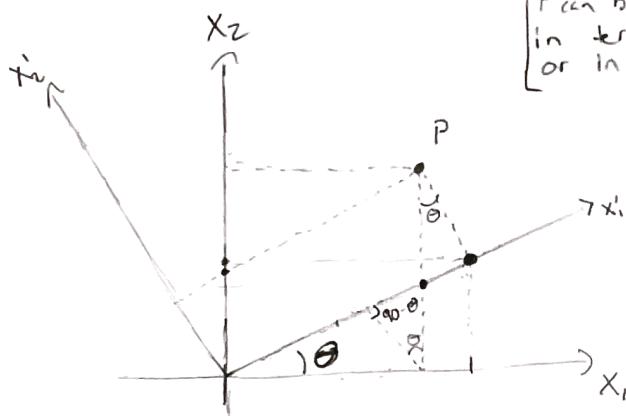


Coordinate Transformations

Class notes



P can be written
in terms of (X_1, X_2)
or in terms of (X'_1, X'_2)

B2

General
Definitions

$$\begin{aligned} X'_1 &= \alpha X_1 + \beta X_2 \rightarrow \vec{X}'_1 = \alpha \hat{X}_1 + \beta \hat{X}_2 \\ X'_2 &= -\alpha X_1 + \beta X_2 \end{aligned}$$

[negative sign prob not needed]

(answer)

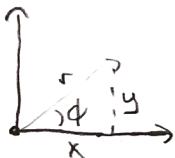
$$\begin{bmatrix} X_1 \\ X_2 \end{bmatrix} = \begin{bmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{bmatrix} \begin{bmatrix} X'_1 \\ X'_2 \end{bmatrix} = \begin{bmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{bmatrix} \begin{bmatrix} X'_1 \\ X'_2 \end{bmatrix}$$

also

$$\begin{aligned} \vec{P} &= a \cdot \hat{X}_1 + b \hat{X}_2 = \begin{bmatrix} a \\ b \end{bmatrix} \\ \vec{P}' &= c \cdot \hat{X}'_1 + d \hat{X}'_2 = \begin{bmatrix} c \\ d \end{bmatrix} \end{aligned}$$

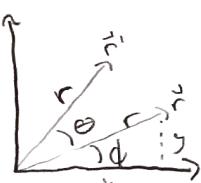
eh

Deriving the rotation matrix



$$\vec{r} = r \cos(\phi) \hat{x} + r \sin(\phi) \hat{y} = \begin{bmatrix} r \cos(\phi) \\ r \sin(\phi) \end{bmatrix}$$

Now rotate it up



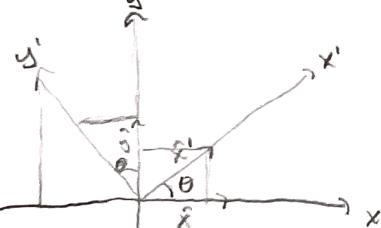
$$\vec{r}' = \begin{bmatrix} x' \\ y' \end{bmatrix} = \begin{bmatrix} r \cos(\phi + \theta) \\ r \sin(\phi + \theta) \end{bmatrix} = \begin{bmatrix} r \cos(\phi) \cos \theta - r \sin(\phi) \sin \theta \\ r \sin(\phi) \cos \theta + r \cos(\phi) \sin \theta \end{bmatrix} = \begin{bmatrix} x \cos \theta - y \sin \theta \\ x \sin \theta + y \cos \theta \end{bmatrix}$$

for me: $\vec{r}' = (x \cos \theta - y \sin \theta) \hat{x} + (x \sin \theta + y \cos \theta) \hat{y}$

$$\vec{r}' = \begin{bmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{bmatrix} \begin{bmatrix} x \\ y \end{bmatrix}$$

This
makes
sense
to
me!

Looking at where \hat{x} and \hat{y} go.



$$e_1 = \begin{bmatrix} 1 \\ 0 \end{bmatrix} \rightarrow e_1 = \begin{bmatrix} \cos \theta & \hat{x} \\ \sin \theta & \hat{y} \end{bmatrix}$$

$$e_2 = \begin{bmatrix} 0 \\ 1 \end{bmatrix} \rightarrow e_2 = \begin{bmatrix} -\sin \theta & \hat{x} \\ \cos \theta & \hat{y} \end{bmatrix}$$

R

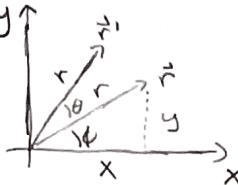
when $R = x \hat{e}_1 + y \hat{e}_2$

$$R = \begin{bmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{bmatrix}$$

$$\begin{bmatrix} x' \\ y' \end{bmatrix} = \begin{bmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{bmatrix} \begin{bmatrix} x \\ y \end{bmatrix}$$

Coordinate Transformations

Rotating a coordinate system



$$\vec{r} = r \cos(\phi) \hat{x} + r \sin(\phi) \hat{y} = \begin{bmatrix} \cos\phi \\ \sin\phi \end{bmatrix} \begin{bmatrix} x \\ y \end{bmatrix}$$

now rotate it *

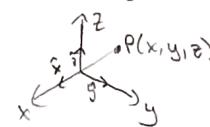
$$\vec{r}' = \begin{bmatrix} x' \\ y' \end{bmatrix} = \begin{bmatrix} r \cos(\phi + \theta) \\ r \sin(\phi + \theta) \end{bmatrix} = \begin{bmatrix} r \cos\phi \cos\theta - r \sin\phi \sin\theta \\ r \sin\phi \cos\theta + r \cos\phi \sin\theta \end{bmatrix} = \begin{bmatrix} x \cos\theta - y \sin\theta \\ x \sin\theta + y \cos\theta \end{bmatrix} = \begin{bmatrix} \cos\theta & -\sin\theta \\ \sin\theta & \cos\theta \end{bmatrix} \begin{bmatrix} x \\ y \end{bmatrix}$$

$$\vec{r}' = \begin{bmatrix} \cos\theta & -\sin\theta \\ \sin\theta & \cos\theta \end{bmatrix} \begin{bmatrix} x \\ y \end{bmatrix}$$

Cartesian

$$\vec{r} = x \hat{e}_1 + y \hat{e}_2 + z \hat{e}_3$$

$$|\vec{r}| = \sqrt{x^2 + y^2 + z^2}$$



Polar/Cylindrical

$$\vec{r} = p \hat{p} + z \hat{z}$$

$$y \uparrow$$

$$\hat{p}$$

$$P$$

$$\phi$$

$$x$$

$$z \text{ out of page } \hat{z}$$

$$y$$

$$P(p, \phi, z)$$

$$x = \vec{r} \cdot \hat{e}_1$$

$$y = \vec{r} \cdot \hat{e}_2$$

$$z = \vec{r} \cdot \hat{e}_3$$

$$p = \sqrt{x^2 + y^2}$$

$$\phi = \tan^{-1}(y/x)$$

Unit vectors $\hat{p}, \hat{\phi}$ in cartesian

$$\hat{p} = ? \rightarrow \vec{r} = x \hat{e}_1 + y \hat{e}_2 + z \hat{e}_3$$

$$x = p \cos\phi \rightarrow \vec{r} = p \cos\phi \hat{e}_1 + p \sin\phi \hat{e}_2 + z \hat{e}_3$$

$$y = p \sin\phi \rightarrow \frac{1}{p \hat{p}}$$

$$\hat{p} = (\cos\phi \hat{e}_1 + \sin\phi \hat{e}_2)$$

Polar

$$1 = \beta^2 \tan^2\phi + \beta^2 \rightarrow 1 = \beta^2 (\tan^2\phi + 1) \rightarrow \beta = \frac{1}{\tan\phi}$$

$$\alpha = -\beta \tan\phi \rightarrow \beta = \pm \cos\phi \quad (\text{take plus so CCW is positive})$$

$$\alpha = -\sin\phi$$

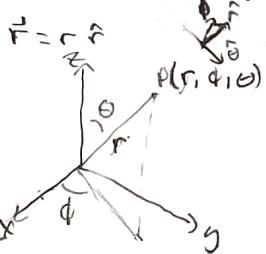
$$\text{THUS! } \hat{\phi} = -\sin\phi \hat{e}_1 + \cos\phi \hat{e}_2$$

Then:

$$\begin{bmatrix} \hat{p} \\ \hat{\phi} \\ \hat{z} \end{bmatrix} = \begin{bmatrix} \cos\phi \sin\theta & 0 \\ -\sin\phi \cos\theta & 0 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} \hat{e}_1 \\ \hat{e}_2 \\ \hat{e}_3 \end{bmatrix}$$

$$\leftarrow \text{from } \begin{array}{l} \hat{p} \rightarrow \begin{bmatrix} \cos\phi \\ -\sin\phi \\ 0 \end{bmatrix} \hat{e}_1 + \begin{bmatrix} \sin\phi \\ \cos\phi \\ 0 \end{bmatrix} \hat{e}_2 + \begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix} \hat{e}_3 \\ \hat{z} \rightarrow \begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix} \end{array}$$

Cartesian to Spherical



From geometry

$$x = r \sin\theta \cos\phi$$

$$y = r \sin\theta \sin\phi$$

$$z = r \cos\theta$$

$$r = \sqrt{x^2 + y^2 + z^2}$$

$$\cos\theta = z/r$$

$$\tan\phi = y/x$$

Looking at Cartesian \vec{r}

$$\vec{r} = x \hat{e}_1 + y \hat{e}_2 + z \hat{e}_3 \quad \text{plug in } x, y, z$$

$$\vec{r} = r \sin\theta \cos\phi \hat{e}_1 + r \sin\theta \sin\phi \hat{e}_2 + r \cos\theta \hat{e}_3$$

$$\textcircled{1} \text{ pull out } r \text{ and then get } \hat{r} \text{ by } \vec{r} = r \hat{r}$$

$$\hat{r} = \sin\theta \cos\phi \hat{e}_1 + \sin\theta \sin\phi \hat{e}_2 + \cos\theta \hat{e}_3$$

Spherical

$\hat{\phi}$: think about $\hat{e}_3 \times \hat{r} = |\hat{e}_3| |\hat{r}| \sin\theta \cdot \hat{\phi} \rightarrow$

$$\hat{\phi} = -\sin\theta \hat{e}_1 + \cos\theta \hat{e}_2$$

$$\hat{\phi} = \frac{\hat{e}_3 \times \hat{r}}{\sin\theta} = \frac{1}{\sin\theta} \begin{vmatrix} \hat{e}_1 & \hat{e}_2 & \hat{e}_3 \\ 0 & 0 & 1 \\ \sec\phi & \tan\phi & 0 \end{vmatrix} = -\sin\theta \hat{e}_1 + \cos\theta \hat{e}_2$$

$$\hat{\theta} = \hat{\phi} \times \hat{r} = \begin{vmatrix} \hat{e}_1 & \hat{e}_2 & \hat{e}_3 \\ -\sin\theta & \cos\theta & 0 \\ \cos\theta & \sin\theta & 0 \end{vmatrix} = \cos\theta \cos\phi \hat{e}_1 + \cos\theta \sin\phi \hat{e}_2 - \sin\theta \hat{e}_3 = \hat{\theta}$$

(In matrix form)

$$\begin{bmatrix} \hat{r} \\ \hat{\theta} \\ \hat{\phi} \end{bmatrix} = \begin{bmatrix} \sin\theta \cos\phi & \sin\theta \sin\phi & \cos\theta \\ \cos\theta \cos\phi & \cos\theta \sin\phi & -\sin\phi \\ -\sin\phi & \cos\phi & 0 \end{bmatrix} \begin{bmatrix} \hat{e}_1 \\ \hat{e}_2 \\ \hat{e}_3 \end{bmatrix}$$

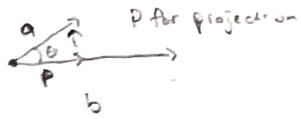
$$\begin{bmatrix} \hat{e}_1 \\ \hat{e}_2 \\ \hat{e}_3 \end{bmatrix} = \begin{bmatrix} \sin\theta \cos\phi & \cos\theta \cos\phi & -\sin\phi \\ \sin\theta \sin\phi & \cos\theta \sin\phi & \cos\phi \\ \cos\theta & -\sin\theta & 0 \end{bmatrix} \begin{bmatrix} \hat{r} \\ \hat{\theta} \\ \hat{\phi} \end{bmatrix}$$

distance formula = $|\vec{r} - \vec{s}|$

dot product intuition $\rightarrow \vec{a} \cdot \vec{b} = |\vec{a}| |\vec{b}| \cos\theta$

and $\cos\theta = \frac{\vec{a} \cdot \vec{b}}{|\vec{a}| |\vec{b}|} \rightarrow \vec{a} \cdot \vec{b} = |\vec{a}| |\vec{b}| \cos\theta$

\hookrightarrow it says how much of \vec{a} 's \vec{b} component is going into \vec{b} (if it's \perp)



if you pull a box 10 meters at an inclined angle there's a horizontal and vertical component to force
dot product gives amount of force going in the direction of displacement
 \hookrightarrow which is what work only cares about $(F \cdot d \hat{d})$

Mechanics Intro Notes

Definition of a vector: $r = x\hat{x} + y\hat{y} + z\hat{z} = x\hat{i} + y\hat{j} + z\hat{k}$ (components x, y, z)

\hookrightarrow Velocity $\rightarrow V_x, V_y, V_z$ components, note $\hat{e}_1 = \hat{x}, \hat{e}_2 = \hat{y}, \hat{e}_3 = \hat{z}$

Linear combo: $r = r_1 \hat{e}_1 + r_2 \hat{e}_2 + r_3 \hat{e}_3 = \sum_{i=1}^3 r_i \hat{e}_i$ or $r = \begin{bmatrix} r_1 \\ r_2 \\ r_3 \end{bmatrix}, \vec{r} = \begin{bmatrix} r_1 \\ r_2 \\ r_3 \end{bmatrix}$

dot product $\vec{r} \cdot \vec{s} = r_1 s_1 + r_2 s_2 + r_3 s_3$ or $\vec{r} \cdot \vec{s} = rs \cos\theta$

$\vec{r} \cdot \hat{e}$ \leftarrow unit vector, $\vec{r} \cdot \vec{s} = \vec{r}^T \vec{s} = [r_1 \ r_2 \ r_3] \begin{bmatrix} s_1 \\ s_2 \\ s_3 \end{bmatrix} = r_1 s_1 + r_2 s_2 + r_3 s_3$

$|\vec{r}| = \sqrt{\vec{r} \cdot \vec{r}} = \sqrt{r_1^2 + r_2^2 + r_3^2}$, unit vector $\frac{\vec{r}}{|\vec{r}|}$, dist between points $\rightarrow ||\vec{r} - \vec{s}||$

$\hookrightarrow \cos\theta = \frac{\vec{r} \cdot \vec{s}}{|\vec{r}| |\vec{s}|}$ (same thing as above) $\rightarrow \vec{r} \cdot \vec{s}$

Cross product $\vec{r} \times \vec{s} = \det \begin{bmatrix} \hat{x} & \hat{y} & \hat{z} \\ r_x & r_y & r_z \\ s_x & s_y & s_z \end{bmatrix}$ or $\vec{p} = \vec{r} \times \vec{s} = \begin{cases} p_x = r_y s_2 - r_z s_y \\ p_y = r_z s_x - r_x s_z \\ p_z = r_x s_y - r_y s_x \end{cases}$

Vector \perp to both \vec{r} and \vec{s} , $|\vec{r} \times \vec{s}| = rs \sin\theta$, $\frac{d}{dt}(axb) = \frac{da}{dt} \times b + a \times \frac{db}{dt}$ (product rule)

Product rule: $\frac{d}{dt}(f(t)r(t)) = f \frac{dr}{dt} + \frac{df}{dt} r$, chain rule: if $y = f(u)$ and $u = g(x)$ then $\frac{dy}{dx} = \frac{dy}{du} \frac{du}{dx}$
example, $F(x) = e^{x^2}$, $f(x) = \text{outer: } e^x, y = f(u) = e^u$, $g(u) = \text{inner: } x^2, u = g(x) = x^2$, $\frac{du}{dx} = 2x$ and $\frac{dy}{du} = e^u$ then $\frac{dy}{dx} = \frac{dy}{du} \frac{du}{dx} = 2x e^{x^2}$

Velocity if $r = x\hat{x} + y\hat{y} + z\hat{z}$ then $\frac{dr}{dt} = \frac{dx}{dt}\hat{x} + \frac{dy}{dt}\hat{y} + \frac{dz}{dt}\hat{z} = v = V_x\hat{x} + V_y\hat{y} + V_z\hat{z}$

Unit vectors $\hat{x}, \hat{y}, \hat{z}$ \rightarrow if unit vectors depend on time like in other coordinate systems its not simple \rightarrow polar coordinates this is not true so you use product rule \rightarrow NOT just derive w/ each of components!

Force Newton, 1 N is magnitude of force that accelerates a kg w/ an acceleration of $\frac{1\text{m}}{\text{s}^2}$

\hookrightarrow direction of Force is the direction of resulting acceleration

Newton's Laws

\hookrightarrow point masses \rightarrow bodies w/ rotation, if 2 frames are moving relative to each other Newton's laws don't hold \rightarrow acceleration for this second frame is non-inertial

First Law \rightarrow absence of forces particle moves w/ constant v , Second Law $\rightarrow F = ma = m\ddot{v} = m\ddot{r}$

Second Law rephrased $\rightarrow p = mv, \dot{p} = m\dot{v} = ma$ so $F = \dot{p}$

Newton's laws for a particle undergoing a constant force \rightarrow \vec{F} constant in x dir, so $\ddot{x}(t) = \frac{F_0}{m}$

Third Law \rightarrow F_{21} is force on 2 by 1 \rightarrow F_{12} is force on 1 by 2

$$\ddot{x}(t) = \int \ddot{x}(t) dt = V_{initial} + \frac{F_0}{m} t, \text{ with } \dot{x}(0) = V_{initial}$$

$$x(t) = \int \dot{x}(t) dt = X_0 + V_0 t + \frac{F_0}{2m} t^2$$

if object 1 exerts a force on object 2 (F_{21}) then object 2 always exerts a reaction force (F_{12}) given by

$$F_{12} = -F_{21}$$

$$F_{21}$$

if you push on a wall you don't fall over so there is a reaction force by the wall due to molecular forces

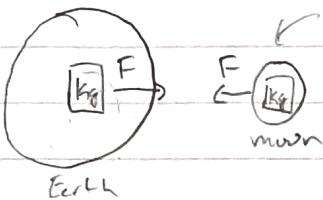
More detailed look into Newton's 3rd Law on back

* a non inertial RF
where first and second law don't hold

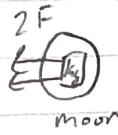
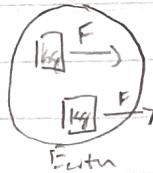
train is moving at constant velocity the ice cube will not move
if the train accelerates the ice cube will fly back even though it experiences no net force! non-inertial RF same goes for a rotating Earth is a non-inertial RF but can be approximated on small time + very low mass length scales to be inertial

Newton's 3rd law or deeper approach

[a couple examples]
gravitational example



equal and opposite



$$-2F = 2F$$

$$-(2ma) = 2(ma)$$

what?

[in 2 by adding an extra 1 kg the force double on the 1 kg for moon but earth also feels double the force but over 2 kg]

↳ the earth will accelerate less but it has more mass, more inertia but feels the same force!

The earth pulls on moon mass but moon pulls on all of earth's mass in such a way its equal and opposite, since earth's bigger its more spread out & has more inertia causing less acceleration

Conservation of momentum from third law

p_1, p_2 for express

(no external forces!)

$$F_{12} = -F_{21} \text{ and } F = \dot{p} \text{ so, } \frac{dp_{12}}{dt} = -\frac{dp_{21}}{dt} \rightarrow \frac{d(p_{12} + p_{21})}{dt} = 0$$

↳ if the rate of change of this sum is zero, which means $p_1 + p_2$ is constant

$$p_1 + p_2 = p_{\text{tot}} = m_1 v_1 + m_2 v_2$$

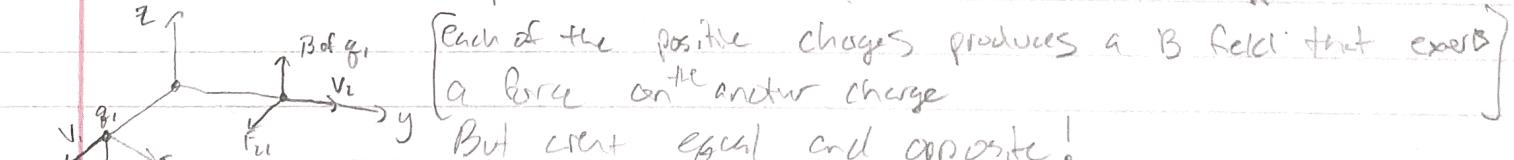
(net force on 1) $\vec{F}_1 = \vec{F}_{12} + \vec{F}_{\text{ext}}$ external force is like the sun or wind by the earth, moon system or skaters

(net force on 2) $\vec{F}_2 = \vec{F}_{21} + \vec{F}_{\text{ext}} = \dot{\vec{p}}_2$ and same for $\vec{F}_1 = \dot{\vec{p}}_1$ special case is $F_{\text{ext}} = 0$ and $p = \text{constant}$
↳ $p_{\text{tot}} = p_1 + p_2$, $\dot{p}_{\text{tot}} = \dot{p}_1 + \dot{p}_2 = \vec{F}_{\text{ext}} + \vec{F}_{\text{ext}}$, ($F_{12} = -F_{21}$), conservation of momentum!
then we can just say $\dot{p}_{\text{tot}} = \vec{F}_{\text{ext}}$ (total external force)

this is true for multi-particle systems, proof on page 316 of Taylor

third law is not true in relativity

also doesn't work in charge particles in B field



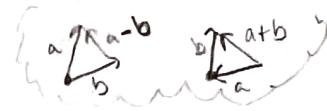
* mechanical momentum isn't only kind of momentum

E fields can also carry momentum (neglect E & V_{CC})

so its ok in low speeds we are good in class 2/3/ mechanics

back to normality

* solving problems in non-rectangular
shifts things very easily

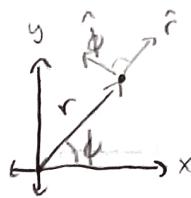


Polar coordinates

unit vectors \hat{r} and $\hat{\phi}$
 direction moves r while ϕ is fixed
 r direction we move when ϕ increases and r is fixed
 \hat{r} are linearly independent

$$\mathbf{F} = F_r \hat{r} + F_\phi \hat{\phi}$$

$$\hat{r} = \frac{\mathbf{r}}{|\mathbf{r}|}$$



(2) throwing a stone in circle.
 F_r is tension and $F_d = air$ resistance

\hat{r} changes as r moves so its not constant like curvilinear. Find $\frac{d\hat{r}}{dt}$

$$\begin{aligned}\hat{r} &= r \hat{i} \\ \frac{d\hat{r}}{dt} &\neq \hat{r}'\end{aligned}$$

doing Δ of vector is connecting heads from tails
 \hat{r} is not vector
 $\hat{r}, \hat{\phi}$ are mag=1 and dont change direction

to find $\frac{d\hat{r}}{dt}$ algebraically use def of polar coord.

$$\begin{aligned}\hat{r} &= \hat{x} \cos\phi + \hat{y} \sin\phi \\ \frac{d\hat{r}}{dt} &= \hat{x} \frac{d}{dt}(\cos\phi) + \hat{y} \frac{d}{dt}(\sin\phi) \\ \frac{d\hat{r}}{dt} &= \hat{x} \left[-\sin\phi \frac{d\phi}{dt} \right] + \hat{y} \left[\cos\phi \frac{d\phi}{dt} \right] \\ \frac{d\hat{r}}{dt} &= \frac{d\phi}{dt} \hat{\phi} \quad \text{wow!}\end{aligned}$$

$$\begin{aligned}\hat{\phi} &= -\hat{x} \sin\phi + \hat{y} \cos\phi \\ \frac{d\phi}{dt} &= -\hat{x} \cos\phi \frac{dd}{dt} - \hat{y} \sin\phi \frac{dd}{dt} \\ \frac{dd}{dt} &= -\frac{d\phi}{dt} (\hat{x} \cos\phi + \hat{y} \sin\phi)\end{aligned}$$

triple product rule

$$\frac{d\hat{r}}{dt} = -\hat{r} \frac{d\phi}{dt}$$

Product rule on $\vec{r} = r\hat{r}$: $\vec{v} = \frac{d\vec{r}}{dt} = \frac{dr}{dt} \hat{r} + r \frac{d\hat{r}}{dt} \Rightarrow \vec{v} = \frac{dr}{dt} \hat{r} + \frac{d\phi}{dt} r \hat{\phi}$

\downarrow $V_r = \dot{r}$ angular velocity [To write Second] $\rightarrow a = \ddot{v} = \ddot{r} = \frac{d}{dt}(\dot{r}\hat{r} + r\dot{\phi}\hat{\phi})$
 $V_\phi = \dot{\phi} r = cor$ Law $\frac{d}{dt} v$

\hookrightarrow Product RULE $\rightarrow a = \frac{d}{dt}(\dot{r}\hat{r}) + \frac{d}{dt}(r\dot{\phi}\hat{\phi}) = (\dot{r}\frac{d\hat{r}}{dt} + \ddot{r}\hat{r}) + ((\dot{r}\dot{\phi} + r\ddot{\phi})\hat{\phi} + r\dot{\phi}\frac{d\hat{\phi}}{dt})$ \hookrightarrow sub in!

Sub in $\frac{d\hat{r}}{dt}$ and collect $\hat{r}, \hat{\phi}$

Holy moly this is cool!

$$\vec{a} = (\ddot{r} - r\omega^2) \hat{r} + (r\ddot{\phi} + 2\dot{r}\dot{\phi}) \hat{\phi}$$

(acceleration) (centrifugal acceleration)
 (as r varies) (tangential acceleration)
 radial stuff

\hookrightarrow $\perp r$ or tangent

$$* r\omega^2 = \frac{v^2}{r}$$

(2nd Law) $\Rightarrow \begin{cases} F_r = m(\ddot{r} - r\omega^2) & * \alpha = \dot{\phi} \\ F_\phi = m(r\ddot{\phi} + 2\dot{r}\dot{\phi}) \end{cases}$ if r is constant like in most problems $\rightarrow a = -r\omega^2 \hat{r} + r\alpha \hat{\phi}$

\vec{r} is constant

* how long will it take to come back to point of release?

$\hookrightarrow \vec{a} = -r\omega^2 \hat{r} + r\alpha \hat{\phi}$

\downarrow $\begin{cases} F_r = m(-r\omega^2) \end{cases} \rightarrow \sum F_r = mg \cos\phi - N = -mR\omega^2$

\downarrow $\begin{cases} F_\phi = m(r\alpha) \end{cases} \rightarrow \sum F_\phi = -mg \sin\phi = mr\alpha$

or α for small angles

$\hookrightarrow \alpha = \dot{\phi} = -\frac{g}{R} \sin\phi$, if you say $\omega = \sqrt{\frac{g}{R}}$ $\rightarrow \dot{\phi} = -\omega^2 \sin\phi = -\omega^2 \phi$

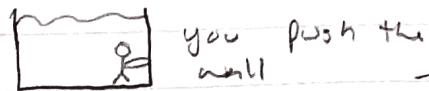
$\dot{\phi} = -\omega^2 \phi$ is a solvable diff eq of form $\dot{\phi}(t) = A \sin(\omega t) + B \cos(\omega t)$ $\dot{\phi} = B$, $B = \phi_0$ and $\phi(t)$ eq implies $\dot{\phi} = \omega A$, $\dot{\phi} = 0$ so $A = 0$ and $\phi(t) = \phi_0 \cos(\omega t)$, returns to ϕ_0 when $\omega t = 2\pi$

$$T = \frac{2\pi}{\omega} = 2\pi \sqrt{\frac{R}{g}}$$

Sorry Newtons
3rd Law Again

ex1: IF I give you \$10, I am \$10 poorer, and you are \$10 better off. So how does anyone get rich?
→ if we sum +\$10 and -\$10 = \$0, mistake is thinking this applies to the same person. They don't.

ex2: sitting on wheelchair chair, if you push on wall, wall pushes on you and you roll backwards

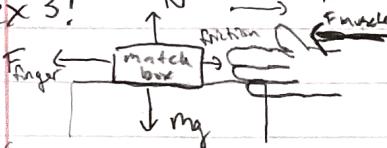


↳ or in a swimming pool, you are at the bottom

By 3rd law you push it equally and it equally pushes on you. The wall/Earth has inertia so it barely moves (reality the wall flexes). But you small so you fly back.

~ equilibrium can only establish itself when forces act on same object (1st Law)

↳ things move because force is applied to the other object.

ex3:  $\sum F = N - mg = 0$ (this first law though)

$$F_{\text{finger}} = -F_{\text{matchbox}} \quad (\text{3rd Law})$$

~ person doesn't move bcc of friction if no friction (and ratio of mass to force) ↳ person would in opposite direction (less box or inertial) $\ddot{m} = a$

$F = ma$ in coordinates

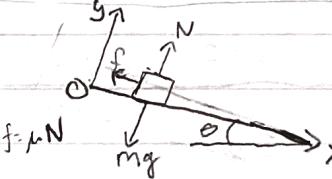
$$(F = m \frac{d^2 r}{dt^2})$$

constant + constant

• finding components ($F = F_x \hat{x} + F_y \hat{y} + F_z \hat{z}$) and ($r = x \hat{x} + y \hat{y} + z \hat{z}$)

↳ $\ddot{r} = \ddot{x} \hat{x} + \ddot{y} \hat{y} + \ddot{z} \hat{z}$ so $F = m \ddot{x} \hat{x} + m \ddot{y} \hat{y} + m \ddot{z} \hat{z}$

$$F = m \ddot{r} \rightarrow \begin{cases} F_x = m \ddot{x} \\ F_y = m \ddot{y} \\ F_z = m \ddot{z} \end{cases}$$



RF: where it starts a.g. y normal to the plane, otherwise x and y would vary, also f and N will have a component = 0

Ex block sliding on incline plane $f \perp N$
How far will it travel in time t?

• no motion in y-direction, $\dot{y} = 0$ therefore $F_y = m \dot{y} = 0$, y direction is in equilibrium

↳ $F_y = N - mg \cos \theta = 0 \rightarrow N = mg \cos \theta$ thus $f = \mu N = \mu mg \cos \theta$

$$\text{Weight } x = f_x = \mu m g \cos \theta$$

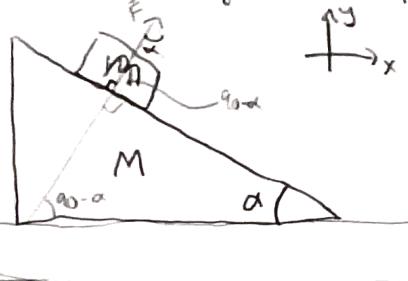
$$\text{Now, } F_x = m \ddot{x} = m g \sin \theta - \mu m g \cos \theta \rightarrow \ddot{x} = g(\sin \theta - \mu \cos \theta)$$

all these are constants, don't depend on x or t

$$x(t) = \frac{1}{2} g (\sin \theta - \mu \cos \theta) t^2$$

↳ for integration and initial cond, $(\dot{x} = 0 \text{ at } t=0)$
 $x = 0 \text{ at } t=0$

a block sliding on a plane (Block slides down and Inclined plane can not move)



$F_N = \text{normal force of } M \text{ on inclined plane (IP)}$

and $SB = \text{small block}$

$$\bullet \vec{F}_L = F_N \sin(\alpha) \hat{i} + F_N \cos(\alpha) \hat{j}$$

$$\bullet \vec{F}_{\text{gravity (SB)}} = -mg \hat{j}$$

[Let \vec{A} be the acceleration of
IP in an inertial ref frame] \rightarrow [\vec{A} has an \hat{i} component = A_x]

(need help out from frames)

Why?

[Let $\vec{a} + \vec{A}$ be the acceleration of
the small block in same ref frame] \rightarrow [\vec{a} is the acceleration of the block
relative to the inclined plane]

I'm guessing
related

↳ (now use Newton's law) $\rightarrow \vec{F}(SB) = \vec{F}_L + \vec{F}_{g(SB)} = m(\vec{a} + \vec{A})$ * now find F_{\parallel} and F_{\perp}

$$F_{\parallel} = mg \sin(\alpha) = ma_{\parallel} + mA_x \cos(\alpha) \leftarrow [\text{parallel to the top of plane}] \leftarrow (\text{simple trig})$$

$$F_{\perp} = F_N - mg \cos(\alpha) = ma_{\perp} + mA_x \sin(\alpha) \leftarrow [\text{perp to the top of the plane}]$$

$$13 \quad \begin{aligned} |F_N| & \\ \text{solve for } F_N & \\ F_N &= mg \cos(\alpha) + mA_x \sin(\alpha) \end{aligned}$$

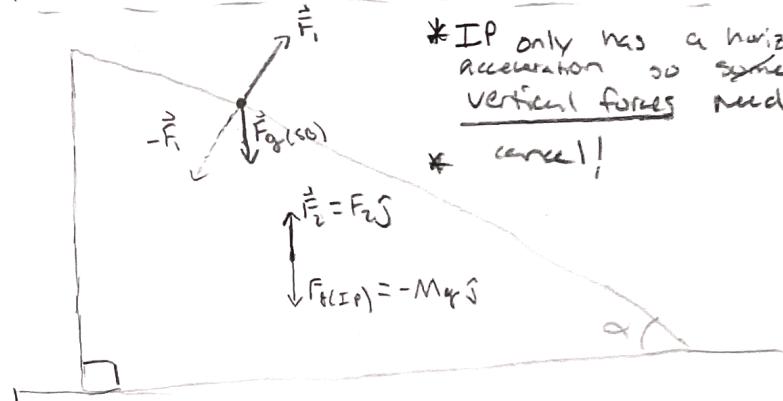


* IP only has a horizontal acceleration so some all vertical forces need to cancel!

Sum forces on inclined plane:

$$\vec{F}_{(IP)} = -\vec{F}_1 + \vec{F}_2 + \vec{F}_{g(IP)} = M\vec{A}$$

3rd law force
IP pushes on m (F_1)
m pushes on IP ($-F_1$)
(normal force of IP on ground)
(just vertical
so $\vec{F}_2 = F_2 \hat{j}$)



$$\begin{aligned} (\text{vertical forces on IP}) \quad F_2 \hat{j} + (-Mg \hat{j}) &= (-F_1 \cos(\alpha)) \rightarrow F_2 \hat{j} - Mg \hat{j} = -F_1 \cos(\alpha) \hat{j} \\ \text{Forces from IP} &= \text{Forces by m} \end{aligned}$$

(determines motion of the inclined plane)

$$(\text{horizontal forces}) \quad MA_x = -F_1 \sin(\alpha)$$

$$\text{Now plug } F_1 \text{ from above into here} \rightarrow MA_x = -(mg \cos(\alpha) + mA_x \sin(\alpha)) \sin(\alpha)$$

$$MA_x = -g \left(\frac{\sin(\alpha) \cos(\alpha)}{\sin^2(\alpha) + \frac{M}{m}} \right)$$

Motion for small block \longrightarrow

(Motion for
S10 is given) $\rightarrow (a_{11}) \rightarrow$ [which we have]
from eq 1.2 $\rightarrow m g \sin(\alpha) = m a_{11} + m A_x \cos(\alpha)$

$$\rightarrow a_{11} = g \sin(\alpha) - A_x \cos(\alpha)$$

(get by plugging in)
 A_x in question)

$$\rightarrow \left[\begin{array}{l} \text{vertical comp of} \\ \text{the acceleration} \end{array} \right] \rightarrow a_y = -a_{11} \sin(\alpha) = -g \sin^2(\alpha) \left(\frac{M+m}{M+m \sin^2 \alpha} \right)$$

may be down

Let's say block starts at height "h"] Using final initial $y_2 - y_1 = a_y \frac{t^2}{2} \rightarrow 0 - h = a_y \frac{t^2}{2} \rightarrow t = \sqrt{-2h/a_y}$

$$\hookrightarrow \text{also } a_x = +a_{11} \cos(\alpha)$$



Notes to take away

1) Started with 2 objects each having their own coordinates $(x_m, y_m), (x_M, y_M)$, 4 coordinates, but with our constraints of the problem (m moves \perp incline plane and incline plane has $A_y = 0$)

2) also notice the constraint forces, normal forces are NOT in our final solution, so it should be reasonable to think of a way to solve the problem w/out them in the first place,

\hookrightarrow the constraints restrict the number of degrees of freedom of the system of rotate 1