

PY 711 Fall 2010
Homework 7: Due Tuesday, October 12

1. (8 points) Let U be the following unitary operator defined in terms of Dirac annihilation and creation operators:

$$U = \exp \left[-\frac{i\pi}{2} \sum_{r=1,2} \int \frac{d^3\vec{p}}{(2\pi)^3} \left(a_{\vec{p}}^{r\dagger} - b_{\vec{p}}^{r\dagger} \right) \left(a_{\vec{p}}^r - b_{\vec{p}}^r \right) \right]. \quad (1)$$

Investigate the effect of this unitary transformation upon the annihilation operators by explicitly calculating $U^\dagger a_{\vec{p}}^r U$ and $U^\dagger b_{\vec{p}}^r U$. Which type of transformation does the unitary operator U produce?

2. (7 points) Using Dirac annihilation and creation operators, explicitly construct a unitary operator P which implements the parity transformation,

$$P^\dagger a_{\vec{p}}^r P = a_{-\vec{p}}^r, \quad (2)$$

$$P^\dagger b_{\vec{p}}^r P = -b_{-\vec{p}}^r. \quad (3)$$

1. LET U BE THE FOLLOWING UNITARY OPERATOR DEFINED IN TERMS OF DIRAC CREATION AND ANNIHILATION OPERATORS:

$$U = \exp \left[-\frac{i\pi}{2} \sum_{\vec{r}} \int \frac{d^3p}{(2\pi)^3} (a_{\vec{p}}^{\dagger} - b_{\vec{p}}^{\dagger})(a_{\vec{p}} - b_{\vec{p}}) \right]$$

INVESTIGATE THE EFFECT OF THIS UNITARY TRANSFORMATION UPON THE ANNIHILATION OPERATORS BY EXPLICITLY CALCULATING $U^{\dagger} a_{\vec{p}} U$ AND $U^{\dagger} b_{\vec{p}} U$. WHICH TYPE OF TRANSFORMATION DOES THE UNITARY OPERATOR U PRODUCE?

$$U^{\dagger} = \exp \left[+\frac{i\pi}{2} \sum_{\vec{r}} \int \frac{d^3p}{(2\pi)^3} ((a_{\vec{p}}^{\dagger} - b_{\vec{p}}^{\dagger})(a_{\vec{p}} - b_{\vec{p}}))^{\dagger} \right]$$

$$= \exp \left[+\frac{i\pi}{2} \sum_{\vec{r}} \int \frac{d^3p}{(2\pi)^3} (a_{\vec{p}}^{\dagger} - b_{\vec{p}}^{\dagger})(a_{\vec{p}} - b_{\vec{p}}) \right]$$

$$U^{\dagger} a_{\vec{p}} U = e^X a_{\vec{p}} e^{-X}$$

$$\text{where } X = \frac{i\pi}{2} \sum_{\vec{q}} \int \frac{d^3q}{(2\pi)^3} (a_{\vec{q}}^{\dagger} - b_{\vec{q}}^{\dagger})(a_{\vec{q}} - b_{\vec{q}})$$

Baker-Hausdorff Lemma (Sakurai pg 96)

$$e^X Y e^{-X} = Y + [X, Y] + \frac{1}{2!} [X, [X, Y]] + \frac{1}{3!} [X, [X, [X, Y]]] + \dots$$

Also need the identity

$$\begin{aligned} [AB, C] &= A\{B, C\} - \{A, C\}B \\ &= ABC + ACB - ACB - CAB \\ &= ABC - CAB = [AB, C] \end{aligned}$$

In addition, recall

$$\{a_{\vec{p}}, a_{\vec{q}}^{\dagger}\} = (2\pi)^3 \delta^3(\vec{p} - \vec{q}) = \{b_{\vec{p}}, b_{\vec{q}}^{\dagger}\}$$

All other combinations have anticommutators equal to zero.

1 CONTINUED

$$\left[\frac{i\pi}{2} \sum_s \int \frac{d^3q}{(2\pi)^3} (a_{\vec{q}}^{s+} - b_{\vec{q}}^{s+})(a_{\vec{q}}^s - b_{\vec{q}}^s), a_{\vec{p}}^r \right]$$

$$= \frac{i\pi}{2} \sum_s \int \frac{d^3q}{(2\pi)^3} \left((a_{\vec{q}}^{s+} - b_{\vec{q}}^{s+}) \{ \cancel{(a_{\vec{q}}^s - b_{\vec{q}}^s)}, a_{\vec{p}}^r \} \right. \\ \left. - \{ a_{\vec{q}}^{s+} - b_{\vec{q}}^{s+}, a_{\vec{p}}^r \} (a_{\vec{q}}^s - b_{\vec{q}}^s) \right)$$

$$= \frac{i\pi}{2} \sum_s \int \frac{d^3q}{(2\pi)^3} \left(- (2\pi)^3 \delta^{rs} \delta^3(\vec{p} - \vec{q}) (a_{\vec{q}}^s - b_{\vec{q}}^s) \right)$$

$$= - \left(\frac{i\pi}{2} \right) (a_{\vec{p}}^r - b_{\vec{p}}^r)$$

$$[X, [X, Y]] = \left[\frac{i\pi}{2} \sum_s \int \frac{d^3q}{(2\pi)^3} (a_{\vec{q}}^{s+} - b_{\vec{q}}^{s+})(a_{\vec{q}}^s - b_{\vec{q}}^s), \left(-\frac{i\pi}{2} \right) (a_{\vec{p}}^r - b_{\vec{p}}^r) \right]$$

$$= - \left(\frac{i\pi}{2} \right)^2 \sum_s \int \frac{d^3q}{(2\pi)^3} \left((a_{\vec{q}}^{s+} - b_{\vec{q}}^{s+}) \{ \cancel{(a_{\vec{q}}^s - b_{\vec{q}}^s)}, (a_{\vec{p}}^r - b_{\vec{p}}^r) \} \right. \\ \left. - \{ a_{\vec{q}}^{s+} - b_{\vec{q}}^{s+}, (a_{\vec{p}}^r - b_{\vec{p}}^r) \} (a_{\vec{q}}^s - b_{\vec{q}}^s) \right)$$

$$= + \left(\frac{i\pi}{2} \right)^2 \sum_s \int \frac{d^3q}{(2\pi)^3} \left((\{ a_{\vec{q}}^{s+}, a_{\vec{p}}^r \} + \{ b_{\vec{q}}^{s+}, b_{\vec{p}}^r \}) (a_{\vec{q}}^s - b_{\vec{q}}^s) \right)$$

$$= \left(\frac{i\pi}{2} \right)^2 \sum_s \int \frac{d^3q}{(2\pi)^3} \left(2 (2\pi)^3 \delta^{rs} \delta^3(\vec{p} - \vec{q}) (a_{\vec{q}}^s - b_{\vec{q}}^s) \right)$$

$$= 2 \left(\frac{i\pi}{2} \right)^2 (a_{\vec{p}}^r - b_{\vec{p}}^r)$$

Each subsequent commutator will have another factor of $2 \left(\frac{i\pi}{2} \right)$ and alternating negative signs, since each is the commutator will always be between the exponent of U^+ , which isn't changing, and $(a_{\vec{p}}^r - b_{\vec{p}}^r)$.

1 CONTINUED

So, we see that

$$U^\dagger a_{\vec{p}}^r U = a_{\vec{p}}^r - \left(\frac{i\pi}{2}\right)(a_{\vec{p}}^r - b_{\vec{p}}^r) + \frac{1}{2!} \left(\frac{i\pi}{2}\right)^2 2(a_{\vec{p}}^r - b_{\vec{p}}^r) \\ - \frac{1}{3!} \left(\frac{i\pi}{2}\right)^3 4(a_{\vec{p}}^r - b_{\vec{p}}^r) + \dots$$

$$= a_{\vec{p}}^r \left(1 - \frac{i\pi}{2} + \frac{1}{2!} \left(\frac{i\pi}{2}\right)^2 2 - \frac{1}{3!} \left(\frac{i\pi}{2}\right)^3 4 + \dots\right) \\ + b_{\vec{p}}^r \left(\frac{i\pi}{2} - \frac{1}{2!} \left(\frac{i\pi}{2}\right)^2 2 + \frac{1}{3!} \left(\frac{i\pi}{2}\right)^3 4 + \dots\right)$$

$$= a_{\vec{p}}^r \left(1 - \frac{1}{2} (i\pi - \frac{1}{2!} (i\pi)^2 + \frac{1}{3!} (i\pi)^3 - \dots)\right) \\ + b_{\vec{p}}^r \left(\frac{1}{2} (i\pi - \frac{1}{2!} (i\pi)^2 + \frac{1}{3!} (i\pi)^3 - \dots)\right)$$

$$= a_{\vec{p}}^r \left(1 - \frac{1}{2} \sum_{n=1}^{\infty} \frac{1}{n!} (i\pi)^n (-1)^{n-1}\right) + b_{\vec{p}}^r \left(\frac{1}{2} \sum_{n=1}^{\infty} \frac{1}{n!} (i\pi)^n (-1)^{n-1}\right)$$

$$\sum_{n=1}^{\infty} \frac{1}{n!} (i\pi)^n (-1)^{n-1} = 2$$

$$= a_{\vec{p}}^r \left(1 - \frac{1}{2} (2)\right) + b_{\vec{p}}^r \left(\frac{1}{2} (2)\right)$$

$$\boxed{U^\dagger a_{\vec{p}}^r U = b_{\vec{p}}^r}$$

I CONTINUED

$U^\dagger b_{\vec{p}}^r U$ will be similar

$$\begin{aligned}
 & \left[\frac{i\pi}{2} \sum_{\vec{s}} \int \frac{d^3 q}{(2\pi)^3} (a_{\vec{q}}^{s\dagger} - b_{\vec{q}}^{s\dagger})(a_{\vec{q}}^s - b_{\vec{q}}^s), b_{\vec{p}}^r \right] \\
 &= \frac{i\pi}{2} \sum_{\vec{s}} \int \frac{d^3 q}{(2\pi)^3} \left((a_{\vec{q}}^{s\dagger} - b_{\vec{q}}^{s\dagger}) \{ (a_{\vec{q}}^s - b_{\vec{q}}^s), b_{\vec{p}}^r \} \right. \\
 &\quad \left. - \{ (a_{\vec{q}}^{s\dagger} - b_{\vec{q}}^{s\dagger}), b_{\vec{p}}^r \} (a_{\vec{q}}^s - b_{\vec{q}}^s) \right) \\
 &= \frac{i\pi}{2} \sum_{\vec{s}} \int \frac{d^3 q}{(2\pi)^3} \left((2\pi)^3 \delta^{rs} \delta^3(\vec{p} - \vec{q}) (a_{\vec{q}}^s - b_{\vec{q}}^s) \right) \\
 &= + \frac{i\pi}{2} (a_{\vec{p}}^r - b_{\vec{p}}^r)
 \end{aligned}$$

So, apparently it will be off by a negative sign from the calculations with $a_{\vec{p}}^r$, since I've already calculated the next commutator.

$$\begin{aligned}
 U^\dagger b_{\vec{p}}^r U &= b_{\vec{p}}^r + \left(\frac{i\pi}{2}\right)(a_{\vec{p}}^r - b_{\vec{p}}^r) - \frac{1}{2!} \left(\frac{i\pi}{2}\right)^2 (a_{\vec{p}}^r - b_{\vec{p}}^r)(2) \\
 &\quad + \frac{1}{3!} \left(\frac{i\pi}{2}\right)^3 (a_{\vec{p}}^r - b_{\vec{p}}^r)(4) - \dots \\
 &= b_{\vec{p}}^r \left(1 - \frac{i\pi}{2} + \frac{1}{2!} \left(\frac{i\pi}{2}\right)^2 (2) - \frac{1}{3!} 4 \left(\frac{i\pi}{2}\right)^3 + \dots \right) \\
 &\quad + a_{\vec{p}}^r \left(\frac{i\pi}{2} - \frac{1}{2!} 2 \left(\frac{i\pi}{2}\right)^2 + \frac{1}{3!} 4 \left(\frac{i\pi}{2}\right)^3 - \dots \right) \\
 &= b_{\vec{p}}^r \left(1 - \frac{1}{2} \sum_{n=1}^{\infty} \frac{1}{n!} (i\pi)^n (-1)^{n-1} \right) + a_{\vec{p}}^r \left(\frac{1}{2} \sum_{n=1}^{\infty} \frac{1}{n!} (i\pi)^n (-1)^{n-1} \right)
 \end{aligned}$$

$$\boxed{U^\dagger b_{\vec{p}}^r U = a_{\vec{p}}^r}$$

Since $U^\dagger a_{\vec{p}}^r U = b_{\vec{p}}^r$ and $U^\dagger b_{\vec{p}}^r U = a_{\vec{p}}^r$, U produces a charge conjugation transformation, turning particles into antiparticles and vice versa.

2. USING DIRAC ANNIHILATION AND CREATION OPERATORS, EXPLICITLY CONSTRUCT A UNITARY OPERATOR P WHICH IMPLEMENTS THE PARITY TRANSFORMATION

$$P^\dagger a_{\vec{p}}^\dagger P = a_{-\vec{p}}^\dagger$$

$$P^\dagger b_{\vec{p}}^\dagger P = -b_{-\vec{p}}^\dagger$$

Let's assume that P is also an exponential, similar to U . We can use the Baker-Hausdorff lemma again to find the form of P .

$$P^\dagger a_{\vec{p}}^\dagger P = e^X a_{\vec{p}}^\dagger e^{-X} = a_{\vec{p}}^\dagger + [X, a_{\vec{p}}^\dagger] + \frac{1}{2!} [X, [X, a_{\vec{p}}^\dagger]] + \dots = a_{-\vec{p}}^\dagger$$

$$P^\dagger b_{\vec{p}}^\dagger P = e^X b_{\vec{p}}^\dagger e^{-X} = b_{\vec{p}}^\dagger + [X, b_{\vec{p}}^\dagger] + \frac{1}{2!} [X, [X, b_{\vec{p}}^\dagger]] + \dots = -b_{-\vec{p}}^\dagger$$

In problem 1, we found that

$$U^\dagger a_{\vec{p}}^\dagger U = a_{\vec{p}}^\dagger \left(1 - \frac{1}{2} \sum_{n=1}^{\infty} \frac{(i\pi)^n}{n!} (-1)^{n-1}\right) + b_{\vec{p}}^\dagger \left(\frac{1}{2} \sum_{n=1}^{\infty} \frac{(i\pi)^n}{n!} (-1)^{n-1}\right)$$

and

$$\left[\frac{i\pi}{2} \sum \frac{d^3 p}{(2\pi)^3} (a_{\vec{p}}^{\dagger\dagger} - b_{\vec{p}}^{\dagger\dagger}) (a_{\vec{p}}^\dagger - b_{\vec{p}}^\dagger), a_{\vec{q}}^\dagger\right] = -\left(\frac{i\pi}{2}\right) (a_{\vec{q}}^\dagger - b_{\vec{q}}^\dagger)$$

Let's try to find an X that gives

$$[X, a_{\vec{p}}^\dagger] = -\left(\frac{i\pi}{2}\right) (a_{\vec{p}}^\dagger - a_{-\vec{p}}^\dagger)$$



$$[X, b_{\vec{p}}^\dagger] = -\left(\frac{i\pi}{2}\right) (b_{\vec{p}}^\dagger + b_{-\vec{p}}^\dagger)$$

Where the difference in sign is necessary to get the correct sign for the transformation.

Since the a and b stay separate, I will keep them separate in the operator (no terms like $a_{\vec{p}}^\dagger b_{\vec{p}}^\dagger$ etc.). This way, the a terms will be ignored when transforming $b_{\vec{p}}^\dagger$ and the b terms will be ignored when transforming $a_{\vec{p}}^\dagger$.



If the commutators follow the same pattern as in problem 1 (extra factor of $-2\left(\frac{i\pi}{2}\right)$ on each subsequent commutator), we'll end up with the same sums as in problem 1 and should have our operator.

2 CONTINUED

The simplest operator to try first would be

$$P = \exp \left[-\frac{i\pi}{2} \sum \int \frac{d^3p}{(2\pi)^3} (a_{\vec{p}}^{r\dagger} (a_{\vec{p}}^r - a_{-\vec{p}}^r) + b_{\vec{p}}^{s\dagger} (b_{\vec{p}}^s + b_{-\vec{p}}^s)) \right]$$

Look at $a_{\vec{p}}^r$ first.

this must be hermitian. see last page.

$$\left[\frac{i\pi}{2} \sum \int \frac{d^3q}{(2\pi)^3} (a_{\vec{q}}^{s\dagger} (a_{\vec{q}}^s - a_{-\vec{q}}^s) + b_{\vec{q}}^{s\dagger} (b_{\vec{q}}^s + b_{-\vec{q}}^s)), a_{\vec{p}}^r \right]$$

$$= \left[\frac{i\pi}{2} \sum \int \frac{d^3q}{(2\pi)^3} (a_{\vec{q}}^{s\dagger} (a_{\vec{q}}^s - a_{-\vec{q}}^s)), a_{\vec{p}}^r \right]$$

$$+ \left[\frac{i\pi}{2} \sum \int \frac{d^3q}{(2\pi)^3} (b_{\vec{q}}^{s\dagger} (b_{\vec{q}}^s + b_{-\vec{q}}^s)), a_{\vec{p}}^r \right]$$

→ 0 (after expansion into anticommutators, $\{a_{\vec{p}}^r, b_{\vec{q}}^{s\dagger}\} = 0$, $\{a_{\vec{p}}^r, b_{\vec{q}}^{s\dagger}\} = 0$)

$$= \left(\frac{i\pi}{2} \right) \sum \int \frac{d^3q}{(2\pi)^3} (a_{\vec{q}}^{s\dagger} \{a_{\vec{q}}^s - a_{-\vec{q}}^s, a_{\vec{p}}^r\} - \{a_{\vec{q}}^{s\dagger}, a_{\vec{p}}^r\} (a_{\vec{q}}^s - a_{-\vec{q}}^s))$$

$$= \left(\frac{i\pi}{2} \right) \sum \int \frac{d^3q}{(2\pi)^3} (- (2\pi)^3 \delta^{rs} \delta^3(\vec{p} - \vec{q}) (a_{\vec{q}}^s - a_{-\vec{q}}^s))$$

$$= - \left(\frac{i\pi}{2} \right) (a_{\vec{p}}^r - a_{-\vec{p}}^r)$$

Now let's check that the next commutator has an extra factor of $-2 \left(\frac{i\pi}{2} \right)$. Since I know the b part falls out, I will ignore it from the start.

$$\left[\frac{i\pi}{2} \sum \int \frac{d^3q}{(2\pi)^3} (a_{\vec{q}}^{s\dagger} (a_{\vec{q}}^s - a_{-\vec{q}}^s)), \left(-\frac{i\pi}{2} \right) (a_{\vec{p}}^r - a_{-\vec{p}}^r) \right]$$

$$= - \left(\frac{i\pi}{2} \right)^2 \sum \int \frac{d^3q}{(2\pi)^3} (a_{\vec{q}}^{s\dagger} \{a_{\vec{q}}^s - a_{-\vec{q}}^s, a_{\vec{p}}^r - a_{-\vec{p}}^r\} - \{a_{\vec{q}}^{s\dagger}, (a_{\vec{p}}^r - a_{-\vec{p}}^r)\} (a_{\vec{q}}^s - a_{-\vec{q}}^s))$$

$$= \left(\frac{i\pi}{2} \right)^2 \sum \int \frac{d^3q}{(2\pi)^3} ((2\pi)^3 \delta^{rs} \delta^3(\vec{p} - \vec{q}) - (2\pi)^3 \delta^{rs} \delta^3(-\vec{p} - \vec{q})) (a_{\vec{q}}^s - a_{-\vec{q}}^s)$$

$$= \left(\frac{i\pi}{2} \right)^2 ((a_{\vec{p}}^r - a_{-\vec{p}}^r) - (a_{-\vec{p}}^r - a_{\vec{p}}^r))$$

$$= 2 \left(\frac{i\pi}{2} \right)^2 (a_{\vec{p}}^r - a_{-\vec{p}}^r)$$

good

So, the pattern does hold.

$$\begin{aligned} P^\dagger a_p^r P &= a_p^r - \frac{1}{2} (in) (a_p^r - a_{-p}^r) + \frac{1}{2!} \frac{1}{2} (in)^2 (a_p^r - a_{-p}^r) + \dots \\ &= a_p^r \left(1 - \underbrace{\frac{1}{2} \sum_{n=1}^{\infty} \frac{(in)^n}{n!} (-1)^{n-1}}_{=+2} \right) - a_{-p}^r \left(\underbrace{\frac{1}{2} \sum_{n=1}^{\infty} \frac{(in)^n}{n!} (-1)^n}_{=-2} \right) \end{aligned}$$

$$P^\dagger a_{\vec{p}}^\dagger P = a_{-\vec{p}}^\dagger$$

Now, we need to check that b_j^r transforms correctly. Again, since I know the a terms will drop out, I will ignore them.

$$\begin{aligned}
 & \left[\frac{i\pi}{2} \Sigma \int \frac{d^3 q}{(2\pi)^3} (b_{\vec{q}}^{s\dagger} (b_{\vec{q}}^s + b_{-\vec{q}}^s), b_{\vec{p}}^{r\dagger}) \right] \\
 &= \left(\frac{i\pi}{2} \right) \Sigma \int \frac{d^3 q}{(2\pi)^3} (b_{\vec{q}}^{s\dagger} \{ \cancel{(b_{\vec{q}}^s + b_{-\vec{q}}^s)} \}, b_{\vec{p}}^{r\dagger}) - \{ b_{\vec{q}}^{s\dagger}, b_{\vec{p}}^{r\dagger} \} (b_{\vec{q}}^s + b_{-\vec{q}}^s)) \\
 &= \left(\frac{i\pi}{2} \right) \Sigma \int \frac{d^3 q}{(2\pi)^3} (-(2\pi)^3 \delta^{rs} \delta^3(\vec{p} - \vec{q}) (b_{\vec{q}}^s + b_{-\vec{q}}^s)) \\
 &= - \left(\frac{i\pi}{2} \right) (b_{\vec{p}}^{r\dagger} + b_{-\vec{p}}^{r\dagger})
 \end{aligned}$$

$$\begin{aligned}
 & \left[\frac{i\pi}{2} \Sigma \int \frac{d^3 q}{(2\pi)^3} (b_{\vec{q}}^{s+} (b_{\vec{q}}^{s+} + b_{-\vec{q}}^{s+})) , - \left(\frac{i\pi}{2} \right) (b_{\vec{p}}^{r+} + b_{-\vec{p}}^{r+}) \right] \\
 &= - \left(\frac{i\pi}{2} \right)^2 \Sigma \int \frac{d^3 q}{(2\pi)^3} (b_{\vec{q}}^{s+} \{ \cancel{(b_{\vec{q}}^{s+} + b_{-\vec{q}}^{s+})}, (b_{\vec{p}}^{r+} + b_{-\vec{p}}^{r+}) \}^0 \\
 &\quad - \{ b_{\vec{q}}^{s+}, (b_{\vec{p}}^{r+} + b_{-\vec{p}}^{r+}) \} (b_{\vec{q}}^{s+} + b_{-\vec{q}}^{s+})) \\
 &= \left(\frac{i\pi}{2} \right)^2 \Sigma \int \frac{d^3 q}{(2\pi)^3} ((2\pi)^3 \delta^{rs} \delta^3(\vec{p} - \vec{q}) + (2\pi)^3 \delta^{rs} \delta^3(-\vec{p} - \vec{q})) (b_{\vec{q}}^{s+} + b_{-\vec{q}}^{s+}) \\
 &= \left(\frac{i\pi}{2} \right)^2 ((b_{\vec{p}}^{r+} + b_{-\vec{p}}^{r+}) + (b_{-\vec{p}}^{r+} + b_{\vec{p}}^{r+})) \\
 &= 2 \left(\frac{i\pi}{2} \right)^2 (b_{\vec{p}}^{r+} + b_{-\vec{p}}^{r+})
 \end{aligned}$$

2 CONTINUED.

Again, our pattern holds

$$P^\dagger b_{\vec{p}}^r P = b_{\vec{p}}^r - \frac{1}{2} (i\pi) (b_{\vec{p}}^r + b_{-\vec{p}}^r) + \frac{1}{2!} \frac{1}{2} (i\pi)^2 (b_{\vec{p}}^r + b_{-\vec{p}}^r) + \dots$$

$$= b_{\vec{p}}^r \left(1 - \frac{1}{2} \sum_{n=1} \frac{(i\pi)^n}{n!} (-1)^{n-1} \right) + b_{-\vec{p}}^r \left(\frac{1}{2} \sum_{n=1} \frac{(i\pi)^n}{n!} (-1)^{n-1} \right)$$

$$P^\dagger b_{\vec{p}}^r P = -b_{\vec{p}}^r$$

So, this operator does perform the parity transformation correctly.

$$P = \exp \left[-\frac{i\pi}{2} \sum \int \frac{d^3 p}{(2\pi)^3} \left(a_{\vec{p}}^{r\dagger} (a_{\vec{p}}^r - a_{-\vec{p}}^r) + b_{\vec{p}}^{r\dagger} (b_{\vec{p}}^r + b_{-\vec{p}}^r) \right) \right]$$

Is $\sum \int \frac{d^3 p}{(2\pi)^3} (a_{\vec{p}}^{r\dagger} (a_{\vec{p}}^r - a_{-\vec{p}}^r) + b_{\vec{p}}^{r\dagger} (b_{\vec{p}}^r + b_{-\vec{p}}^r))$ Hermitian?

$$\left(\sum \int \frac{d^3 p}{(2\pi)^3} (a_{\vec{p}}^{r\dagger} (a_{\vec{p}}^r - a_{-\vec{p}}^r) + b_{\vec{p}}^{r\dagger} (b_{\vec{p}}^r + b_{-\vec{p}}^r)) \right)^\dagger$$

$$= \sum \int \frac{d^3 p}{(2\pi)^3} \left((a_{\vec{p}}^{r\dagger})^\dagger (a_{\vec{p}}^r - a_{-\vec{p}}^r)^\dagger + (b_{\vec{p}}^{r\dagger})^\dagger (b_{\vec{p}}^r + b_{-\vec{p}}^r)^\dagger \right)$$

$$= \sum \int \frac{d^3 p}{(2\pi)^3} \left(a_{\vec{p}}^r a_{\vec{p}}^{r\dagger} - \underbrace{a_{\vec{p}}^{r\dagger} a_{-\vec{p}}^r}_{\substack{\text{Since we are integrating over all } \vec{p}, \text{ the signs on} \\ \text{these momenta can both switch and everything} \\ \text{remains the same.}}} + b_{\vec{p}}^r b_{\vec{p}}^{r\dagger} + \underbrace{b_{-\vec{p}}^{r\dagger} b_{\vec{p}}^r}_{\substack{\text{Since we are integrating over all } \vec{p}, \text{ the signs on} \\ \text{these momenta can both switch and everything} \\ \text{remains the same.}}} \right)$$

$$= \sum \int \frac{d^3 p}{(2\pi)^3} \left(a_{\vec{p}}^r a_{\vec{p}}^{r\dagger} - a_{\vec{p}}^{r\dagger} a_{-\vec{p}}^r + b_{\vec{p}}^r b_{\vec{p}}^{r\dagger} + b_{\vec{p}}^{r\dagger} b_{-\vec{p}}^r \right)$$

$$= \sum \int \frac{d^3 p}{(2\pi)^3} \left(a_{\vec{p}}^{r\dagger} (a_{\vec{p}}^r - a_{-\vec{p}}^r) + b_{\vec{p}}^{r\dagger} (b_{\vec{p}}^r + b_{-\vec{p}}^r) \right)$$

Since taking the Hermitian conjugate left this unchanged, it is indeed Hermitian, and P is unitary.

PY111 Solutions #7

1. We start with some general observations. Consider two operators S, T such that $[S, T] = cT$ for some constant c . Then

$$ST = TS + cT = T(S + c)$$

$$S^2T = SST = ST(S + c) = T(S + c)^2$$

In general we can prove by induction that $S^nT = T(S + c)^n$.

So

$$\begin{aligned} e^{iS}T &= T e^{i(S+c)} \\ \Rightarrow e^{iS}T e^{-iS} &= e^{+ic}T \end{aligned}$$

For our problem $S = \frac{\pi}{2} \sum_{r=1,2} \int \frac{d^3p}{(2\pi)^3} (a_p^{r\dagger} - b_p^{r\dagger})(a_p^r - b_p^r)$.

Let $C_p^r = \frac{1}{\sqrt{2}}(a_p^r - b_p^r)$ and $d_p^r = \frac{1}{\sqrt{2}}(a_p^r + b_p^r)$. Note that

$$\{C_p^r, C_{p'}^{r'}\} = (2\pi)^3 \delta_{\vec{p}, -\vec{p}'} \delta_{r, r'} \quad \{d_p^r, d_{p'}^{r'}\}$$

and all other commutators vanish. Notice that

$$S = \pi \sum_{r=1,2} \int \frac{d^3p}{(2\pi)^3} C_p^{r\dagger} C_p^r,$$

which is π times the number of " c "-type quanta.

We have

$$\begin{aligned} [S, C_p^r] &= \pi \sum_{r'=1,2} \int \frac{d^3p'}{(2\pi)^3} (C_{p'}^{r'\dagger} C_{p'}^{r'} C_p^r - C_p^r C_{p'}^{r'\dagger} C_{p'}^{r'}) \\ &= \pi \sum_{r'=1,2} \int \frac{d^3p'}{(2\pi)^3} \left[(2\pi)^3 \delta_{\vec{p}', -\vec{p}} \delta_{r', r} (-C_{p'}^{r'}) + \underbrace{(C_{p'}^{r'\dagger} C_{p'}^{r'} C_p^r)}_{\text{cancel}} + \underbrace{(C_p^r C_{p'}^{r'\dagger} C_{p'}^{r'})}_{\text{cancel}} \right] \\ &= -\pi C_p^r \end{aligned}$$

Also $[S, d_p^r] = 0$. Therefore $e^{iS} C_p^r e^{-iS} = e^{-i\pi} C_p^r = -C_p^r$, and $e^{iS} d_p^r e^{-iS} = d_p^r$. Since $a_p^r = \frac{1}{\sqrt{2}}(C_p^r + d_p^r)$, $b_p^r = \frac{1}{\sqrt{2}}(-C_p^r + d_p^r)$,

$$\begin{aligned} e^{iS} a_p^r e^{-iS} &= \frac{1}{\sqrt{2}}(-C_p^r + d_p^r) = b_p^r \\ e^{iS} b_p^r e^{-iS} &= \frac{1}{\sqrt{2}}(C_p^r + d_p^r) = a_p^r. \end{aligned}$$

The transformation corresponds with charge conjugation.

2. Let us define the operators

$$c_p^r = \frac{1}{\sqrt{2}}(a_p^r - a_{-p}^r), \quad d_p^r = \frac{1}{\sqrt{2}}(a_p^r + a_{-p}^r) \\ f_p^r = \frac{1}{\sqrt{2}}(b_p^r - b_{-p}^r), \quad g_p^r = \frac{1}{\sqrt{2}}(b_p^r + b_{-p}^r)$$

$$\text{Then } \{c_p^r, c_{p'}^{r'}\} = \{g_p^r, g_{p'}^{r'}\} = \delta_{pp'} \delta_{rr'} [\delta^{(1)}(p-p') - \delta^{(1)}(p+p')] \\ \{d_p^r, d_{p'}^{r'}\} = \{f_p^r, f_{p'}^{r'}\} = (\pi)^3 \delta_{pp'} \delta_{rr'} [\delta^{(1)}(p-p') + \delta^{(1)}(p+p')]$$

Note also that $c_p^r = -c_p^r$, $d_p^r = d_p^r$, $f_p^r = -f_p^r$, $g_p^r = g_p^r$.
All other anticommutators vanish.

$$\text{Let } S_c = \frac{\pi}{2} \sum_{r=1,2} \int \frac{d^3p}{(2\pi)^3} c_p^{r\dagger} c_p^r \quad \text{and} \quad S_f = \frac{\pi}{2} \sum_{r=1,2} \int \frac{d^3p}{(2\pi)^3} f_p^{r\dagger} f_p^r.$$

Then S_c and S_f commute with each other and also

$$[S_c, c_p^r] = \frac{\pi}{2}(-c_p^r + c_p^r) = -\pi c_p^r \\ [S_c, d_p^r] = [S_c, f_p^r] = [S_c, g_p^r] = 0 \\ [S_f, f_p^r] = \frac{\pi}{2}(-f_p^r + f_p^r) = -\pi f_p^r \\ [S_f, c_p^r] = [S_f, d_p^r] = [S_f, g_p^r] = 0$$

So let $U = e^{-iS_c - iS_f}$. Then

$$U^\dagger c_p^r U = e^{-i\pi} c_p^r = -c_p^r, \quad U^\dagger d_p^r U = d_p^r, \\ U^\dagger f_p^r U = e^{-i\pi} f_p^r = -f_p^r, \quad U^\dagger g_p^r U = g_p^r,$$

This produces the parity transformation required,

$$U^\dagger a_p^r U = a_{-p}^r, \quad U^\dagger b_p^r U = -b_{-p}^r$$

which are explicitly

$$U = \exp \left[-\frac{i\pi}{4} \sum_{r=1,2} \int \frac{d^3p}{(2\pi)^3} (a_p^{r\dagger} - a_{-p}^{r\dagger})(a_p^r - a_{-p}^r) \right. \\ \left. - \frac{i\pi}{4} \sum_{r=1,2} \int \frac{d^3p}{(2\pi)^3} (b_p^{r\dagger} + b_{-p}^{r\dagger})(b_p^r + b_{-p}^r) \right]$$