

Homework# 4

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P87 Problem12

(a)

$$\left. \begin{aligned} \frac{\partial u_\lambda}{\partial t} &= \lambda^2 \frac{\partial u}{\partial t}(\lambda x, \lambda^2 t) \\ \frac{\partial^2 u_\lambda}{\partial x_i^2} &= \lambda^2 \frac{\partial^2 u}{\partial x_i^2}(\lambda x, \lambda^2 t) \end{aligned} \right\} \Rightarrow \frac{\partial u_\lambda}{\partial t} - \Delta u_\lambda = \lambda^2 \left(\frac{\partial u}{\partial t} - \Delta u \right) = 0. \quad (1)$$

It means that u_λ solves the heat equation.

(b)

$$\frac{du_\lambda}{d\lambda} = x \cdot Du(\lambda x, \lambda^2 t) + 2\lambda t u_t(x, t) \xrightarrow{\lambda \rightarrow 1} v(x, t). \quad (2)$$

As $\forall \lambda$, u_λ satisfies the heat equation, we can see $\frac{u_\lambda + \Delta \lambda - u_\lambda}{\Delta \lambda}$ satisfies heat equation for all $\Delta \lambda$. It means that v satisfies the heat equation.

P87 Problem13

(a) We can see:

$$\begin{aligned} \frac{\partial u}{\partial t} &= v'(z) \frac{\partial(\frac{x}{\sqrt{t}})}{\partial t} = -\frac{1}{2} x t^{-\frac{3}{2}} v'(z) \\ \frac{\partial u}{\partial x} &= v'(z) \frac{1}{\sqrt{t}} \\ \frac{\partial^2 u}{\partial x^2} &= v''(z) \frac{1}{t} \end{aligned} \quad (3)$$

So:

$$\begin{aligned} u_{xx} &= u_t \\ \Leftrightarrow -\frac{1}{2} v'(z) x t^{-\frac{3}{2}} &= v''(z) t^{-1} \\ \Leftrightarrow v''(z) t^{-1} + \frac{1}{2} v'(z) x t^{-\frac{3}{2}} &= 0 \\ \Leftrightarrow v''(z) + \frac{z}{2} v'(z) &= 0 \end{aligned} \quad (4)$$

To get the general solution for 4, we can see:

$$\begin{aligned}
 v'' + \frac{z}{2}v' &= 0 \Rightarrow \frac{dv'}{dz} + \frac{z}{2}v' = 0 \\
 &\Rightarrow \frac{dv'}{v'} + \frac{zdz}{2} = 0 \\
 &\Rightarrow d \log(v') + d \frac{z^2}{4} = 0 \\
 &\Rightarrow \frac{z^2}{4} + \log(v') \equiv C \\
 &\Rightarrow v' = e^{C - \frac{z^2}{4}} \\
 &\Rightarrow v(z) = c \int_0^z e^{-\frac{s^2}{4}} ds + d.
 \end{aligned} \tag{5}$$

(b) As (a) suggests, for $t > 0$, we can see:

$$u(x, t) = v(z) = c \int_0^{\frac{x}{\sqrt{t}}} e^{-\frac{s^2}{4}} ds + d = c \int_0^x e^{-\frac{w^2}{4t}} \frac{1}{\sqrt{t}} dw + d. \tag{6}$$

So:

$$\frac{\partial u}{\partial x} = \frac{c}{\sqrt{t}} e^{-\frac{x^2}{4t}}. \tag{7}$$

Just set $c = \frac{1}{\sqrt{4\pi}}$, we get the fundamental solution Φ .

Then I should explain the reason of this phenomenon. By the definition $u(x, t) = v(\frac{x}{\sqrt{t}})$, we can see:

$$u(x, 0) = \begin{cases} \lim_{z \rightarrow +\infty} v(z) = \sqrt{\pi}c + d, x > 0 \\ v(0) = d, x = 0 \\ \lim_{z \rightarrow -\infty} v(z) = d - \sqrt{\pi}c, x < 0 \end{cases} \tag{8}$$

So: $\frac{\partial u(x, 0)}{\partial x} = \delta_0(x)$. It shows that this solution is the fundamental solution.

P87 Problem14

Define $v(x, t) = u(x, t)e^{ct}$, then:

$$\begin{cases} v_t - \Delta v = e^{ct} f \text{ in } \mathbb{R}^n \times (0, \infty) \\ v = g \text{ on } \mathbb{R}^n \times \{t = 0\} \end{cases} \tag{9}$$

So:

$$v(x, t) = \int_{\mathbb{R}^n} \Phi(x - y, t) g(y) dy + \int_0^t \int_{\mathbb{R}^n} \Phi(x - y, t - s) f(y, s) e^{cs} dy ds. \tag{10}$$

And $u(x, t) = e^{-ct} v(x, t)$.

P87 Problem15

Set $v(x, t) = u(x, t) - g(t)$, then we can see:

$$\begin{cases} v_t - v_{xx} = -g'(t) \text{ in } \mathbb{R}_+ \times (0, \infty) \\ v(x, 0) = 0 \\ v(0, t) = 0 \end{cases} \quad (11)$$

Then extend v to $x < 0$: just set $v(x, t) = -v(-x, t)$ on $x < 0$, as the equation satisfies the compatible condition, this extension is okay. We can see v satisfies:

$$\begin{cases} v_t - v_{xx} = f(x, t) \\ v(x, 0) = 0 \end{cases} \quad (12)$$

where $f(x, t) = -g'(t)$ if $x \geq 0$, and $f(x, t) = g'(t)$ if $x < 0$. Solve this equation, we can see

$$u(x, t) = v(x, t) + g(t) = \frac{x}{\sqrt{4\pi}} \int_0^t \frac{1}{(t-s)^{\frac{3}{2}}} e^{-\frac{x^2}{4(t-s)}} g(s) ds. \quad (13)$$

P89 Problem22

$$\begin{cases} u_t + u_x = d(v - u) \\ v_t - v_x = d(u - v) \end{cases} \Rightarrow \begin{cases} u_{tt} + u_{tx} = d(v_t - u_t), u_{tx} + u_{xx} = d(v_x - u_x) \\ v_{tt} - v_{tx} = d(u_t - v_t), v_{tx} - v_{xx} = d(u_x - v_x) \end{cases}$$

$$\Rightarrow u_{tt} - u_{xx} = d(v_t - u_t - v_x + u_x) \quad (14)$$

What's more, $v_t - v_x = -(u_t + u_x)$, so $u_{tt} - u_{xx} = -2du_t$. In the same way, $v_{tt} + 2dv_t - v_{xx} = 0$.

P90 Problem24

(a)

$$\begin{aligned} \frac{dk+p}{dt} &= \int_{-\infty}^{+\infty} (u_t u_{tt} + u_x u_{xt}) dx \\ &= \int_{-\infty}^{+\infty} (u_t u_{xx} + u_x u_{xt}) dx \\ &= \int_{-\infty}^{+\infty} \frac{\partial u_t u_x}{\partial x} dx \\ &= 0. \end{aligned} \quad (15)$$

The final equation is true for g and h both have compact support.

(b) By D.Almbert's formula, we can see:

$$u(x, t) = \frac{g(x+t) + g(x-t)}{2} + \frac{1}{2} \int_{x-t}^{x+t} h(y) dy. \quad (16)$$

It means that:

$$\begin{aligned} u_x &= \frac{g'(x+t) + g'(x-t)}{2} + \frac{1}{2}(h(x+t) - h(x-t)). \\ u_t &= \frac{g'(x+t) - g'(x-t)}{2} + \frac{1}{2}(h(x+t) + h(x-t)). \end{aligned} \quad (17)$$

As g, h both have compact support, we can choose large t such that $u_x^2 = u_t^2 \forall x$, which means $k(t) = p(t)$.