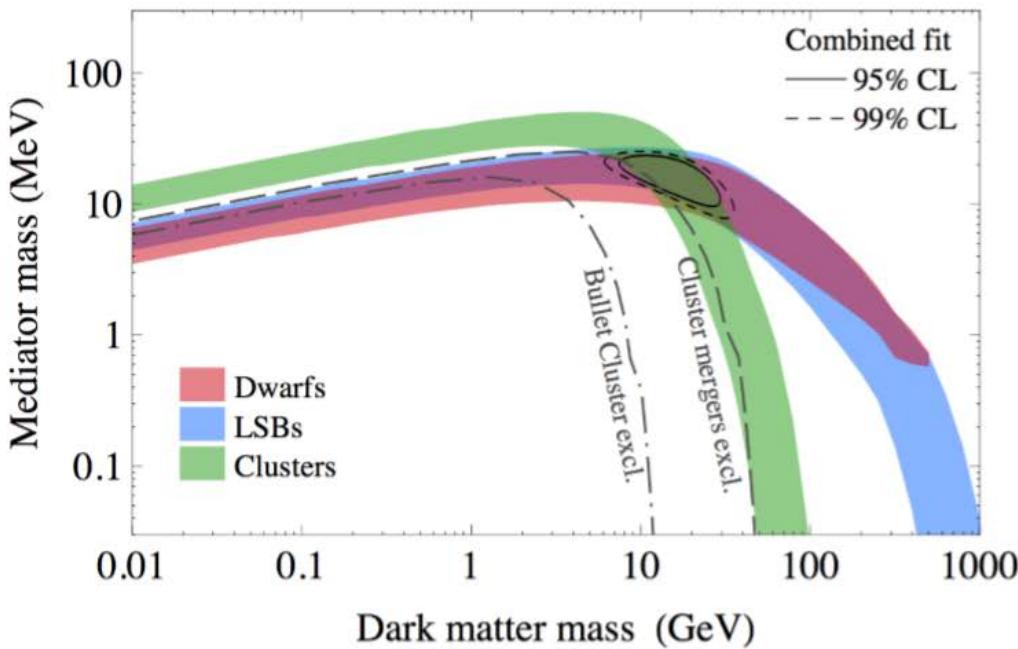
Core size in clusters: $\sim 10 \text{ kpc}$

If it were $\sim 1 \text{ cm}^2/\text{g}$ in clusters,
the core size would be $\sim 100 \text{ kpc}$

The strongest limit!

Measuring Dark Matter Mass

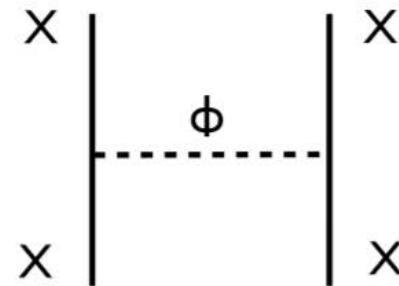
- Self-scattering kinematics determines SIDM mass



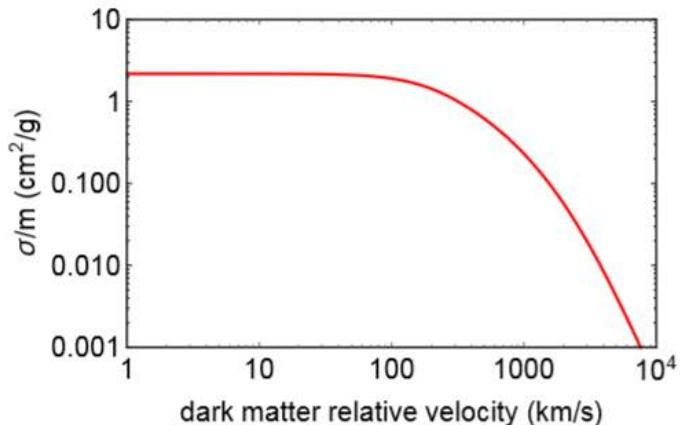
$$\alpha_X = 1/137$$

$$m_X \sim 15 \text{ GeV}, m_\phi \sim 17 \text{ MeV}$$

with Kaplinghat, Tulin (PRL 2015)

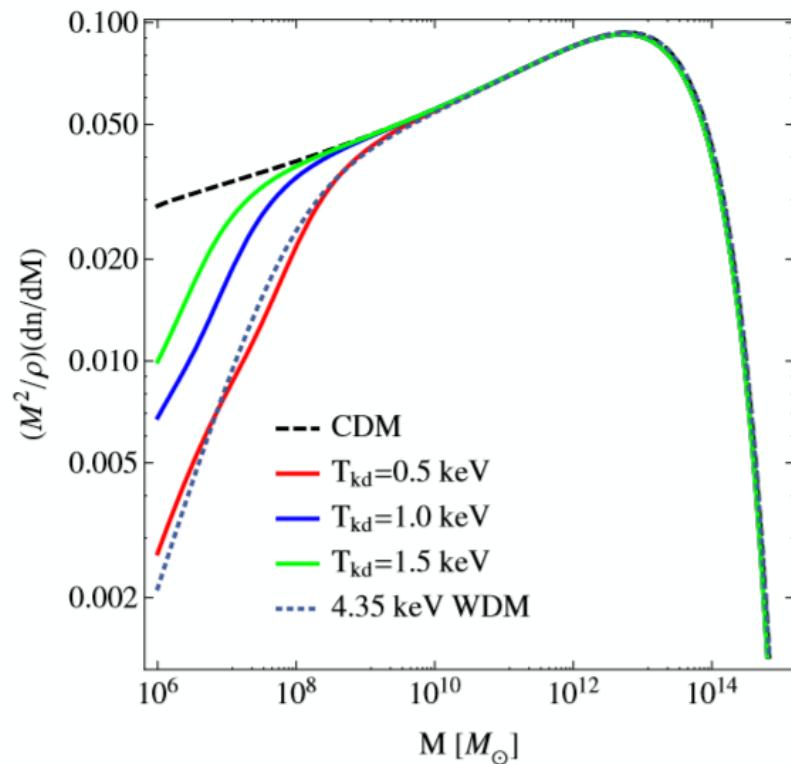
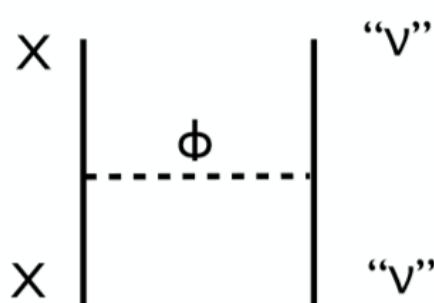
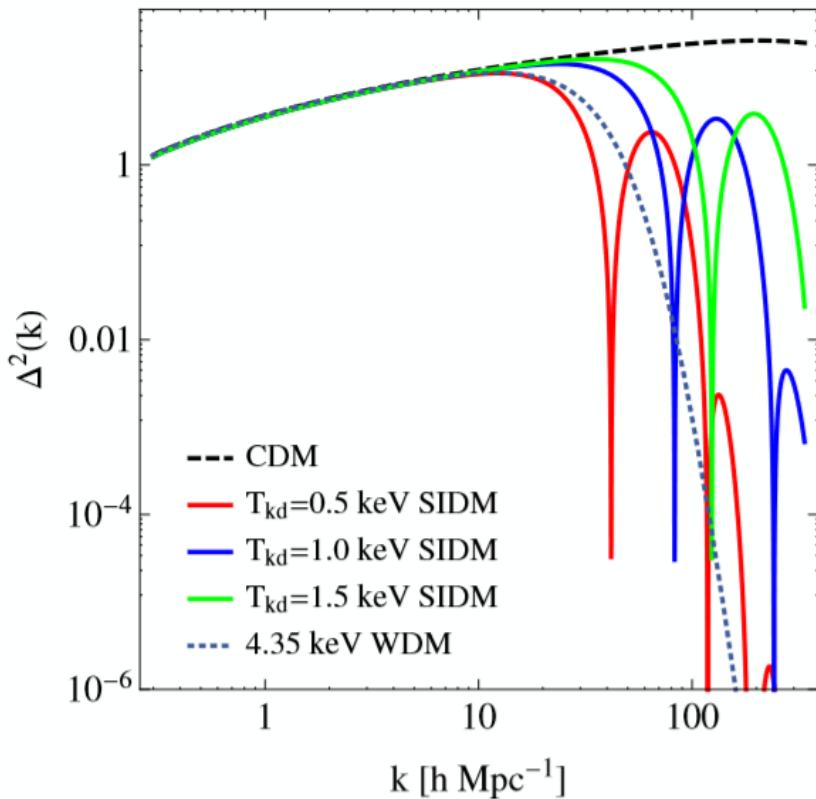


$$V(r) = \frac{\alpha_X}{r} e^{-m_\phi r}$$



Dark Acoustic Oscillation

- Roles of dark radiation, damped SIDM power spectrum



With Ran, Kaplinghat (in prep)

CDM (WIMP): $m_\phi \sim 1 \text{ TeV}$, $T_{kd} \sim 30 \text{ MeV}$
SIDM: $m_\phi \sim 10 \text{ MeV}$, $T_{kd} \sim 1 \text{ keV}$

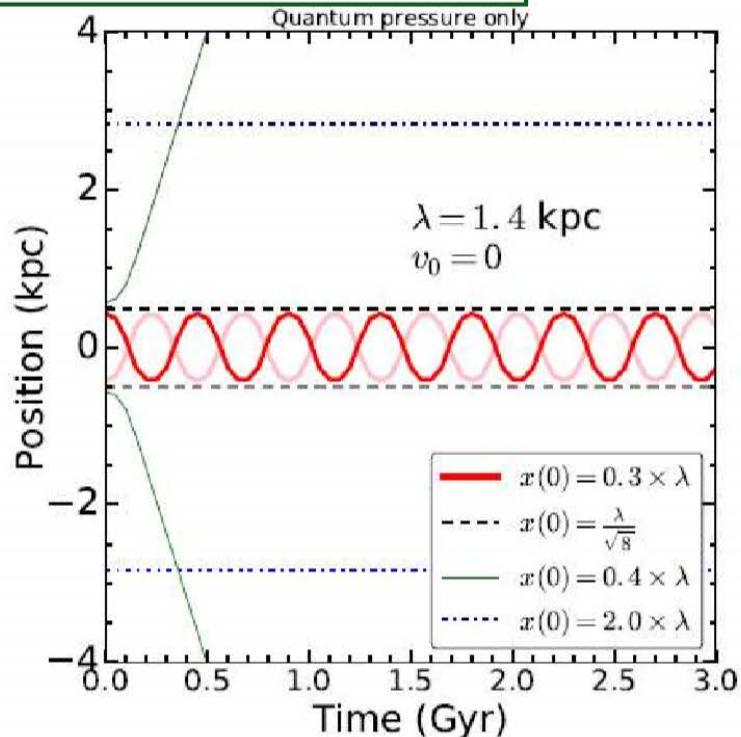
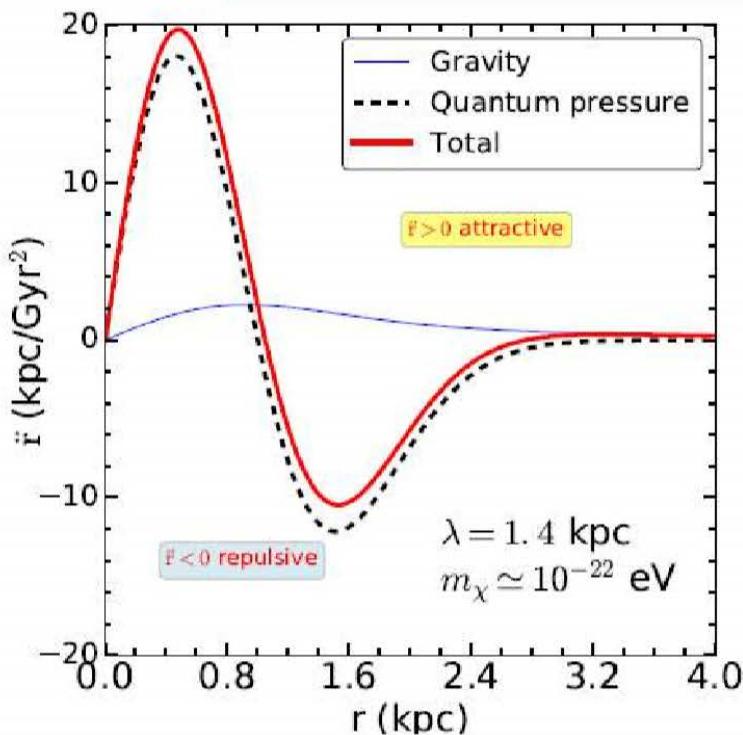
Ultra Light Axion

- T at $H^{\sim}m$ is 500 eV, $z=2e6$ (after BBN) behaves like matter, no contribution to Dark Energy.
- It is unnatural for just looking the mass within the order $1e-22$ eV but such mass can be natural to have correct relic density. In string theory, some mechanism can generate exponentially small mass.
- It is fluid and its mass has some tension with Lyman-alpha forest. (Feedback processes, such as galactic winds or outflows, are assumed to have negligible impact on the forest.)
- The matter wave length $\sim kpc$ if the velocity dispersion $\sim 100 \text{ km s}^{-1}$.

$$\frac{\lambda}{2\pi} = \frac{\hbar}{mv} = 1.92 \text{ kpc} \left(\frac{10^{-22} \text{ eV}}{m} \right) \left(\frac{10 \text{ km s}^{-1}}{v} \right)$$

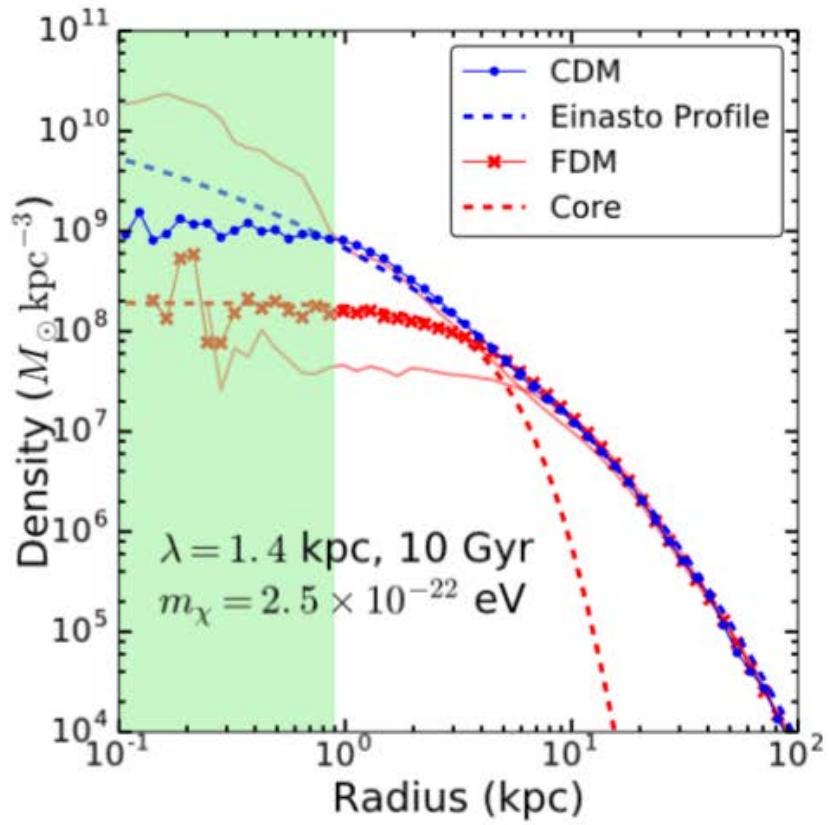
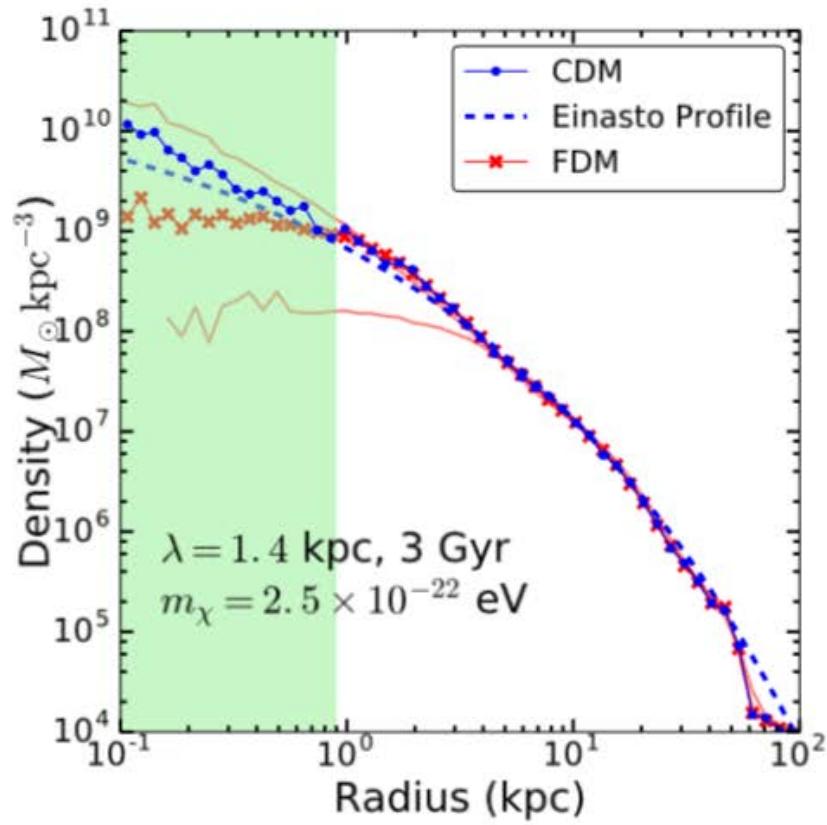
Quantum pressure in the N-body simulation

$$\ddot{r} = \frac{4M\hbar^2}{M_0 m_\chi^2 \lambda^4} \sum_j \exp\left[-\frac{2|r - r_j|^2}{\lambda^2}\right] \left(1 - \frac{2|r - r_j|^2}{\lambda^2}\right) (r_j - r).$$



The quantum pressure as a short-range interaction in the exponentially decay term.

Let's see the N-body simulation Movie



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