The Hydrogen Atom and Harmonic Oscillator(s)

Huan Bui

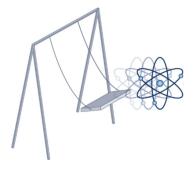
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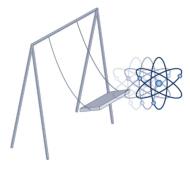
Harmonic oscillator universe?



Harmonic oscillator in physics:

- Hooke's law
- QHO
- Einstein solid
- Atom-radiation interaction
- 2nd quantization of EM fields
- QFT
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- Gravity? Inverse-square law?

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A particle in a central potential V(r) = -k/r:

$$H=\frac{\vec{p}^2}{2m}-\frac{k}{r}.$$

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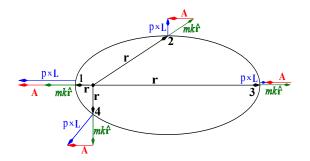
$$H=\frac{\vec{p}^2}{2m}-\frac{k}{r}.$$

Constants of motion: H, $\vec{L} = \vec{r} \times \vec{p}$, and \vec{A} , the Laplace-Runge-Lenz vector:

$$\vec{A} = \vec{p} \times \vec{L} - mk\frac{\vec{r}}{r}.$$

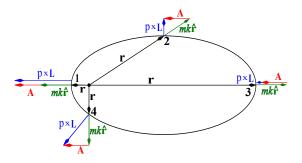
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Brief review:



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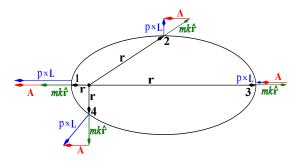
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 \vec{A} is in the plane of the orbit (so $\vec{A} \cdot \vec{L} = 0$), with $A^2 = m^2 k^2 + 2mEL^2$.

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• Shape: from

$$\vec{A} \cdot \vec{r} = \vec{r} \cdot (\vec{p} \times \vec{L}) - mk = (\vec{r} \times \vec{p}) \cdot \vec{L} - mk = L^2 - mk = Ar \cos \theta$$

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we get the orbit equation

$$\frac{1}{r} = \frac{mk}{L^2} \left(1 + \epsilon \cos \theta \right), \quad \text{eccentricity } \epsilon = \frac{A}{|mk|} = \sqrt{1 + \frac{2EL^2}{mk^2}} \ge 0$$

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• Orientation: \vec{A} points from source to periapsis

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• Used \vec{A} to derive the spectrum of hydrogen (pre-SE!)

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Further readings (so fun):

- History: [1], [2], [3], [4], [5]
- "Discoveries" and application: [6], [7], [8], [9], [10], [11]

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The Hydrogen Atom

The energy levels and wavefunctions for the bound states of hydrogen are gotten by solving the Schrödinger equation:

$$\left\{-\frac{\hbar^2}{2m}\nabla^2 - \frac{e^2}{r}\right\}\psi = E\psi, \qquad E < 0.$$

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With

$$\lambda = \frac{8}{a} \qquad \alpha^4 = -\frac{8E}{e^2 a}, \qquad a = \frac{\hbar^2}{me^2}, \tag{1}$$

the SE becomes

$$\left\{4\nabla^2 + \frac{\lambda}{r} - \alpha^4\right\}\psi = 0. \tag{2}$$

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Following [12], introduce coordinates $\zeta_A, \zeta_B \in \mathbb{C}$ and demand

$$x + iy = 2\zeta_A \overline{\zeta_B} \qquad z = \zeta_A \overline{\zeta_A} - \zeta_B \overline{\zeta_B}.$$

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$$x + iy = 2\zeta_A \overline{\zeta_B} \qquad z = \zeta_A \overline{\zeta_A} - \zeta_B \overline{\zeta_B}.$$

With this,

$$r = \sqrt{x^2 + y^2 + z^2} = \zeta_A \overline{\zeta_A} + \zeta_B \overline{\zeta_B}$$

Note:

- Each pair (ζ_A, ζ_B) gives a unique point (x, y, z)
- ullet Converse is true up to arbitrary but equal arguments of ζ_A,ζ_B

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Let $\sigma = 2 \arg(\zeta_A) = 2 \arg(\zeta_B)$. Can write ζ_A, ζ_B in spherical coordinates:

$$\zeta_A = r^{1/2} e^{i(\sigma + \varphi)/2} \cos \frac{\theta}{2} \qquad \zeta_B = r^{1/2} e^{i(\sigma - \varphi)/2} \sin \frac{\theta}{2}$$
 (3)

 $\implies (x, y, z)$ determines (ζ_A, ζ_B) up to $e^{i\sigma}$.

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Can show that

$$r\nabla^2\psi=(\partial_A\partial_{\bar{A}}+\partial_B\partial_{\bar{B}})\psi.$$

 \implies Can now write SE in terms of $\zeta_A, \zeta_B, \overline{\zeta_A}, \overline{\zeta_B}$.

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SE in terms of $\zeta_A, \zeta_B, \overline{\zeta_A}, \overline{\zeta_B}$:

$$\left\{4\nabla^2 + \frac{\lambda}{r} - \alpha^4\right\}\psi = \left\{4\left(\partial_A\partial_{\bar{A}} + \partial_B\partial_{\bar{B}}\right) + \lambda - \alpha^4\left(\zeta_A\overline{\zeta_A} + \zeta_B\overline{\zeta_B}\right)\right\}\psi = 0 \quad (4)$$

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SE in terms of $(\Delta, \zeta_R, \overline{\zeta_\Delta}, \overline{\zeta_R})$:

$$\left\{4\nabla^2 + \frac{\lambda}{r} - \alpha^4\right\}\psi = \left\{4(\partial_A\partial_{\bar{A}} + \partial_B\partial_{\bar{B}}) + \lambda - \alpha^4(\zeta_A\overline{\zeta_A} + \zeta_B\overline{\zeta_B})\right\}\psi = 0 \quad (4)$$

Since $\psi(x, y, z)$ independent of σ ,

$$\frac{\partial \psi}{\partial \sigma} = 0 \quad \Longleftrightarrow \quad (\overline{\zeta}_A \partial_{\bar{A}} - \zeta_A \partial_A) \psi = -(\overline{\zeta}_B \partial_{\bar{B}} - \zeta_B \partial_B) \psi. \tag{5}$$

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Together, (4) and (5) are equivalent to SE (2).

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Let $\zeta_A = q_1 + iq_2$ and $\zeta_B = q_3 + iq_4$, then (4) is the equation for a 4D HO

$$\left[\partial_1^2 + \partial_2^2 + \partial_3^2 + \partial_4^2 + \lambda - \alpha^4 (q_1^2 + q_2^2 + q_3^2 + q_4^2)\right] \psi = 0$$
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Where are the harmonic oscillators?

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with frequency ω and energy ϵ given by (1):

$$\alpha^2 \equiv \sqrt{-\frac{8E}{e^2 a}} = \frac{m\omega}{\hbar}$$
 $\lambda \equiv \frac{8}{a} = \frac{2m\epsilon}{\hbar^2}$.

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Condition (5) becomes

$$(q_1\partial_2 - q_2\partial_1)\psi = -(q_3\partial_4 - q_4\partial_3)\psi. \tag{7}$$

(6)+(7): Two 2D HO's with equal and opposite angular momenta

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Separating variables $\psi = \psi(q_1, q_2)\psi(q_3, q_4)$,

$$[\partial_1^2 + \partial_2^2 + \lambda_A - \alpha^4 (q_1^2 + q_2^2)] \psi_A = 0,$$

with $\lambda_A = 2m\epsilon_A/\hbar^2$.

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with $\lambda_A = 2m\epsilon_A/\hbar^2$. Solution for A:

$$\psi_{An_Am_A} = C_{n_Am_A} \left(\frac{\zeta_A}{\overline{\zeta_A}}\right)^{m_A/2} \left(\alpha^2 \zeta_A \overline{\zeta_A}\right)^{|m_A|/2} e^{-\frac{\alpha^2 \zeta_A \overline{\zeta_A}}{2}} L_{n_A+|m_A|}^{|m_A|} \left(\alpha^2 \zeta_A \overline{\zeta_A}\right)$$

$$n_A = 0, 1, 2, \dots \qquad m_A = 0, \pm 1, \pm 2, \dots$$

Energy:
$$\epsilon_{An_Am_A} = \hbar\omega(2n_A + |m_A| + 1) = \frac{\hbar^2\lambda_{An_Am_A}}{2m}$$

Angular momentum: $L_{An_Am_A} = m_A\hbar$

Similar solution for B. $\lambda_A + \lambda_B = \lambda$ and $m_A = -m_B = m$ due to (7).

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Full solution

$$\psi_{n_A n_B m} = \psi_{A n_A m} \left(\zeta_A, \overline{\zeta_A} \right) \psi_{B n_B - m} \left(\zeta_B, \overline{\zeta_B} \right).$$

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Can relate this back to the hydrogen atom.

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Full solution

$$\psi_{n_{A}n_{B}m}=\psi_{An_{A}m}\left(\zeta_{A},\overline{\zeta_{A}}\right)\psi_{Bn_{B}-m}\left(\zeta_{B},\overline{\zeta_{B}}\right).$$

Can relate this back to the hydrogen atom. From

$$\lambda = \lambda_A + \lambda_B = 4\alpha^2(n_A + n_B + |m| + 1) = \frac{8}{a}$$

can get energy in terms of n_A , n_B , m:

$$E = \frac{-\alpha^4 e^2 a}{8} = -\frac{\alpha^4 e^2}{\lambda} = \frac{-e^2}{2a(n_A + n_B + |m| + 1)^2} \equiv \frac{-e^2}{2aN^2}.$$

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How about the wavefunctions?

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How about the wavefunctions? Going to parabolic coordinates (ξ, η, φ) :

$$x = \sqrt{\xi \eta} \cos \varphi \qquad y = \sqrt{\xi \eta} \sin \varphi \qquad z = (\xi - \eta)/2$$

$$\iff \xi = 2r \cos^2(\theta/2) = 2|\zeta_A|^2 \qquad \eta = 2r \sin^2(\theta/2) = 2|\zeta_B|^2.$$

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we get

$$\psi_{n_{A}n_{B}m} = K_{n_{A}n_{B}m}e^{im\varphi}(\xi\eta)^{|m|/2}e^{-\frac{\alpha^{2}(\xi^{2}+\eta^{2})}{4}}L_{n_{A}+|m|}^{|m|}\left(\frac{\alpha^{2}\xi}{2}\right)L_{n_{B}+|m|}^{|m|}\left(\frac{\alpha^{2}\eta}{2}\right).$$

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These are simultaneous eigenfunctions of H, L_z , and M_z where

$$\mathbf{M} = \frac{1}{2m}(\mathbf{p} \times \mathbf{L} - \mathbf{L} \times \mathbf{p}) - \frac{e^2}{r}\mathbf{r}.$$

is the Laplace-Runge-Lenz operator, symmetrized by Pauli, 1926.

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From how (ζ_A, ζ_B) is defined:

$$\mathbf{M}_{z} = \frac{e^{2}a}{r} \left[\left| \zeta_{B} \right|^{2} \partial_{A} \partial_{\bar{A}} - \left| \zeta_{A} \right|^{2} \partial_{B} \partial_{\bar{B}} - \frac{1}{a} \left(\left| \zeta_{A} \right|^{2} + \left| \zeta_{B} \right|^{2} \right) \right].$$

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CSCO is $\{H, L_z, M_z\}$ instead of $\{H, L^2, L_z\}$. Eigenvalue equations:

$$\begin{aligned} \mathbf{H}\psi_{n_An_Bm} &= \frac{-e^2}{2aN^2}\psi_{n_An_Bm} \\ \mathbf{L}_z\psi_{n_An_Bm} &= mh\psi_{n_An_Bm} \\ \mathbf{M}_z\psi_{n_An_Bm} &= \frac{e^2(n_B-n_A)}{N}\psi_{n_An_Bm}. \end{aligned}$$

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$$\mathbf{L}_z\psi_{n_An_Bm} = mh\psi_{n_An_Bm}$$

$$\mathbf{M}_z\psi_{n_An_Bm} = \frac{e^2(n_B - n_A)}{N}\psi_{n_An_Bm}.$$

$$\label{eq:main_section} \left[\boldsymbol{\mathsf{M}}_{z}, \boldsymbol{\mathsf{L}}^{2} \right] \neq 0, \qquad \left[\boldsymbol{\mathsf{M}}_{z}, \boldsymbol{\mathsf{M}}^{2} \right] \neq 0, \qquad \left[\boldsymbol{\mathsf{M}}^{2}, \boldsymbol{\mathsf{H}} \right] = \left[\boldsymbol{\mathsf{M}}^{2}, \boldsymbol{\mathsf{L}}^{2} \right] = \left[\boldsymbol{\mathsf{M}}^{2}, \boldsymbol{\mathsf{L}}_{z} \right] = 0.$$

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Aside: Hydrogen wavefunctions in parabolic coordinates

What do eigenfunctions of M_z look like?

$$\psi_{n_{A}n_{B}m} = K_{n_{A}n_{B}m}e^{im\varphi}(\xi\eta)^{|m|/2}e^{-\frac{\alpha^{2}(\xi^{2}+\eta^{2})}{4}}L_{n_{A}+|m|}^{|m|}\left(\frac{\alpha^{2}\xi}{2}\right)L_{n_{B}+|m|}^{|m|}\left(\frac{\alpha^{2}\eta}{2}\right).$$

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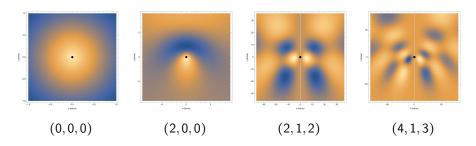


Figure: $|\psi_{n_A n_B m}(x, 0, z)|^2$ for different values of (n_A, n_B, m) [13]

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Following [14], use the Kustaanheimo-Stiefel transformation $S : \mathbb{R}^4 \to \mathbb{R}^3$:

$$x_1 = 2(s_1s_3 - s_2s_4)$$

$$x_2 = 2(s_1s_4 + s_2s_3)$$

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$$\begin{cases} x_1 = r \sin \theta \cos \phi \\ x_2 = r \sin \theta \sin \phi \\ x_3 = r \cos \theta \end{cases} \begin{cases} s_1 = s \cos \alpha \cos \beta \\ s_2 = s \cos \alpha \sin \beta \\ s_3 = s \sin \alpha \cos \gamma \\ s_4 = s \sin \alpha \sin \gamma \end{cases}$$

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Constraint on coordinates: $r = s^2$, $\theta = 2\alpha$, $\phi = \beta + \gamma$ Constraint on velocities: $s_2\dot{s}_1 - s_1\dot{s}_2 - s_4\dot{s}_3 + s_3\dot{s}_4 = 0$

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The equivalence: With $4|E| = m\omega^2/2$ and $\epsilon = 4k$,

$$\frac{1}{2}mv^2 - \frac{k}{r} = -|E| \longrightarrow \frac{1}{2}m\dot{s}^2 + \frac{1}{2}m\omega^2 s^2 = \epsilon \quad (4D \text{ H.O.})$$

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In polar coordinates: $u = s \cos \alpha$, $v = s \sin \alpha$

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The constraint on velocities $\iff \vec{L} \cdot \vec{M} = 0$ and implies

$$-u^2\dot{\beta} + v^2\dot{\gamma} = 0 \implies mu^2\dot{\beta} = mv^2\dot{\gamma} \implies L_{\beta} = L_{\gamma}$$

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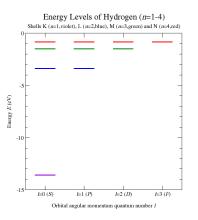
$$-u^2\dot{\beta}+v^2\dot{\gamma}=0 \implies mu^2\dot{\beta}=mv^2\dot{\gamma} \implies L_\beta=L_\gamma$$

→ Two coupled 2D H.O.'s with equal angular momenta

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Something deeper?

Symmetry implies degeneracy. $[\mathbf{M}, \mathbf{H}] = 0$ explains the n^2 degeneracy in H.



More info: SO(4) symmetry of H, etc. See Chapter 14 of [15].

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