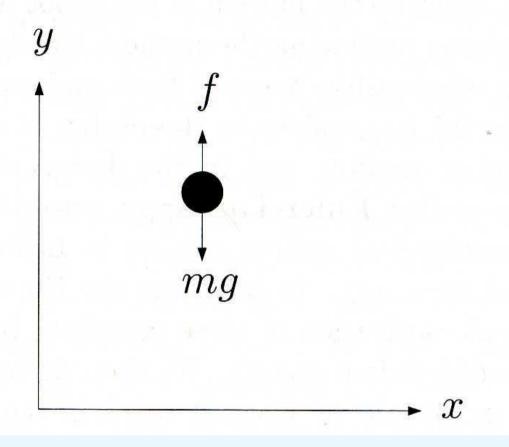
Lecture 10: Dynamics: Euler-Lagrange Equations

Examples

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- Examples
- Holonomic Constraints and Virtual Work



The second Newton law says that the equation of motion of the particle is

$$mrac{d^2}{dt^2}y=\sum_i F_i=f-mg$$

- f is an external force;
- mg is the force acting on the particle due to gravity.

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ight]
ight) = rac{d}{dt} \left(rac{\partial}{\partial \dot{y}} \mathcal{K}
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$$\begin{array}{ll} m\frac{d^2}{dt^2}y & = & \frac{d}{dt}\left(m\frac{d}{dt}y\right) = \frac{d}{dt}\left(m\frac{\partial}{\partial\dot{y}}\left[\frac{1}{2}\,\dot{y}^2\right]\right) = \frac{d}{dt}\left(\frac{\partial}{\partial\dot{y}}\mathcal{K}\right) \\ mg & = & \frac{\partial}{\partial y}\left[mgy\right] = \frac{\partial}{\partial y}\mathcal{P} \end{array}$$

with kinetic/potential energies defined by $\mathcal{K} = \frac{1}{2}m\dot{y}^2, \ \mathcal{P} = mgy$

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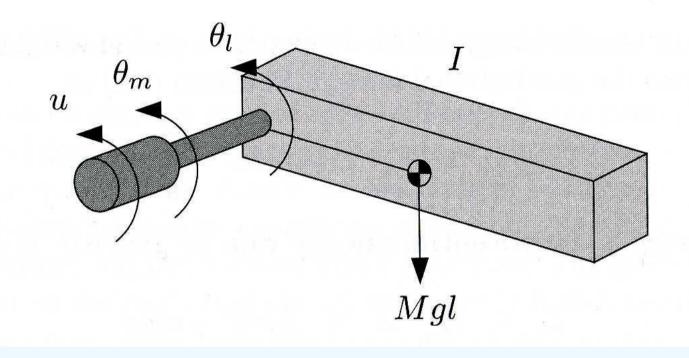
$$mg = \frac{\partial}{\partial y} \left[mgy \right] = \frac{\partial}{\partial y} \mathcal{P}$$

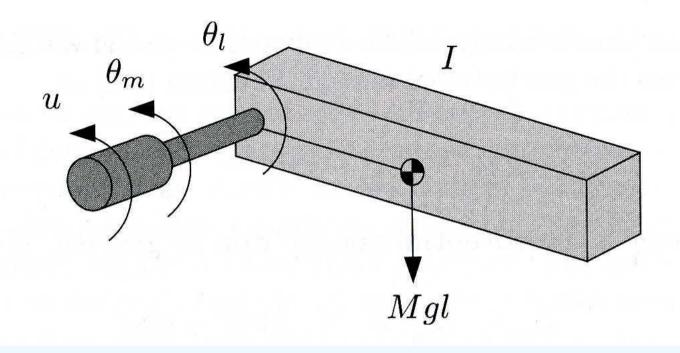
with kinetic/potential energies defined by $\mathcal{K} = \frac{1}{2}m\dot{y}^2, \ \mathcal{P} = mgy$

Then the second Newton law can be rewritten as

$$ilde{m{ iny}} rac{d}{dt} \left(rac{\partial}{\partial \dot{y}} {m{\mathcal{L}}}
ight) - rac{\partial}{\partial y} {m{\mathcal{L}}} = f$$
 with ${m{\mathcal{L}}} = {m{\mathcal{K}}} - {m{\mathcal{P}}}$

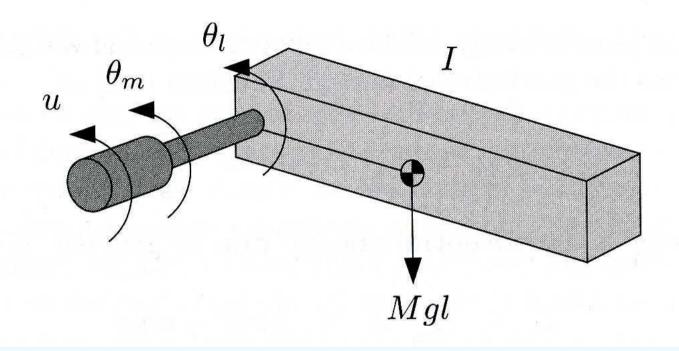
where the function $\mathcal{L}(y, \dot{y})$ is called the Lagrangian.



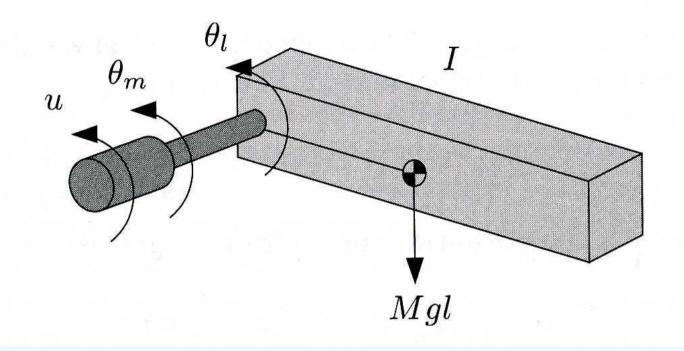


A rigid link (θ_l) coupled through a gear to DC motor $(\theta_m = r\theta_l)$:

ullet Kinetic energy: ${\cal K}=rac{1}{2}J_m\dot heta_m^2+rac{1}{2}J_l\dot heta_l^2=rac{1}{2}\left({m r}^2J_m+J_l
ight)\dot heta_l^2$

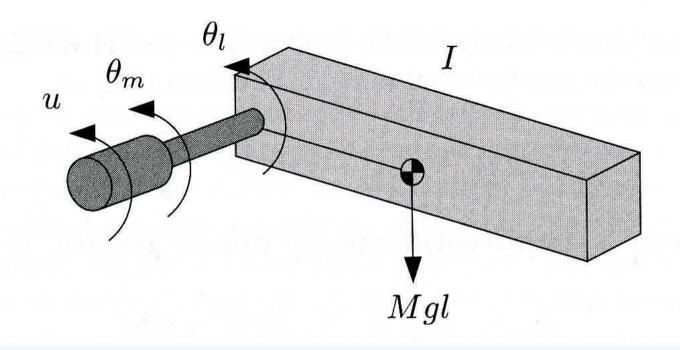


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$$rac{d}{dt}\left(rac{\partial}{\partial\dot{ heta}_l}\mathcal{L}
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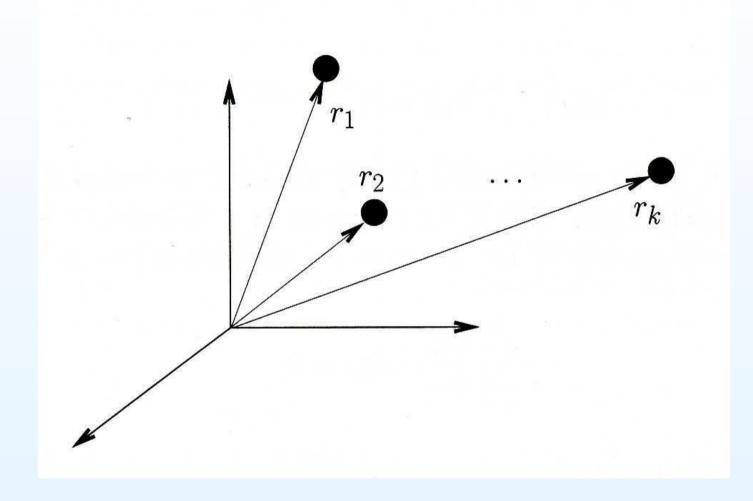


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$$(\mathbf{r}^2 J_m + J_l) \ddot{\theta}_l + Mgl \sin \theta_l = \mathbf{r}u$$

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- Examples
- Holonomic Constraints and Virtual Work



Unconstrained system of k particles has 3k degrees of freedom. The number of Dof is less if the particles are constrained

A constraint imposed on k particles (with coordinates $r_1, r_2, \ldots, r_k \in \mathbb{R}^3$) is called holonomic, if it is of the form

$$g_i(r_1, r_2, \ldots, r_k) = 0, \qquad i = 1, 2, \ldots, l$$

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For example, given two particles joined by massless rigid wire of length l, then

$$\left\Vert r_{1},\,r_{2}\in\mathbb{R}^{3}:\,\left\Vert r_{1}-r_{2}
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Presence of constraint implies presence a force

(called constraint force), that forces this constraint to hold.

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Differentiating the constraint function $g_i(\cdot)$ with respect to time, we obtain new constraint

$$rac{d}{dt}g_i(r_1,\,r_2,\,\ldots,\,r_k)=rac{\partial g_i}{\partial r_1}rac{d}{dt}r_1+\cdots+rac{\partial g_i}{\partial r_k}rac{d}{dt}r_k=0$$
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$$\frac{d}{dt}g_i(r_1, r_2, \dots, r_k) = \frac{\partial g_i}{\partial r_1} \frac{d}{dt} r_1 + \dots + \frac{\partial g_i}{\partial r_k} \frac{d}{dt} r_k = 0$$

$$\frac{\partial g_i}{\partial r_1} dr_1 + \dots + \frac{\partial g_i}{\partial r_k} dr_k = 0$$

or

$$rac{\partial g_i}{\partial r_1} dr_1 + \dots + rac{\partial g_i}{\partial r_k} dr_k = 0$$

The constraint of the form

$$\omega_1(r_1,\ldots r_k)dr_1+\cdots+\omega_k(r_1,\ldots r_k)dr_k=0$$

is called non-holonomic if it cannot be integrated back.

Concept of Generalized Coordinates

If the system is subject to holonomic constraint then

 If system consists of k particles, then it may be possible to express their coordinates as functions of fewer than 3k variables

$$r_1 = r_1(q_1, \dots, q_n), \ r_2 = r_2(q_1, \dots, q_n), \dots,$$
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- The smallest set of variables is called generalized coordinates
- This smallest number is called a number of degrees of freedom
- If the system consists of an infinite number of particles, then it might have finite number of degrees of freedom

Given a system of k-particles and a holonomic constraint

$$g_i(r_1, r_2, \ldots, r_k) = 0, \qquad i = 1, 2, \ldots, l$$

or the same

$$\frac{\partial g_i}{\partial r_1} dr_1 + \dots + \frac{\partial g_i}{\partial r_k} dr_k = 0, \quad i = 1, 2, \dots, l$$

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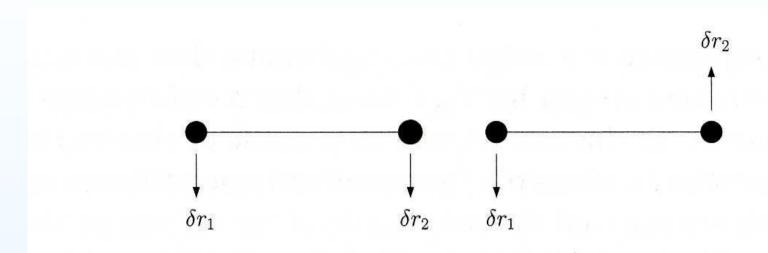
By definition a set of infinitesimal displacements

$$\delta r_1, \ \delta r_2, \ \ldots, \ \delta r_k$$

that are consistent with the constraint, i.e.

$$\frac{\partial g_i}{\partial r_1} \frac{\delta r_1}{\delta r_1} + \dots + \frac{\partial g_i}{\partial r_k} \frac{\delta r_k}{\delta r_k} = 0, \quad i = 1, 2, \dots, l$$

are called virtual displacements



Virtual displacements of a rigid bar. Such infinitesimal motions do not destroy the constraint

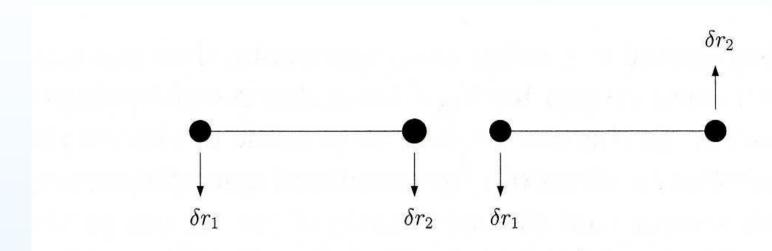
$$\left(r_1-r_2
ight)^{\scriptscriptstyle T}\left(r_1-r_2
ight)=l^2$$

if r_1 and r_2 are perturbed

$$m{r_1}
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that is

$$\left(\left(r_1+\boldsymbol{\delta r_1}\right)-\left(r_2+\boldsymbol{\delta r_2}\right)\right)^{\mathrm{\scriptscriptstyle T}}\left(\left(r_1+\boldsymbol{\delta r_1}\right)-\left(r_2+\boldsymbol{\delta r_2}\right)\right)=l^2$$



Virtual displacements of a rigid bar. Such infinitesimal motions do not destroy the constraint

$$(r_1 - r_2)^{^{\mathrm{\scriptscriptstyle T}}} (r_1 - r_2) = l^2$$

if r_1 and r_2 are perturbed

$$r_1 \rightarrow (r_1 + \boldsymbol{\delta r_1}) \qquad r_2 \rightarrow (r_2 + \boldsymbol{\delta r_2})$$

that is

$$(r_1-r_2)^{{\scriptscriptstyle T}} ({\color{blue}\delta r_1-\delta r_2})=l^2$$

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Then the total sum of all forces applied to i^{th} -particle is zero

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Then the work done by all forces applied to i^{th} -particle along each set of virtual displacement is zero, i.e.

$$oxed{0} = \sum_{oldsymbol{i}} \left(f_{oldsymbol{i}}^c + f_{oldsymbol{i}}^e
ight) oldsymbol{\delta r_i}} = \sum_{oldsymbol{i}} f_{oldsymbol{i}}^c oldsymbol{\delta r_i} + \sum_{oldsymbol{i}} f_{oldsymbol{i}}^e oldsymbol{\delta r_i}$$