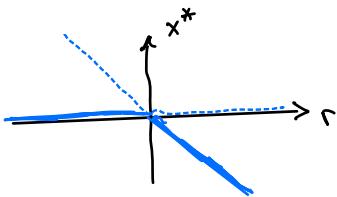


13

1. (a) transcritical bifurcation diagram
 (b)



4. (b) • Dimples on a golf ball reduce the drag force
 2. (or increase the range, also OK)

- They do this by allowing turbulent boundary layer to form for smaller R (smaller velocity)

2. (c) $\dot{x} = \mu x + x^3 - x^5$, varying μ

5

$$H = H_0 + \Delta H.$$

- Assume H_0 exhibits periodic behavior with period τ
- ΔH causes constant variables under H_0 to change in time
- variable that returns to initial value after $t = \tau$ is periodic
- variable that shifts after $t = \tau$ is secular

or

e.g. H_0 = Harmonic Osc with $\Delta H \propto x^4$. Here amplitude J is periodic & phase shift β is secular.

2. 16

- (a)
 (b)

Variables:	r	v	f	ϵ	n	g
dimension:	m	m/s	$1/s$	kg/m^3	$\frac{kg}{m^3 s^2} = \frac{kg}{m \cdot s}$	m/s^2

$$\epsilon^0 = n$$

$6-3 = 3$ dimensionless ratios

2 are arguments of F : $f = \frac{v}{r} F\left[\frac{vr\epsilon}{n}, \frac{gr}{v^2}\right]$

$$r \rightarrow \frac{r}{\sqrt{10}}, v \rightarrow \frac{v}{\sqrt{10}}, \epsilon \rightarrow 10\sqrt{10}\epsilon \quad \therefore f \rightarrow \sqrt{10} f \quad \epsilon_R$$

(b) $H = \frac{p^2}{2m} + k|x| = E$ turning points $x = \pm E/k \equiv \pm x_0$ -2-

(8) $J = \oint p dx = 2 \int_{-x_0}^{+x_0} dx \sqrt{2m} \sqrt{E - k|x|}$

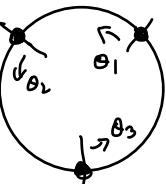
$$= 4 \int_0^{x_0} dx \sqrt{2m} (E - kx)^{1/2} = -\frac{8}{3k} \sqrt{2m} (E - kx)^{3/2} \Big|_0^{x_0} = \frac{8\sqrt{2m}}{3k} E^{3/2}$$

$$\therefore E = \left(\frac{3k}{8\sqrt{2m}} \right)^{2/3} J^{2/3}$$

$$\omega = \frac{\partial E}{\partial J} = \left(\frac{3k}{8\sqrt{2m}} \right)^{2/3} \frac{2}{3} J^{-1/3}$$

$$= \frac{3k}{8\sqrt{2m}} \cancel{\frac{2}{3}} \cancel{J^{-1/2}} = \frac{k}{4\sqrt{2m}} E^{-1/2}$$

$\delta_0 \quad \gamma = \frac{1}{\omega} = \frac{4\sqrt{2m}E}{k}$

3 [30] 

 $L = T - V$

(a)

 $T = \frac{m}{2} R^2 (\dot{\theta}_1^2 + \dot{\theta}_2^2) + \frac{(2m)}{2} R^2 \dot{\theta}_3^2 \quad 3.$

(8)

displacements small angle, springs stretched by $s_{ij} = R(\theta_j - \theta_i)$

 $V = \frac{k}{2} (s_{12}^2 + s_{23}^2 + s_{31}^2) = \frac{kR^2}{2} [(\theta_1 - \theta_2)^2 + (\theta_2 - \theta_3)^2 + (\theta_3 - \theta_1)^2]$

5. $= \frac{kR^2}{2} [2\theta_1^2 + 2\theta_2^2 + 2\theta_3^2 - 2\theta_1\theta_2 - 2\theta_1\theta_3 - 2\theta_2\theta_3]$

(b) $L = \frac{1}{2} \vec{\eta}^\top \hat{T} \cdot \hat{T} \cdot \vec{\eta} - \frac{1}{2} \vec{\eta}^\top \hat{V} \cdot \vec{\eta} \quad \text{with } \vec{\eta} = \begin{pmatrix} \theta_1 \\ \theta_2 \\ \theta_3 \end{pmatrix}$

(8) $\hat{T} = mR \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 2 \end{pmatrix}, \quad \hat{V} = kR^2 \begin{pmatrix} 2 & -1 & -1 \\ -1 & 2 & -1 \\ -1 & -1 & 2 \end{pmatrix}$

$\vec{\eta} = \vec{a} e^{-i\omega t} \quad (\hat{V} - \omega^2 \hat{T}) \cdot \vec{a} = 0 = \left(\frac{\hat{V}}{kR^2} - \frac{\omega^2}{kR^2} \hat{T} \right)$

Let $\omega_0^2 = \frac{k}{m}$

$$\left(\begin{array}{ccc} 2 - \omega^2/\omega_0^2 & -1 & -1 \\ -1 & 2 - \omega^2/\omega_0^2 & -1 \\ -1 & -1 & 2 - 2\omega^2/\omega_0^2 \end{array} \right) \cdot \vec{a} = 0$$

$\underbrace{\quad}_{\equiv \hat{B}}$

(c) 10 pts solution & 4 pts pictures

Either Physical Intuition + bit math or Math

Intuition

- $\theta_1 = \theta_2 = \theta_3$, everyone rotates together

$$\vec{a} = \begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix}$$

← no need to $-^3-$
normalize,
but could

3. $\hat{B} \begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix} = \left(-\frac{\omega^2}{\omega_0^2} \right) \begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix}$ so take $\boxed{\omega=0}$

- should have solution with $\theta_3 = 0$, $\theta_1 = -\theta_2$

3. $\vec{a} = \begin{pmatrix} 1 \\ -1 \\ 0 \end{pmatrix}$, $\hat{B} \begin{pmatrix} 1 \\ -1 \\ 0 \end{pmatrix} = \begin{pmatrix} 3 - \omega^2/\omega_0^2 \\ -3 + \omega^2/\omega_0^2 \\ 0 \end{pmatrix} = 0 \therefore \boxed{\omega = \sqrt{3} \omega_0 = \sqrt{\frac{3k}{m}}}$

- final solution, use orthogonality $\vec{a}^{(k)} \cdot \vec{a}^{(\ell)} = 0 \quad k \neq \ell$

4. $\vec{a} = \begin{pmatrix} 1 & 1 & 0 \\ 0 & 1 & 2 \end{pmatrix}$ so $\vec{a} = \begin{pmatrix} 1 \\ 1 \\ \alpha \end{pmatrix}$ orthogonal $\begin{pmatrix} 1 \\ -1 \\ 0 \end{pmatrix}$

$$(111) \begin{pmatrix} 1 \\ 1 \\ \alpha \end{pmatrix} \begin{pmatrix} 1 \\ -1 \\ 0 \end{pmatrix} = 0 = (112) \begin{pmatrix} 1 \\ 1 \\ \alpha \end{pmatrix} = 2 + 2\alpha \therefore \alpha = -1$$

$$\vec{a} = \begin{pmatrix} 1 \\ 1 \\ -1 \end{pmatrix}, \hat{B} \cdot \vec{a} = \begin{pmatrix} 2 - \omega^2/\omega_0^2 \\ 2 - \omega^2/\omega_0^2 \\ -4 + 2\omega^2/\omega_0^2 \end{pmatrix} = 0 \Rightarrow \boxed{\omega = \sqrt{2} \omega_0 = \sqrt{\frac{2k}{m}}}$$

or could guess $\begin{pmatrix} 1 \\ 1 \\ \alpha \end{pmatrix}$ & try it $\hat{B} \begin{pmatrix} 1 \\ 1 \\ \alpha \end{pmatrix} = \begin{pmatrix} 1 - \omega^2/\omega_0^2 - \alpha \\ 1 - \omega^2/\omega_0^2 - \alpha \\ -2 + 2\alpha - 2\omega^2/\omega_0^2 \alpha \end{pmatrix} = 0$

$$\Rightarrow \alpha = 1 - \omega^2/\omega_0^2, 0 = -2 + 2(1 - \omega^2/\omega_0^2) - 2\frac{\omega^2}{\omega_0^2}(1 - \omega^2/\omega_0^2) \\ = (-\omega^2/\omega_0^2)(4 - 2\omega^2/\omega_0^2) \therefore \omega = \sqrt{2} \omega_0 = \sqrt{\frac{2k}{m}}$$

Math

$$\det \hat{B} = 0 = (2 - \omega^2/\omega_0^2) \left[(2 - \frac{\omega^2}{\omega_0^2})(2 - 2\frac{\omega^2}{\omega_0^2}) - 1 \right] + \underbrace{\left[-2 + 2\frac{\omega^2}{\omega_0^2} - 1 \right]}_{\neq \sqrt{3}\omega_0} - \left[1 + 2 - \frac{\omega^2}{\omega_0^2} \right]$$

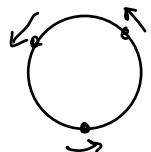
4. $0 = (2 - \omega^2/\omega_0^2) \left(\frac{2\omega^2}{\omega_0^2} \right) \left(\frac{\omega^2}{\omega_0^2} - 3 \right) \Rightarrow \omega = 0, \sqrt{2} \omega_0, -\sqrt{3} \left(2 - \omega^2/\omega_0^2 \right)$

2. $\boxed{\omega=0} \quad \begin{pmatrix} 2 & -1 & -1 \\ -1 & 2 & -1 \\ -1 & -1 & 2 \end{pmatrix} \vec{a} = 0 \therefore \vec{a} = \begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix} \quad \boxed{\omega=\sqrt{2}\omega_0} \quad \begin{pmatrix} 0 & -1 & -1 \\ -1 & 0 & -1 \\ -1 & -1 & 2 \end{pmatrix} \vec{a} = 0 \therefore \vec{a} = \begin{pmatrix} 1 \\ 1 \\ -1 \end{pmatrix}$

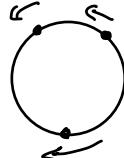
$\boxed{\omega=\sqrt{3}\omega_0}$ $\begin{pmatrix} -1 & -1 & -1 \\ -1 & -1 & -1 \\ -1 & -1 & -4 \end{pmatrix} \vec{a} = 0 \therefore \vec{a} = \begin{pmatrix} 1 \\ -1 \\ 0 \end{pmatrix}$

Either method:

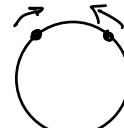
4. $\omega=0$



$\omega=\sqrt{2}\omega_0$



$\omega=\sqrt{3}\omega_0$



4. Dumbbell $\rho = \frac{M}{\pi R^2}$ $\{r, \theta, z\}$ coords

(a) (12) 

$$I_{zz} = \rho \int_0^R r dr \int_0^{2\pi} d\theta \left(\underbrace{x^2 + y^2}_{r^2} \right) = \rho \frac{R^4}{4} (2\pi) = \frac{MR^2}{2}$$

$$I_{xx} = I_{yy} = \rho \int_0^R r dr \int_0^{2\pi} d\theta (x^2 + y^2) = \rho \int_0^R r dr \int_0^{2\pi} r^2 \cos^2 \theta$$

$$= \rho \frac{R^4}{4} \pi = \frac{MR^2}{4}$$

5. $I_{disk} = \frac{MR^2}{2} \begin{pmatrix} \frac{1}{2} & 0 \\ 0 & 1 \end{pmatrix}$ ← full points for just writing it down

For one disk use // axes to translate $I_{ab}^{(a)} = I_{ab}^{(cm)} + M(s_{ab}\vec{R}^2 - R_a R_b)$

$$I^{(a)} = I_{disk} + ML^2 \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix}$$

For two disk double it

$$I_{dumbbell} = 2I_{disk} + 2ML^2 \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} = \begin{pmatrix} \frac{MR^2}{2} + 2ML^2 & 0 & 0 \\ 0 & \frac{MR^2}{2} + 2ML^2 & 0 \\ 0 & 0 & MR^2 \end{pmatrix}$$

or $I_{xx} = I_{yy} = \frac{MR^2}{2} + 2ML^2$ & $I_{zz} = MR^2$

(b) 3. $\vec{\omega} = \omega (0, \sin\theta, \cos\theta)$

4. $\hat{\vec{L}} = \hat{\vec{I}}_{dumbbell} \cdot \vec{\omega} = \omega \sin\theta \left(\frac{MR^2}{2} + 2ML^2 \right) \hat{y} + \omega \cos\theta MR^2 \hat{z}$

(c) Use Euler Equations with $(1, 2, 3) = (x, y, z)$

5. $\omega_1 = 0, \omega_2 = \omega \sin\theta, \omega_3 = \omega \cos\theta, \dot{\omega} = \alpha t, \ddot{\omega} = \alpha$

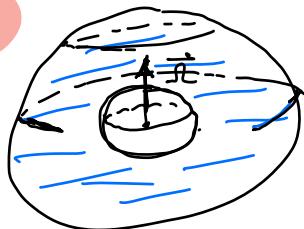
$$\begin{aligned} 3. \quad \tau_1 &= I_1 \dot{\omega}_1 - (I_2 - I_3) \omega_2 \omega_3 = 0 - \left(\frac{MR^2}{2} + 2ML^2 - MR^2 \right) \omega^2 \sin\theta \cos\theta \\ &= \left(\frac{MR^2}{2} - 2ML^2 \right) (\alpha t)^2 \sin\theta \cos\theta \end{aligned}$$

$$2. \quad \tau_2 = I_2 \dot{\omega}_2 - (I_3 - I_1) \underbrace{\omega_3 \omega_1}_0 = \left(\frac{MR^2}{2} + 2ML^2 \right) \alpha \sin\theta$$

$$2. \quad \tau_3 = I_3 \dot{\omega}_3 - (I_1 - I_2) \underbrace{\omega_1 \omega_2}_0 = MR^2 \alpha \cos\theta$$

5.

35



a) 4)

$$\vec{v} = \vec{r} \times \vec{r}_1 = r r_1 \hat{z} \times \hat{r}$$

$$\hat{z} = \cos\theta \hat{r} - \sin\theta \hat{\phi}, \quad \hat{z} \times \hat{r} = -\sin\theta (-\hat{\phi})$$

so

$$\vec{v} = -r r_1 \sin\theta \hat{\phi}$$

b) 4)

No dependence on ϕ :

$$\vec{v} = v_\phi(r, \theta) \hat{\phi}$$

$$p = p(r, \theta)$$

c) 2)

$$R = \frac{uL}{\eta} = \frac{uL\rho}{\eta}$$

$$L = r_1, \quad u = r r_1$$

$$R = \frac{r r_1^2 \rho}{\eta}$$

[OK if $r_1 \rightarrow r_2$, or left with $\omega = n/e$]

d) 8)

$$\text{Look at } \nabla^2 \vec{v} = \left[\underbrace{\frac{1}{r^2} \frac{\partial}{\partial r} r^2 \frac{\partial}{\partial r} + \frac{1}{r^2 \sin\theta} \frac{\partial}{\partial \theta} \sin\theta \frac{\partial}{\partial \theta}}_{\text{only acts on } v_\phi} + \underbrace{\frac{1}{r^2 \sin^2\theta} \frac{\partial^2}{\partial \phi^2}}_{\text{only acts on } \hat{\phi}} \right] (v_\phi \hat{\phi})$$

$$\therefore \nabla^2 (v_\phi \hat{\phi}) = (\nabla^2 v_\phi(r, \theta)) \hat{\phi} + \frac{v_\phi}{r^2 \sin^2\theta} \frac{\partial^2}{\partial \phi^2} \hat{\phi}$$

$$= \frac{1}{r^2} \left(-\cos\theta \hat{r} - \sin\theta \hat{\phi} \right)$$

$$\nabla^2 (v_\phi \hat{\phi}) = \left[\nabla^2 v_\phi - \frac{v_\phi}{r^2 \sin^2\theta} \right] \hat{\phi}$$

$$= -\cos^2\theta \hat{\phi} - \sin^2\theta \hat{\phi}$$

$$= -\hat{\phi}$$

so purely in $\hat{\phi}$ direction

$$\vec{\nabla} p = \eta \nabla^2 \vec{v}$$

$$\left(r \frac{\partial}{\partial r} + \hat{\phi} + \frac{\partial}{\partial \theta} \right) p(r, \theta) = \eta \nabla^2 \vec{v} \propto \hat{\phi} \quad \therefore \frac{\partial p}{\partial r} = 0$$

$$\frac{\partial p}{\partial \theta} = 0$$

so pressure $p = \text{constant}$

e) 10)

$$\text{Direction } \vec{v} = \vec{r} \times [\vec{\nabla} f(r)] = r \hat{z} \times [\hat{r} f'(r)]$$

$$\hat{z} \times \hat{r} = \hat{\phi} \sin\theta \Rightarrow \vec{v} = r f'(r) \sin\theta \hat{\phi} \quad \text{correct direction}$$

$$\nabla^2 \vec{v} = 0 = \vec{r} \times \underbrace{[\vec{\nabla} \nabla^2 f(r)]}_{\text{must be 0}}$$

Equation

$$\therefore \nabla^2 f(r) = \text{constant} = K$$

$$K = \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial f}{\partial r} \right) f(r) = \frac{1}{r^2} \frac{\partial}{\partial r} r^2 f'(r) = \boxed{f''(r) + \frac{2}{r} f'(r) = K} \quad -6-$$

Check Solution $f(r) = Ar^2 + B/r$

$$f'(r) = 2Ar - B/r^2, \quad \frac{2}{r} f'(r) = 4A - \frac{2B}{r^3}$$

$$f''(r) = 2A + 2B/r^3,$$

$$f'' + \frac{2}{r} f' = 6A = \text{constant} \quad \checkmark$$

f ⁷ Boundary Conditions ① $\vec{v} = 0$ for $\vec{r} = \vec{r}_2$ outer sphere
 ② $\vec{v} = \hat{\theta} r_1 \sin \theta \hat{\phi}$ for $\vec{r} = \vec{r}_1$ inner sphere

$$\vec{v} = \hat{\theta} r \sin \theta (2Ar - B/r^2) \hat{\phi}$$

$$\text{for ① need } 2Ar_2^3 = B, \text{ ② needs } 2Ar_1 - \frac{B}{r_1^2} = r_1$$

$$\text{so } 2A \left(1 - \frac{r_2^3}{r_1^3} \right) = 1, \quad A = \frac{1}{2 \left(1 - \frac{r_2^3}{r_1^3} \right)}, \quad B = \frac{r_2^3}{1 - \frac{r_2^3}{r_1^3}}$$

$$\Rightarrow \boxed{\vec{v} = \hat{\theta} r \sin \theta \frac{\left(1 - \frac{r_2^3}{r_1^3} \right)}{\left(1 - \frac{r_2^3}{r_1^3} \right)}}$$

$$\lim_{r_2 \rightarrow \infty} \vec{v} = \hat{\theta} r \sin \theta \frac{r_1^3}{r^3} = \hat{\theta} r \sin \theta \frac{r_1^3}{r^2}$$

falls off $\propto r^{-2}$ as expected

6. 30

"Non-linear Attractions"

-7-

(a)

$$\dot{x} = xy - 1$$

Fixed pts

$$xy = +1$$

 \leftarrow can't have $y=0$

(b)

$$\dot{y} = y(1-y^2)$$

$$y(1-y^2) = 0$$

$$\text{so } y = \pm 1$$

$$= y(1-y)(1+y)$$

$$\begin{aligned} \text{if } x &= \frac{y}{y} \\ &= \pm 1 \end{aligned}$$

2

$$(x^*, y^*) = (1, 1) \text{ or } (-1, -1)$$

(1, 1)

$$\text{let } x = 1+u$$

$$\dot{u} = (1+u)(1-u) - 1 = u + u^2 + \dots$$

$$y = 1+v$$

$$\dot{v} = (1+v)(-v)(2+v) = -2v + \dots$$

5.

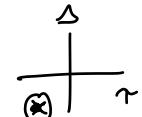
$$M = \begin{pmatrix} 1 & 1 \\ 0 & -2 \end{pmatrix}$$

$$\tau = \text{tr} M = -1$$

$$\tau^2 - 4\Delta = 1 + 8 = 9$$

$$\Delta = \det M = -2$$

$$\lambda_{\pm} = \frac{-1 \pm 3}{2} = \boxed{\pm 1, -2}$$

Saddle node

$$\lambda_+ = +1$$

$$\begin{pmatrix} 0 & 1 \\ 0 & -3 \end{pmatrix} \begin{pmatrix} a \\ b \end{pmatrix} = 0$$

$$\begin{pmatrix} 1 \\ 0 \end{pmatrix}$$

eigenvector

$$\lambda_- = -2$$

$$\begin{pmatrix} 3 & 1 \\ 0 & 0 \end{pmatrix} \begin{pmatrix} a \\ b \end{pmatrix} = 0$$

$$\frac{1}{\sqrt{10}} \begin{pmatrix} 1 \\ -3 \end{pmatrix}$$

eigenvector

(-1, -1)

$$\text{let } x = -1+u$$

$$\dot{u} = (-1+u)(-1+u) - 1 = -u - u^2 + \dots$$

$$y = -1+v$$

$$\dot{v} = (-1+v)(2-v)(v) = -2v + \dots$$

5.

$$M = \begin{pmatrix} -1 & -1 \\ 0 & -2 \end{pmatrix}, \tau = -3, \Delta = +2$$

$$\tau^2 - 4\Delta = 9 - 8 = +1 > 0$$

Stable node

$$\lambda_{\pm} = -\frac{3 \pm 1}{2} = \boxed{-1, -2}$$

$$\lambda_+ = -1$$

$$\begin{pmatrix} 0 & -1 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} a \\ b \end{pmatrix} = 0$$

$$\begin{pmatrix} 1 \\ 0 \end{pmatrix}$$

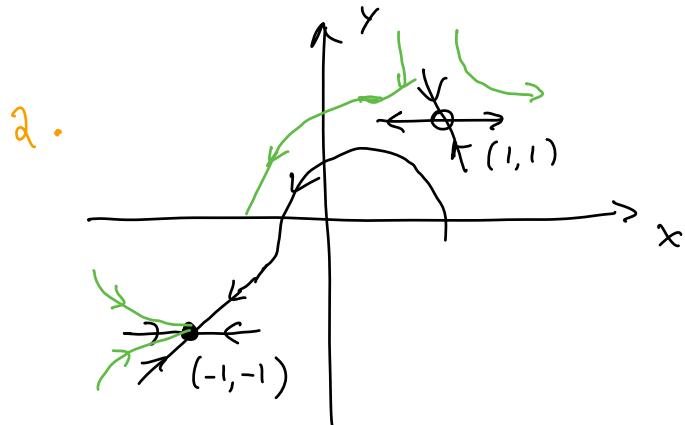
eigenvector

$$\lambda_- = -2$$

$$\begin{pmatrix} 1 & 1 \\ 0 & 0 \end{pmatrix} \begin{pmatrix} a \\ b \end{pmatrix} = 0$$

$$\frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix}$$

\rightarrow eigenvector



(b)
6

$$\dot{r} = \mu r - r^3 = r(\mu - r^2) = r(\sqrt{\mu} - r)(-\sqrt{\mu} - r)$$

$$\dot{\theta} = \omega + b r^2 \quad \omega > 0, b > 0$$

fixed pt $r=0$

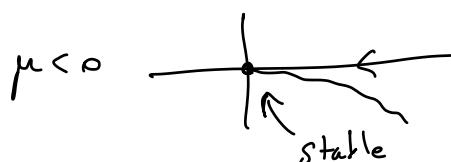
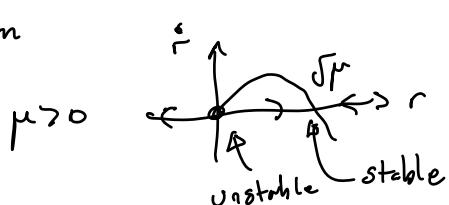
$\mu > 0$
unstable

$\mu < 0$
stable

2

limit cycle $r = \sqrt{\mu}$ needs $\mu > 0$, & stable 1

from



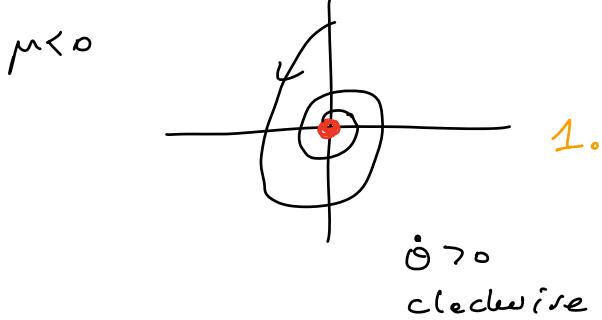
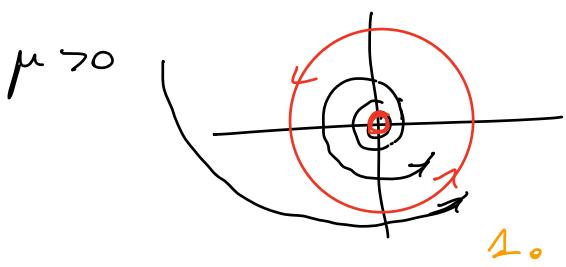
OR linear analysis in r , $\dot{r} = \mu r$ stable $\mu < 0$, unstable $\mu > 0$

$$r = \sqrt{\mu} + \Delta r, \quad \dot{\Delta r} = \sqrt{\mu}(2\sqrt{\mu} + \cancel{\Delta r})(-\Delta r) = -\sqrt{\mu}\Delta r$$

stable

5(6) cont

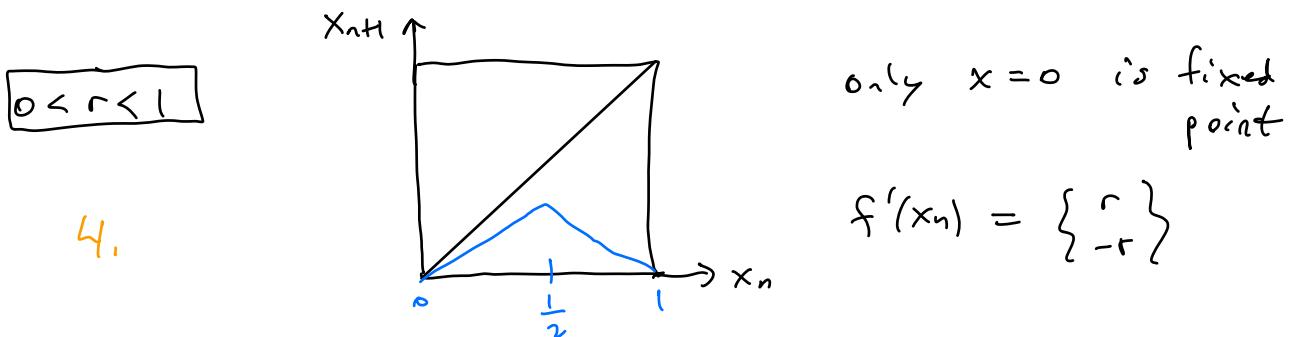
-9-



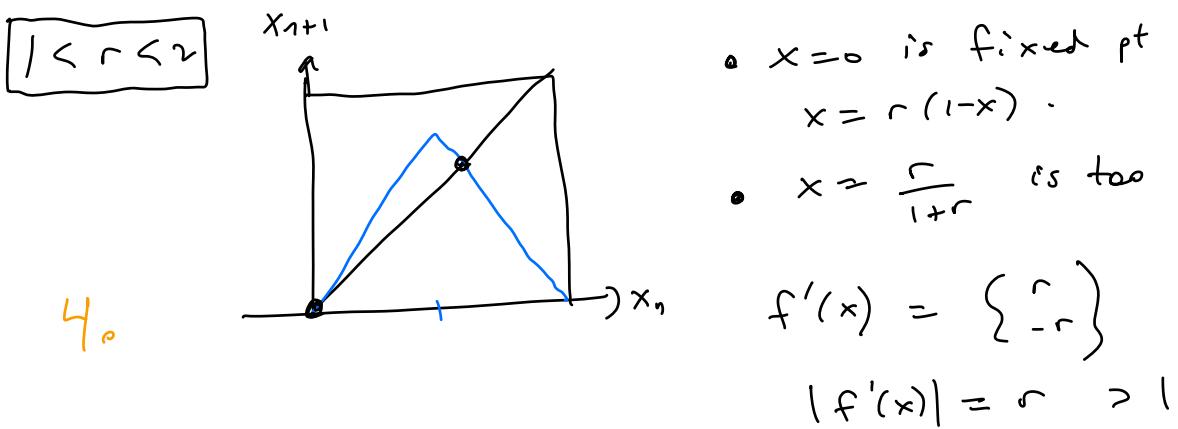
1. $\therefore \mu = 0$ is supercritical Hopf bifurcation

(c) 8

$$x_{n+1} = \begin{cases} rx_n & 0 \leq x_n \leq y_1 \\ r(1-x_n) & y_2 \leq x_n \leq 1 \end{cases} = f(x_n)$$



$$\text{so } |f'(0)| = r < 1 \Rightarrow \text{stable}$$



both unstable