

20)

- a) Let $\frac{dx}{dt}$ be the rate of change of the amount of chemical per unit time. By the input-output principle the accumulation of chemical is $\alpha - \beta$, therefore a differential equation satisfied by x would be

$$\frac{dx}{dt} = \alpha - \beta$$

- b) At steady-state $\frac{dx}{dt} = 0$, so the model predicts $\alpha = \beta$, i.e. the rates of input and decay are equal.

21)

a) Rearranging into the form $\frac{dy}{dx} = g(x)h(y)$

$$\frac{dy}{dx} = \frac{1}{x^2} \times \frac{1}{y}$$

applying the separation of variables method

$$\int y \, dy = \int x^{-2} \, dx$$

$$\frac{1}{2} y^2 = -\frac{1}{x} + C$$

$$y^2 = 2C - \frac{2}{x}$$

$$y = \sqrt{D - \frac{2}{x}}$$

where D is an arbitrary constant

b) Substituting $y = 1$ and $x = 1$ into the solution gives

$$1 = \sqrt{D - 2}$$

$$1 = D - 2$$

$$D = 3$$

So the particular solution satisfying $y(1) = 1$ is

$$y = \sqrt{3 - \frac{2}{x}}$$

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a) In matrix form this system is

$$\begin{pmatrix} \dot{x} \\ \dot{y} \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 1 & -1 \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix}$$

b) The coefficient matrix is triangular, so the eigenvalues are the diagonal elements, 1 and -1. An eigenvector for $\lambda = 1$ is a solution of

$$\begin{pmatrix} 1-1 & 0 \\ 1 & -1-1 \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \underline{\underline{0}}$$

which gives

$$\begin{aligned} x - 2y &= 0 \\ x &= 2y \end{aligned}$$

so $(2 \ 1)^T$ is an eigenvector for $\lambda = 1$

An eigenvector for $\lambda = -1$ is a solution of

$$\begin{pmatrix} 1-(-1) & 0 \\ 1 & -1-(-1) \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \underline{\underline{0}}$$

which gives

$$\begin{aligned} 2x &= 0 \\ x &= 0 \end{aligned}$$

So $(0 \ 1)^T$ is an eigenvector for $\lambda = -1$

c) The general solution is therefore

$$\begin{pmatrix} \dot{x} \\ \dot{y} \end{pmatrix} = \alpha \begin{pmatrix} 2 \\ 1 \end{pmatrix} e^t + \beta \begin{pmatrix} 0 \\ 1 \end{pmatrix} e^{-t}$$

where α and β are arbitrary constants.

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a) $X(x)$ must satisfy the boundary conditions

$$X(0) = 0 \quad \text{and} \quad \frac{\partial}{\partial x} X(\pi) = 0$$

b) If K is negative then the solution takes the form

$$X(x) = A \cos(\sqrt{-K} x) + B \sin(\sqrt{-K} x)$$

Using the boundary condition $X(0) = 0$:
 $0 = A$

Therefore

$$X(x) = B \sin(\sqrt{-K} x)$$

And

$$\frac{\partial}{\partial x} X(x) = \sqrt{-K} B \cos(\sqrt{-K} x)$$

Using the second boundary condition

$$0 = \sqrt{-K} B \cos(\sqrt{-K} \pi)$$

As $K \neq 0$ and $B \neq 0$, $\cos(\sqrt{-K} \pi)$ must equal zero for a non-trivial solution.

$$\cos(\sqrt{-K} x) = 0$$

$$\Rightarrow \sqrt{-K} x = \frac{1}{2} \pi (2n-1)$$

$$\sqrt{-K} = \frac{1}{2x} \pi (2n-1)$$

24

The volume integral in cylindrical coordinates is

$$\int_D f dV = \int_D f(\rho, \phi, z) \rho dz d\phi d\rho$$

In this example, $f(\rho, \phi, z) = \cos \phi$. The limits of integration are

$$z = 0, \quad z = h + x$$

$$\phi = -\pi, \quad \phi = \pi$$

$$\rho = 0, \quad \rho = a$$

Hence the volume integral is

$$\int_{\rho=0}^a \left(\int_{z=0}^{h+x} \left(\int_{\phi=-\pi}^{\pi} \rho \cos \phi d\phi \right) dx \right) d\rho$$

$$= \int_{\rho=0}^a \left(\int_{z=0}^{h+x} \right)$$

$$\left[\rho \sin \right]$$



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a) The linear momentum of the two particles before collision is

$$\underline{p}_1 = m_1 \underline{\dot{r}}_1 = m u \underline{i}$$

$$\underline{p}_2 = m_2 \underline{\dot{r}}_2 = m u \underline{j}$$

After collision the particles have mass $(m_1 + m_2)$ and velocity \underline{w} . The principle of conservation of linear momentum tells us that

$$m u \underline{i} + m u \underline{j} = 2m \underline{w}$$

$$\therefore \underline{u} \underline{i} + u \underline{j} = 2 \underline{w}$$

$$\underline{w} = \frac{u}{2} \underline{i} + \frac{u}{2} \underline{j}$$

The speed after collision is

$$|\underline{w}| = \sqrt{\left(\frac{u}{2}\right)^2 + \left(\frac{u}{2}\right)^2}$$

$$= \sqrt{\frac{u^2}{2}}$$

b) The kinetic energy before collision is

$$E_1 = \frac{1}{2} m u^2 + \frac{1}{2} m u^2$$

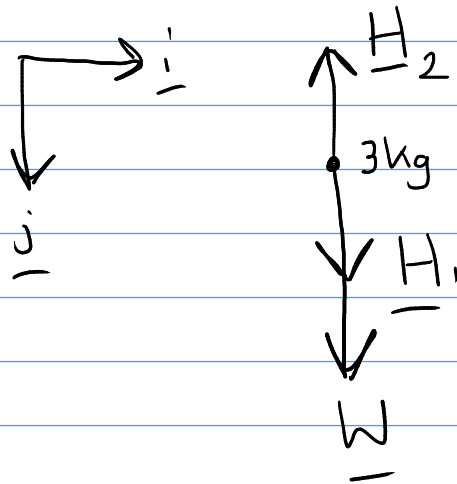
The kinetic energy after collision is

$$E_2 = \frac{1}{2} (2m) \sqrt{\frac{u^2}{2}}$$

$$= m u \sqrt{\frac{1}{2}}$$

The ratio of kinetic energy is $\frac{m u \sqrt{\frac{1}{2}}}{2 m u} = \frac{\sqrt{2}}{4}$

a) 26



where \underline{W} is the weight of P, \underline{H}_1 is the spring force of BP and \underline{H}_2 is the spring force of AP.

b) $\underline{W} = 3g \underline{j}$, where g is acceleration due to gravity. \underline{H}_1 and \underline{H}_2 are found in the table below.

Spring	L	L_0	$L - L_0$	K	\underline{j}	\underline{H}
\underline{H}_1	$\frac{3}{2} - x$	2	$-\frac{1}{2} - x$	30	\underline{j}	$30(-\frac{1}{2} - x)\underline{j}$
\underline{H}_2	x	1	$x - 1$	18	$-\underline{j}$	$-18(x - 1)\underline{j}$

Therefore

$$\underline{H}_1 = 30(-\frac{1}{2} - x)\underline{j}$$

$$\underline{H}_2 = -18(x - 1)\underline{j}$$

c) Applying Newton's second law gives

$$3\ddot{x} = 3g\mathbf{j} + 30(-\frac{1}{2} - x)\mathbf{j} - 18(x-1)\mathbf{j}$$

If the system is in equilibrium the $\ddot{x} = 0$

$$0 = 3g\mathbf{j} + 30(-\frac{1}{2} - x)\mathbf{j} - 18(x-1)\mathbf{j}$$

Resolving in the \mathbf{j} direction and rearranging gives

$$0 = 3g - 15 - 30x - 18x + 18$$
$$48x = 3g + 3$$
$$x = \frac{3g + 3}{48}$$

$$= \frac{g+1}{16}$$

d) When the particle is in motion we have

$$3\ddot{x} = 3g\mathbf{j} + 30(-\frac{1}{2} - x)\mathbf{j} - 18(x-1)\mathbf{j}$$

Resolving in the \mathbf{j} direction and rearranging

$$3\ddot{x} = 3g - 15 - 30x - 18x + 18$$

$$3\ddot{x} + 48x = 3g + 3$$

$$\ddot{x} + 16x = g + 1$$

e) Let $\omega^2 = 16$ then the equation of motion can be rewritten as

$$\ddot{x} + \omega^2 x = \omega^2 x_{eq}$$

$$\text{with } x_{eq} = \frac{g+1}{16}$$

The general solution is

$$\begin{aligned} x(t) &= B \cos(\omega t) + C \sin(\omega t) + x_{eq} \\ &= B \cos(4t) + C \sin(4t) + \frac{g+1}{16} \end{aligned}$$

where B and C are constants. At $t=0$ we have $x(0) = \frac{g+1}{16} - \frac{1}{4}$

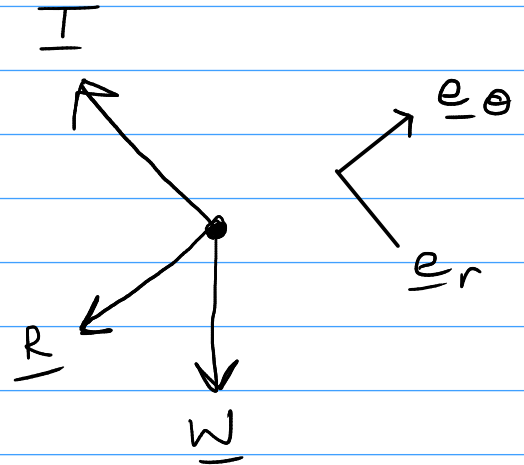
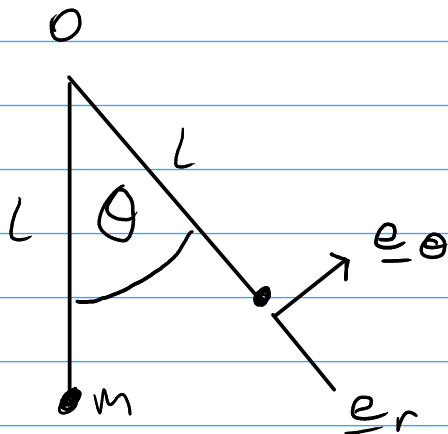
$$\frac{g+1}{16} - \frac{1}{4} = B \cos(0) + C \sin(0) + \frac{g+1}{16}$$

$$-\frac{1}{4} = B$$



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a)



b) The tension in the rod is

$$\underline{T} = -|T|\underline{e}_r$$

The resistive force is

$$\underline{R} = -kv\underline{e}_\theta$$

The weight of the bob is

$$\underline{W} = mg\cos\theta\underline{e}_r - mg\sin\theta\underline{e}_\theta$$

The position of the bob in polar coordinates is

$$\underline{r} = L\underline{e}_r$$

c) The total torque about O is the sum of the torques of all forces acting on the bob. As the lines of action of \underline{T} and \underline{R} pass through O , they contribute no torque. The torque due to \underline{W} is

$$\begin{aligned}\tau &= l \underline{e}_r \times (mg \cos \theta \underline{e}_r - mg \sin \theta \underline{e}_\theta) \\ &= mgl \sin \theta \underline{u}\end{aligned}$$