

# Notes 8.370/18.435 Fall 2022

## Lecture 5 Prof. Peter Shor

In the last few lectures, we saw how to operate on one qubit—or more generally, one  $d$ -dimensional quantum system. We studied how to transform quantum systems with a unitary matrix, and how to measure them via a complete set of projection matrices. In this lecture, we explain how the state space of a joint quantum system, composed of two individual quantum systems, is constructed.

Suppose we have two quantum systems. Each of these has a state space which is a complex vector space of dimensions  $d_1$  and  $d_2$ , respectively. When we consider these two quantum systems together, we get a state space which is the tensor product of their individual state spaces, and which has dimension  $d_1 d_2$ .

Some of you probably haven't seen tensor products before. We will show how it works by example. If you have two qubits, each of which has a state space with basis  $\{|0\rangle, |1\rangle\}$ , then the state space of the joint system has basis  $\{|0\rangle \otimes |0\rangle, |0\rangle \otimes |1\rangle, |1\rangle \otimes |0\rangle, |1\rangle \otimes |1\rangle\}$ . The convention is to write the basis in lexicographical order. Thus, if you have two qubits in state  $\alpha_0 |0\rangle + \alpha_1 |1\rangle$  and  $\beta_0 |0\rangle + \beta_1 |1\rangle$ , the system of both qubits is in state

$$\alpha_0 \beta_0 |0\rangle \otimes |0\rangle + \alpha_0 \beta_1 |0\rangle \otimes |1\rangle + \alpha_1 \beta_0 |1\rangle \otimes |0\rangle + \alpha_1 \beta_1 |1\rangle \otimes |1\rangle,$$

so the normal distributed law applies to tensor products. Often, rather than writing  $|0\rangle \otimes |1\rangle$ , we will write this as  $|01\rangle$ .

To illustrate tensor products using a more usual vector notation,

$$\begin{pmatrix} \beta_0 \\ \beta_1 \end{pmatrix} \otimes \begin{pmatrix} \alpha_0 \\ \alpha_1 \end{pmatrix} = \begin{pmatrix} \alpha_0 \beta_0 \\ \alpha_0 \beta_1 \\ \alpha_1 \beta_0 \\ \alpha_1 \beta_1 \end{pmatrix}$$

What if you have  $n$  qubits? The state space has dimension  $2^n$ , and has basis  $|00 \dots 00\rangle, |00 \dots 01\rangle, |00 \dots 10\rangle, \dots$  (Recall that by notation,  $|00 \dots 00\rangle = |0\rangle \otimes |0\rangle \otimes \dots \otimes |0\rangle \otimes |0\rangle$ .)

Now let us count up (real) degrees of freedom. A qubit has a state  $\alpha |0\rangle + \beta |1\rangle$ , and since  $\alpha$  and  $\beta$  are complex numbers, this would give four real degrees of freedom. However, it also satisfies  $\alpha^2 + \beta^2 = 1$ , which means that it only has three real degrees of freedom (if we use the fact that multiplying by a global phase leaves the state essentially unchanged, there are only two degrees of freedom). However, the joint state of two qubits has four complex coefficients, leaving 7 (or 6) degrees of freedom. Thus, because  $7 > 2 \cdot 3$ , there are some states of the joint system which cannot be the tensor product of states of the individual systems. These states are called *entangled*. In fact, because the number of degrees of freedom of tensor product states is less than that of entangled states, the tensor product states form a lower-dimensional manifold in the entangled states, and we see that most quantum states are entangled.

For example, the state

$$\frac{1}{\sqrt{3}} |00\rangle + \frac{1}{\sqrt{3}} |01\rangle + \frac{1}{\sqrt{6}} |10\rangle + \frac{1}{\sqrt{6}} |11\rangle = \left( \frac{\sqrt{2}}{\sqrt{3}} |0\rangle + \frac{1}{\sqrt{3}} |1\rangle \right) \left( \frac{1}{\sqrt{2}} |0\rangle + \frac{1}{\sqrt{2}} |1\rangle \right)$$

is a tensor product state, while the state

$$\frac{1}{\sqrt{3}}|00\rangle + \frac{1}{\sqrt{3}}|01\rangle + \frac{1}{\sqrt{3}}|10\rangle$$

is entangled. You can see that the second one is not a tensor product state, and thus entangled, because if we could represent it as a tensor product state

$$\alpha_0\beta_0|00\rangle + \alpha_0\beta_1|01\rangle + \alpha_1\beta_0|10\rangle + \alpha_1\beta_1|11\rangle,$$

then we would need  $\alpha_1\beta_1 = 0$ . But if that holds, either  $\alpha_1$  or  $\beta_1$  is 0, so one of  $\alpha_0\beta_1$  and  $\alpha_1\beta_0$  must be 0, which is not the case.

What about unitary transformations on a joint state space. Suppose we have two unitary matrices  $U$  and  $V$ , of dimensions  $k$  and  $\ell$ . The tensor product of them is the  $k\ell \times k\ell$  matrix

$$U \otimes V = \begin{pmatrix} u_{11}V & u_{12}V & \dots & u_{1k}V \\ u_{21}V & u_{22}V & \dots & u_{2k}V \\ \vdots & \vdots & \ddots & \vdots \\ u_{k1}V & u_{k2}V & \dots & u_{kk}V \end{pmatrix},$$

where  $u_{ij}V$  is the  $(i, j)$  entry of  $U$  multiplied by  $V$ .

Let's do an example and figure out what  $\sigma_x \otimes \sigma_z$  is. We have

$$\begin{aligned} \sigma_x \otimes \sigma_z &= \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \otimes \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \\ &= \left( \begin{array}{cc|cc} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \\ \hline 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \end{array} \right) \end{aligned}$$

If the input is in a tensor product state  $|\phi_1\rangle \otimes |\phi_2\rangle$ , and you apply a tensor product unitary  $U_1 \otimes U_2$ , then the output is also in a tensor product state, namely  $U_1|\phi_1\rangle \otimes U_2|\phi_2\rangle$ . So to get an entangled output from a tensor product input, you need to apply a unitary that is a tensor product. One non-tensor-product unitary that we will be using extensively in this class is the CNOT.

We have

$$\begin{aligned} CNOT &= \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{pmatrix} \\ &= |0\rangle\langle 0| \otimes I + |1\rangle\langle 1| \otimes \sigma_x \end{aligned}$$

The last description says that if the first qubit is a  $|0\rangle$ , then we apply an identity to the second qubit, and if the first qubit is a 1, we apply a  $\sigma_x$  (or NOT gate) to the second qubit. This is why it's called a controlled Not — depending on the state of the first qubit, we apply a NOT gate to the second one (or we don't).

Note that we don't measure the first qubit — we we don't get a classical record of the state of the first qubit, and a state that was in a superposition of the first qubit being  $|0\rangle$  and  $|1\rangle$  remains in a superposition. Let's do an example. Suppose we start with the state  $|+\rangle|0\rangle = \frac{1}{\sqrt{2}}(|00\rangle + |10\rangle)$ . When we apply the CNOT gate, the  $|00\rangle$  remains unchanged, because the first qubit is  $|0\rangle$ , but the  $|10\rangle$  becomes  $|11\rangle$ . This means

$$\text{CNOT } |+\rangle|0\rangle = \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle),$$

so we have used the CNOT gate to entangle two qubits.