



李改道研究所
TSUNG-DAO LEE INSTITUTE

Ultralight dark matter: From small-scale structure to dynamic neutrino mass

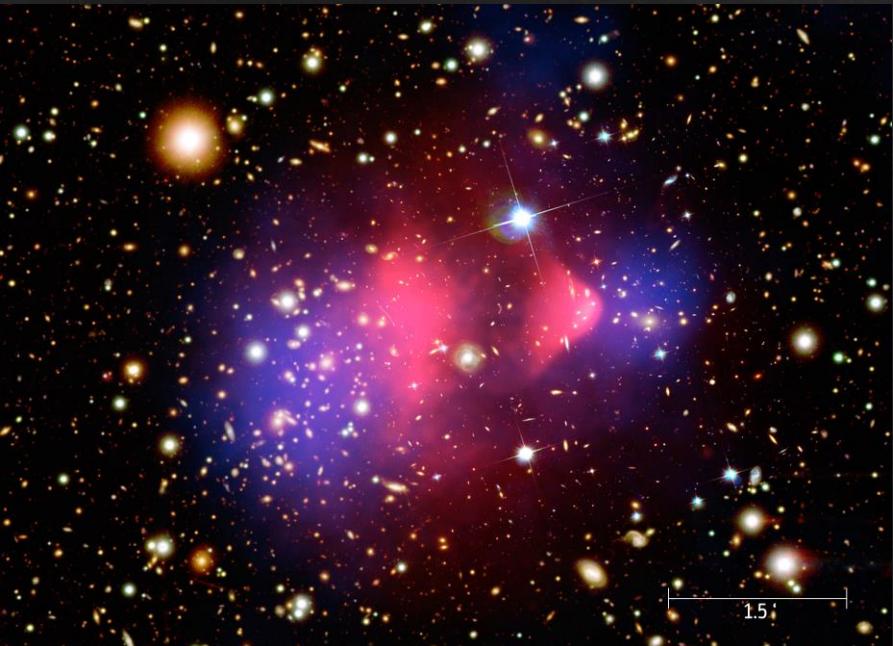
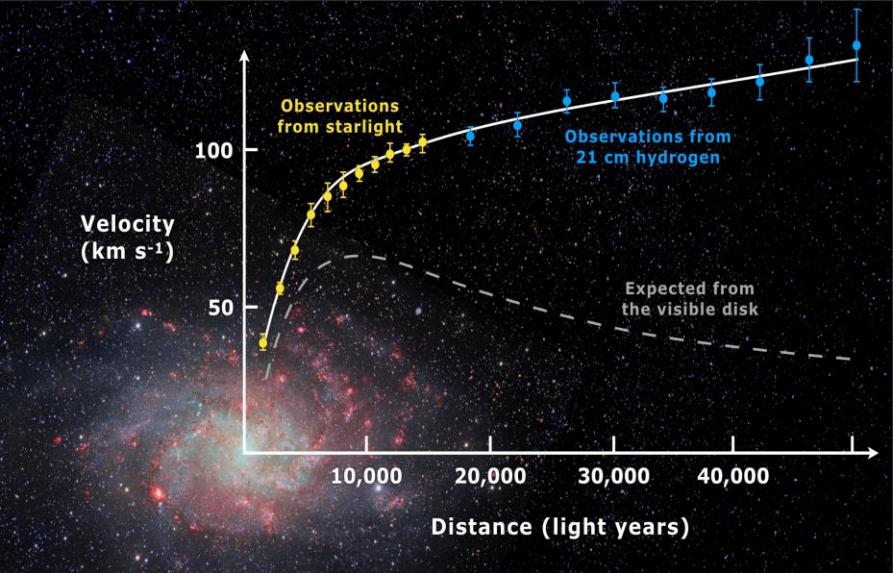
Hong-Yi Zhang

Tsung-Dao Lee Institute, Shanghai Jiao Tong University

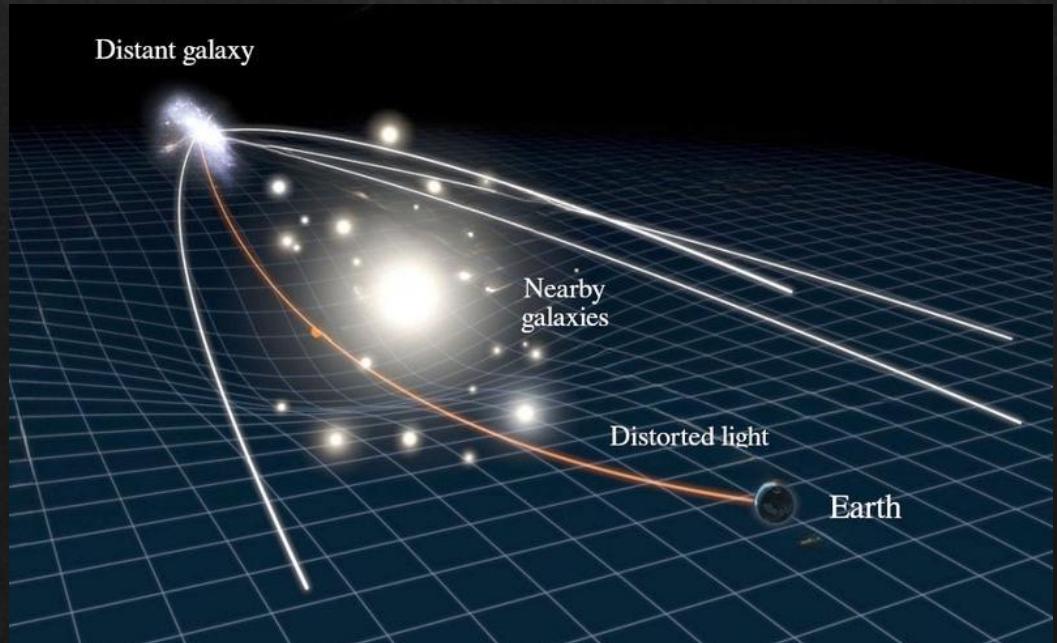
@ Università degli Studi di Salerno

Jun 10, 2025

Evidence of dark matter



Velocity dispersion of galaxy clusters
Galaxy rotation curves
Gravitational lensing
Bullet cluster
Cosmic microwave background
Large scale structure
Baryon acoustic oscillations
Type Ia supernovae, ...



How light could dark matter particles be?

Uncertainty principle

$$\Delta x \Delta p \geq \frac{\hbar}{2}$$

$$m \gtrsim 4.8 \times 10^{-21} \text{ eV} \left(\frac{10 \text{ km/s}}{v} \right) \left(\frac{0.02 \text{ kpc}}{\Delta x} \right)$$

Willman 1 (dSph, discovered in 2018)

$$\lambda_{\text{dB}} = 0.25 \text{ kpc} \left(\frac{4.8 \times 10^{-21} \text{ eV}}{m} \right) \left(\frac{10 \text{ km/s}}{v} \right)$$

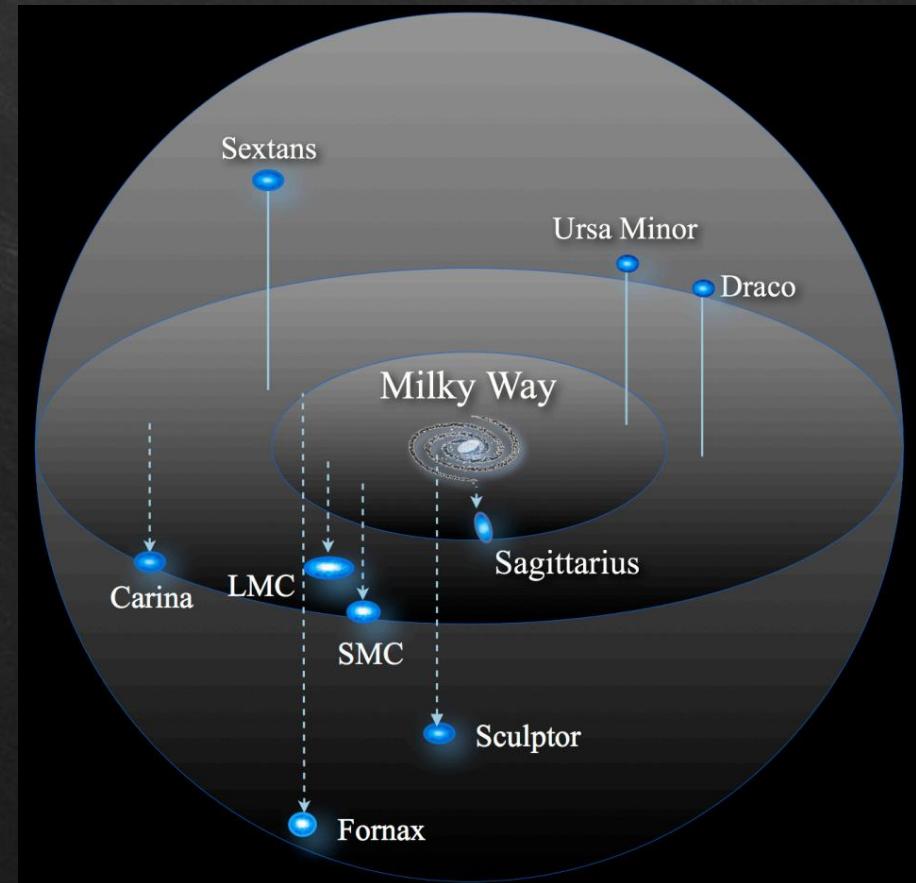


Figure from D. H. Weinberg et al. (PNAS, 2013)

If dark matter particles are fermions ...

Pauli exclusion principle

$$n = g \int \frac{d^3 p}{(2\pi)^3} f \lesssim \frac{g}{(2\pi)^3} \frac{4\pi}{3} (mv)^3$$

$$m \gtrsim 24 \text{ eV} \left(\frac{2}{g} \right)^{\frac{1}{4}} \left(\frac{\rho_{DM}}{0.4 \text{ GeV/cm}^3} \right)^{\frac{1}{4}} \left(\frac{200 \text{ km/s}}{v} \right)^{\frac{3}{4}}$$

Solar neighborhood

$$\lambda_{dB} = 7.7 \times 10^{-5} \text{ m} \left(\frac{24 \text{ eV}}{m} \right) \left(\frac{200 \text{ km/s}}{v} \right)$$

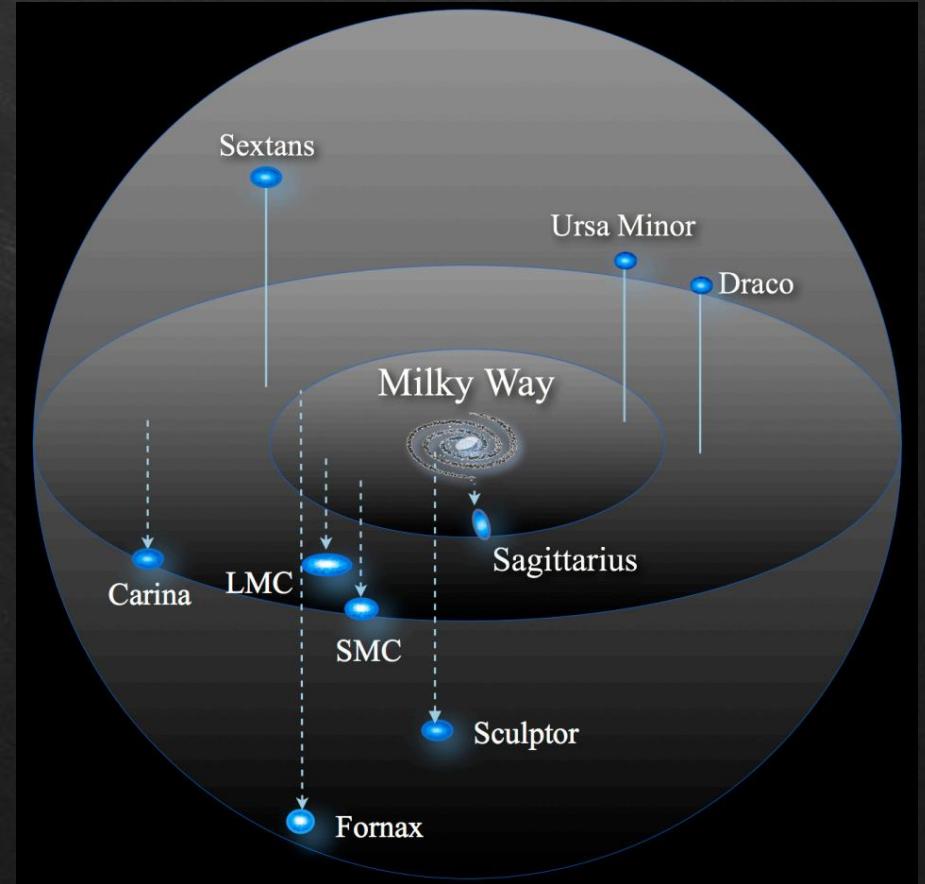
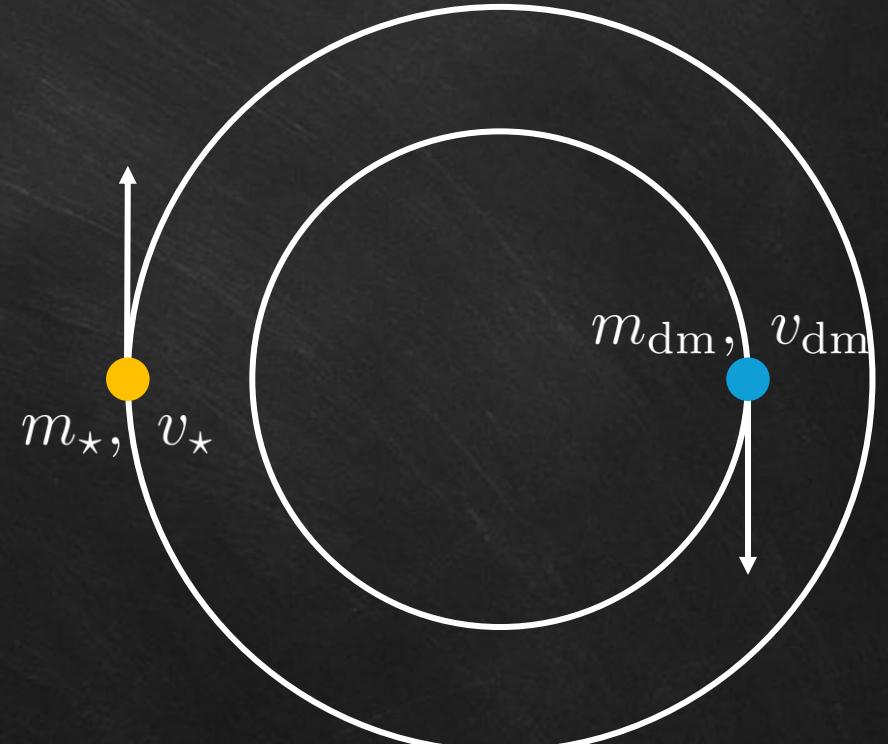


Figure from D. H. Weinberg et al. (PNAS, 2013)

How heavy could dark matter “particles” be?

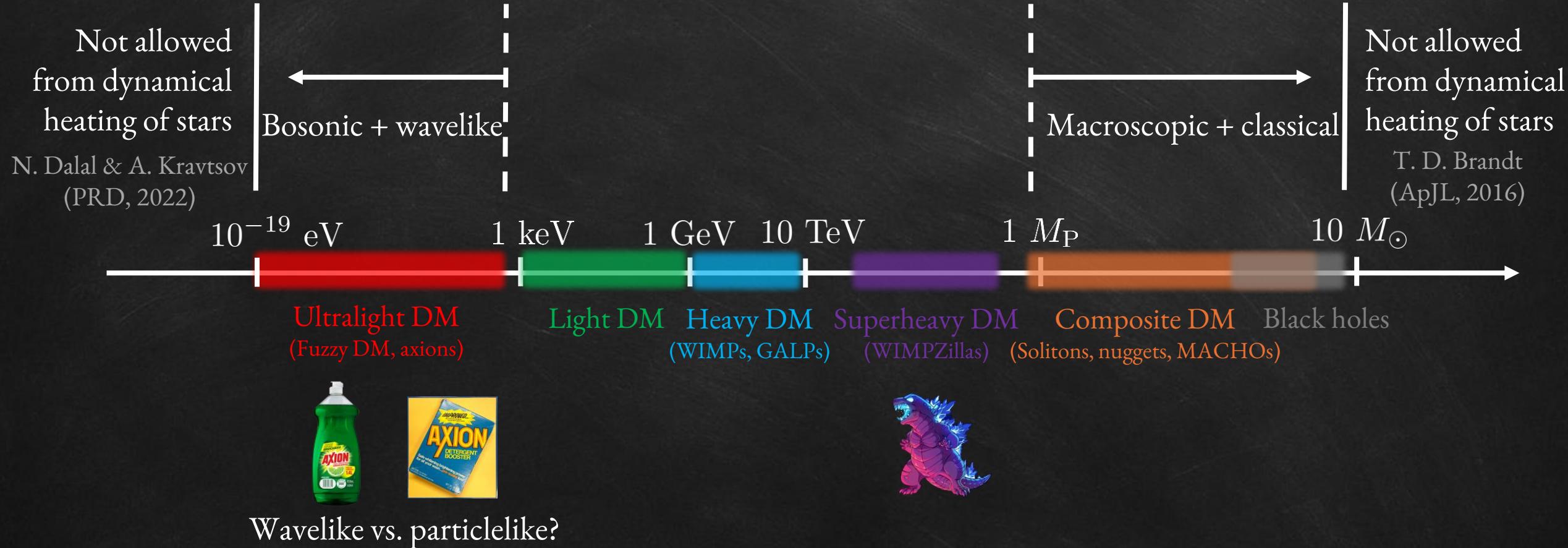
$$v_\star = \frac{m_{\text{dm}}}{m_\star} v_{\text{dm}} \sim \frac{m_{\text{dm}}}{m_\star} \sqrt{\frac{GM_{\text{gal}}}{R_{\text{gal}}}}$$

↓
Virial/escape velocity



Too heavy dark matter ($\gg M_\odot$) → Stellar heating

Dark matter mass landscape



Ultralight dark matter



- Large occupation number → Classical fields

$$n\lambda_{dB}^3 \sim \left(\frac{40 \text{ eV}}{m}\right)^4 \sim 3 \times 10^{82} \left(\frac{10^{-19} \text{ eV}}{m}\right)^4$$

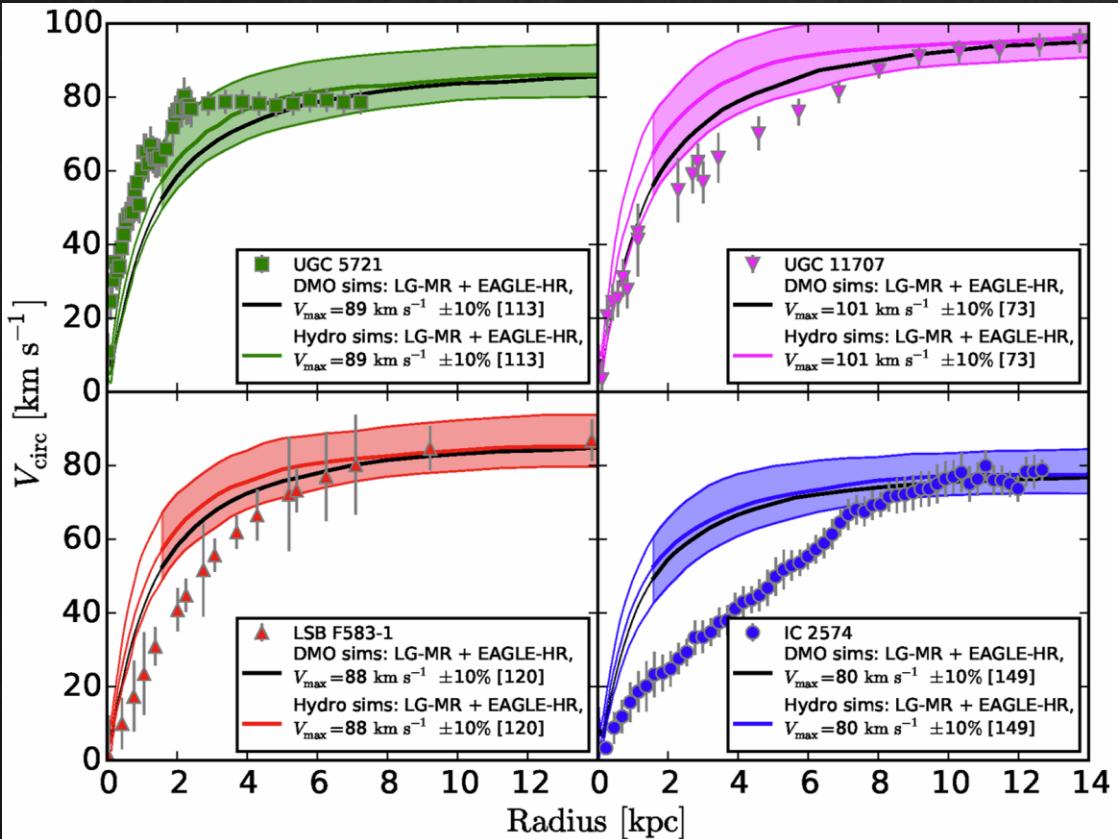
- Macroscopic/astrophysical scales

$$\lambda_{dB} \sim 50 \text{ } \mu\text{m} \left(\frac{40 \text{ eV}}{m}\right) \sim 0.6 \text{ pc} \left(\frac{10^{-19} \text{ eV}}{m}\right)$$

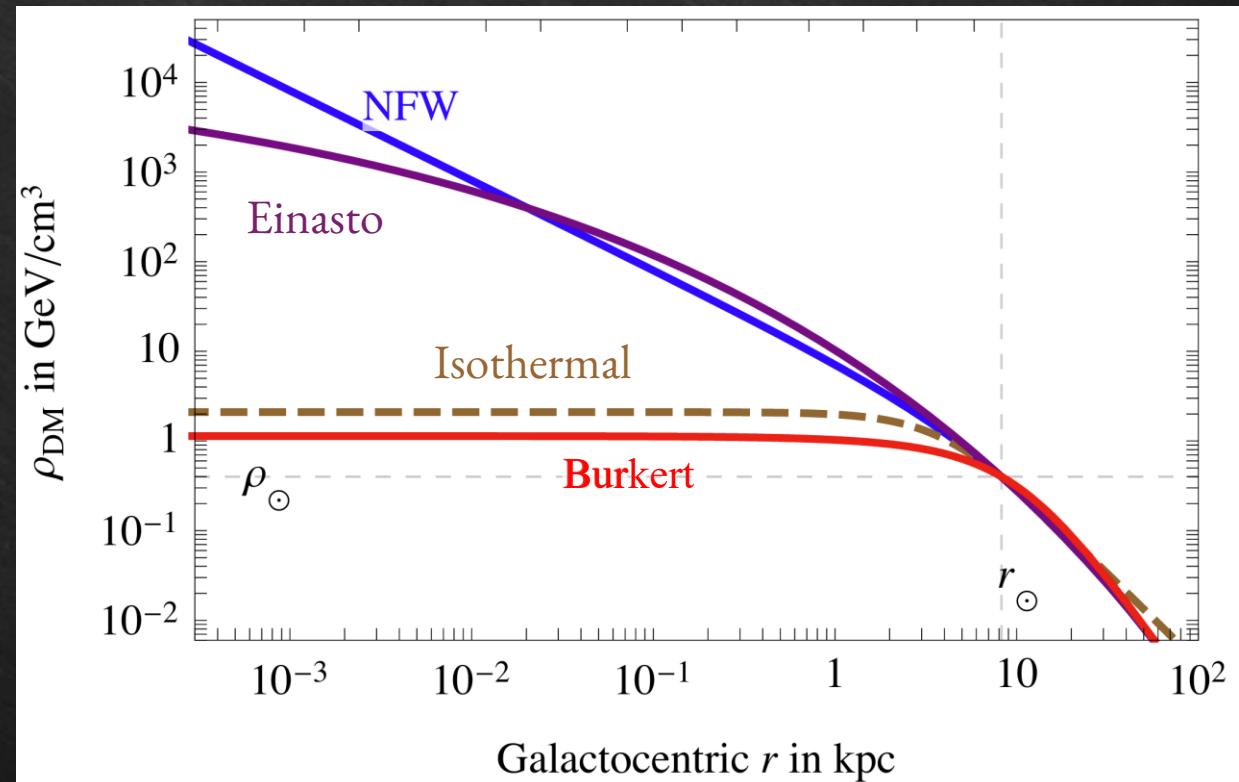
- Wave dynamics, rich phenomenology

Interference, Bose-Einstein condensation, polarization,
modulation of standard model constants, etc.

Diversity of dark matter profiles



K. A. Oman et al. (MNRAS, 2015)



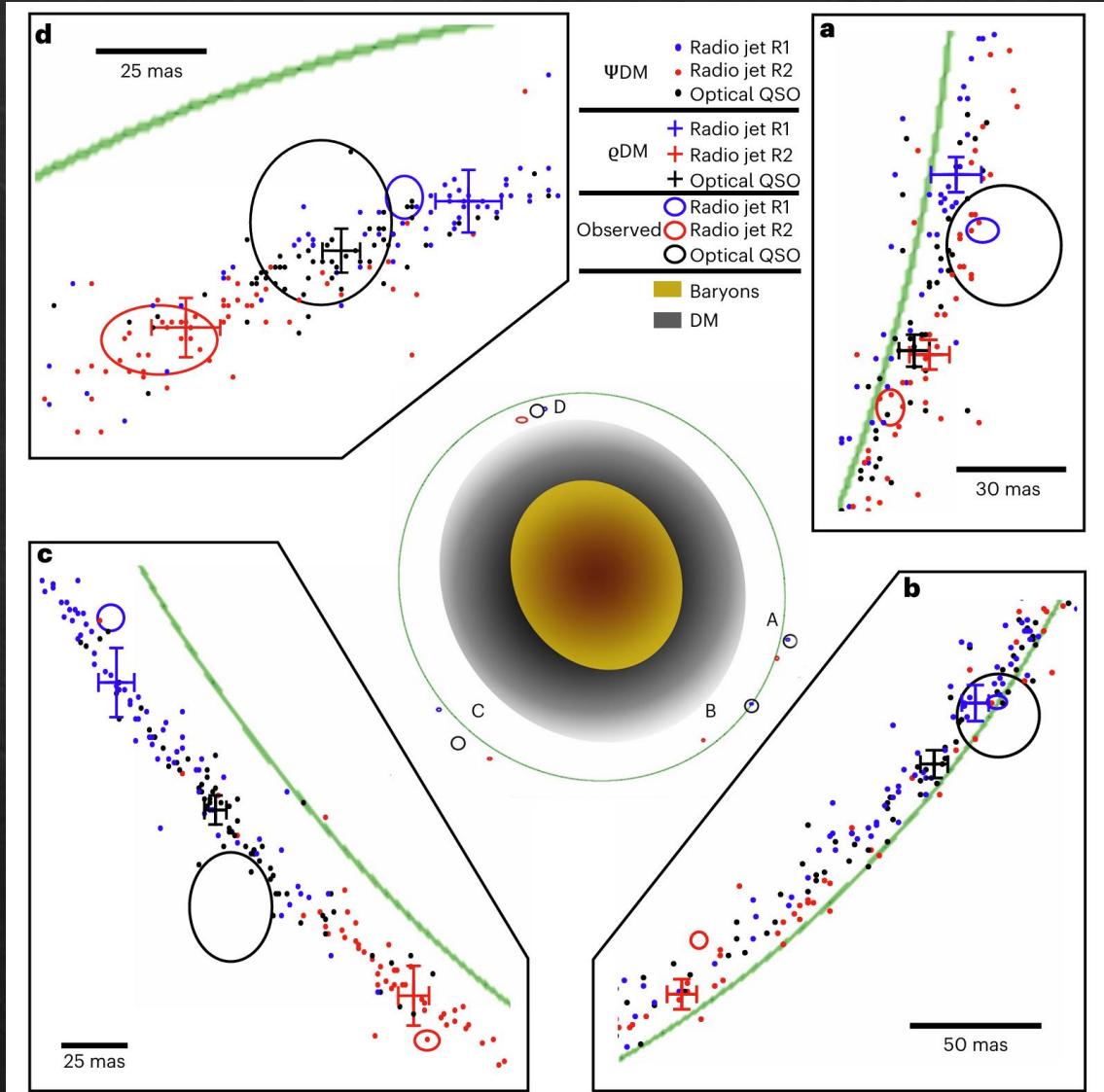
M. Cirelli et al. (2024)

Strong lensing anomalies in HS 0810+2554

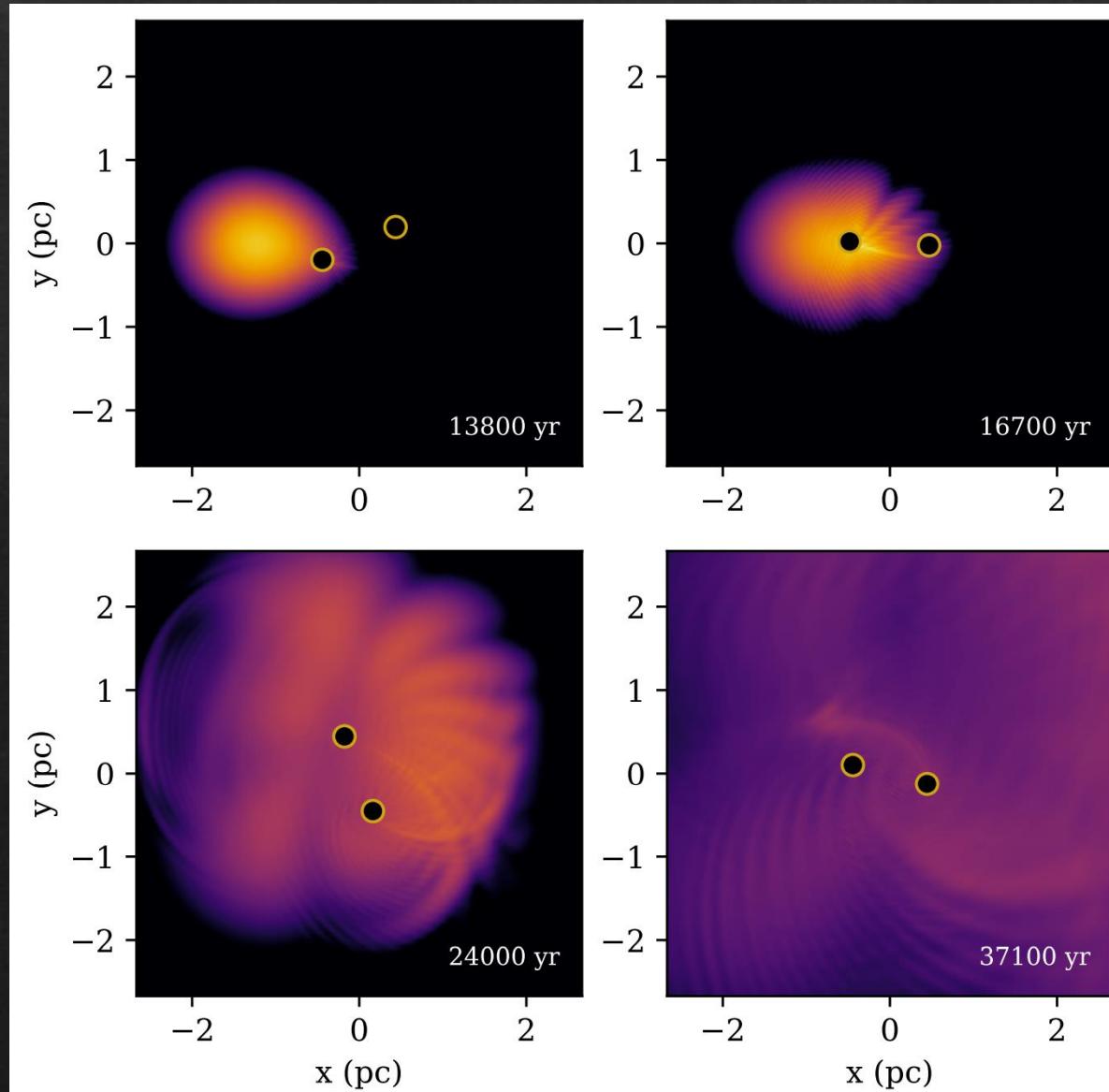
Big circles: Observed locations of a quasi-stellar object and two radio jets

Cross points: NFW profile

Points: 75 Gaussian realizations of fuzzy DM fluctuations



Final parsec problem



Mathematical prescriptions

$$i\partial_t \psi = -\frac{\nabla^2}{2ma^2} \psi + \frac{m}{a} \Phi \psi$$

$$\nabla^2 \Phi = \frac{1}{2M_P^2} (\rho - \bar{\rho})$$

$$\rho = m|\psi|^2$$

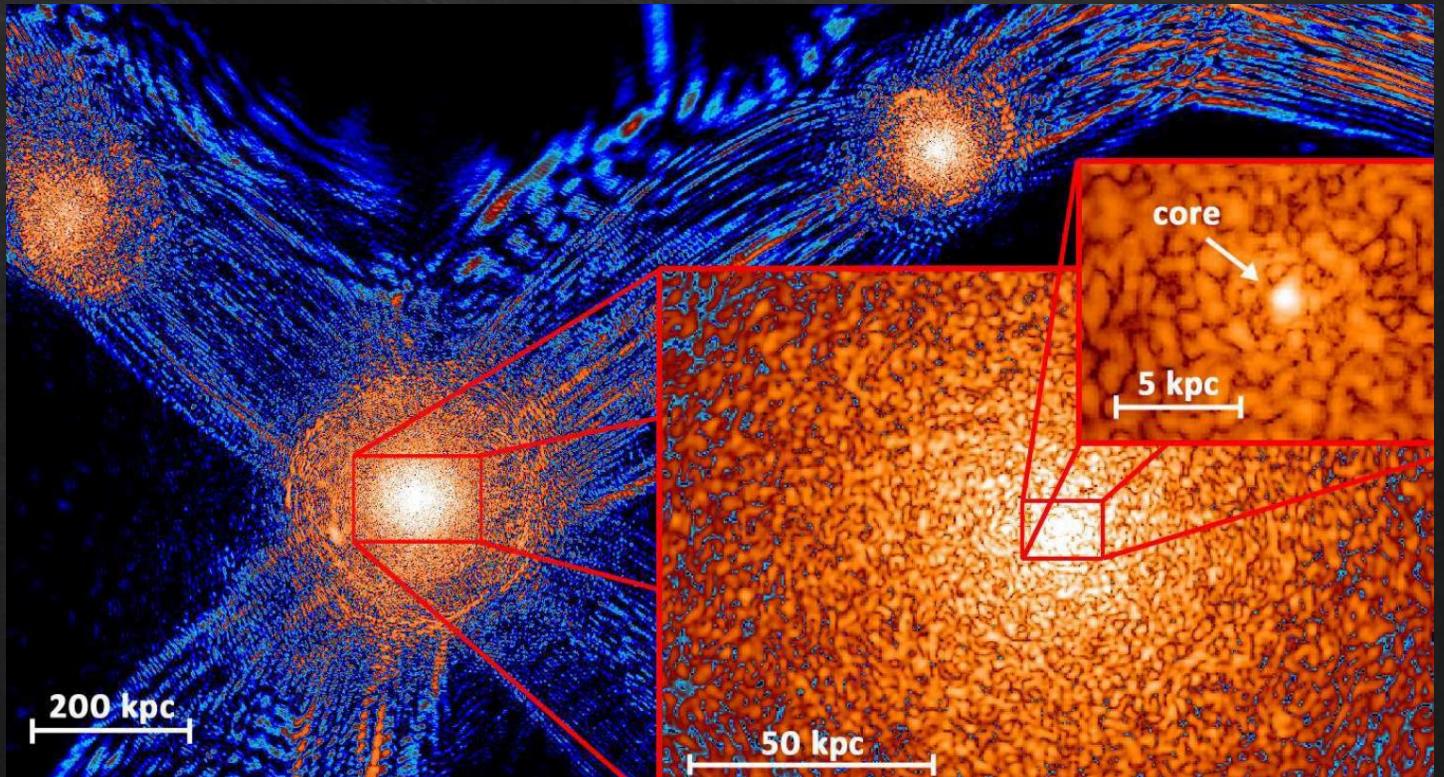
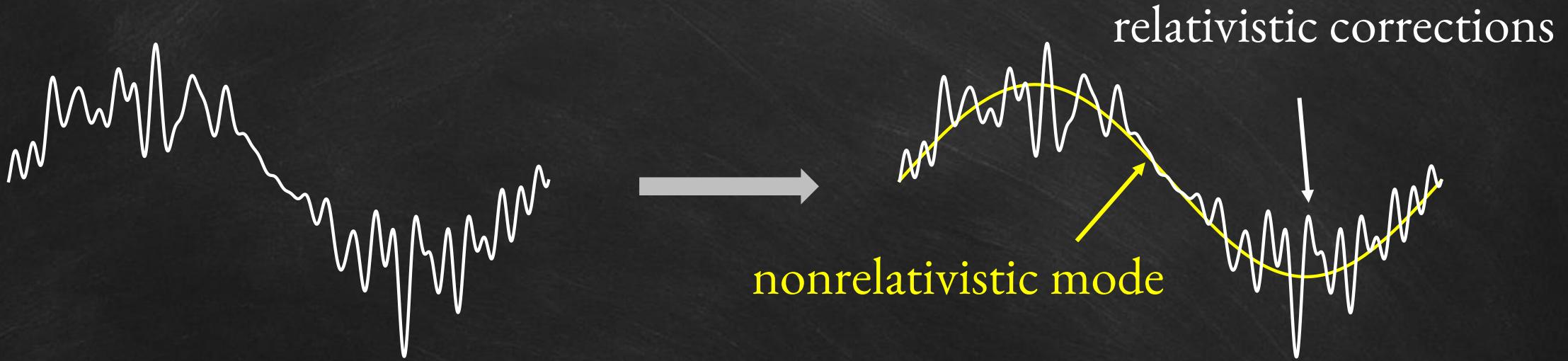


Figure from H.-Y. Schive et al. (Nature Physics, 2014)

Nonrelativistic effective field theory



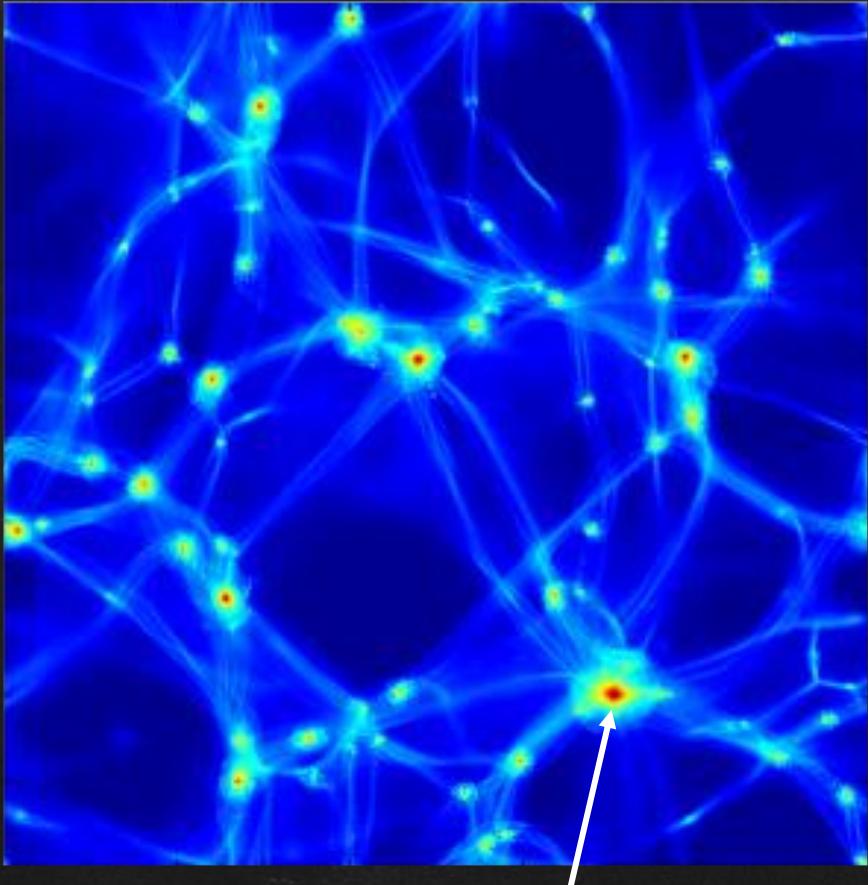
$$\phi(t, \mathbf{x}) = \frac{1}{\sqrt{2ma^3}} [\psi(t, \mathbf{x}) e^{-imt} + \text{c.c.}]$$

B. Salehian et al. (JHEP, 2020)

B. Salehian, H.-Y. Zhang et al (JHEP, 2021)

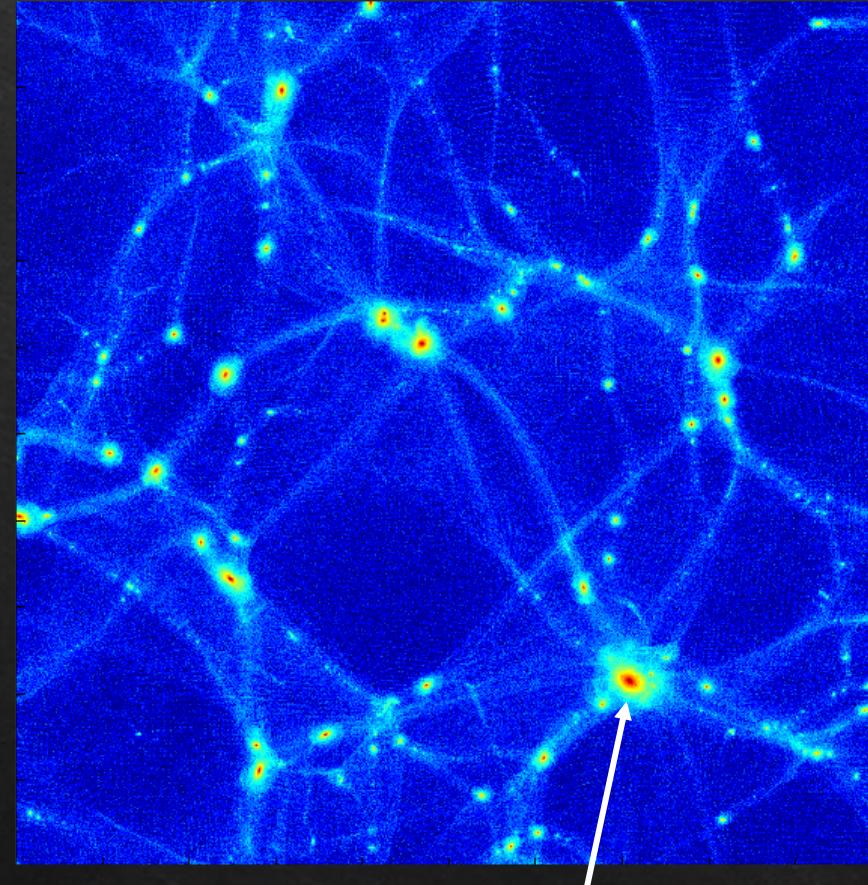
Large- and small-scale structure

Wavelike dark matter



Solitons, cored profiles

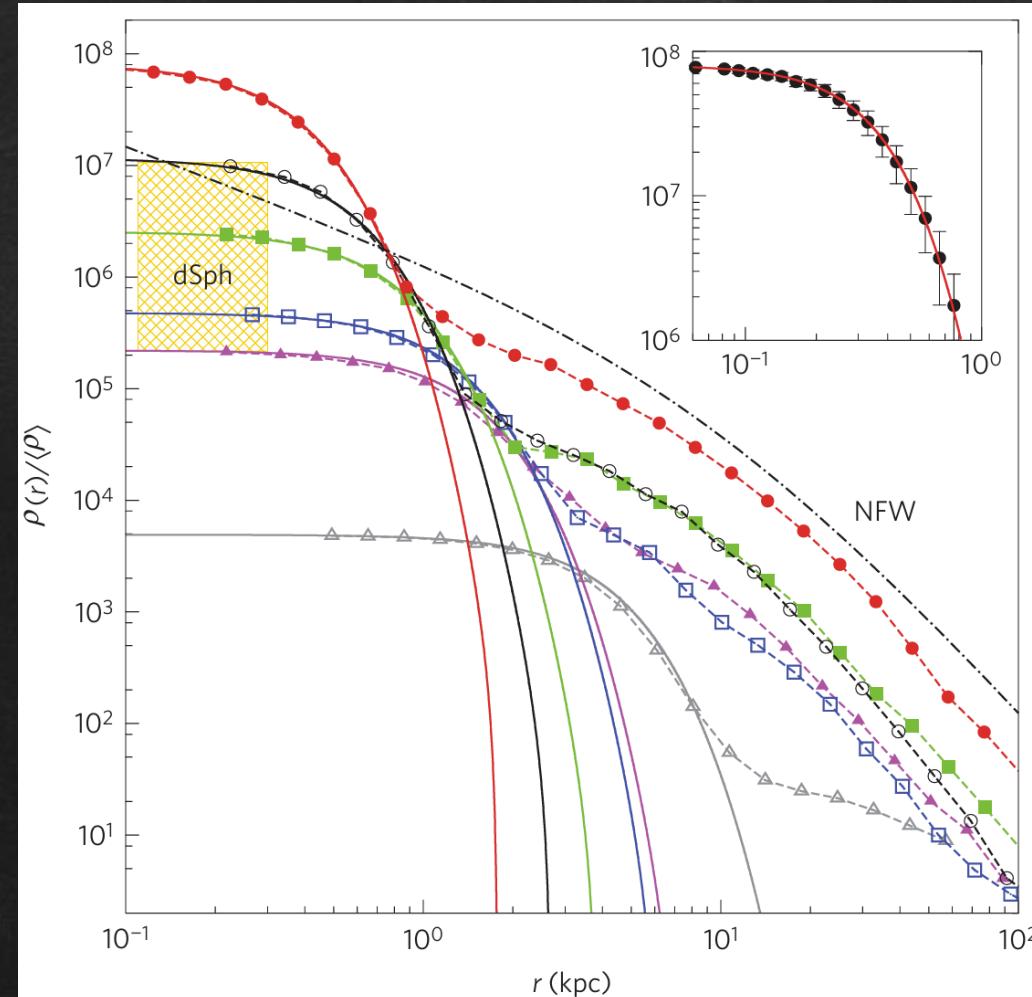
Particlelike dark matter



NFW, cuspidal profiles

Density profiles of ultralight dark matter halos

Soliton core + NFW profile



Solitons

$$\psi(t, \mathbf{x}) = f(r) e^{i\mu t} , \quad \mu \ll m$$

$$\phi(t, \mathbf{x}) \approx \sqrt{\frac{2}{m}} f(r) \cos(\omega t) , \quad \omega = m - \mu$$



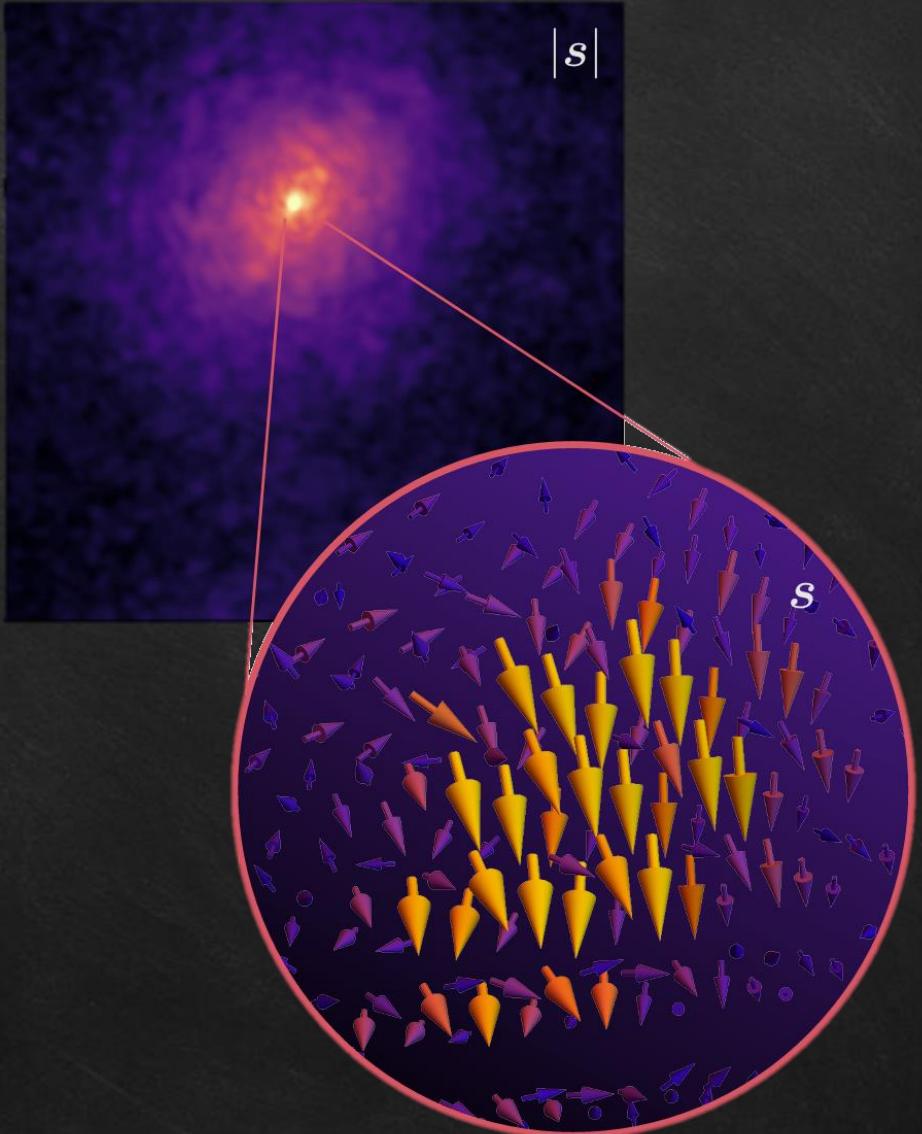
Ground state of the SP equations

Spin in vector solitons

$$\psi_i(t, \mathbf{x}, \sigma) = f(r)e_i(\sigma)e^{i\mu t}$$

$$e_i(0) = \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix} \quad , \quad e_i(\pm 1) = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ \pm i \\ 0 \end{pmatrix}$$

Linearly polarized Circularly polarized

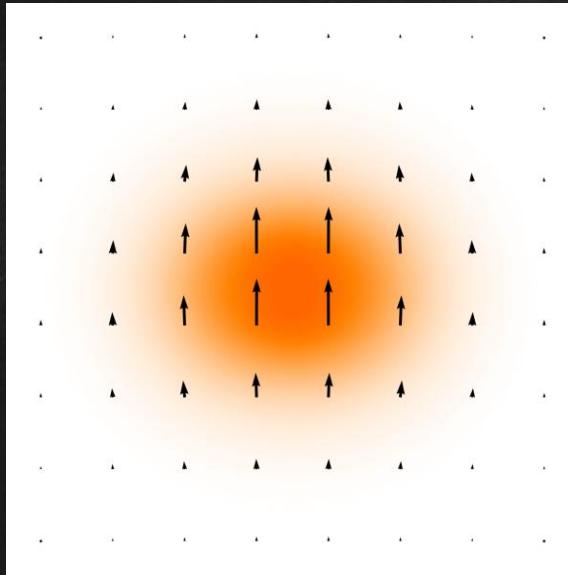


M. Jain, and M. A. Amin (PRD, 2022)
H.-Y. Zhang et al. (PRD, 2022)

Figure from M. A. Amin et al. (JCAP, 2022)

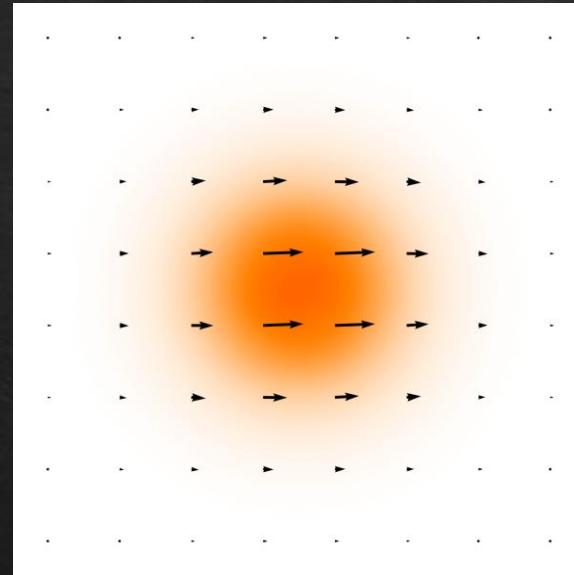
Vector solitons

$$X_i \approx \sqrt{\frac{2}{m}} f(r) \begin{pmatrix} 0 \\ 0 \\ \cos(\omega t) \end{pmatrix}$$



Linearly polarized

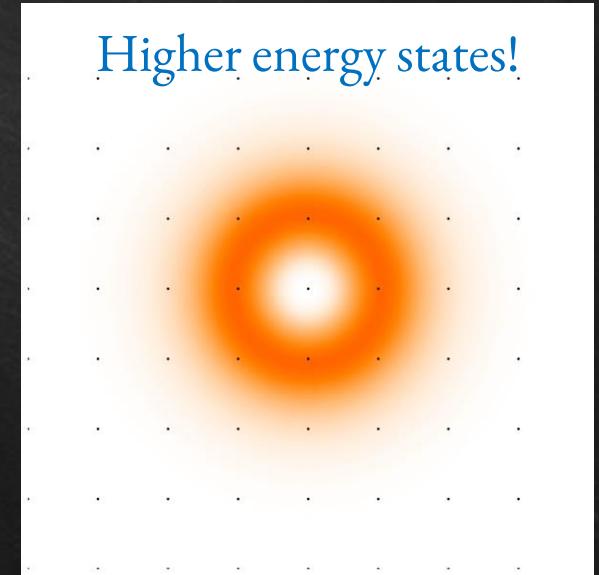
$$X_i \approx \frac{1}{\sqrt{m}} f(r) \begin{pmatrix} \cos(\omega t) \\ \sin(\omega t) \\ 0 \end{pmatrix}$$



Circularly polarized

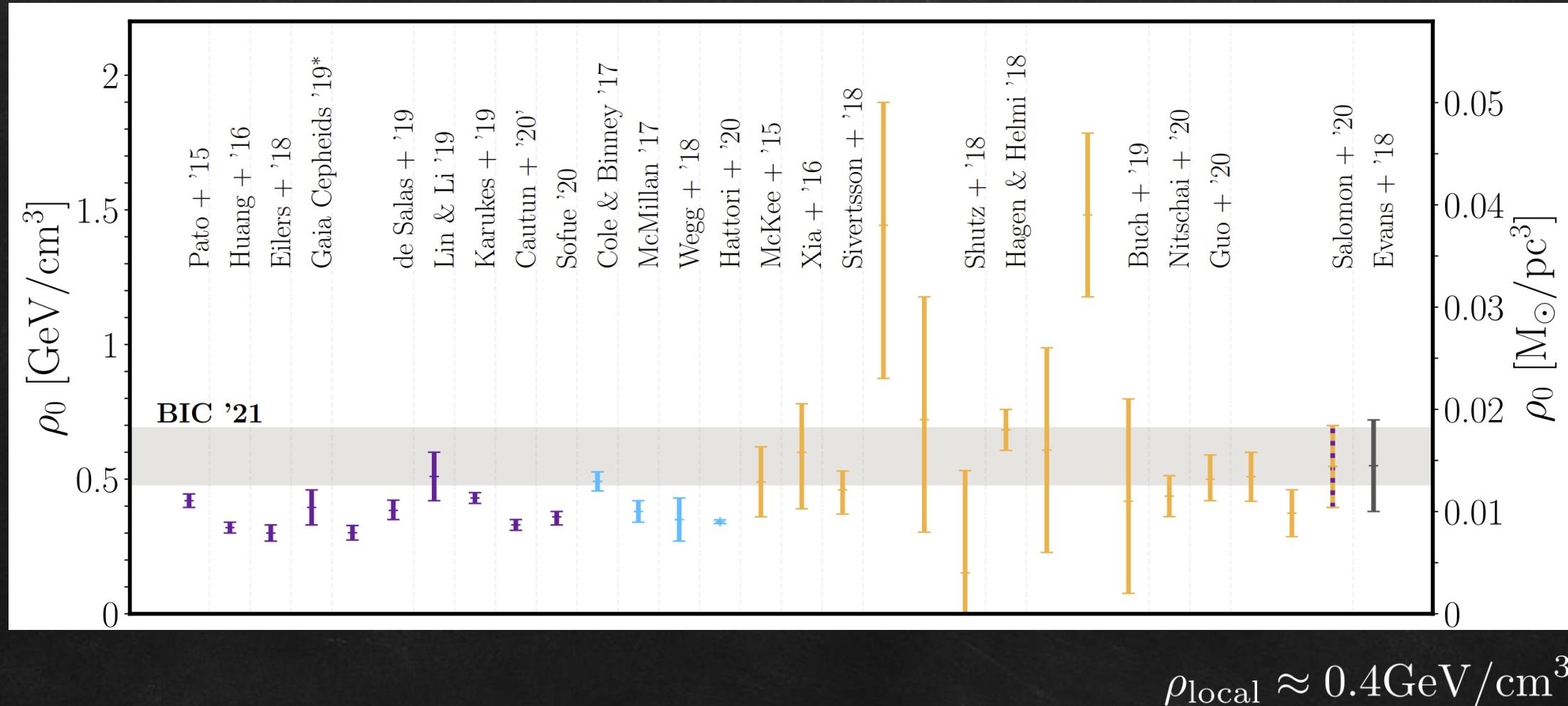
$$X_i \propto g(r) \cos(\omega t) \hat{r}$$

Higher energy states!



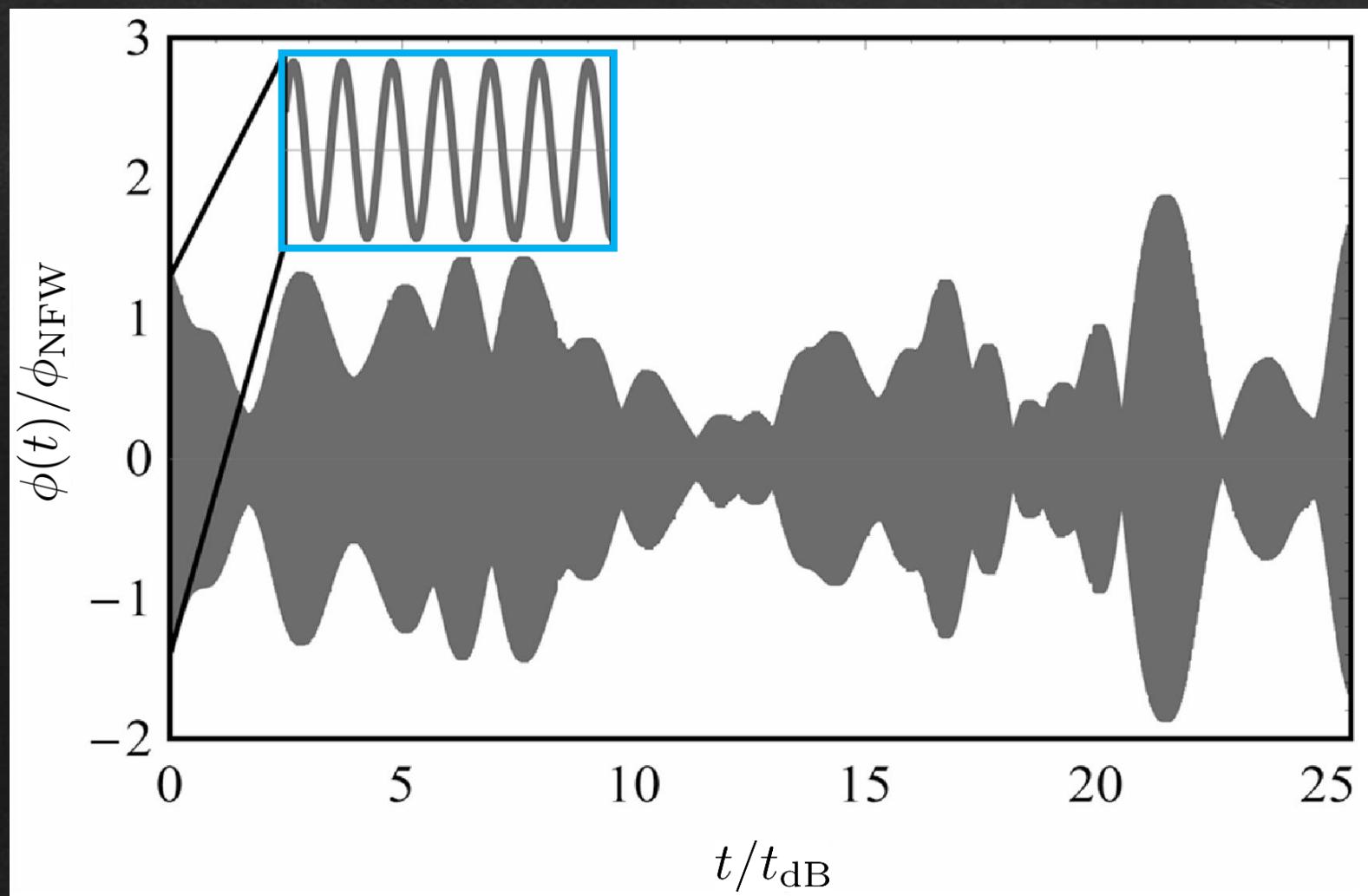
Spherically symmetric
(Solutions with a node)

Local dark matter density



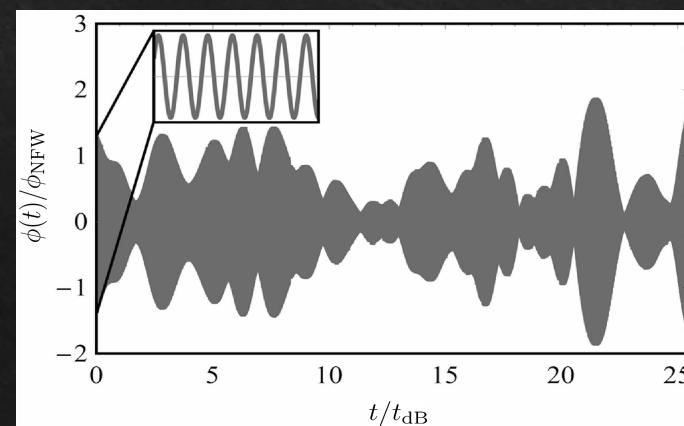
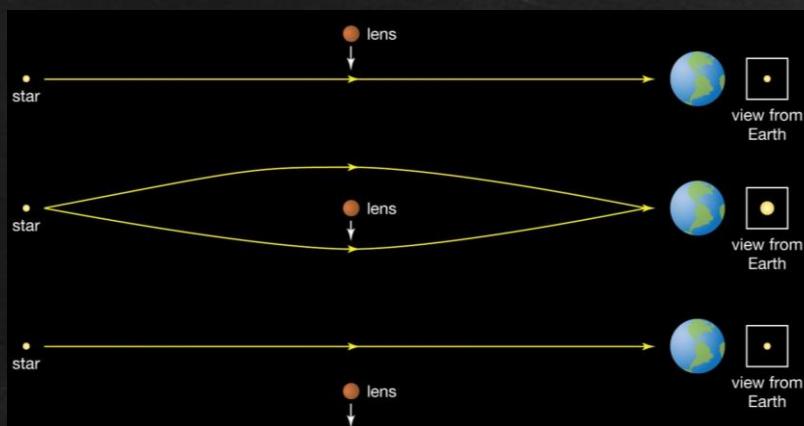
Stochastic fluctuations

Constant amplitude within a de Broglie time



$$t_{\text{dB}} \sim \frac{\lambda_{\text{dB}}}{v}$$

Rich phenomenology

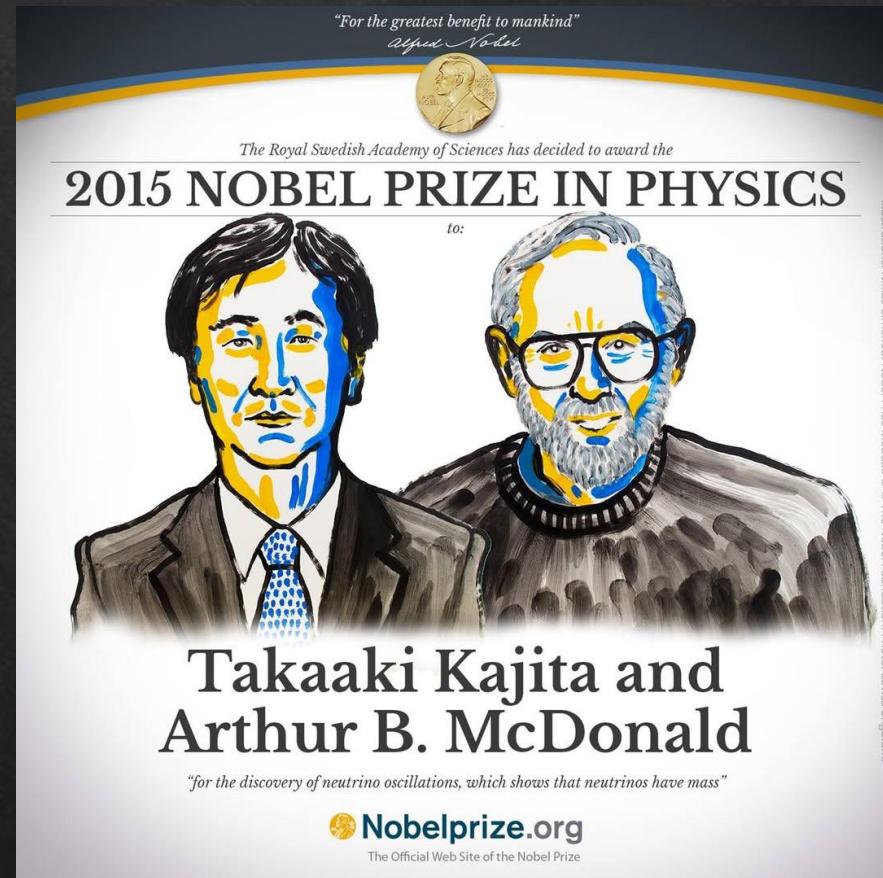


Neutrino mass?

The standard model and neutrino mass

The standard model:

- Minimal lepton sector
- No right-handed neutrino
- Massless neutrinos
- Accidental lepton number symmetry



“for the discovery of neutrino oscillations,
which shows that neutrinos have mass”

The seesaw mechanism

Mass matrix for ν_L, ν_R

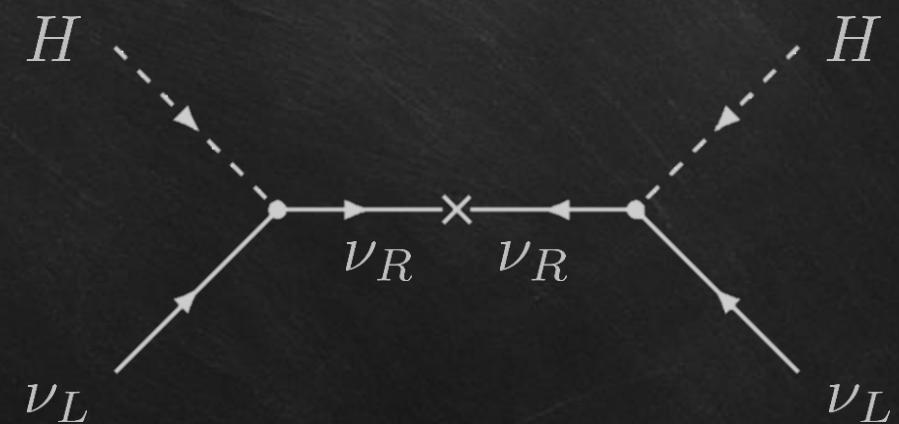
$$\mathcal{M}_\nu = \begin{pmatrix} 0 & m_D \\ m_D & m_R \end{pmatrix}$$

$$\downarrow m_D \ll m_R$$

$$m_l \simeq \frac{m_D^2}{m_R}, \quad m_h \simeq m_R$$

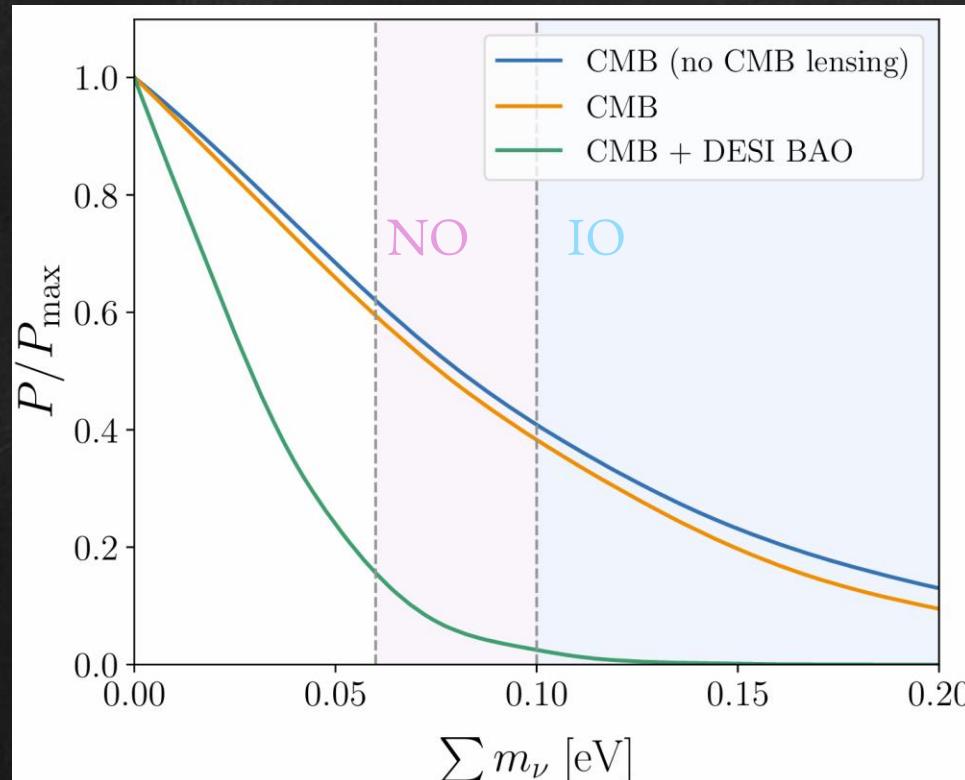
$$\mathcal{L} \supset -m_D \bar{\nu}_D \nu_D - \frac{1}{2} m_R \overline{\nu_M^c} \nu_M + h.c.$$

$$\nu_D = \nu_L + \nu_R, \quad \nu_M = \nu_R + \nu_R^c$$



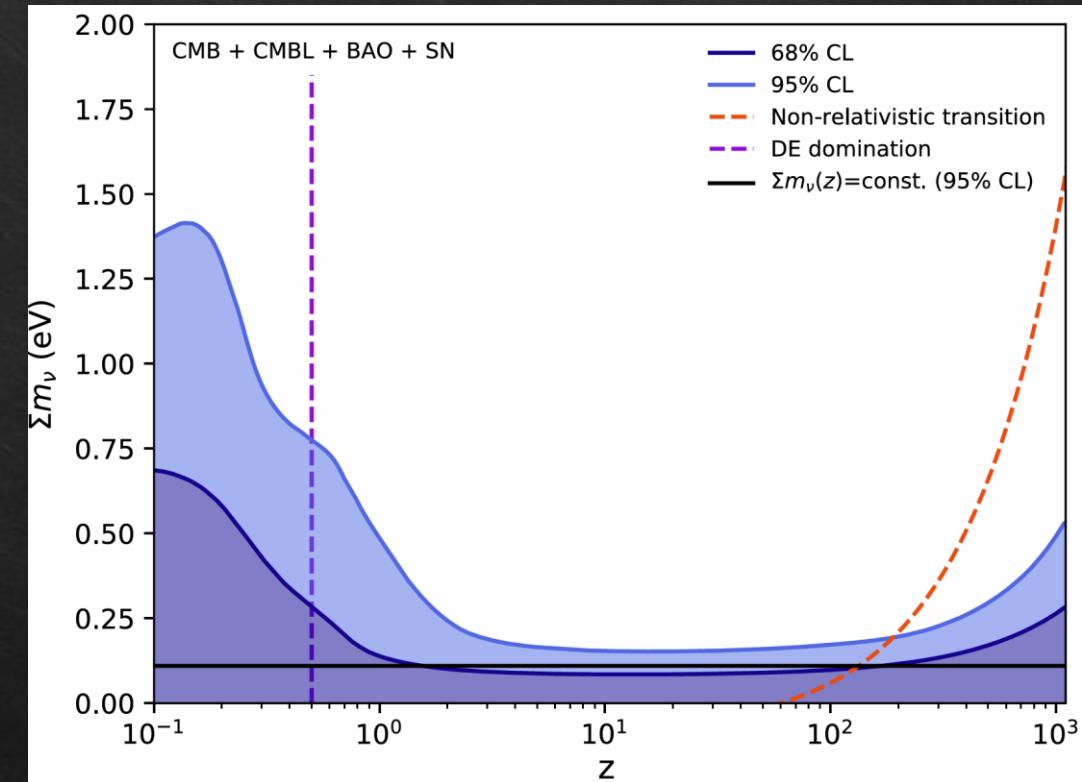
Dynamic neutrino mass? Some motivations

Some tension in $\sum m_\nu$?



DESI Collaboration (JCAP, 2025)

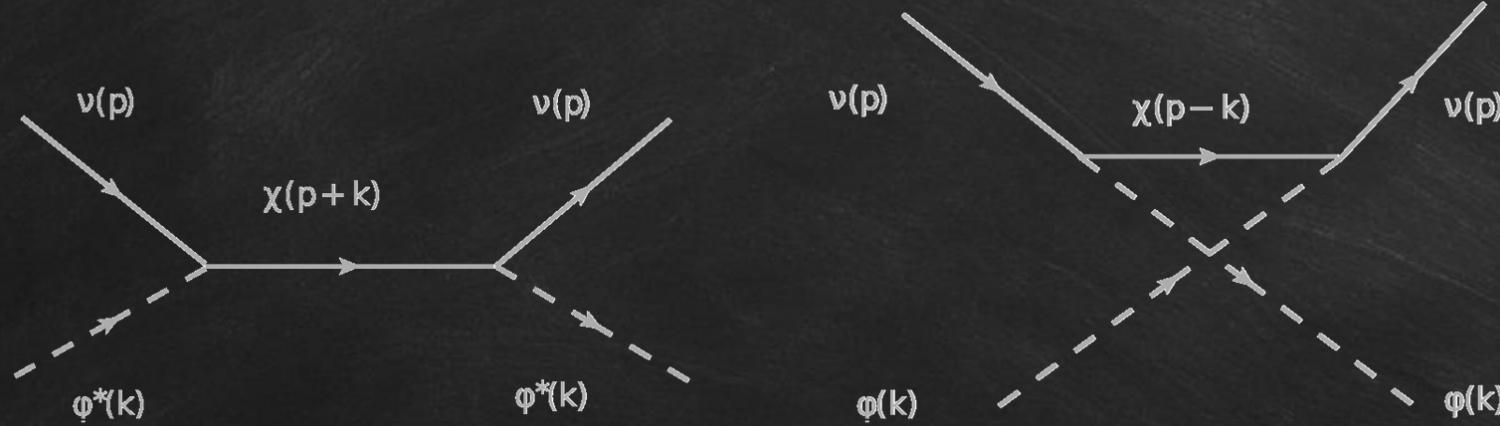
Redshift-dependent $\sum m_\nu$?



C. S. Lorenz et al. (PRD, 2021)

Can DM explain m_ν ?

A specific realization of “dark” neutrino mass



ν : Neutrinos
 χ : Fermionic mediators
 φ : Scalar dark matter
 ν, χ have zero bare mass

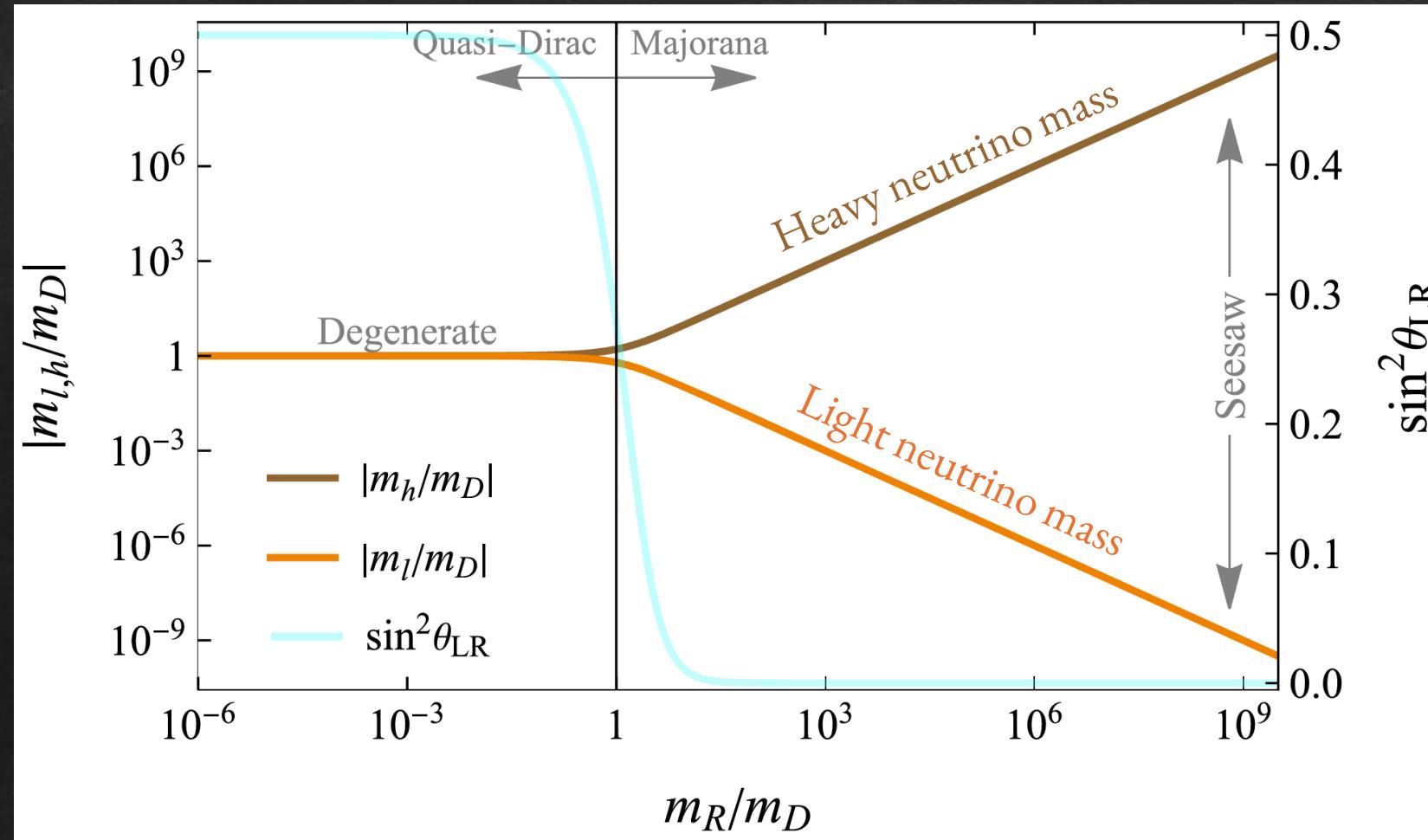
Cold gas of dark matter particles: $m_\nu^2 \propto \frac{\rho_\phi}{m_\phi^2} \frac{y(y-\epsilon)}{y^2 - 1}$, $y = \frac{2E_\nu m_\phi}{m_\chi^2}$, $\epsilon = \frac{n_\phi - \bar{n}_\phi}{n_\phi + \bar{n}_\phi}$
 (Forward scattering)

Classical scalar field background: $m_\nu^2 \propto \frac{\rho_\phi}{m_\phi^2} \cos^2(m_\phi t)$ (Relevant to ultralight dark matter)

Another realization of “dark” neutrino mass



If the Majorana mass is $m_R = g\phi_0 \cos(m_\phi t)$ in the seesaw



$$\begin{pmatrix} \nu_l \\ \nu_h \end{pmatrix} = \begin{pmatrix} \cos \theta_{LR} & \sin \theta_{LR} \\ -\sin \theta_{LR} & \cos \theta_{LR} \end{pmatrix} \begin{pmatrix} \nu_L^c \\ \nu_R \end{pmatrix} + \text{H.c.}$$



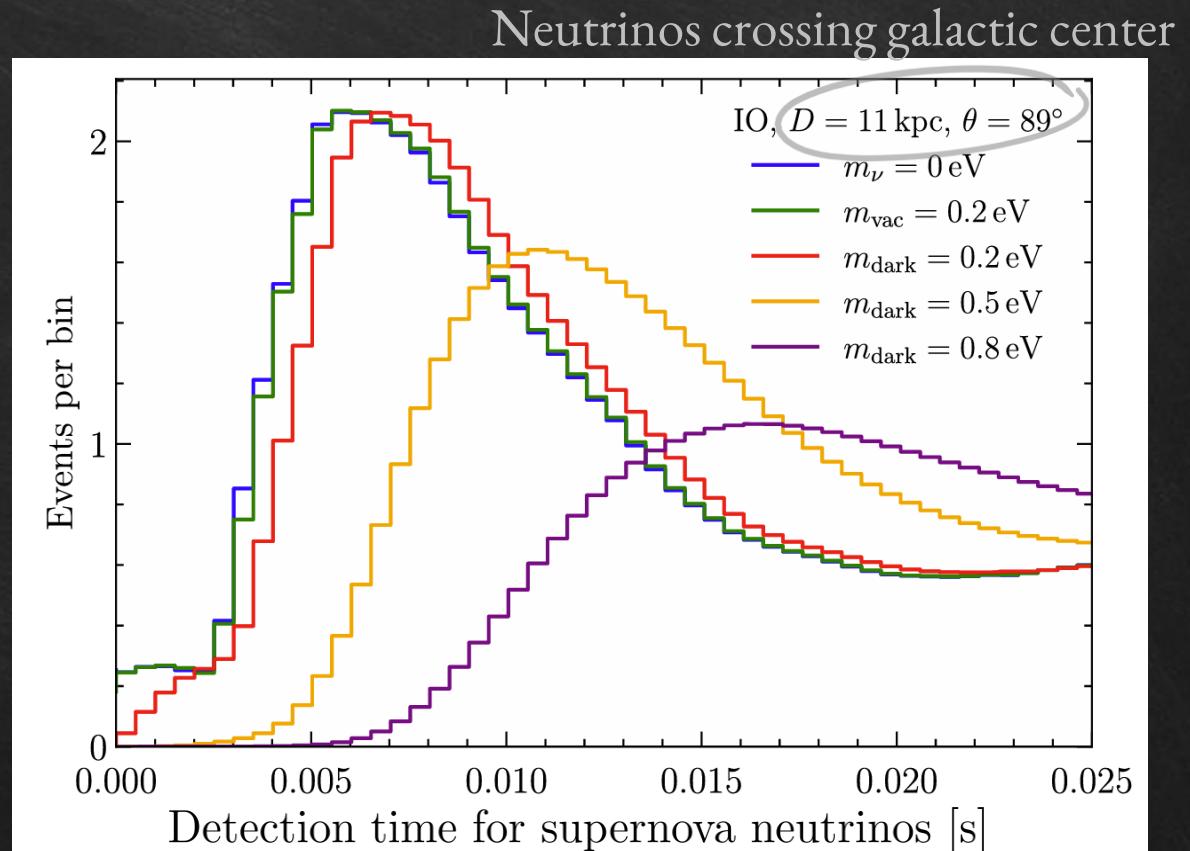
Testing the mass origin with supernova neutrinos

Galactic core-collapse supernovae rate $\sim \mathcal{O}(1)$ /century

S. M. Adams et al. (ApJ, 2013)

Arrival time delay effect is pronounced for:

- Large “dark” mass
 - Supernovae near galactic center
(even lower rate)



S.-F. Ge, C.-F. Kong and A. Y. Smirnov (PRL, 2024)

Assuming NFW profile for dark matter, neglecting time dependence and density fluctuations

Tests with oscillation experiments?



For ultralight DM $10^{-19} \lesssim m_\phi \ll 10\text{eV}$

$$\Delta m_{ij}^2 \sim \Delta m_{ijD}^2(x) \cos^2(m_\phi t)$$

↑ ↑
DM density- Time
dependent modulation

assuming relativistic neutrinos

Several time and length scales



Modulation period

$$T_\phi = \frac{\pi}{m_\phi} = 5.7\text{hr} \left(\frac{10^{-19}\text{eV}}{m_\phi} \right)$$

Time scale for experiments

$$T_{\text{exp}} \gtrsim \mathcal{O}(10)\text{days}$$

→ Time-averaged probabilities

Coherence length

$$\lambda_{\text{dB}} = 1.24\text{au} \left(\frac{10^{-14}\text{eV}}{m_\phi} \right) \left(\frac{200\text{km/s}}{v} \right)$$

Earth crossing distance

$$l_\oplus = 1.16\text{au} \left(\frac{T_{\text{exp}}}{10\text{days}} \right) \left(\frac{v_\oplus}{200\text{km/s}} \right)$$

→ Space-averaged probabilities for $m_\phi \gg 10^{-14}\text{eV}$

Strategy



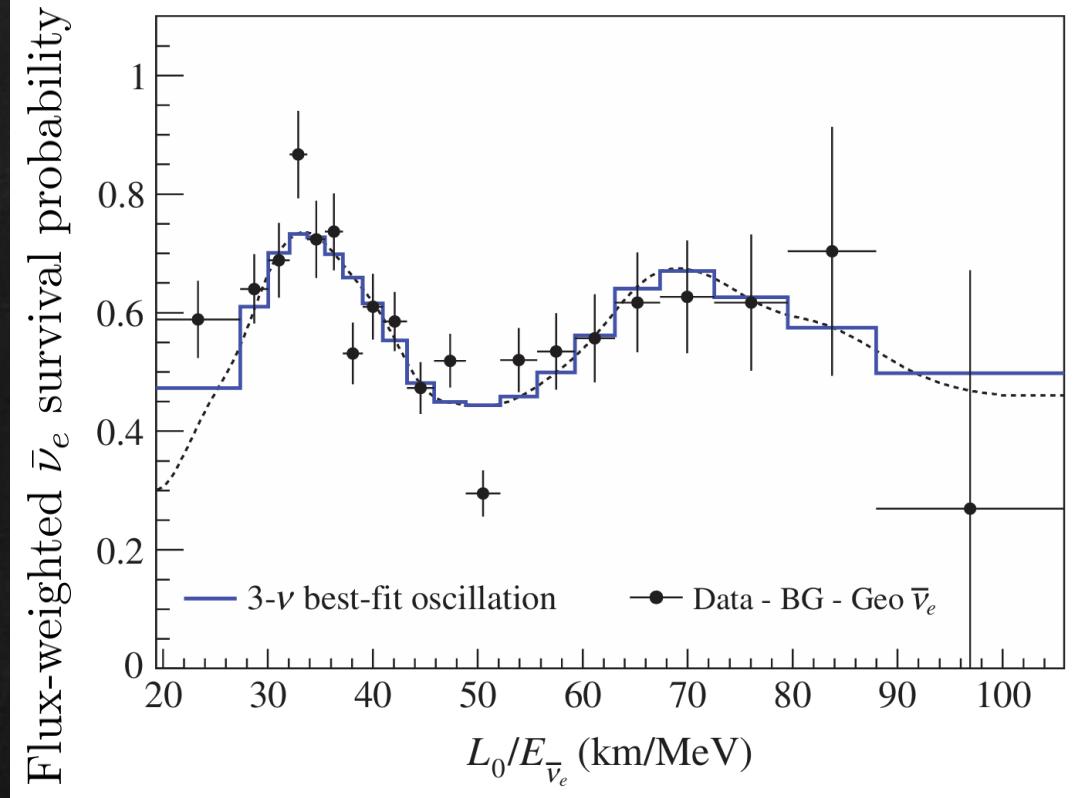
For $m_\phi \ll 10^{-14}\text{eV}$, constant m_ν :

1. Take time average
2. Compare with data

For $m_\phi \gg 10^{-14}\text{eV}$, varying m_ν :

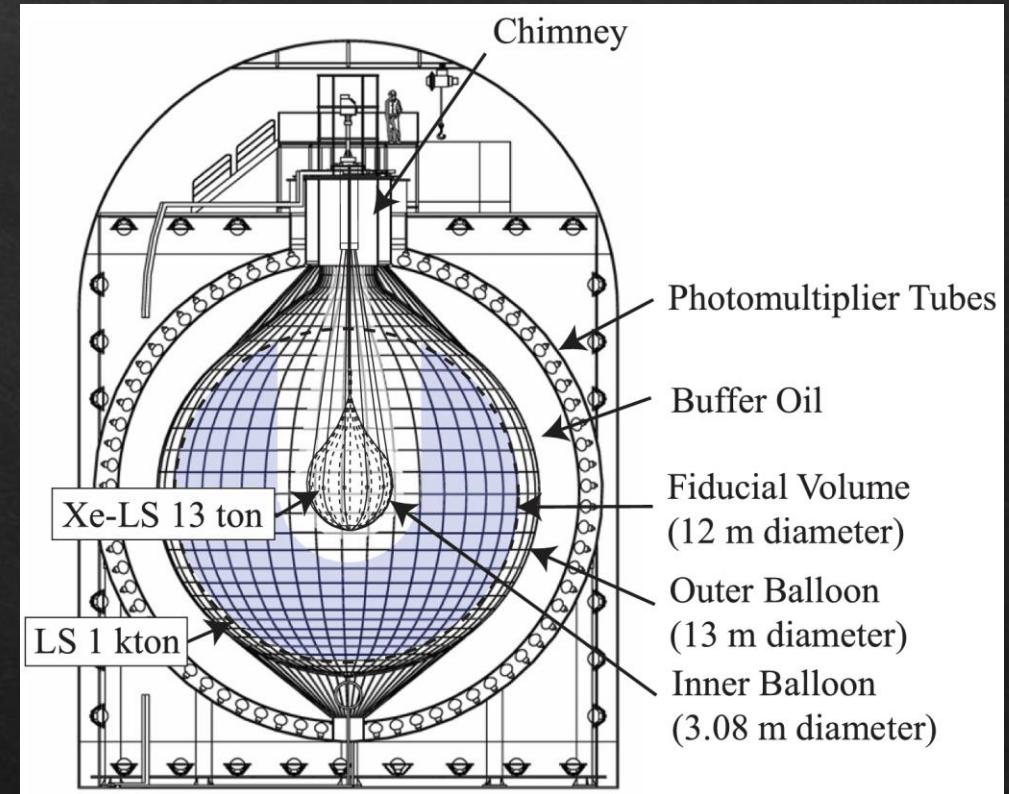
1. Model DM density fluctuations
2. Take time + spatial average
3. Compare with data

Long-baseline reactor experiment: KamLAND



$L_0 = 180\text{km}$ (flux-weighted average baseline)

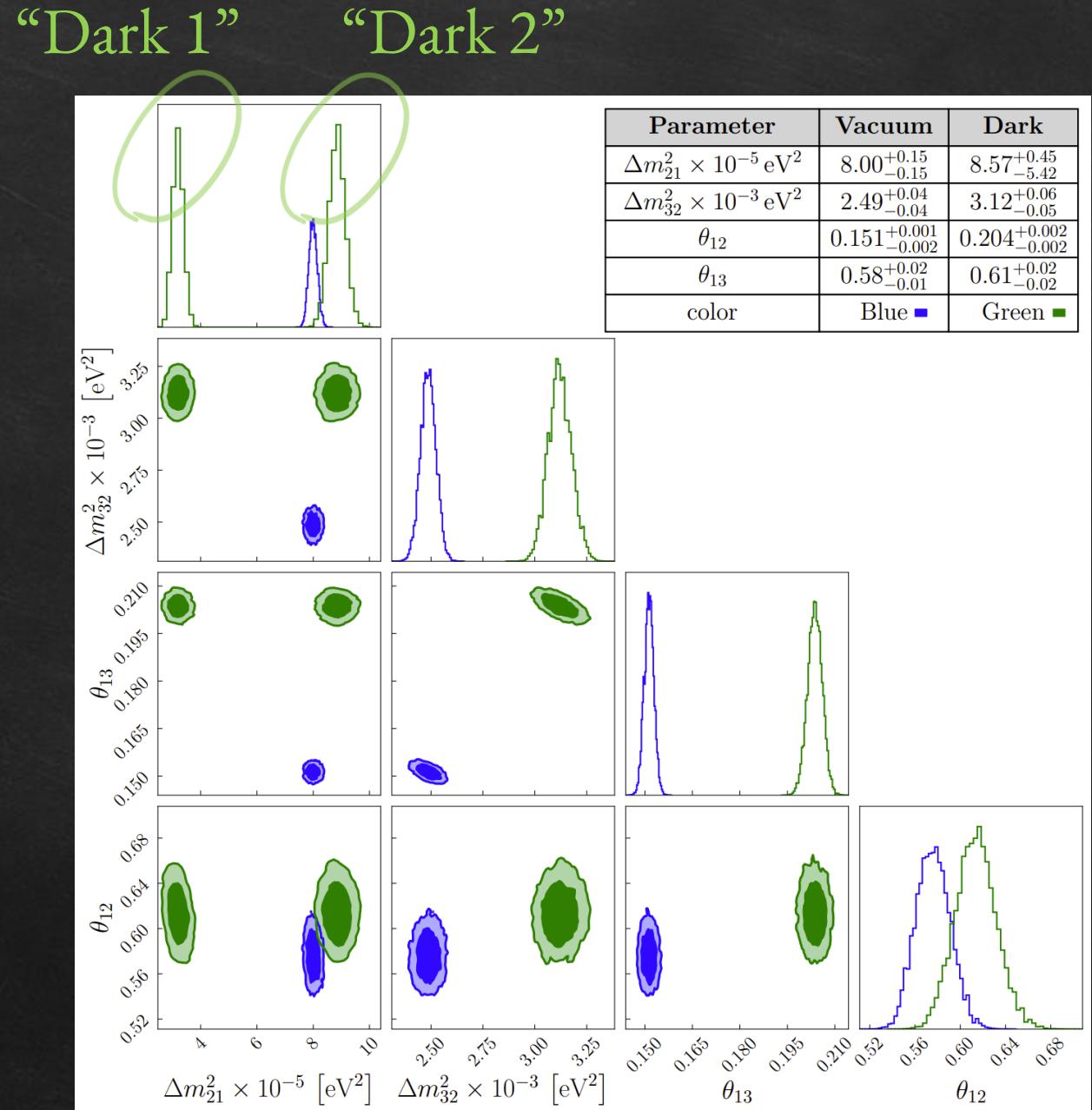
KamLAND Collaboration (2013)



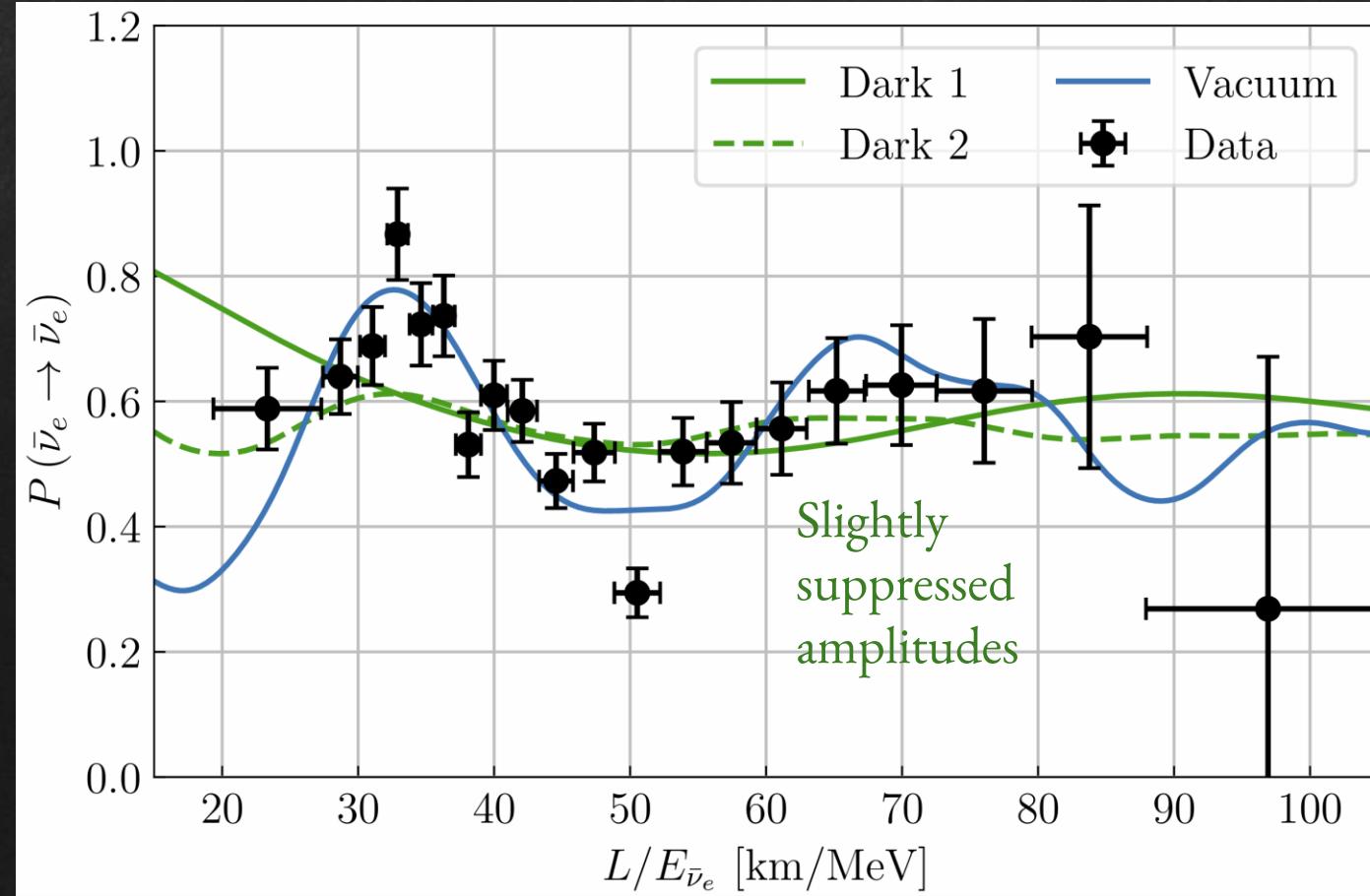
Located at 1km underground, Hida, Japan
Detected antineutrinos from >50 reactors (before 2013)
Sensitive to Δm_{21}^2 , θ_{12} , θ_{13}

Chi-square analysis

KamLAND (main dataset)
+ RENO + Double Chooz
(short baseline experiments)
+ Solar θ_{12}

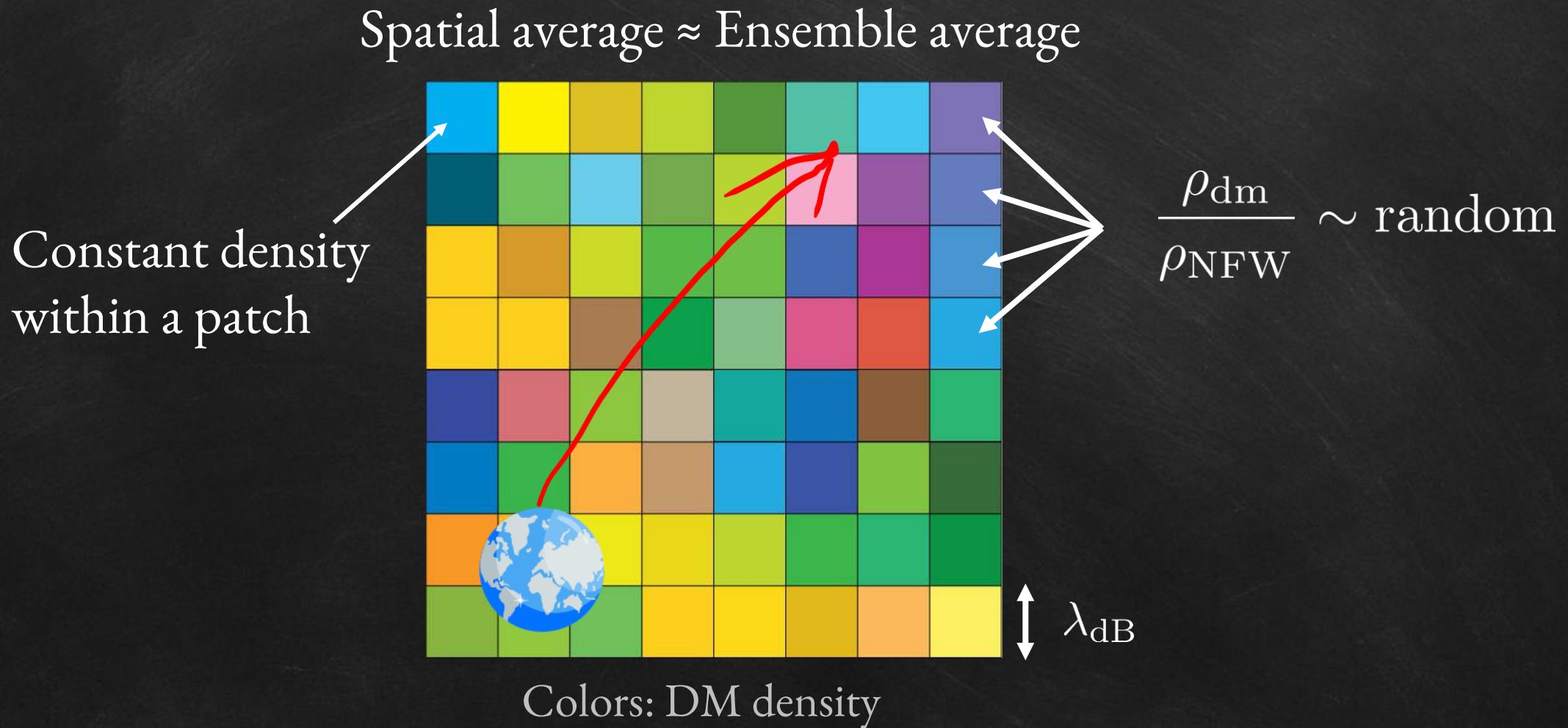


Survival probabilities with best-fit parameters



“Dark” mass is disfavored at $> 4\sigma$.

Earth crossing through different de Broglie patches

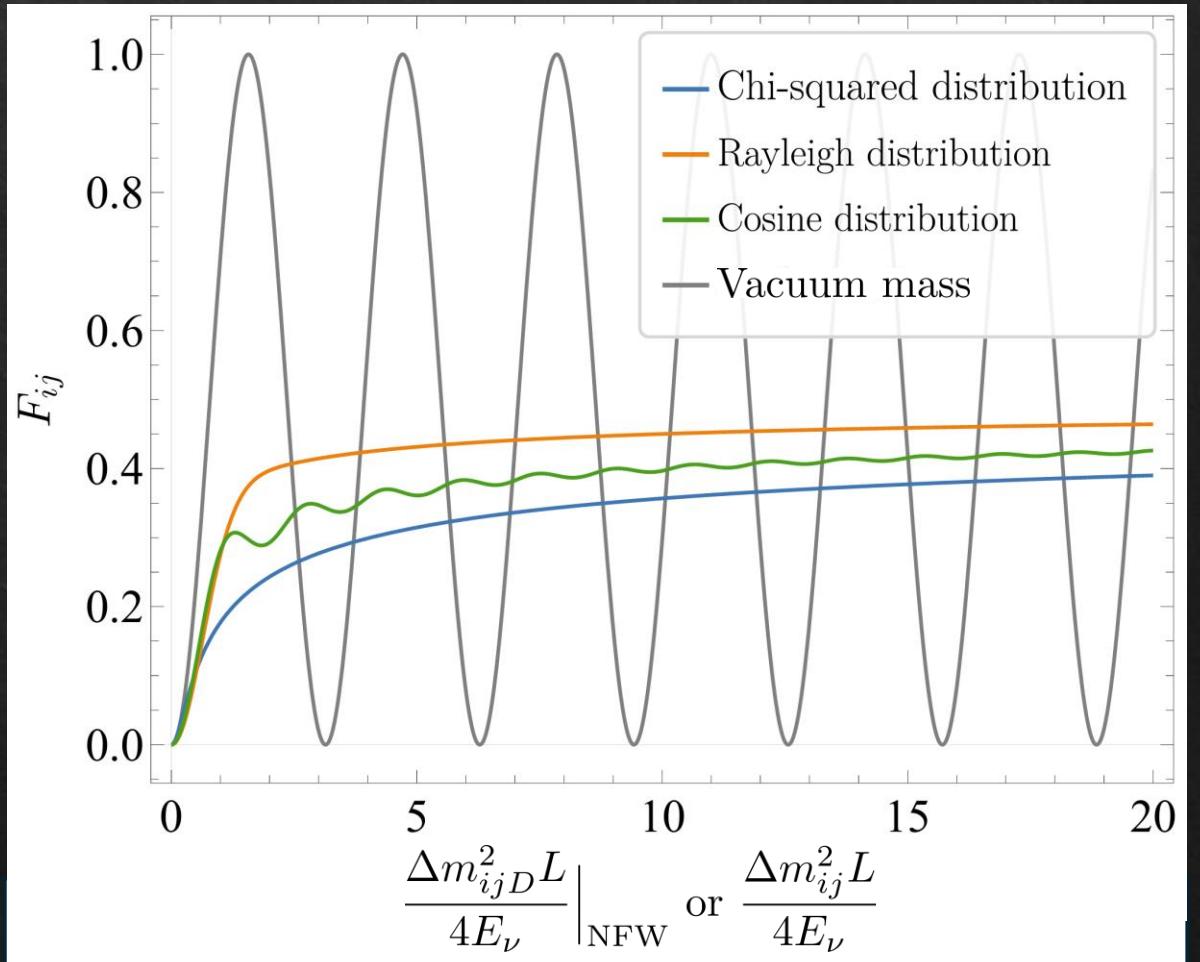


Suppressed oscillation behaviors

Distributions of mass-squared differences in different patches:

$$\frac{\Delta m_{ijD}^2}{\Delta m_{ijD}^2|_{\text{NFW}}} \sim \chi^2(1), \text{ Rayleigh}$$
$$\sim 1 - \cos(k_\theta \theta)$$

Flavor oscillations in terms of distances are suppressed in a model-independent way!



For vacuum mass: $F_{ij} = \sin^2 \frac{\Delta m_{ij}^2 L}{4E_\nu}$

Summary

- Ultralight DM density on small scales
 - Soliton + NFW
 - Stochastic fluctuations
- Dynamic m_ν due to ultralight DM?
 - For $10^{-19} \lesssim m_\phi \ll 10^{-14}$ eV
KamLAND disfavors “dark” mass by $> 4\sigma$
 - For 10^{-14} eV $\ll m_\phi \ll 10$ eV
Stochastic fluctuations suppress neutrino oscillations
- Ultralight DM is unlikely to account for m_ν

