

Today, we start our unit on quantum error-correcting codes.

I'm going to start with a digression into classical information theory. In 1948, Claude Shannon published a paper, "A Mathematical Theory of Communication", which started the field of information theory. In it, he showed that noisy communication channels had a *capacity*, and he derived *Shannon's formula* for the capacity of a channel. If you try to send information over them at a rate less than their capacity, in the limit as the length of the message goes to infinity, you can almost always succeed. On the other hand, if you try to send information over them at a rate larger than the capacity, it is overwhelmingly likely that the message the receiver gets has errors—i.e., it is different from the message that the sender transmitted.

Claude Shannon's theory was not constructive. While he showed that an algorithm existed that would succeed in transmitting information at nearly the channel capacity existed, it was a randomized construction that didn't explicitly produce such an algorithm. Worse, the algorithm itself would take exponential time to implement in the most straightforward way. It wasn't until 60 years later that a coding method was found that provably approached Shannon's bound in the limit of long messages (*polar codes*).

However, only two years after Shannon published his paper, Richard Hamming discovered the first error-correcting codes, and in the next few decades, computer technology advanced greatly and error correcting codes were developed to the point where they could be used in practice for the reliable transmission of information over noisy channels. Hamming codes are one of a class of codes called *linear codes*, and these have nice properties. They are the class of codes that is most often used in practice, and they have received extensive study. We will look at the code that Hamming discovered much more closely in future classes, when we construct the quantum analog of this code.

Before Hamming, the only error-correcting codes engineers knew about were *repetition codes*. In these, you just repeat the message bit  $k$  times. If  $k = 2t + 1$ , the code can correct  $t$  errors. These codes can be put into the framework of linear codes described above. For a repetition code, the generator matrix is

$$G = [1, 1, 1, \dots, 1].$$

Let's look at the three-bit repetition more carefully. Assume that each bit has a  $p$  probability of having an error, and a  $1 - p$  probability of being correct. Then there will be an error in an encoded bit if and only if at least two of the encoding bits have errors, so instead of a probability  $p$  of having an error, the probability is  $3p^2(1 - p) + p^3$ , which is an improvement as long as  $p < \frac{1}{2}$ , and is around  $3p^3$  if  $p$  is small.

Today, we will look at the quantum analog of repetition codes. Since it is impossible to clone a qubit, you can't actually implement a repetition code  $|\psi\rangle \rightarrow |\psi\rangle |\psi\rangle |\psi\rangle$ . One thing you could do instead is use the unitary:

$$\begin{aligned} U |0\rangle |00\rangle &= |000\rangle \\ U |1\rangle |00\rangle &= |111\rangle \end{aligned}$$

Here, we first adjoin the qubits  $|00\rangle$  to the qubit we want to encode and then perform the unitary.

This is a three-qubit code that protects against bit-flip errors. It was first investigated by Asher Peres in 1985. The code maps

$$\alpha |0\rangle + \beta |1\rangle \rightarrow \alpha |000\rangle + \beta |111\rangle ,$$

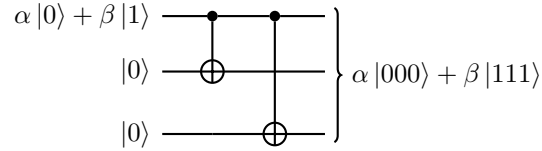
so you can see it does not perform quantum cloning; if you cloned the original qubit, you would get  $(\alpha |0\rangle + \beta |1\rangle)^{\otimes 3}$ .

This code can correct one bit-flip, or  $\sigma_x$ , error. How does this work? We measure the answer to the question “which bit is different?” More specifically, we project the three qubits onto one of the following four subspaces:

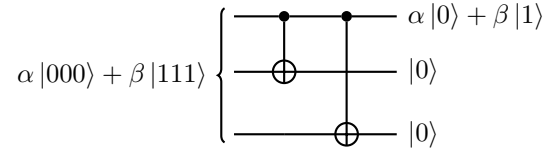
$$\begin{aligned} &|000\rangle\langle 000| + |111\rangle\langle 111| \\ &|100\rangle\langle 100| + |011\rangle\langle 011| \\ &|010\rangle\langle 010| + |101\rangle\langle 101| \\ &|001\rangle\langle 001| + |110\rangle\langle 110|. \end{aligned}$$

Once we know which subspace are in, we can correct the error. For example, if the state was projected onto the third subspace above, we would apply a  $\sigma_x$  to the second qubit to correct it.

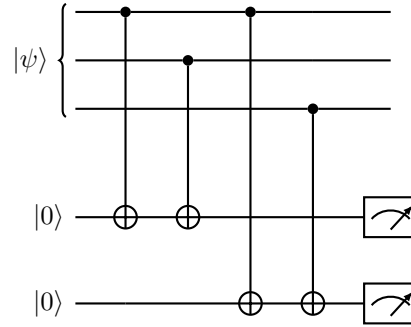
How do we encode a qubit into this code? We use the following circuit:



How do we decode? We basically reverse the circuit above, which gives exactly the same circuit since the two CNOT gates commute with each other:



How do we correct errors? We use the following quantum circuit, that projects onto one of the four subspaces described above:



The two measurement results here are called the *syndrome* and tell us what the error is.

If the two measurements are both  $|0\rangle$ , then we know that  $|\psi\rangle$  was in the code subspace. If the first one is  $|1\rangle$  and the second is  $|0\rangle$ , we know that the first and third qubits are equal, but the first and second qubits are different. This means that the second qubit must be different. Thus, if there is only one error, it is on the second qubit, and we can correct the state by applying  $\sigma_x$  to the second qubit. Similarly, if we get the measurement results  $(0, 1)$ , we know the third qubit is different, and if the measurement results are  $(1, 1)$ , we know the first qubit is different.

For bit-flip errors, the probabilities work exactly the same way as they did for the classical three-bit code. If the probability of a bit-flip error on a single qubit is  $p$ , then the probability of a bit-flip error on the encoded qubit is  $3p^2 + p^3$ , which is approximately  $3p^2$  for small  $p$ .

But what about phase-flip errors? What happens if we apply a  $\sigma_z$  to one of the three qubits in the code? We will call an encoded  $|0\rangle$  and  $|1\rangle$  a *logical*  $|0\rangle$  and  $|1\rangle$ , and we will represent them by

$$|0\rangle_L = |000\rangle \quad |1\rangle_L = |111\rangle$$

If we apply a phase-flip error in any of three encoding qubits, it will take  $|0\rangle$  to  $|0\rangle$  and  $|1\rangle$  to  $-|1\rangle$ . That is, it will apply a phase error to the logical qubit. So if the probability of a phase error on a single qubit is  $p$ , the probability of a phase error on the encoded qubit is  $3p + p^3$ , or approximately  $3p$  for small  $p$ .

So we can reduce the error rate for bit flip errors at the cost of increasing the error rate for phase flip errors. Is there anything we can do to protect phase-flip errors?

There is. Recall that a Hadamard gate took a  $\sigma_x$  to a  $\sigma_z$  and vice versa. Thus, we can find a code that interchanges the role of bit-flip and phase-flip errors. This code is simply

$$\begin{aligned} |0\rangle &\rightarrow |+++ \rangle \\ |1\rangle &\rightarrow |-- - \rangle. \end{aligned}$$

or

$$\begin{aligned} |0\rangle &\rightarrow \frac{1}{\sqrt{8}}(|0\rangle + |1\rangle)^{\otimes 3}, \\ |1\rangle &\rightarrow \frac{1}{\sqrt{8}}(|0\rangle - |1\rangle)^{\otimes 3}. \end{aligned}$$

This code protects against phase-flip errors, but any bit-flip error in an encoding qubit results in a bit-flip error on the logical qubit, so bit-flip errors are roughly three times as likely as on unencoded qubits. This can be seen directly from the fact that a bit-flip error takes a bit string with an odd number of 1's to a bit string with an even number of 1's.

We will be representing this code slightly differently later in the course, as

$$\begin{aligned} |0\rangle &\rightarrow \frac{1}{4}(|000\rangle + |011\rangle + |101\rangle + |111\rangle) \\ |1\rangle &\rightarrow \frac{1}{4}(|111\rangle + |100\rangle + |010\rangle + |001\rangle) \end{aligned}$$

We obtain this alternate representation by applying a Hadamard gate to the qubit before encoding it. This code encodes a  $|0\rangle$  as the superposition of all states with an even number of 0's, and a  $|1\rangle$  as the superposition of all states with an odd number of 1s.

So now we have two codes, and we have a choice: we can protect against one bit-flip error, but only by making phase-flip errors more likely, or we can protect against phase-flip errors but only by making bit-flip errors more likely. Is there any way around this problem.

It turns out there is. The answer comes from a technique of classical coding theory: you *concatenate* the two codes. This means you first encode using one code, and then you encode using the other. Let's see how this works.

$$\begin{aligned} |0\rangle &\rightarrow |+++ \rangle \rightarrow (|000\rangle + |111\rangle)^{\otimes 3} \\ |1\rangle &\rightarrow |-- - \rangle \rightarrow (|000\rangle - |111\rangle)^{\otimes 3} \end{aligned}$$

Now, what happens? if you have a bit-flip error, it gets corrected immediately by the inner code. If you have a phase-flip error, the inner code turns this into a phase-flip on the logical qubit of the inner code, which gets corrected by the outer code. Thus, any single  $\sigma_x$  or  $\sigma_z$  can be corrected.

How about  $\sigma_y$  errors? They can be corrected as well. A  $\sigma_y$  error can be viewed as a  $\sigma_x$  error and a  $\sigma_z$  error acting on the same qubit, since  $\sigma_y = i\sigma_x\sigma_z$ . Thus, any single  $\sigma_y$  error can be corrected as well—the inner code will correct the  $\sigma_x$  component and the outer code the  $\sigma_z$  component.

But what about more general errors? It turns out that the 9-qubit code given in this lecture can correct these as well, as long as they are restricted to one qubit. We will see this in the next lecture.