

# Characteristics of pinned polymer loop in external field

Wenwen Huang

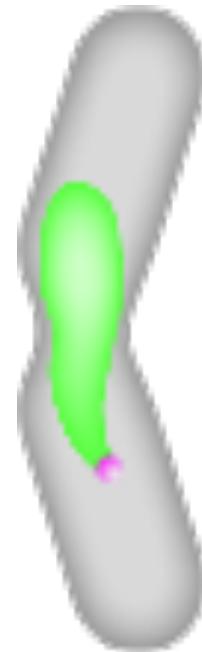
05.01.2015

For PhD seminar

# Outline

- Background Reminder
- Results
- TODO

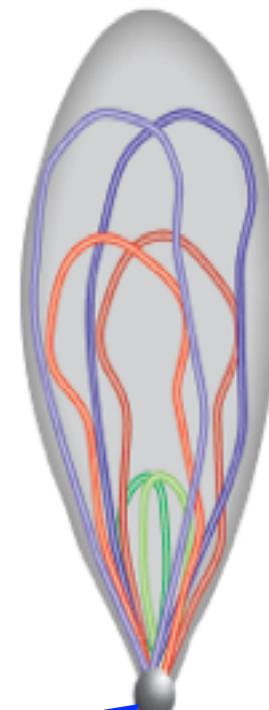
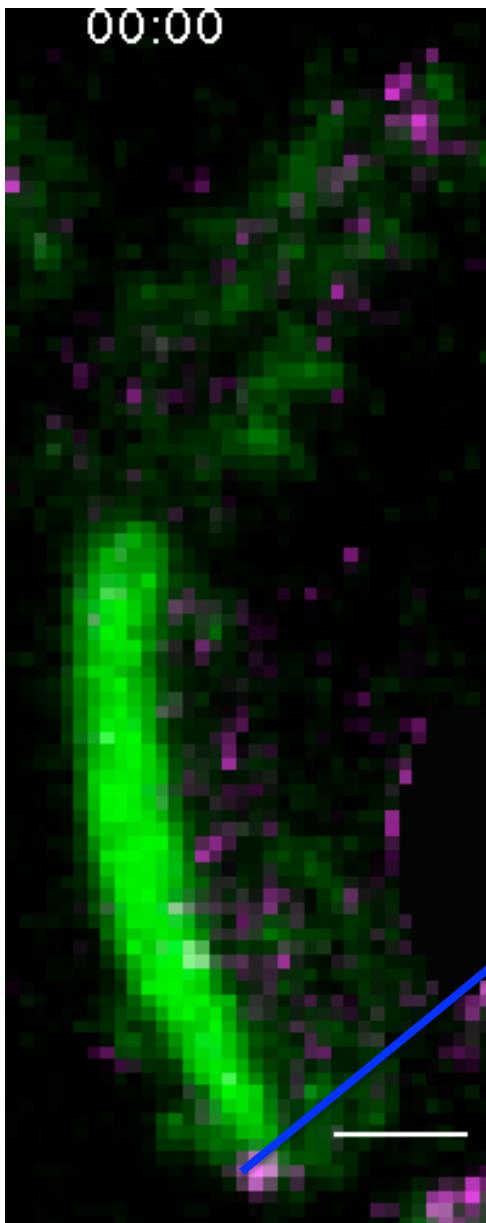
# Fission Yeast (*s. pombe*)



Sketch of fission yeast  
during cell division

Rhind N et al. Science. April 21, 2011

# Chromosome Movement (*S. pombe*)



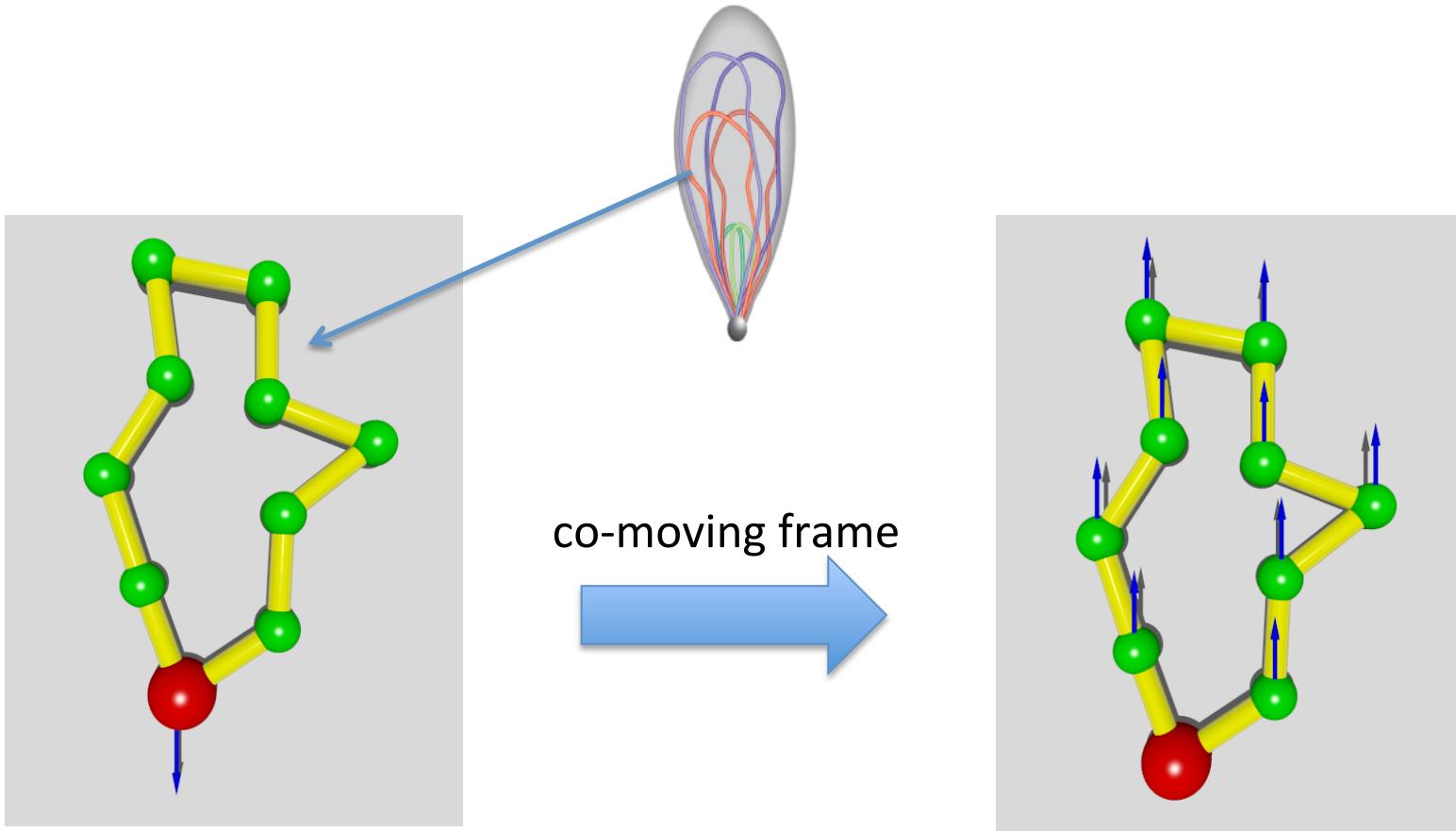
- 3 pairs of chromosomes
- Telomere bond to SPB
- Oscillation is driven by SPB

In the meantime:  
Homologous chromosomes  
pairing and recombination

Chikashige et al., Science 1994  
Yamamoto et al, JCB 1999  
Ding et al., DevCell 2004

# Polymer Ring

- One Chromosome is represented by a polymer ring with the driven bead be the **SPB**(spindle pole body).



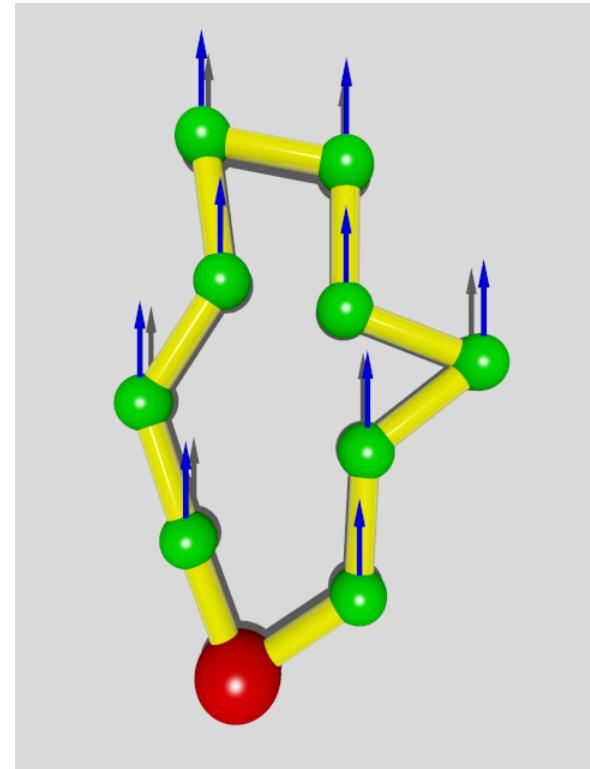
# Freely Joined Bead-Rod Ring

- Bead:  
position ----  $\mathbf{r}_i$
- Rod:  
length ---  $a$

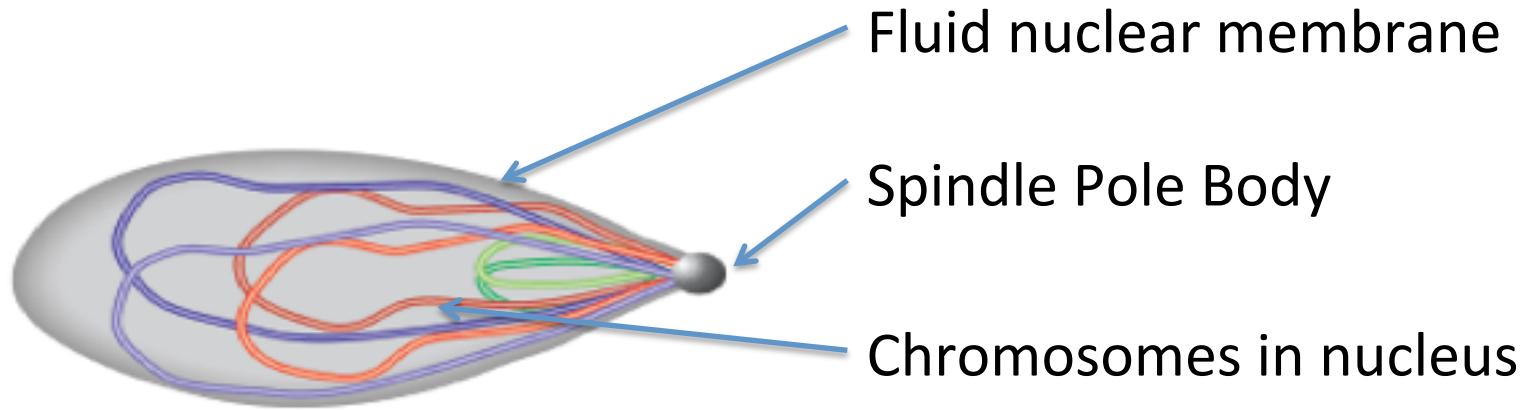
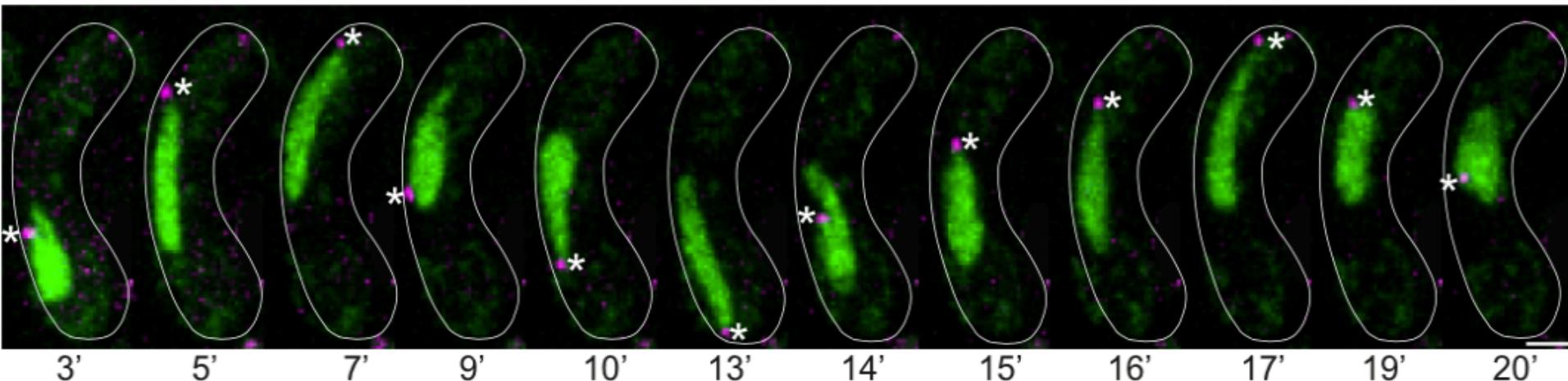
Results we already have:

Mean position:  $\langle \mathbf{r}_i \rangle$

Position Variance:  $Var[\mathbf{r}_i]$



# Shape of nucleus



Assume the shape of nucleus is determined by chromosomes in it

# Characteristic quantities

- Gyration radius:  $R_g^2 = \left\langle \frac{1}{N} \sum_{i=1}^N (\mathbf{r}_i - \mathbf{r}_{CM})^2 \right\rangle$
- Gyration tensor:  $Q_{xy} = \frac{1}{N} \sum_{i=1}^N \mathbf{r}_x^i \mathbf{r}_y^i$   
$$Q = \begin{bmatrix} Q_{xx} & Q_{xy} & Q_{xz} \\ Q_{yx} & Q_{yy} & Q_{yz} \\ Q_{zx} & Q_{zy} & Q_{zz} \end{bmatrix} = \begin{bmatrix} \lambda_x^2 & 0 & 0 \\ 0 & \lambda_y^2 & 0 \\ 0 & 0 & \lambda_z^2 \end{bmatrix}$$
$$\lambda_x^2 \leq \lambda_y^2 \leq \lambda_z^2$$

Asphericity

$$b = \lambda_z^2 - \frac{1}{2}(\lambda_x^2 + \lambda_y^2)$$

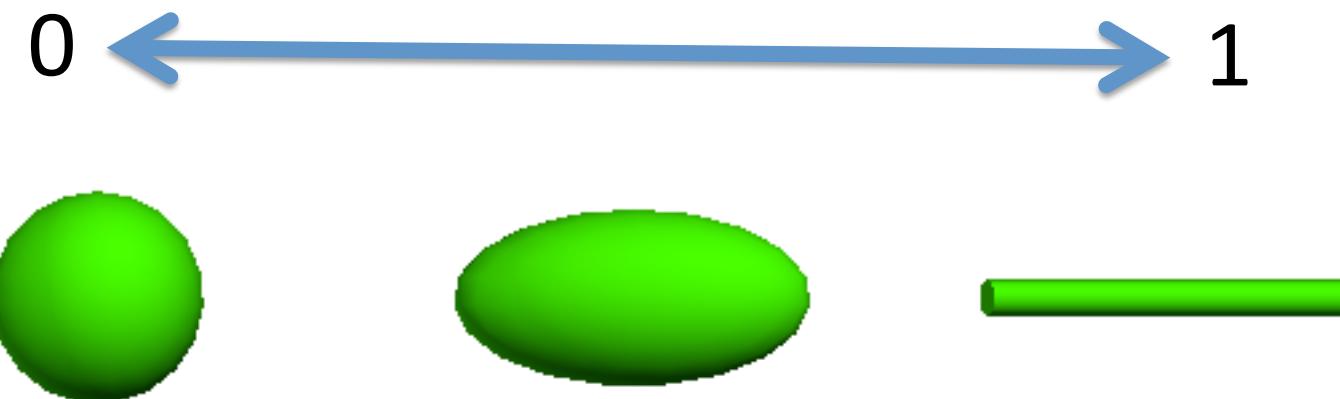
Acylicity

$$c = \lambda_y^2 - \lambda_x^2$$

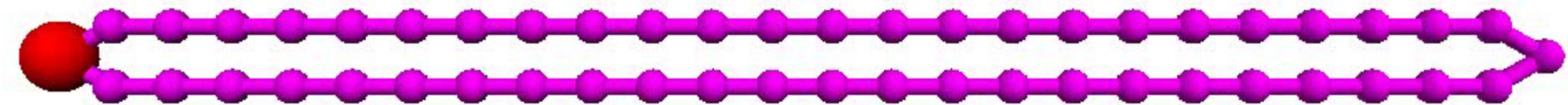
# Relative shape anisotropy

$$\kappa^2 = \frac{3}{2} \frac{\lambda_x^4 + \lambda_y^4 + \lambda_z^4}{(\lambda_x^2 + \lambda_y^2 + \lambda_z^2)^2} - \frac{1}{2}$$

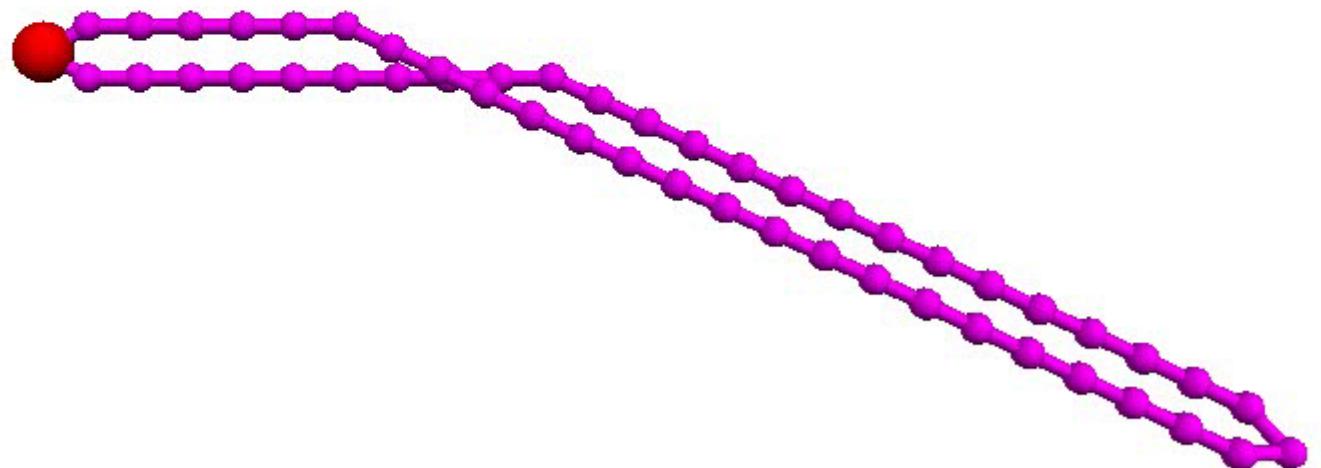
$$0 \leq \kappa \leq 1$$



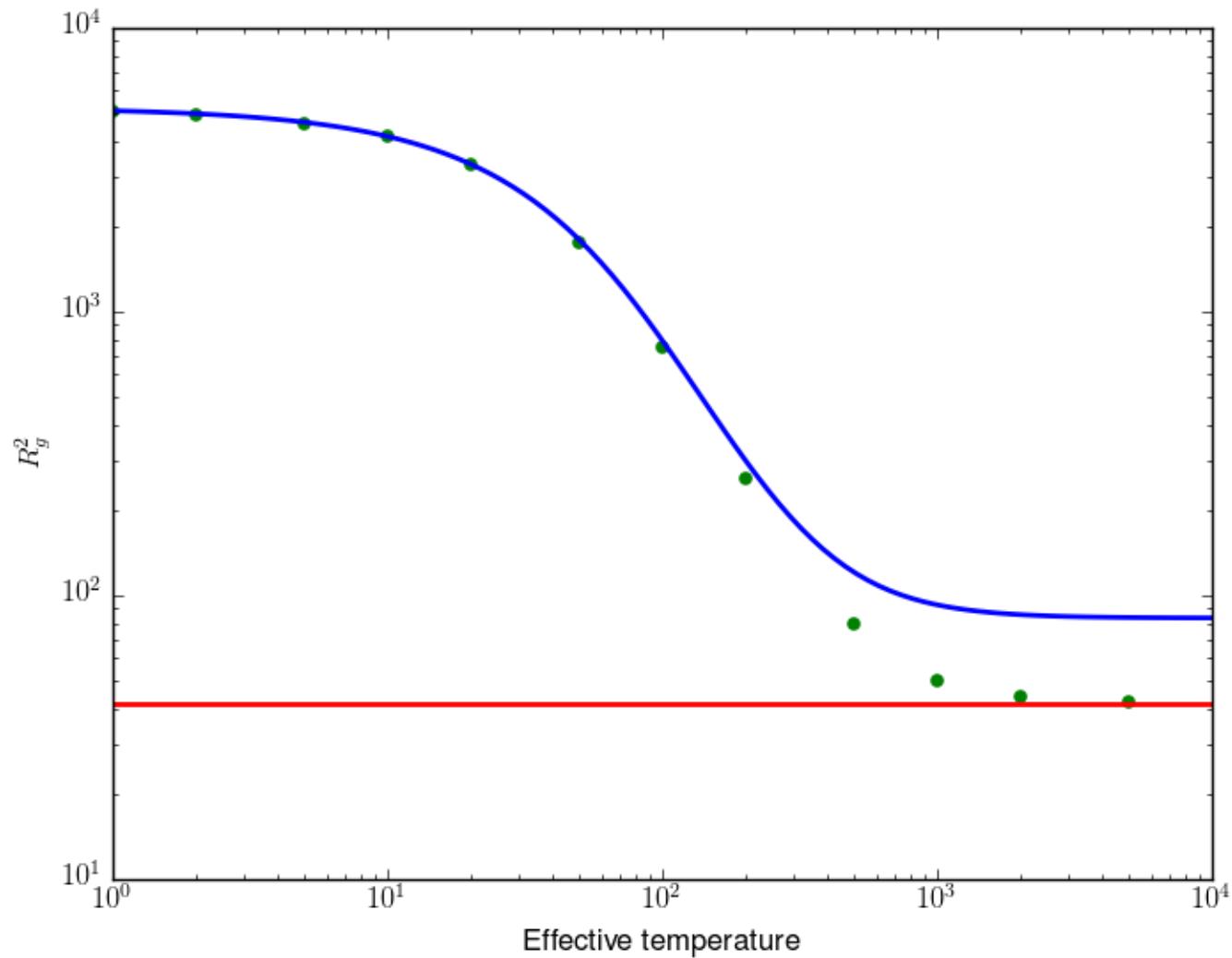
Results ( $T = 1$ ,  $N = 50$ )



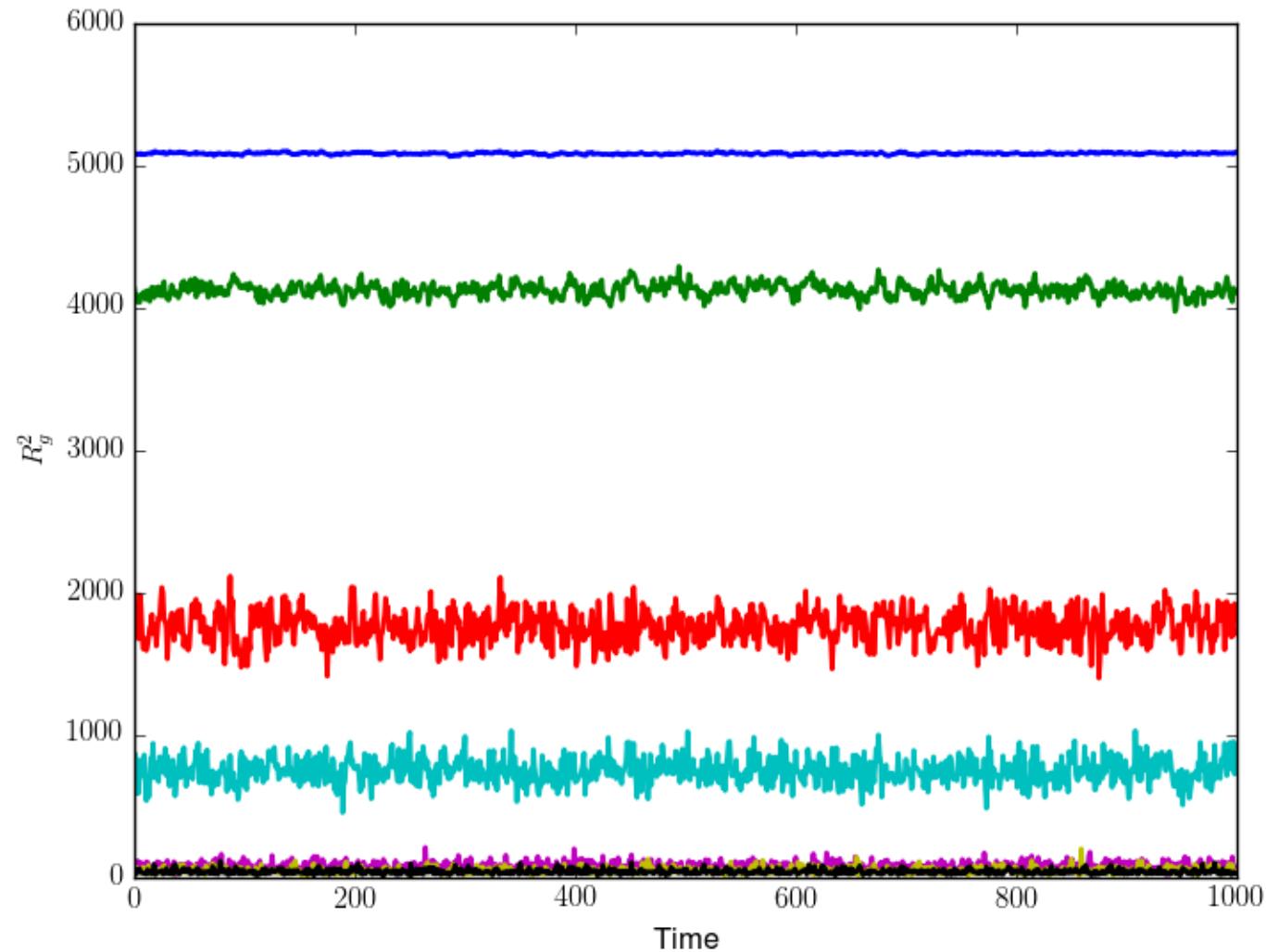
Results ( $T=10000$ ,  $N = 50$ )



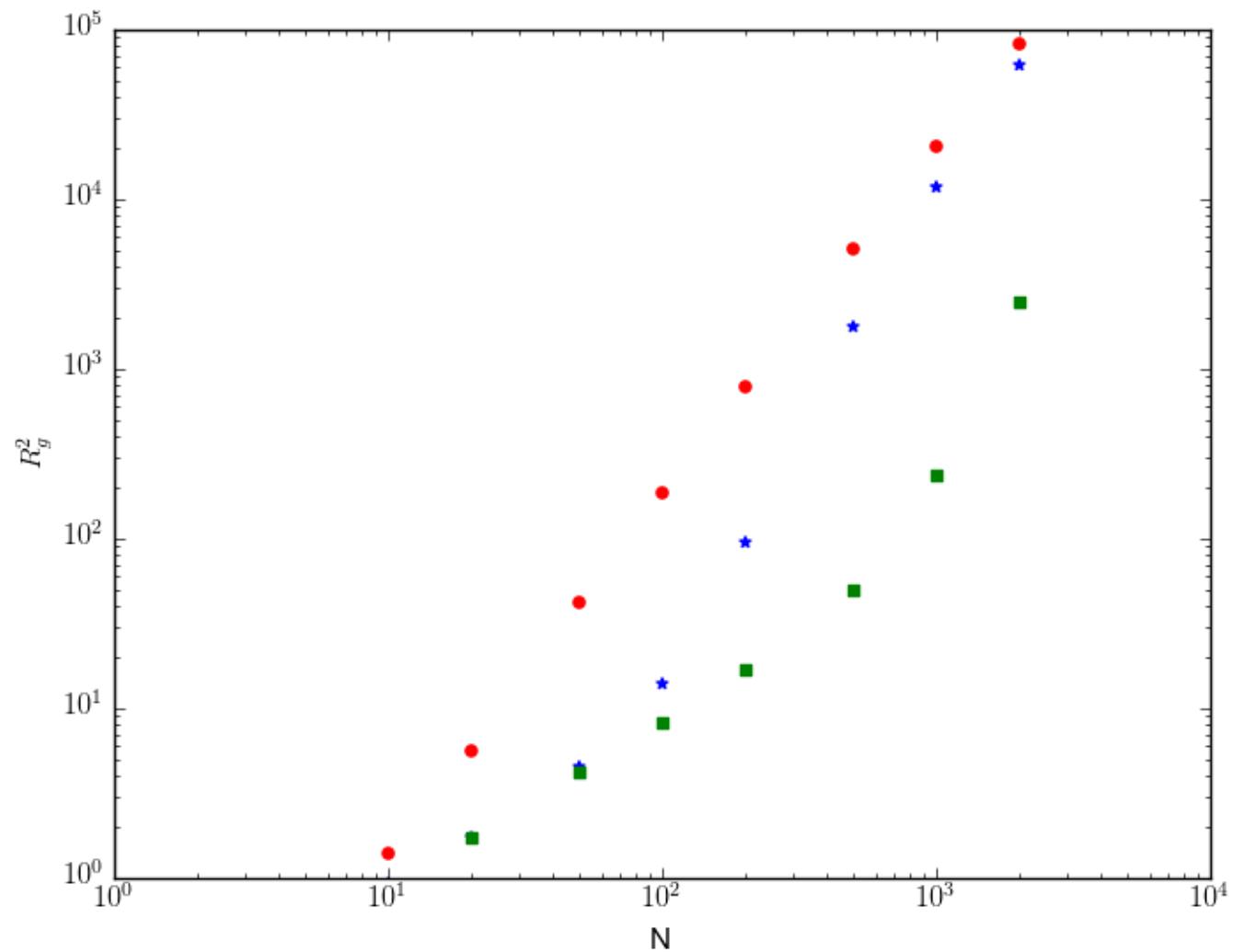
# Results (N=500)



# Results (N=500)

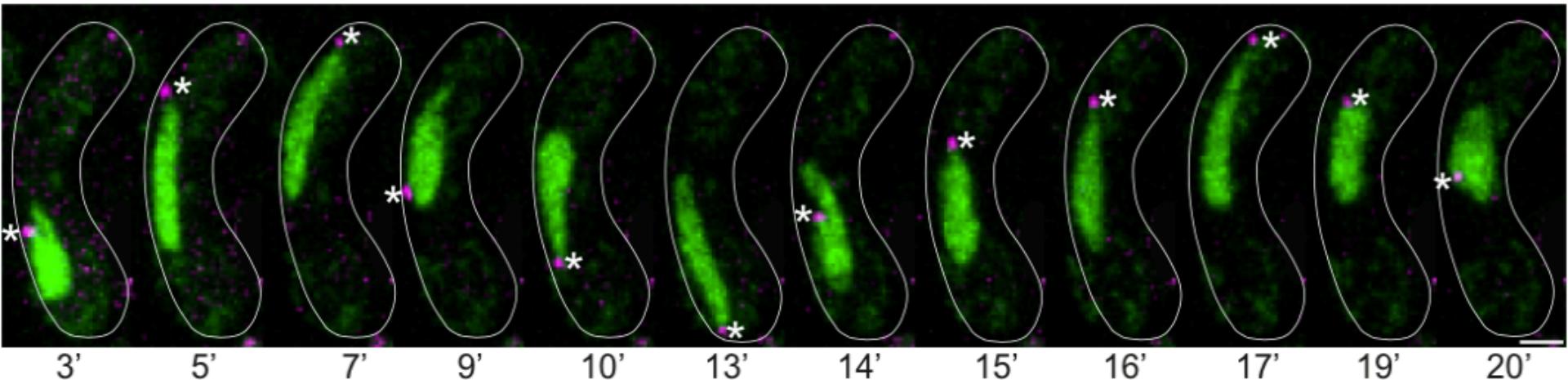


# Results (r: T=1, b: T=50, g: T=1000)



# TODO

- Gyration tensor: both numerically and analytically
- Multiple polymer rings ---- Lennard-Jones interaction?
- Shape at turning point  Tumbling dynamics  
?



# Tumbling dynamics

Philipp S. Lang et al. Phys. Rev. E 89, 022606 (2014)



All suggestions and comments  
are greatly appreciated!

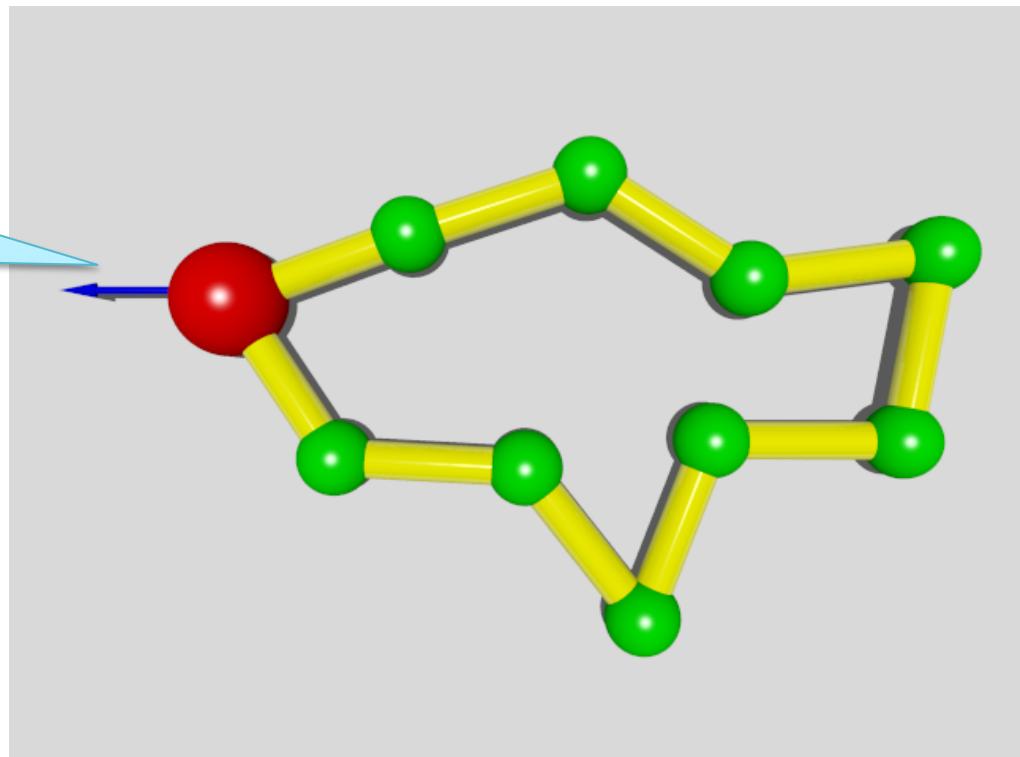
Thank you!

# Bead-Rod Ring Model

A Chromosome is represented by a bead-rod ring with the driven bead be the **SPB**(spindle pole body).

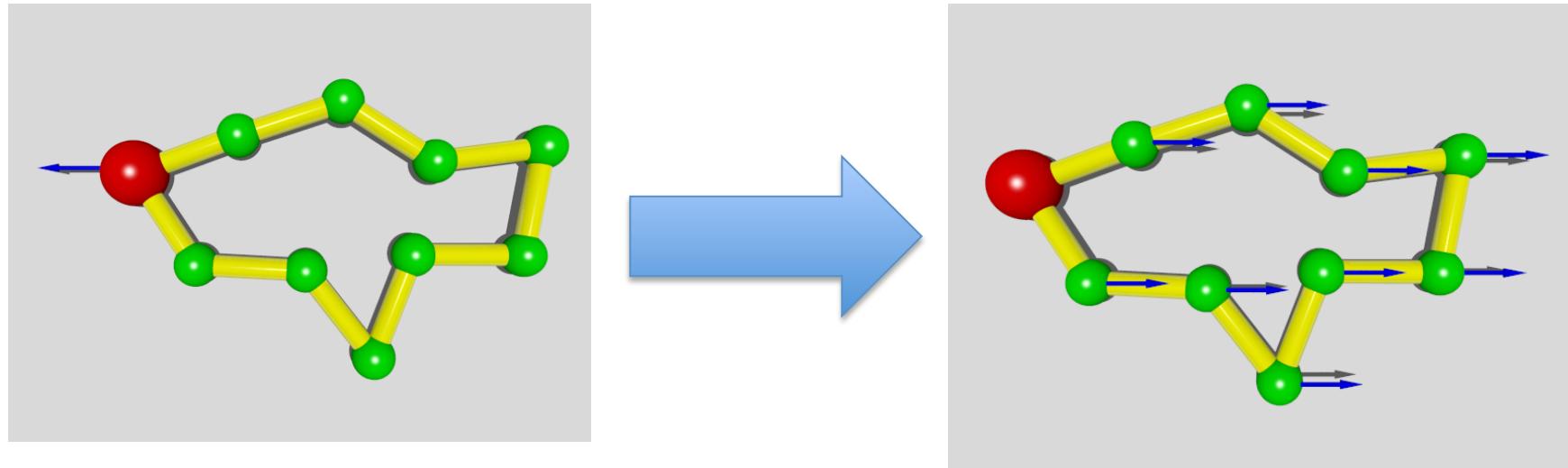
Ring driven by a  
**Constant force**

- Bead:  
position ----  $\mathbf{r}_i$
- Rod:  
length ---  $a$



# Theory for 1D model

Transfer to co-moving frame:



Constant driven force on SPB

Constant force field and pinned SPB

For 1D system:

$$H = E_0 - V_0 \sum_{j=1}^N j \Delta x_j \xrightarrow{z_j = (\Delta x_j + 1)/2} H = E'_0 - 2V_0 \sum_{j=1}^N j n_j$$

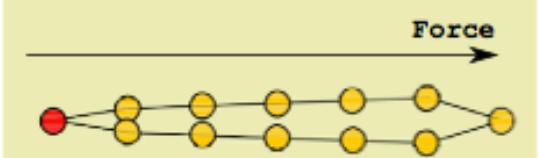
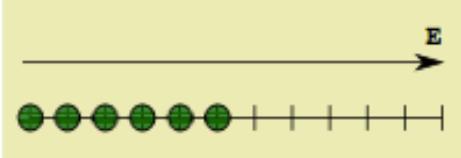
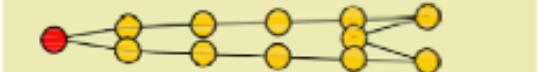
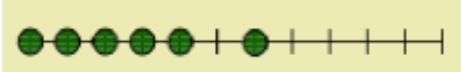
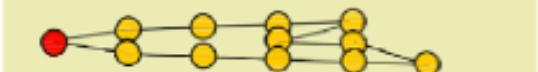
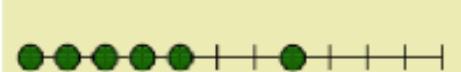
Loop Constraint:

$$\sum_{j=1}^N \Delta x_j = 0 \rightarrow \sum_{j=1}^N n_j = \frac{N}{2}$$

# Theory for 1D model

Hamiltonian:  $H = E_0' - 2V_0 \sum_{j=1}^N j n_j$

$n_j \in \{0,1\}$   Equivalent to  $N/2$  Fermions in  $N$  energy levels

	Polymer Loop	Rod config	Corresponding Fermi
Ground state		RRRRRRRLLLLLL	
First Excitation		RRRRRLRLLLLL	
Second Excitation	 	RRRRLRRRLLLL RRRRRLLRLLLL	 

# Theory for 1D model

- Partition function

$$P \propto \exp\left[-\frac{1}{T_{eff}} \sum_{j=1}^N j n_j\right]$$

By setting natural temperature unit  $2V_0 / k_B \equiv 1$

- Fermi-Dirac distribution

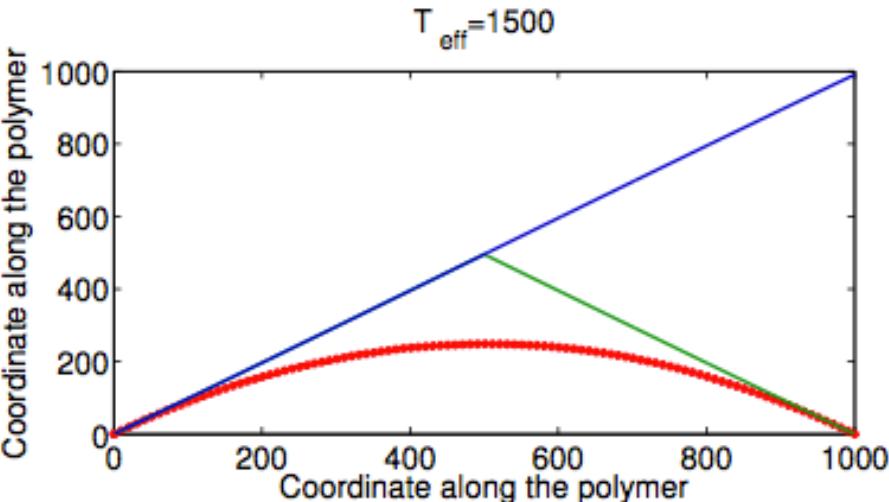
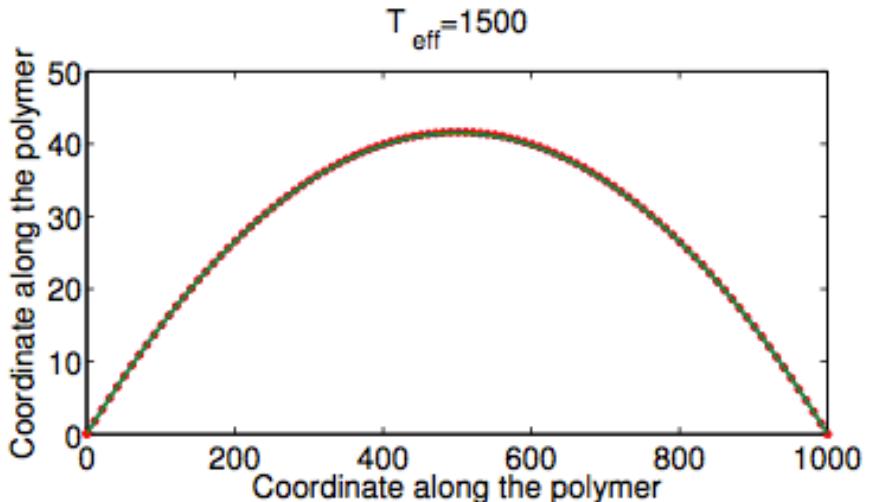
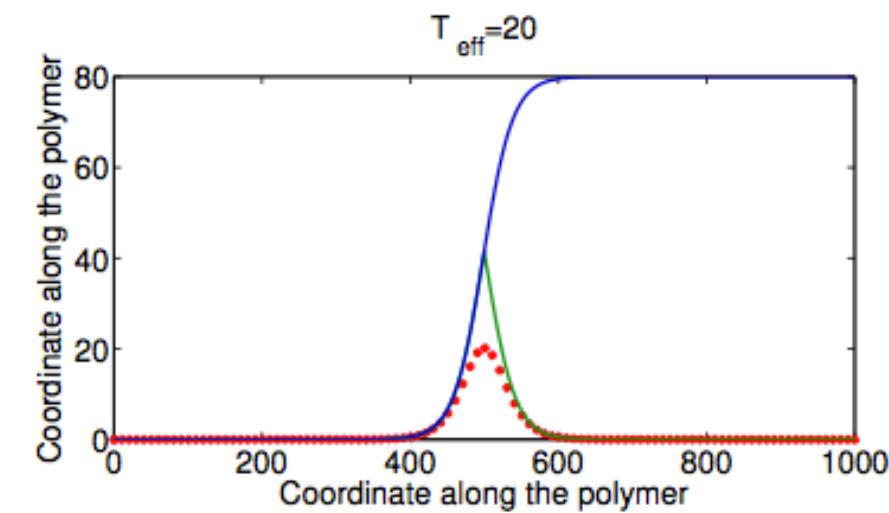
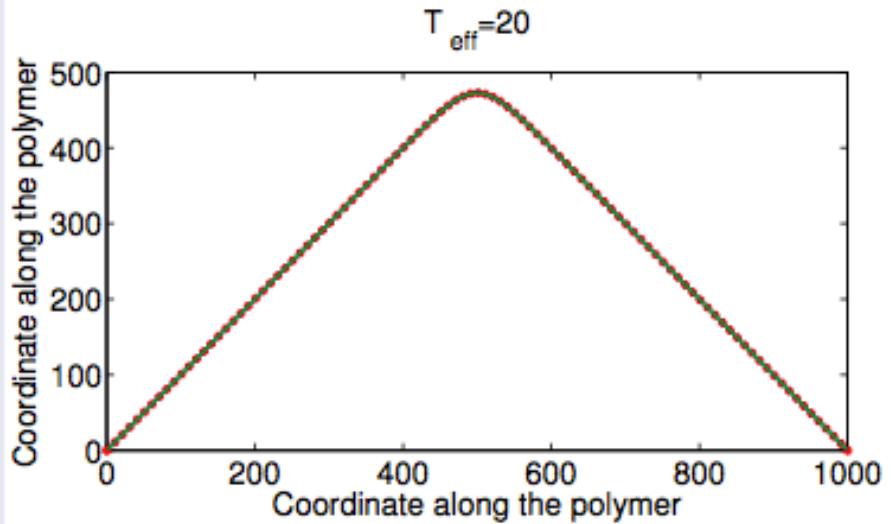
$$P\{n_j = 1\} = \frac{1}{1 + e^{j - \frac{N}{2}}}$$

- Statistical Results

$$\langle x_i \rangle = \sum_{j=1}^i \langle \Delta x_j \rangle = 2 \sum_{j=1}^i \langle n_j \rangle - i = 2T_{eff} \ln \frac{1 + e^{N/2T_{eff}}}{e^{i/2T_{eff}} + e^{(N-i)/2T_{eff}}}$$

$$\text{var}[x_i] = 4 \sum_{j=1}^N \text{var}[n_j]$$

# 1D Theory



# Brownian Bridge Approach

--- devise an effective stochastic process satisfies  $Z_N = Z_0$

Step 1: construct effective process:

- $Z_0 = 0$ ;
- $Z_1$  is almost surely continuous;
- $Z_1$  has independent increments;
- $Z_n - Z_m = N(\bar{z}_{m \rightarrow n}, \bar{\sigma}_{m \rightarrow n}^2)$

Where the mean and the variance is determined by:

$$\begin{aligned}\mathbb{E}[z_p - z_q] &= \sum_{j=p+1}^q \mathbb{P}\{r_i = 1\} - \mathbb{P}\{r_i = -1\} \\ &= \sum_{j=p+1}^q \left[ \frac{2}{1 + e^{\frac{j-\mu}{k_B T}}} - 1 \right],\end{aligned}$$

$$\begin{aligned}\text{var}[z_p - z_q] &= 4 \sum_{j=m+1}^n \text{var}[r_i] \\ &= 4 \sum_{j=p+1}^q \frac{e^{\frac{j-\mu}{k_B T}}}{\left(1 + e^{\frac{j-\mu}{k_B T}}\right)^2}.\end{aligned}$$

$$\begin{aligned}\bar{z}_{m \rightarrow n} &:= k_B T \int_m^n \left[ \frac{2}{1 + e^{\zeta - \frac{N}{2k_B T}}} - 1 \right] d\zeta \\ \bar{\sigma}_{m \rightarrow n; z}^2 &= 4k_B T \int_m^n \frac{e^{\zeta - \frac{N}{2k_B T}}}{\left(1 + e^{\zeta - \frac{N}{2k_B T}}\right)^2} d\zeta\end{aligned}$$

# Brownian Bridge Approach

--- devise an effective stochastic process satisfies  $Z_N = Z_0$

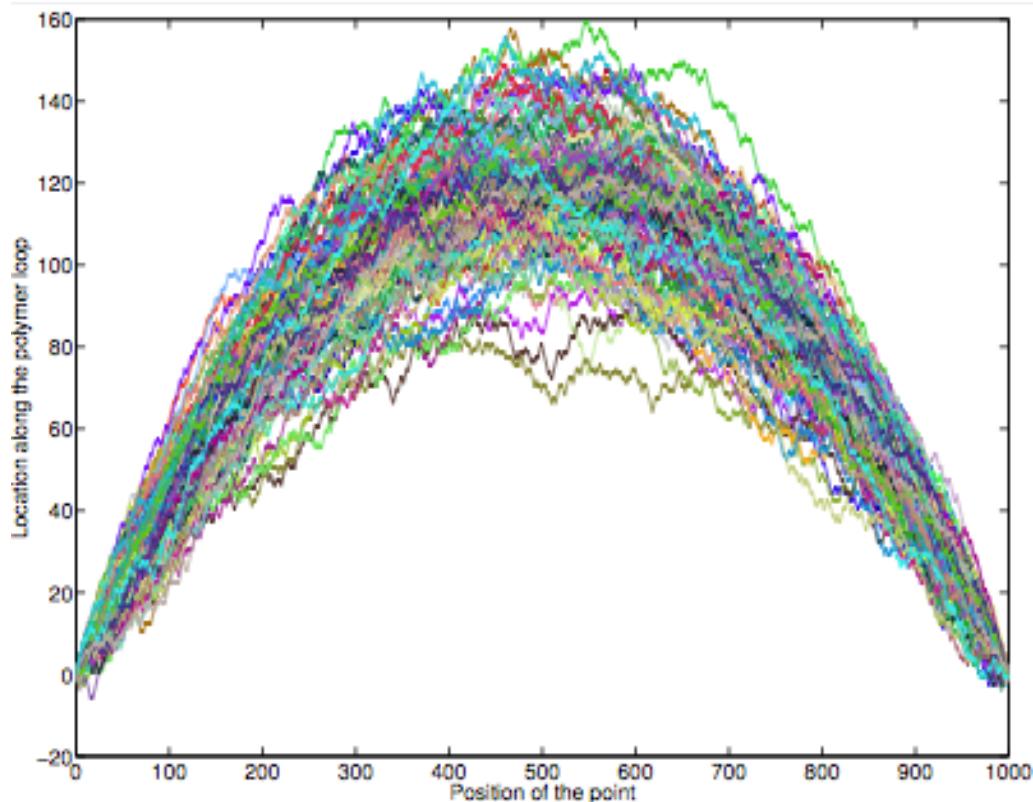
Step 2: introduce loop condition

Construct conditional process:

$$Z_l^0 = (Z_l \mid Z_l = 0), l \in (0, N)$$

where

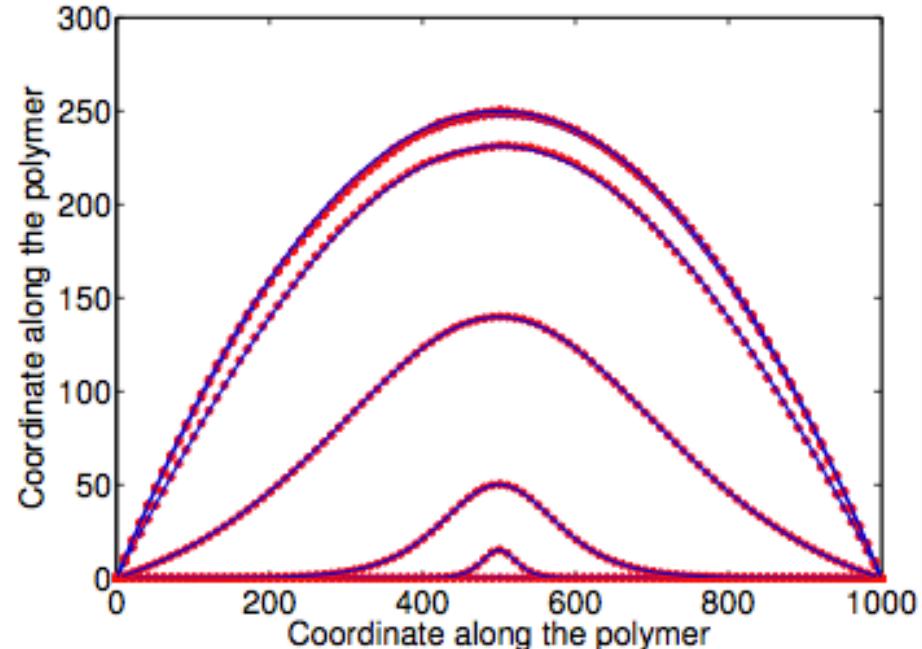
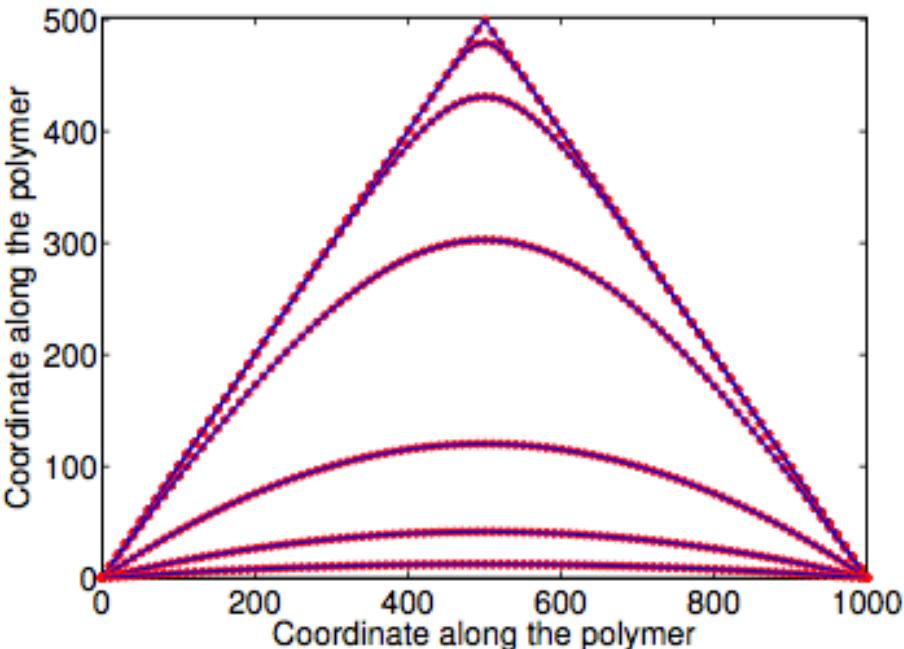
$$Z_l^0 = Z_l - \frac{\bar{\sigma}_{0 \rightarrow l; z}^2}{\bar{\sigma}_{0 \rightarrow N; z}^2} Z_N$$



# Numerical Verification of 1D model

Our analysis yields the following conclusion:

$$\mathbb{E}[x_i] = 2T_{\text{eff}} \ln \frac{1 + e^{N/T_{\text{eff}}}}{e^{i/2T_{\text{eff}}} + e^{(N-i)/2T_{\text{eff}}}}, \quad \text{var}[x_i] = \frac{\sigma_{0i}^2 \sigma_{iN}^2}{\sigma_{0N}^2}$$



# 3D Continuous Model

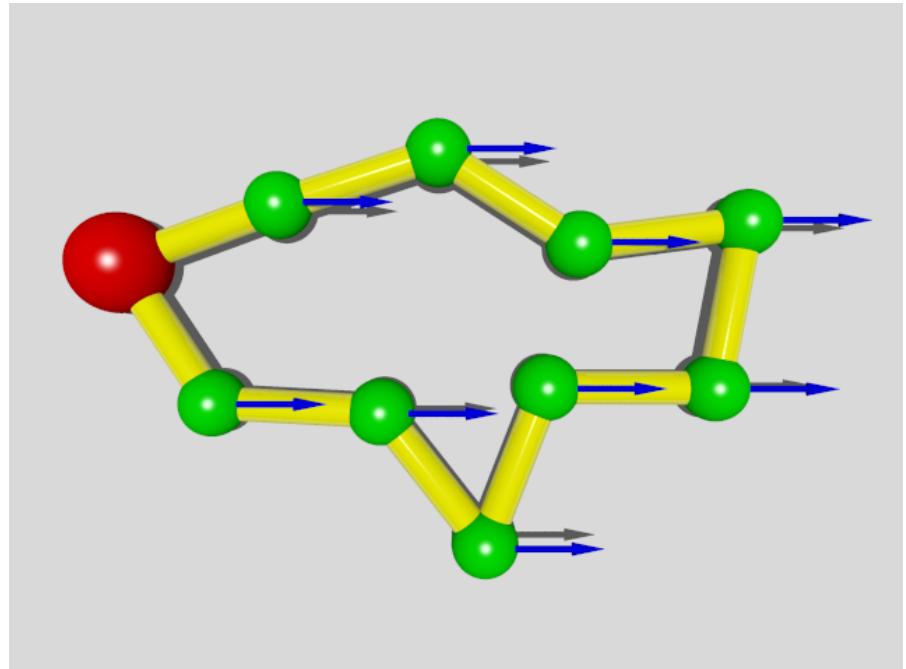
- Hamiltonian

$$H = E_0 - V_0 \sum_{i=1}^N \sum_{j=1}^i \cos \theta_j$$

Heisenberg model

- Partition function

$$\begin{aligned} \mathcal{P}_C &= \prod_{j=1}^N \mathcal{P}_j \\ &= \prod_{j=1}^N \int_0^{2\pi} d\phi_j \int_0^\pi \sin \theta_j \, d\theta_j e^{-\frac{-(j-\mu) \cos \theta_j}{k_B T}} \\ &= \frac{4\pi k_B T \sinh \frac{\mu-j}{k_B T}}{\mu-j} \end{aligned}$$



$$\begin{aligned} \mathbb{E} [\cos \theta_j] &= k_B T \partial_\mu \mathcal{P}_j \\ &= \coth \frac{\mu-j}{k_B T} - \frac{k_B T}{\mu-j}, \\ \text{var} [\cos \theta_j] &= (k_B T \partial_\mu)^2 \mathcal{P}_j \\ &= \left( \frac{k_B T}{\mu-j} \right)^2 - \operatorname{csch}^2 \frac{\mu-j}{k_B T}. \end{aligned}$$

# Adopt Brownian Bridge Approach

Define Effective 3D stochastic process:

- $W_0 = 0,$
- $W_l$  is almost surely continuous,
- $W_l$  has independent increments,
- $W_n - W_m \sim \mathcal{N}(\bar{w}_{m \rightarrow n}, \bar{\sigma}_{m \rightarrow n; w}^2), 0 < n \leq m < N,$

$$\bar{x}_{m \rightarrow n} = 0,$$

$$\bar{y}_{m \rightarrow n} = 0,$$

$$\bar{x}_{m \rightarrow n} = k_B T \left[ \log \frac{l - \frac{N}{2}}{k_B T \sinh \left( \frac{l - \frac{N}{2}}{k_B T} \right)} \right]_{l=m}^n$$

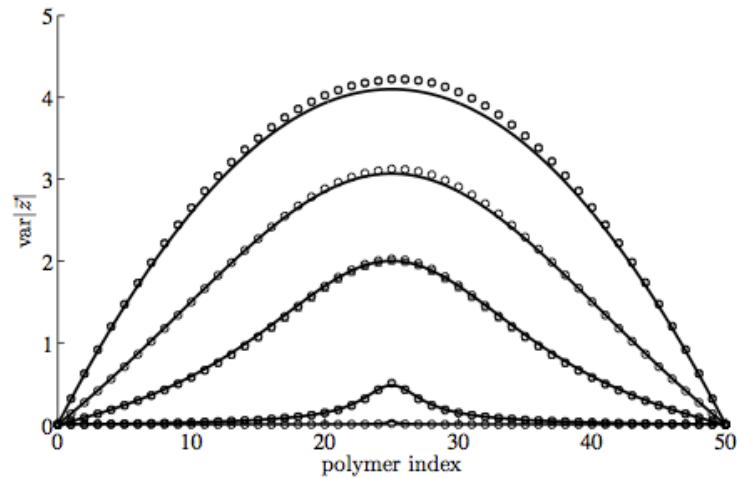
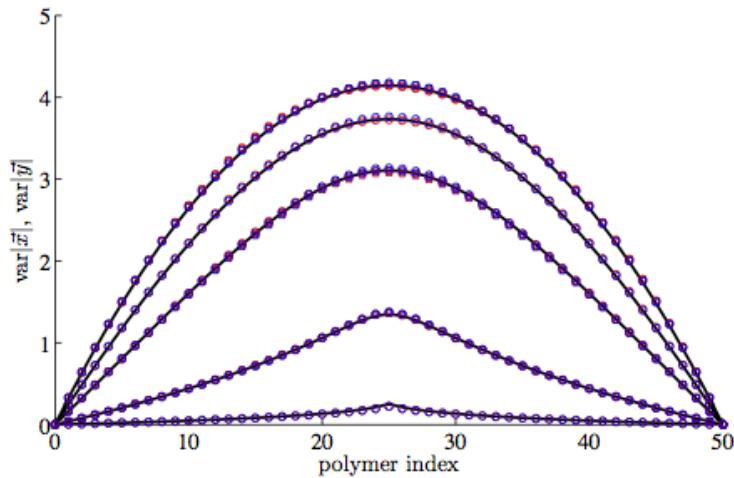
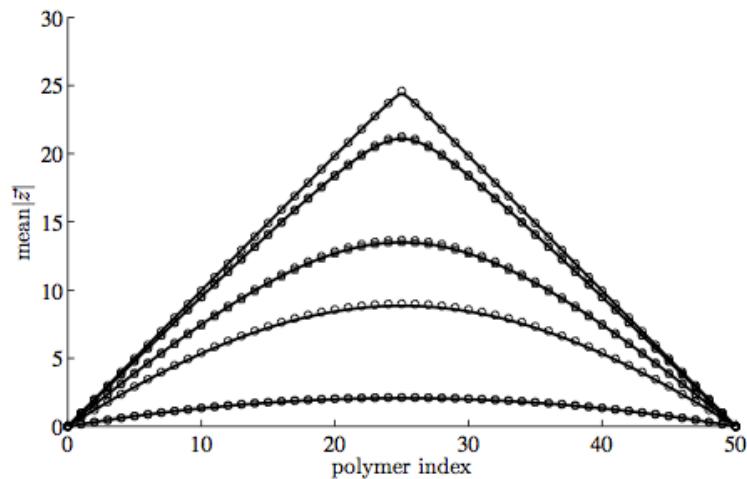
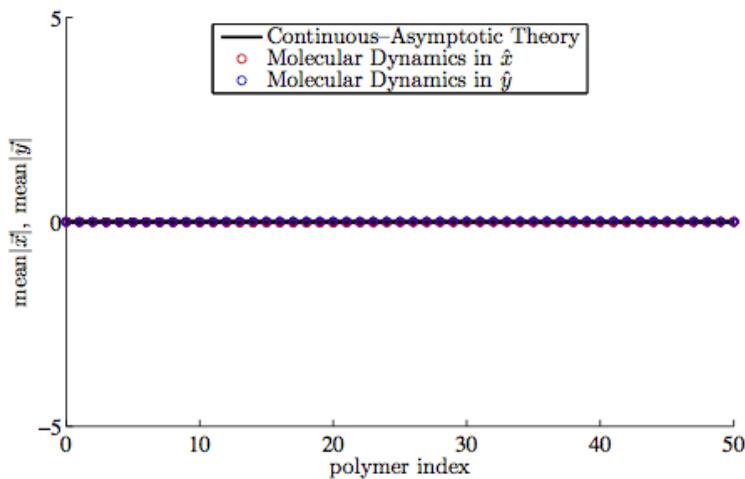
$$\bar{\sigma}_{m \rightarrow n; x}^2 = k_B T \int_m^n \text{var} [\vec{s}_{k_B T \zeta}] \cdot \hat{x} d\zeta,$$

$$\bar{\sigma}_{m \rightarrow n; y}^2 = k_B T \int_m^n \text{var} [\vec{s}_{k_B T \zeta}] \cdot \hat{y} d\zeta,$$

$$\bar{\sigma}_{m \rightarrow n; z}^2 = k_B T \left[ \coth \frac{l - \frac{N}{2}}{k_B T} - \frac{1}{l - \frac{N}{2}} \right]_{l=m}^n$$

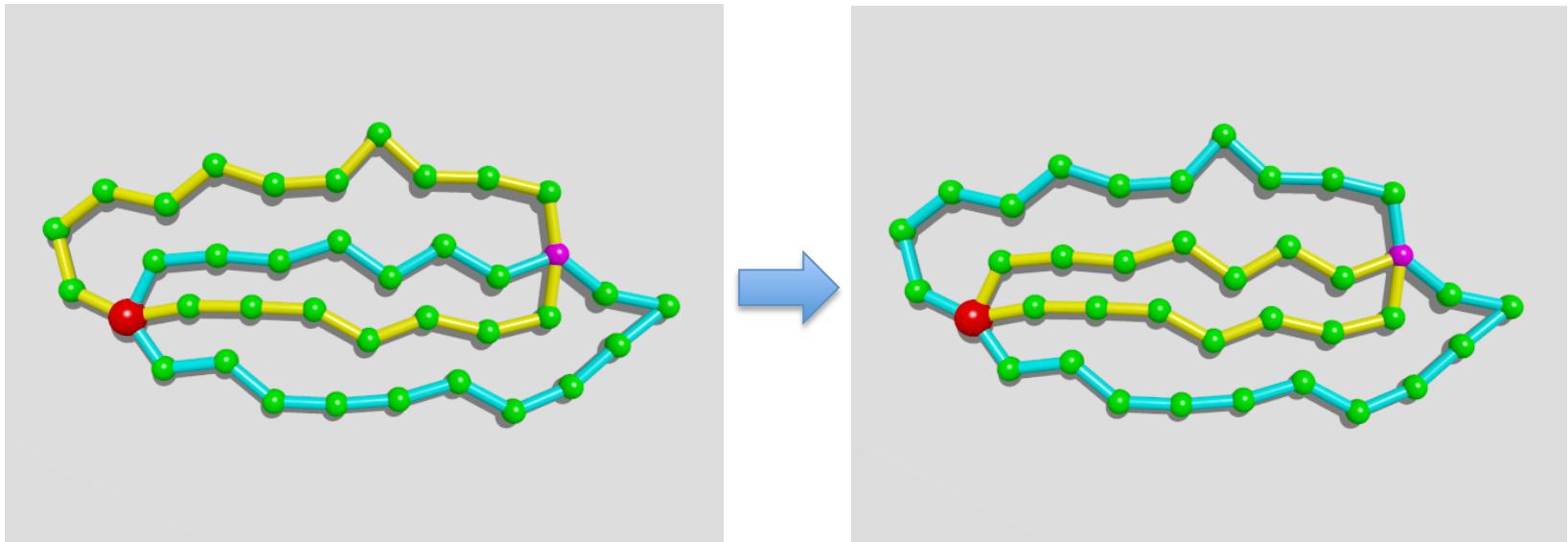
Conditional effective process:  $W_l^0 = (W_l | W_l = 0), l \in (0, N)$

# Numerical Verification of 3D model



# A Pair of Rings with Centromere

- Re-label the loops



- Brownian Bridge

$$\bar{w}_{m \rightarrow n; \mathcal{C}_1} := \bar{w}_{0 \rightarrow n; \mathcal{C}_1} - \bar{w}_{0 \rightarrow m; \mathcal{C}_1}$$
$$\bar{\sigma}_r^2 := \bar{\sigma}_n^2 - \bar{\sigma}_m^2$$

- $W_{0; \mathcal{C}_1} = 0,$

$\bar{w}_{l; \mathcal{C}_1}^0 \equiv W_{l; \mathcal{C}_1} - \frac{\bar{\sigma}_{0 \rightarrow l; w, \mathcal{C}_1}^2}{\bar{\sigma}_{0 \rightarrow N; w, \mathcal{C}_1}^2} W_{N; \mathcal{C}_1}$

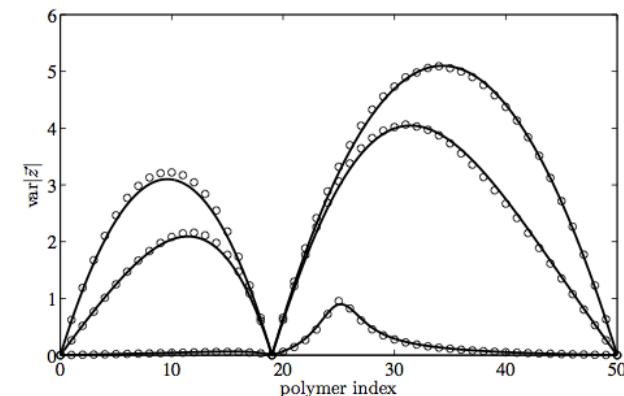
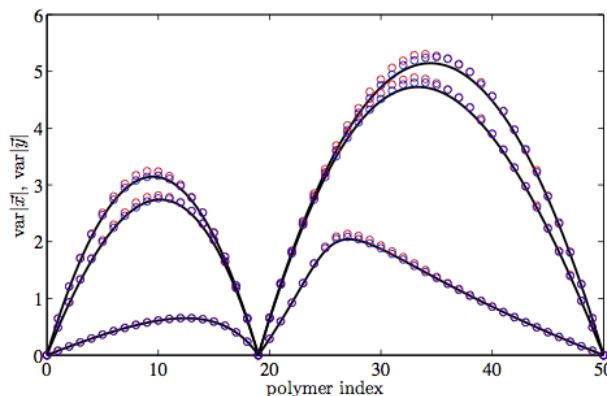
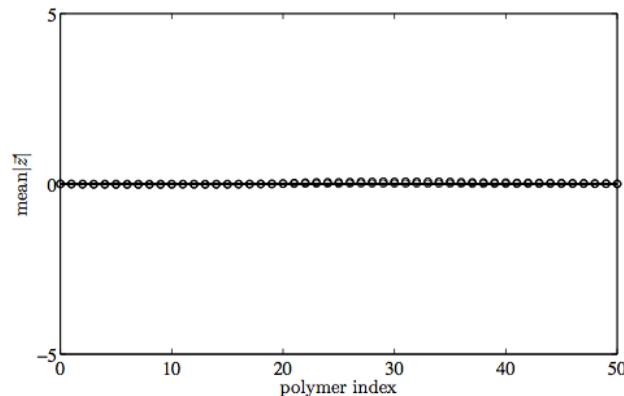
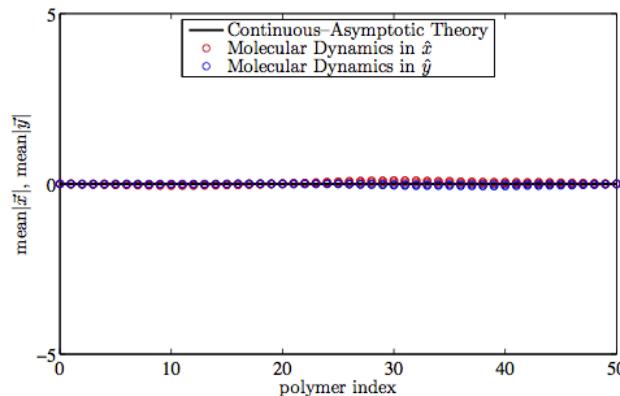
Bridge Condition:  $W_{l; \mathcal{C}_1}^0 = W_{l; \mathcal{C}_1} - \frac{\bar{\sigma}_{0 \rightarrow l; w, \mathcal{C}_1}^2}{\bar{\sigma}_{0 \rightarrow N; w, \mathcal{C}_1}^2} W_{N; \mathcal{C}_1}$

$\sigma_0^2 < \sigma_n^2 \leq \sigma_m^2 < \sigma_r^2$

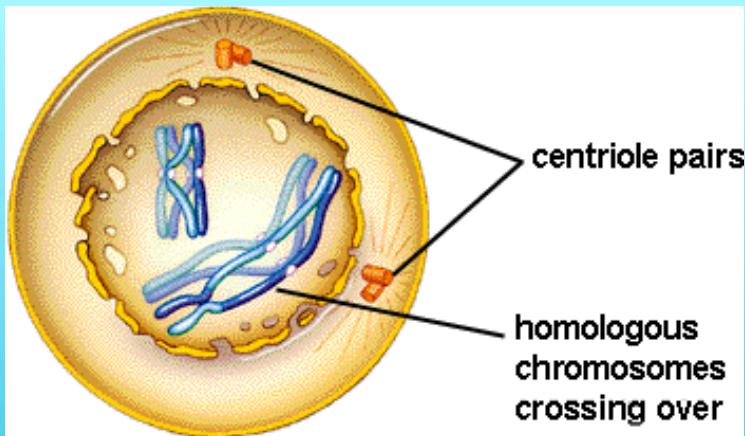
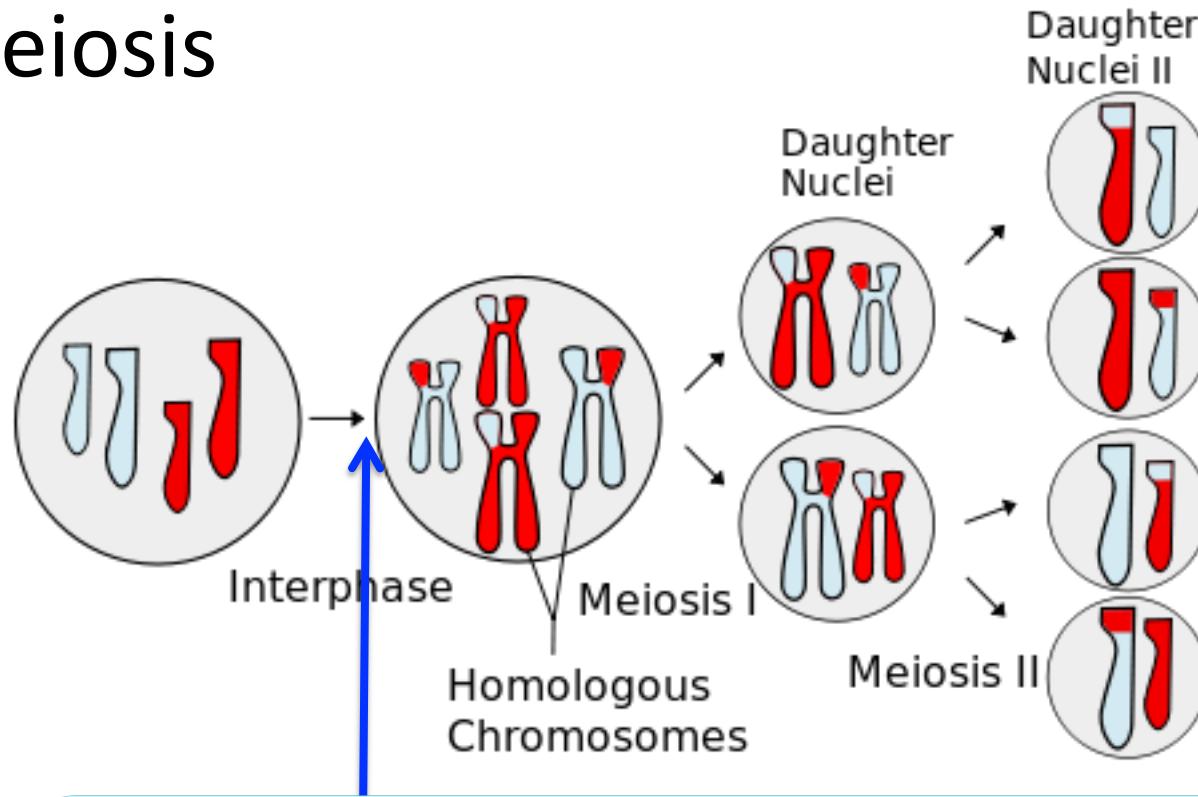
# Numerical Verification

- Theoretical Result

$$\text{var} [W_{\mu;C_1}^0 - W_{\mu';C_1}^0] = \mathbb{E} \left[ (W_{\mu;C_1}^0 - W_{\mu';C_1}^0)^2 \right] = \frac{2\bar{\sigma}_{0 \rightarrow \mu; w}^2 \bar{\sigma}_{\mu \rightarrow \lambda; w}^2}{\bar{\sigma}_{0 \rightarrow \lambda; w}^2}$$



# Meiosis



## Prophase I

- Chromosome movements
- Pairing
- Recombination

# Brownian Dynamics:

Inertia-less assumption

$$\mathbf{F}_i = \mathbf{0}$$

$$\mathbf{F}_i = \mathbf{F}_i^h + \mathbf{F}_i^\phi + \mathbf{F}_i^c + \mathbf{F}_i^b + \mathbf{F}_i^e$$

$i$  is the index of bead

$\mathbf{F}_i^h$  is hydrodynamic force

$\mathbf{F}_i^\phi$  represents force generate from potential

$\mathbf{F}_i^c$  is constraint force to keep the rod length

$\mathbf{F}_i^b$  is brownian force

$\mathbf{F}_i^e$  is external force

# Brownian Dynamics:

$$\mathbf{F}_i = \mathbf{F}_i^h + \mathbf{F}_i^\phi + \mathbf{F}_i^c + \mathbf{F}_i^b + \mathbf{F}_i^e = 0$$

$$\mathbf{F}_i^h = -\zeta \dot{\mathbf{r}}_i$$

$$\mathbf{F}_i^\phi = -\nabla U(\mathbf{r})$$

$$\mathbf{F}_i^c = T_i \mathbf{u}_i - T_{i-1} \mathbf{u}_{i-1}; \mathbf{u}_i = (\mathbf{r}_{i+1} - \mathbf{r}_i) / a$$

$$\langle \mathbf{F}_i^b(t) \rangle = \mathbf{0}; \langle \mathbf{F}_i^b(t) \mathbf{F}_j^b(t + \Delta t) \rangle = 2k_B T \zeta \delta_{ij} \delta(\Delta t)$$

$$\mathbf{F}_i^e = f(\mathbf{r}_i, t)$$

➤ Dynamical differential equation:

$$\frac{d\mathbf{r}_i}{dt} = \zeta^{-1} (\mathbf{F}_i^\phi + \mathbf{F}_i^c + \mathbf{F}_i^b + \mathbf{F}_i^e)$$

# Numerical Scheme

- Predictor-corrector algorithm:

Step 1: predict using known forces

$$\mathbf{r}_i^*(t + \Delta t) = \mathbf{r}_i(t) + \zeta^{-1}(\mathbf{F}_i^\phi + \mathbf{F}_i^b + \mathbf{F}_i^e)\Delta t$$

Step 2: correct using constraint force

$$\mathbf{r}_i(t + \Delta t) = \mathbf{r}_i^*(t + \Delta t) + \zeta^{-1}\mathbf{F}_i^c\Delta t \quad (*)$$

Step 3: substitute eq. (\*) to constraint equations

$$(\mathbf{r}_{i+1} - \mathbf{r}_i)^2 - a^2 = 0$$

Solve a set of nonlinear algebraic equations   $\mathbf{F}_i^c$

Step 4: re-substitute  $\mathbf{F}_i^c$  to (\*) obtain the final  $\mathbf{r}_i(t + \Delta t)$

# Parameter Estimation

## ➤ Length of Chromosomes in base pairs:

There are **3 pairs** of chromosomes in prophase fission yeast

Chromosome I: 5.579.133 bp ~ 5,6Mbp

Chromosome II: 4.539.804 bp ~ 4,5Mbp

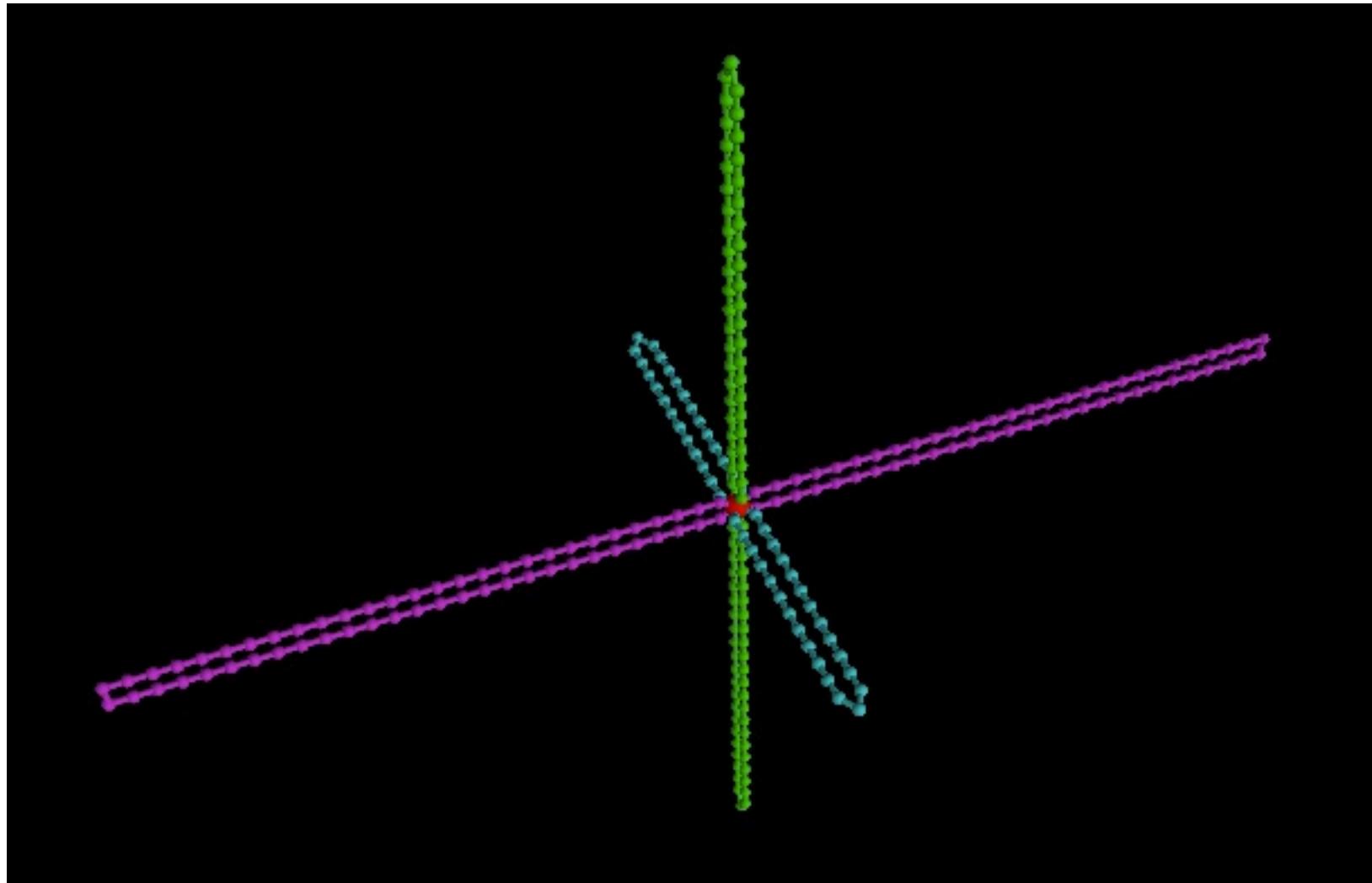
Chromosome III: 2.452.833 bp ~ 2,5Mbp

## ➤ Compaction ratio of chromosomes: ~ 100bp/nm

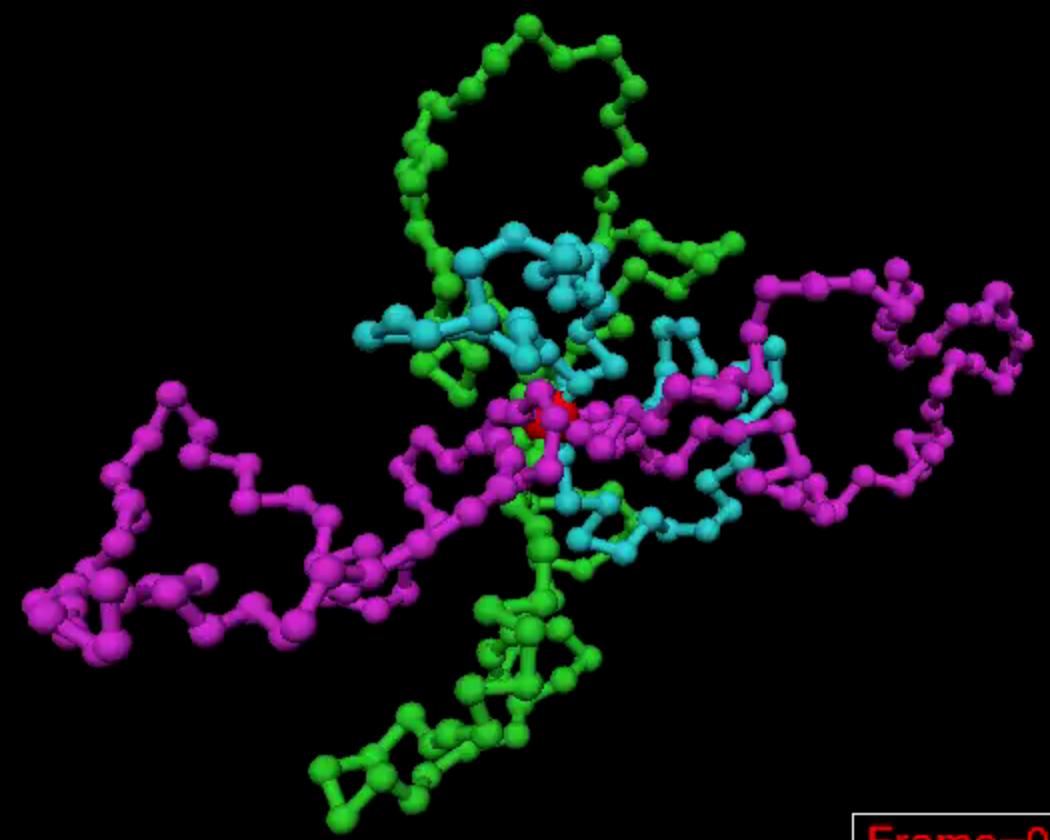
## ➤ Kuhn length(rod length): ~ 100nm

## ➤ System size for 3 pairs of chromosomes: ~ 2000-3000 monomers in one ring : 200 ~ 600

# Initial configuration

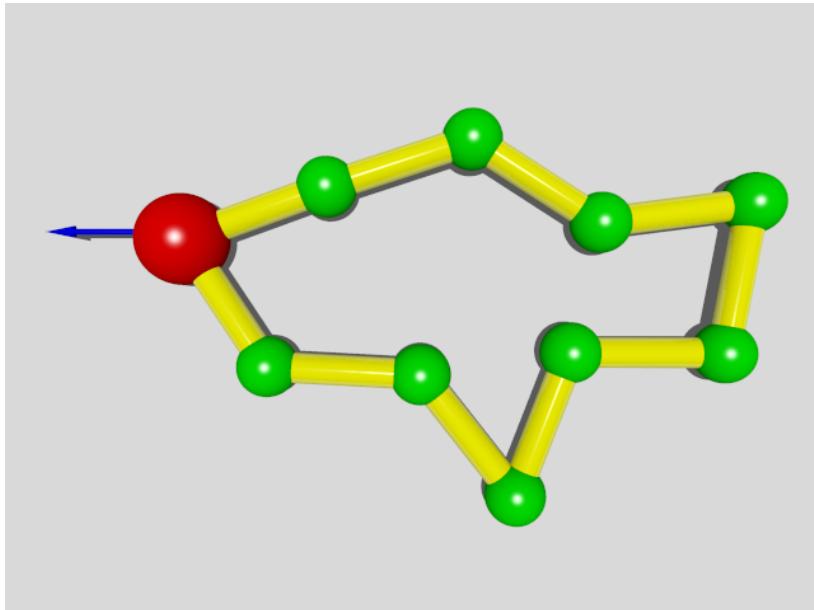


# Animation Movie

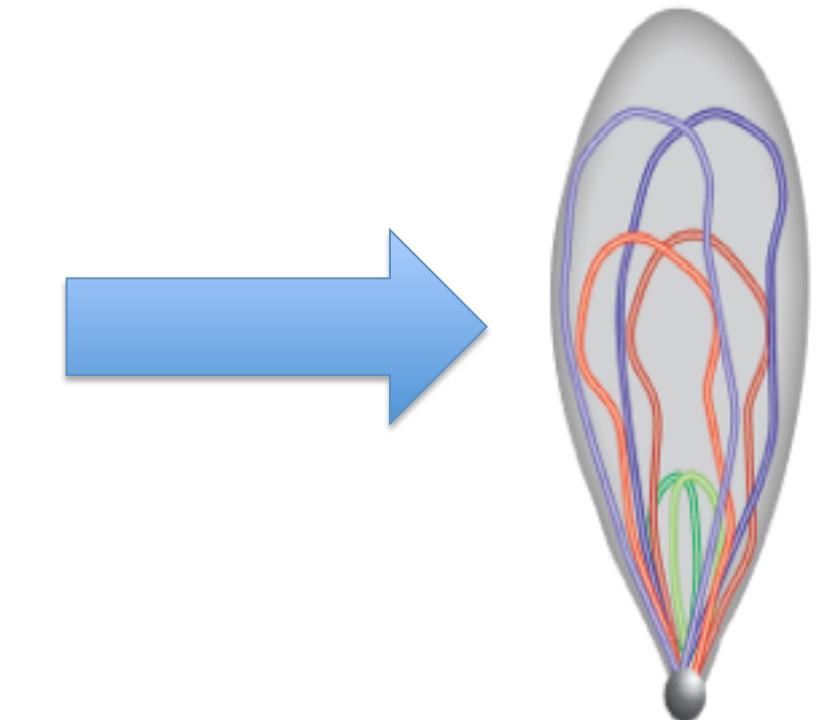


Frame-0

# Problem



Toy model with  
100 beads



Real System

Solution: Parallel computation

# Ways to go

- Use some kinds of sophisticated MD packages

**Advantage:**

easier to go, parallelize automatically

**Disadvantage:**

cannot control some details, e.g. random force

not so flexible

- Write a parallel program

Totally under control but need more time on technique details

## My Strategy

Try MD packages

Doesn't work

Write a parallel program

# List of MD packages

[http://en.wikipedia.org/wiki/  
List\\_of\\_software\\_for\\_molecular\\_mechanics\\_modeling](http://en.wikipedia.org/wiki/List_of_software_for_molecular_mechanics_modeling)

## Rules of selection

- No commercial, should be free
- Need GPU acceleration
- Widely used

## Selected packages:

- [Gromacs](#)
- [LAMMPS](#)
- [ESPResSo](#)