



Due on 25 Nov. 2025

Discussion on 25 Nov. 2025

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**Setup** In this exercise sheet we will explore the basic properties of two coupled superconductors. The order parameter  $\psi$  of each superconductor can be written as

$$\psi = |\psi| e^{i\phi} \equiv \sqrt{\rho_s} e^{i\phi} \quad (1)$$

where  $\rho_s$  is interpreted as the density of Cooper pairs (two-electron pairs) and  $\phi$  is the phase of the order parameter. This order parameter behaves as a wave function and satisfies the time-dependent Schrödinger equation

$$i\dot{\psi} = H\psi$$

where we set  $\hbar = 1$  for simplicity,  $\dot{\psi} \equiv \partial\psi/\partial t$ , and  $H$  is the Hamiltonian.

## 1 1 DC Josephson effect ★ New

If the two superconductors are separated, they evolve independently. The wave function of the two-superconductor system can be written as a vector  $\Psi = (\psi_1, \psi_2)$ , where  $\psi_i$  is the order parameter of the  $i$ -th superconductor. The Schrödinger equation for the two system is given by

$$i \begin{pmatrix} \dot{\psi}_1 \\ \dot{\psi}_2 \end{pmatrix} = \begin{pmatrix} E_1 & 0 \\ 0 & E_2 \end{pmatrix} \begin{pmatrix} \psi_1 \\ \psi_2 \end{pmatrix}$$

Now we bring the two superconductors closer and put a thin insulating layer between the two superconductors, which creates a coupling effect. As we discussed in exercise sheet 7, the coupling is characterized by the off-diagonal elements in the Hamiltonian. Let us add in an extra coupling constant  $\kappa$  to the Hamiltonian matrix:

$$i \begin{pmatrix} \dot{\psi}_1 \\ \dot{\psi}_2 \end{pmatrix} = \begin{pmatrix} E_1 & \kappa \\ \kappa & E_2 \end{pmatrix} \begin{pmatrix} \psi_1 \\ \psi_2 \end{pmatrix} \quad (2)$$

This off-diagonal term also encodes the tunneling effect.<sup>1</sup>

- (a) Using Eq. (1), prove

$$\frac{\dot{\psi}_i}{\psi_i} = \frac{1}{2} \frac{\dot{\rho}_i}{\rho_i} + i\dot{\phi}_i, \quad i = 1, 2 \quad (3)$$

- (b) Rewrite Eq. (2) and prove

$$i \frac{\dot{\psi}_1}{\psi_1} = E_1 + \kappa \sqrt{\frac{\rho_2}{\rho_1}} e^{i\delta}, \quad \delta \equiv \phi_2 - \phi_1 \quad (4)$$

similar for  $\psi_2$ .

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<sup>1</sup>It is good to review the concept of coupling and the physical interpretation of the off-diagonal term in the Hamiltonian.

(c) Combine Eq. (3) and Eq. (4). Prove that from the real part of the equation, we have

$$-\dot{\phi}_1 = E_1 + \kappa \sqrt{\frac{\rho_2}{\rho_1}} \cos \delta \quad (5)$$

and from the imaginary part, we have

$$\dot{\rho}_1 = 2\kappa \sqrt{\rho_1 \rho_2} \sin \delta \quad (6)$$

Similar for  $\psi_2$  and  $\rho_2$ .

- (d) Use Eq. (6) for  $\rho_1$  and  $\rho_2$  to prove that the Cooper pair density is conserved in the entire system, i.e.  $\rho_1 + \rho_2 = \text{const.}$
- (e) Prove that the current density is<sup>2</sup>

$$J = 4e\kappa \sqrt{\rho_1 \rho_2} \sin \delta \equiv J_c \sin \delta \quad (7)$$

This means that as long as there is a phase difference, there is a current flowing between the two superconductors.

## 2 2 AC Josephson effect ★ New

Now we assume that the two superconductors have the same properties, i.e. same material, same transition temperature, and same Cooper pair density:  $\rho_1 = \rho_2$ .

- (a) Use Eq. (2) to show that

$$\dot{\delta}(t) = E_1 - E_2 \equiv 2eV, \quad (8)$$

where  $V$  is the voltage across the junction. The extra factor of 2 again comes from the fact that each Cooper pair carries 2 electrons. Prove that this further leads to

$$J = J_c \sin(2eVt) \quad (9)$$

- (b) Analyze the dimension, and resume the  $\hbar$  in Eq. 9.
- (c) Calculate the frequency of the AC Josephson effect when the voltage  $V = 1 \text{ V}$ . How is this compared to the frequency of (1) the visible light, soft X-ray, microwave; (2) plasma frequency; (3) electron cyclotron frequency; (4) phonons?

## 3 3 SQUID ★ New

In exercise sheet 9, we derive the London equation when the order parameter  $\psi$  is uniform. However, if the phase is not uniform, the current density in the superconductor will not only depend on the vector potential  $A$ , but also depend on the gradient of the phase  $\nabla\phi$ :

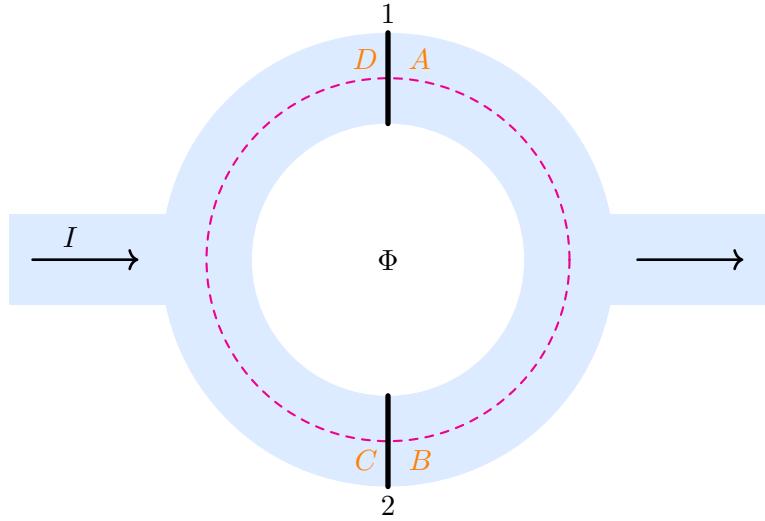
$$J = \frac{\rho e}{m} (\nabla\phi - 2eA) \quad (10)$$

where  $m$  is the mass of the electron, and  $\hbar = 1$ .

Let us consider a superconducting ring with two junctions denoted by 1 and 2, shown in the figure below.

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<sup>2</sup>The extra factor of 2 comes from the fact that each Cooper pair carries 2 electrons.



- (a) Choose a path deep inside the ring (dashed line) where the current density is zero. Calculate the loop integral of both sides of Eq. (10). Prove that

$$2\pi \frac{\Phi}{\Phi_0} = \delta_1 - \delta_2, \quad (11)$$

where  $\Phi$  is the magnetic flux,  $\Phi_0 = \frac{h}{2e}$  is the flux quantum,  $\delta_1 \equiv \phi_A - \phi_D$  and  $\delta_2 \equiv \phi_B - \phi_C$ .

- (b) Assume that two junctions are identical, for simplicity. The current density flowing through the junction  $i$  is denoted by  $J_i$ . Use the relation ( $J = J_c \sin \delta_1 + J_c \sin \delta_2$ ) to prove that<sup>3</sup>

$$J = 2J_c \sin \left( \frac{\delta_1 + \delta_2}{2} \right) \cos \left( \frac{\delta_1 - \delta_2}{2} \right) \quad (12)$$

which further leads to

$$J = 2J_c \sin \left( \frac{\delta_1 + \delta_2}{2} \right) \cos \left( \frac{\pi \Phi}{\Phi_0} \right)$$

- (c) Depending on different circuit setup,  $\delta_1$  and  $\delta_2$  could evolve differently with time. In some cases,  $\delta_1 + \delta_2$  leads to a high-frequency oscillation, while  $\delta_1 - \delta_2$  gives a low-frequency oscillation. Therefore, the current density in Eq. 12 represents a modulated oscillation (high-frequency oscillation modulated by a low-frequency oscillation). Plot the function

$$f(t) = \sin(100t) \cos(5t)$$

and describe what it looks like.

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<sup>3</sup>Hint:

$$\delta_1 = \frac{\delta_1 + \delta_2}{2} + \frac{\delta_1 - \delta_2}{2}, \quad \delta_2 = \frac{\delta_1 + \delta_2}{2} - \frac{\delta_1 - \delta_2}{2}$$