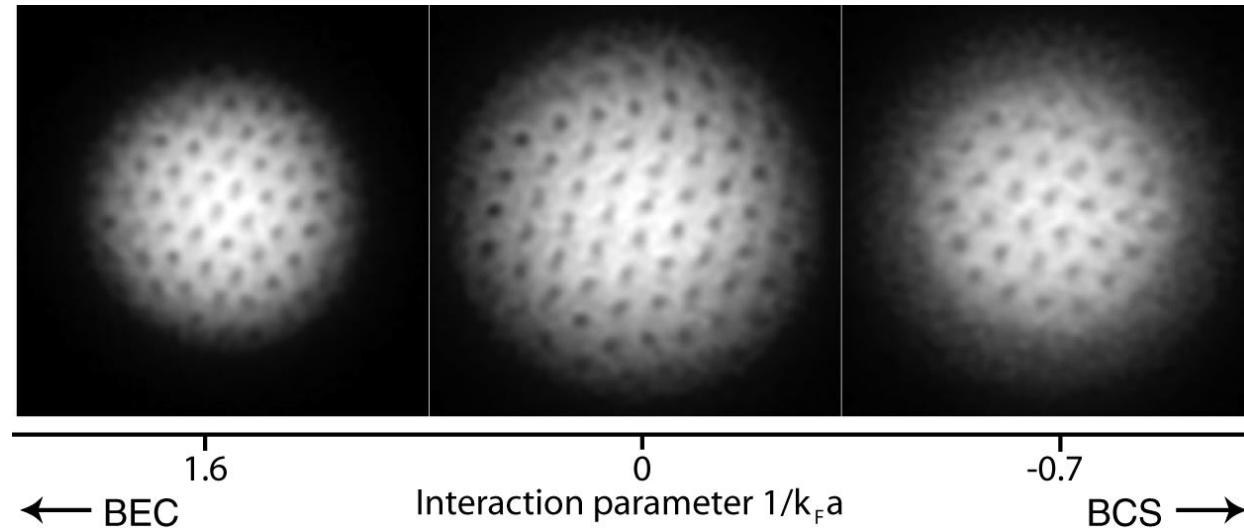


Superfluid Fermi Gases

Martin Zwierlein

Massachusetts Institute of Technology
Center for Ultracold Atoms



GORDON AND BETTY
MOORE
FOUNDATION

Some reviews

Varenna Notes:

on Fermi Gases:

Ketterle, Zwierlein, *Making, Probing and Understanding Ultracold Fermi Gases*

<http://arxiv.org/abs/0801.2500>

Zwierlein, *Thermodynamics of Strongly Interacting Fermi Gases*,

Varenna School of Physics “Enrico Fermi” 2014, vol 191

Stefano Giorgini, Lev P. Pitaevskii, Sandro Stringari

The theory of Fermi gases

<http://arxiv.org/abs/0706.3360>

Immanuel Bloch, Jean Dalibard, Wilhelm Zwerger

Many-Body Physics with Ultracold Gases:

<http://arxiv.org/abs/0704.3011>

Lecture Notes “The BEC-BCS crossover and the Unitary Fermi Gas”

Edited by W. Zwerger, Springer, 2012

How cold is ultracold?

Atoms move at:

$\sim 1 \text{ mm/s}$

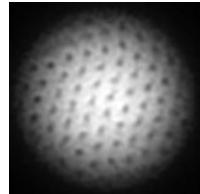
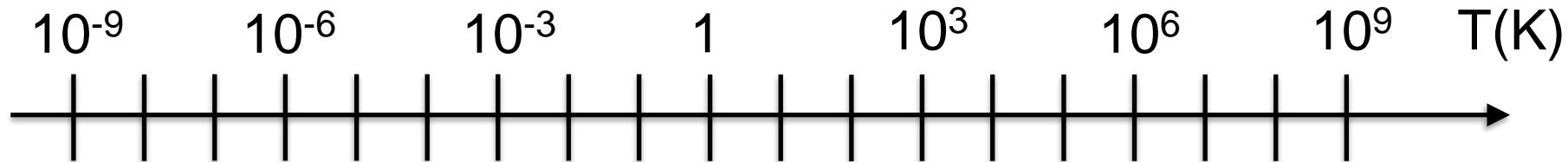


$\sim 100 \text{ m/s}$



$\sim 10^6 \text{ m/s}$

NY – Paris in 10
seconds



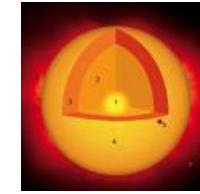
Ultracold atom
experiments



Outer
space



Your
living
room



Center
of the
sun



Supernova
explosion

Particles behave as waves

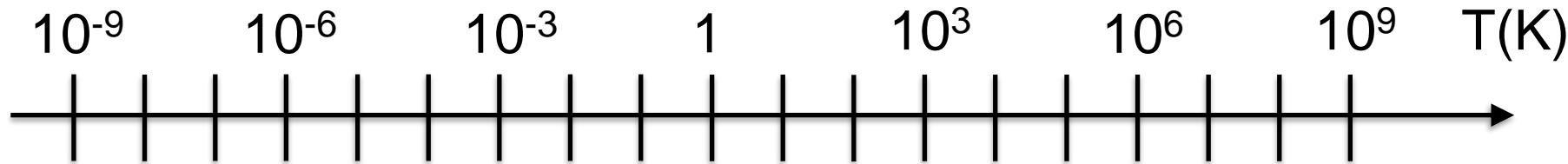


*Louis
de Broglie*

$$\lambda = \frac{h}{mv}$$

Planck's constant
Mass
Velocity

A yellow box contains the de Broglie wavelength formula. Arrows point from the text labels to the corresponding terms in the equation: 'Planck's constant' to 'h', 'Mass' to 'm', and 'Velocity' to 'v'.



*Werner
Heisenberg*

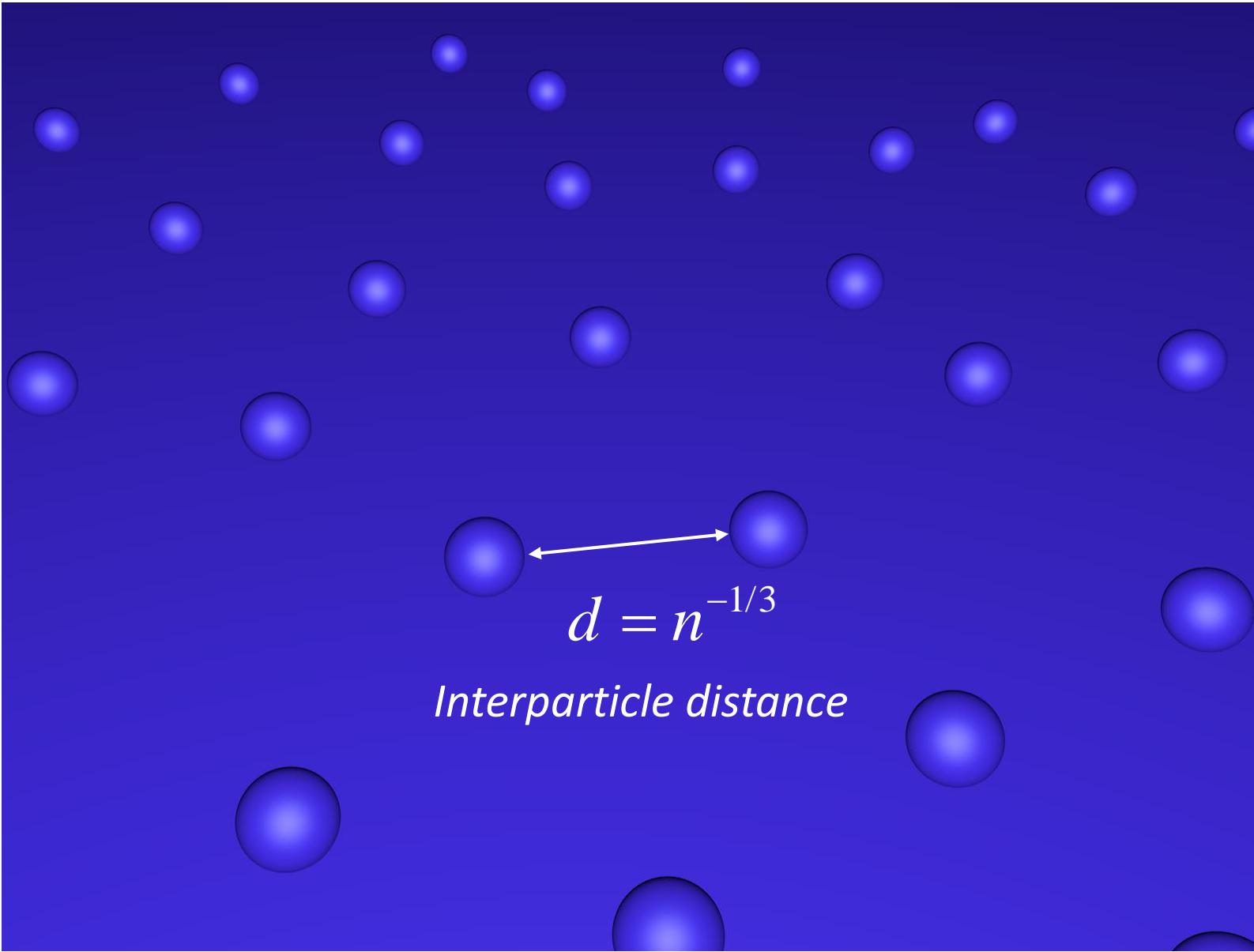
Temperature = Uncertainty of velocity²

Size of a wave packet = Uncertainty of position

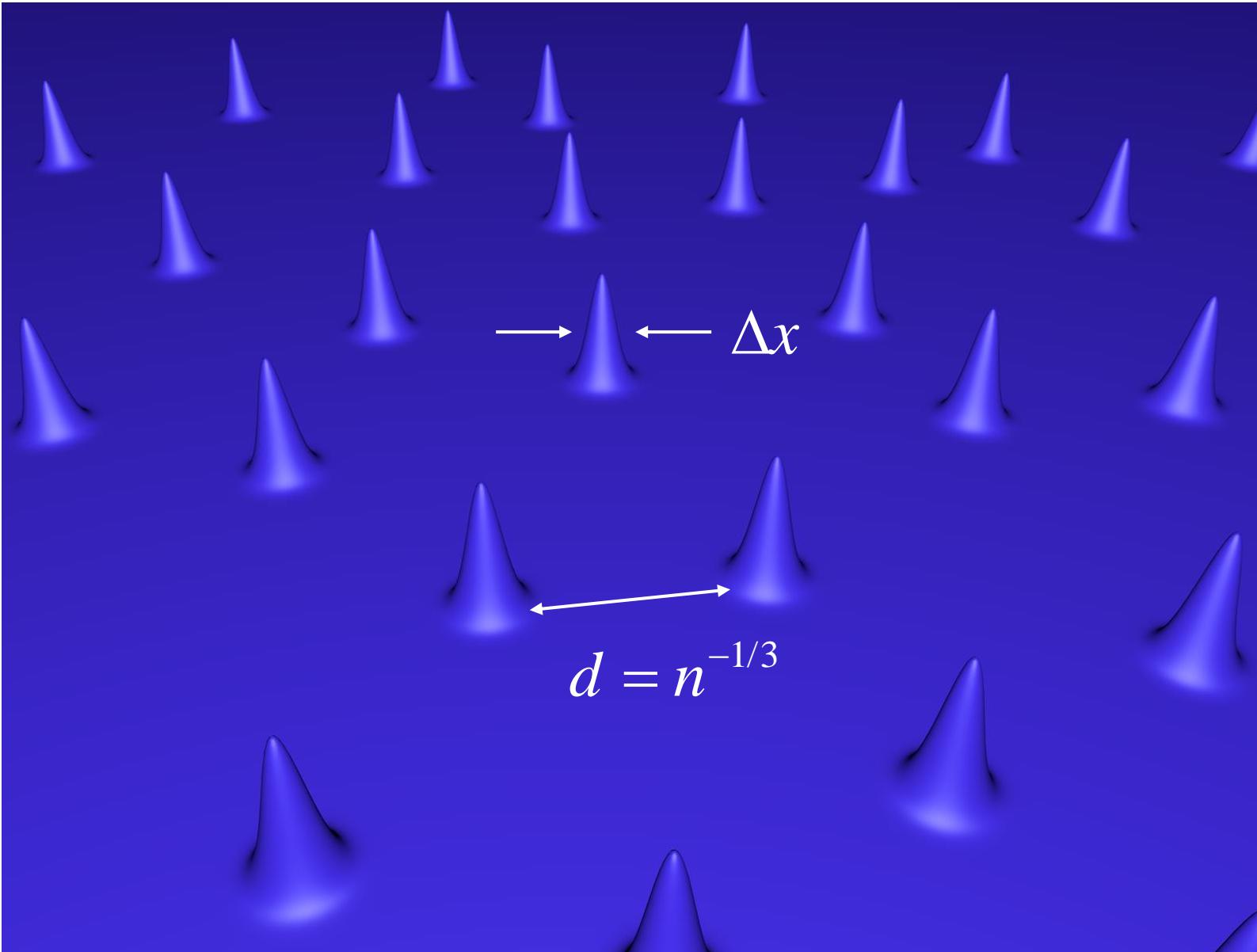
A 3D surface plot of a Gaussian function, representing the size of a wave packet. Two arrows point from the text labels to the plot: one pointing right along the horizontal axis and one pointing down along the vertical axis.

$$\Delta x \approx \frac{\hbar}{\Delta p} \approx \frac{\hbar}{\sqrt{mT}}$$

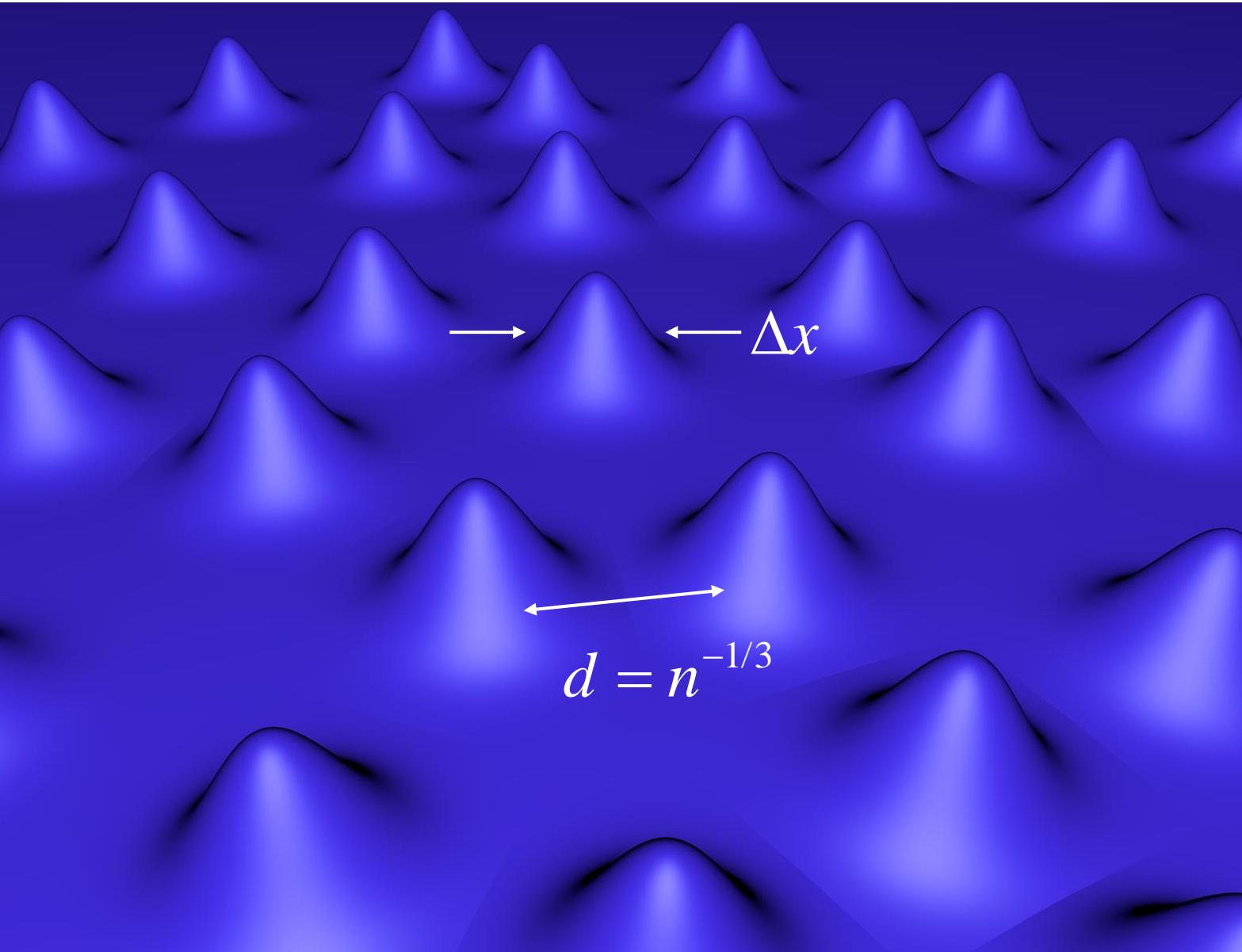
Particles...



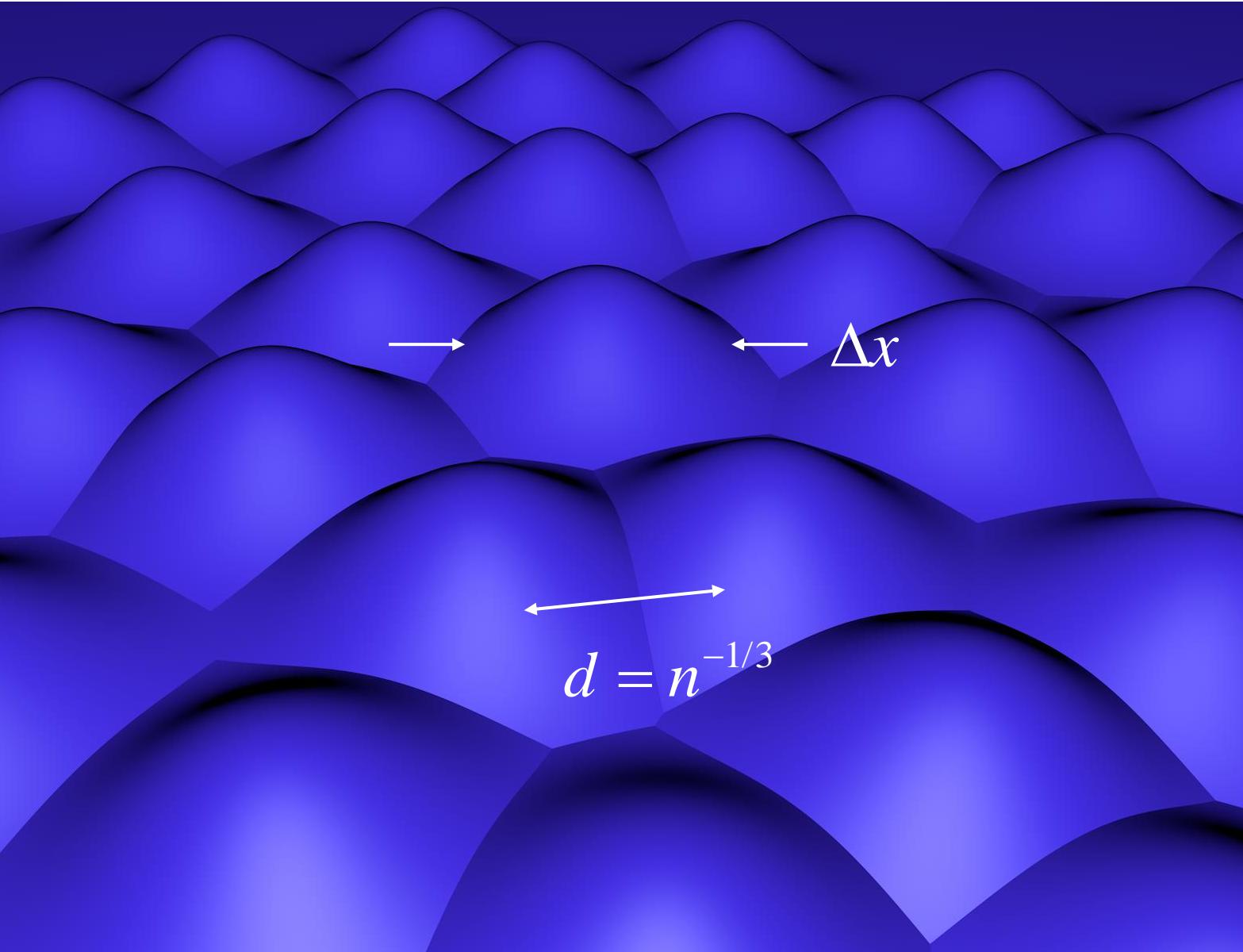
...behave as waves



When does wave mechanics matter?



When does wave mechanics matter?



Bosons versus Fermions

Fermions (unsociable):

Half-Integer Spin

Pauli blocking → Form Fermi sea

No phase transition at low Temperature



Fermi Feb. 1926

Dirac Oct. 1926

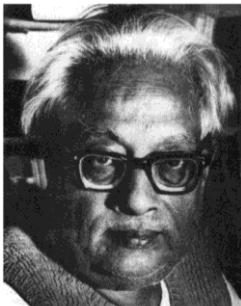
Bosons (sociable):

Integer Spin

Can share quantum states

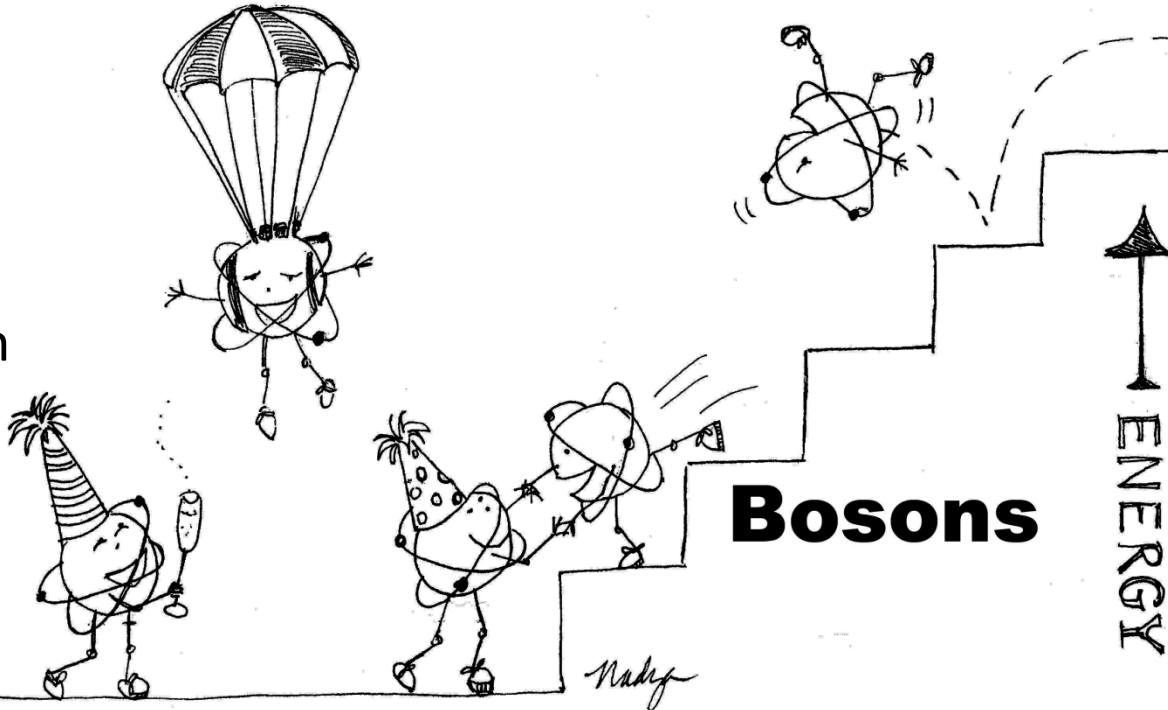
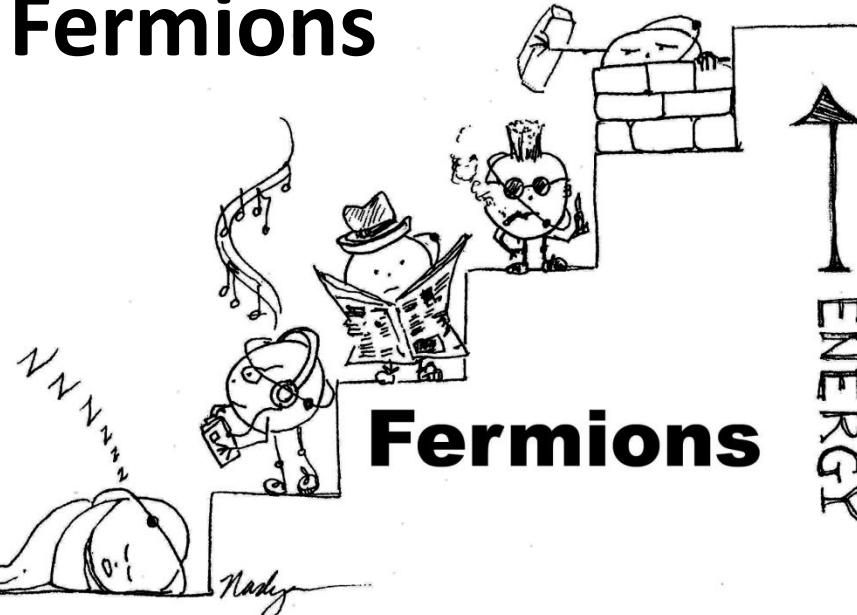
At low temperatures:

Bose-Einstein condensation



Bose 1924

Einstein 1924/25



Bosons



N bosons sharing one and the same macroscopic matter wave

(Artist's conception)

Fermions

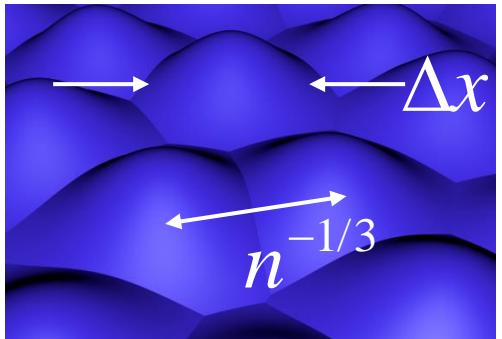


N fermions avoiding each other

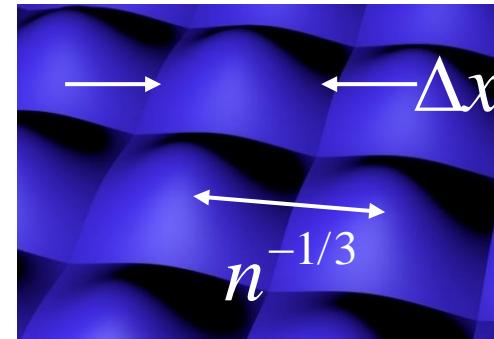
(Artist's conception)

Condition for quantum degeneracy

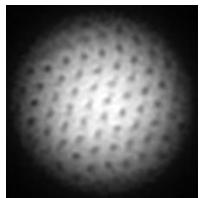
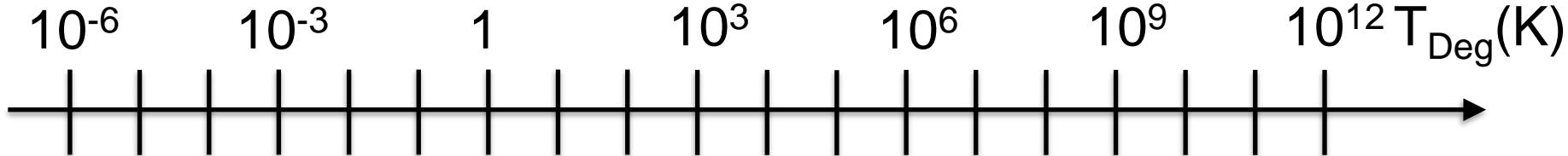
Position uncertainty \sim Interparticle spacing



$$\Delta x = \frac{\hbar}{\Delta p} \approx n^{-1/3}$$



$$k_B T_{\text{Degeneracy}} \approx \frac{\hbar^2}{m} n^{2/3} \approx E_F \quad \text{Fermi energy}$$



Ultracold
atomic gases



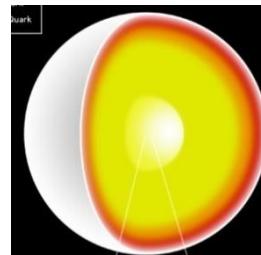
Liquid Helium



Metals

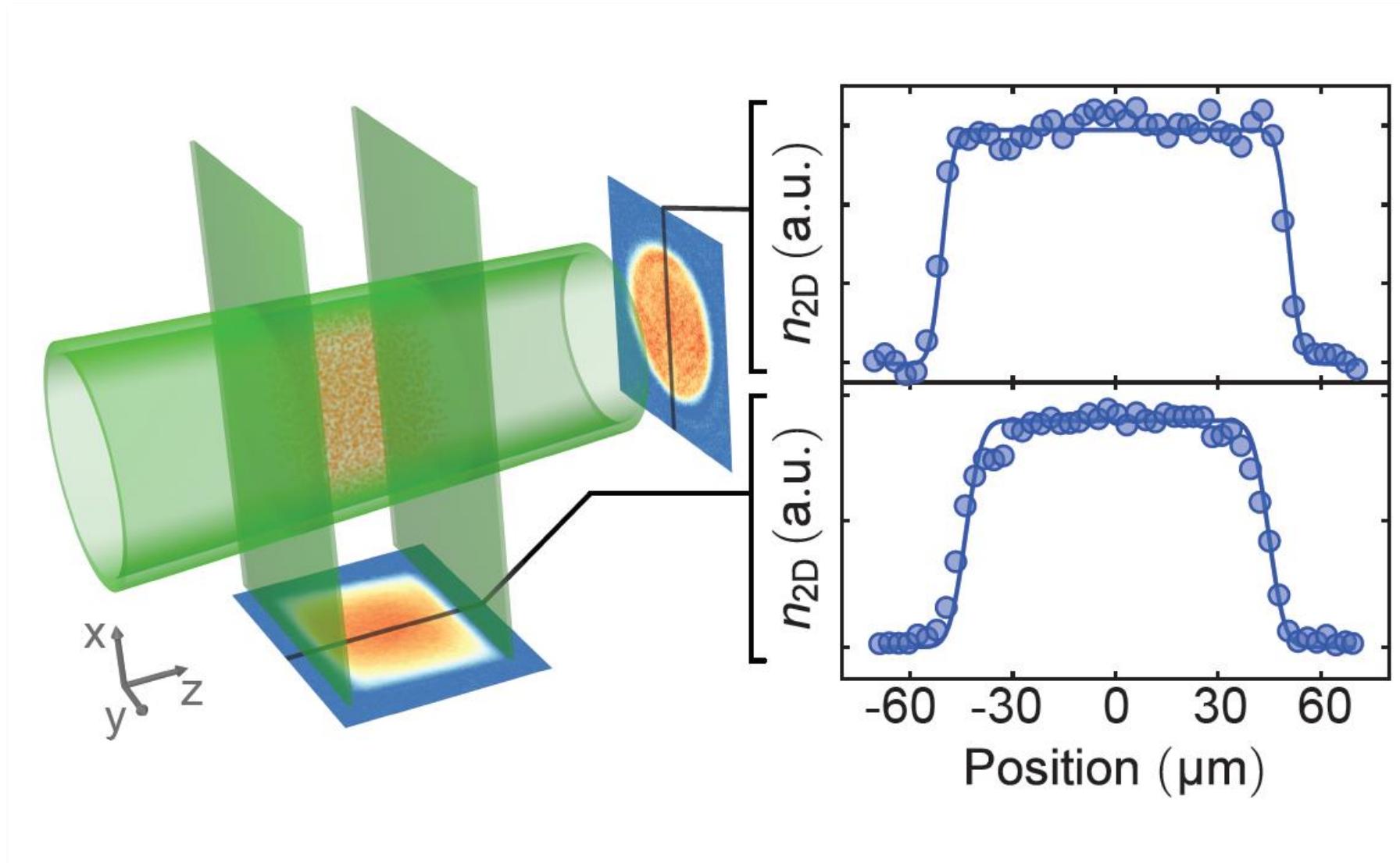


White dwarf

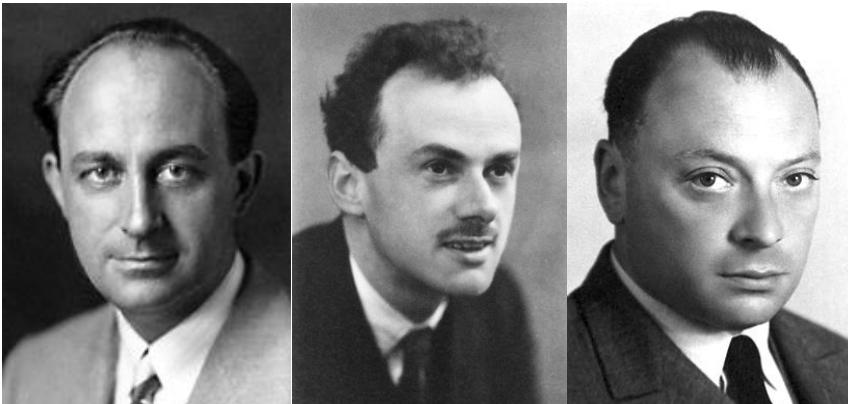
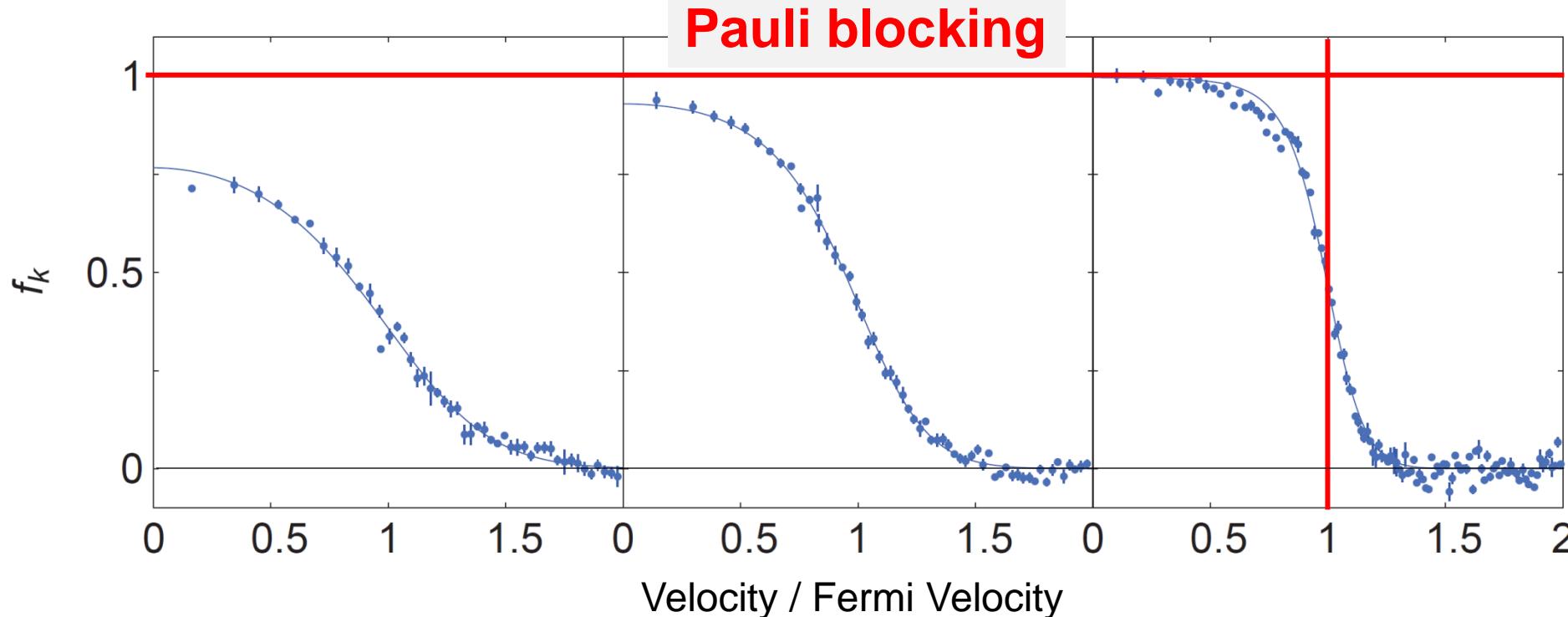


Neutron star

Fermions in a Box



Measuring the Fermi-Dirac distribution



$$f_k = \frac{1}{e^{\beta(\varepsilon_k - \mu)} + 1}$$

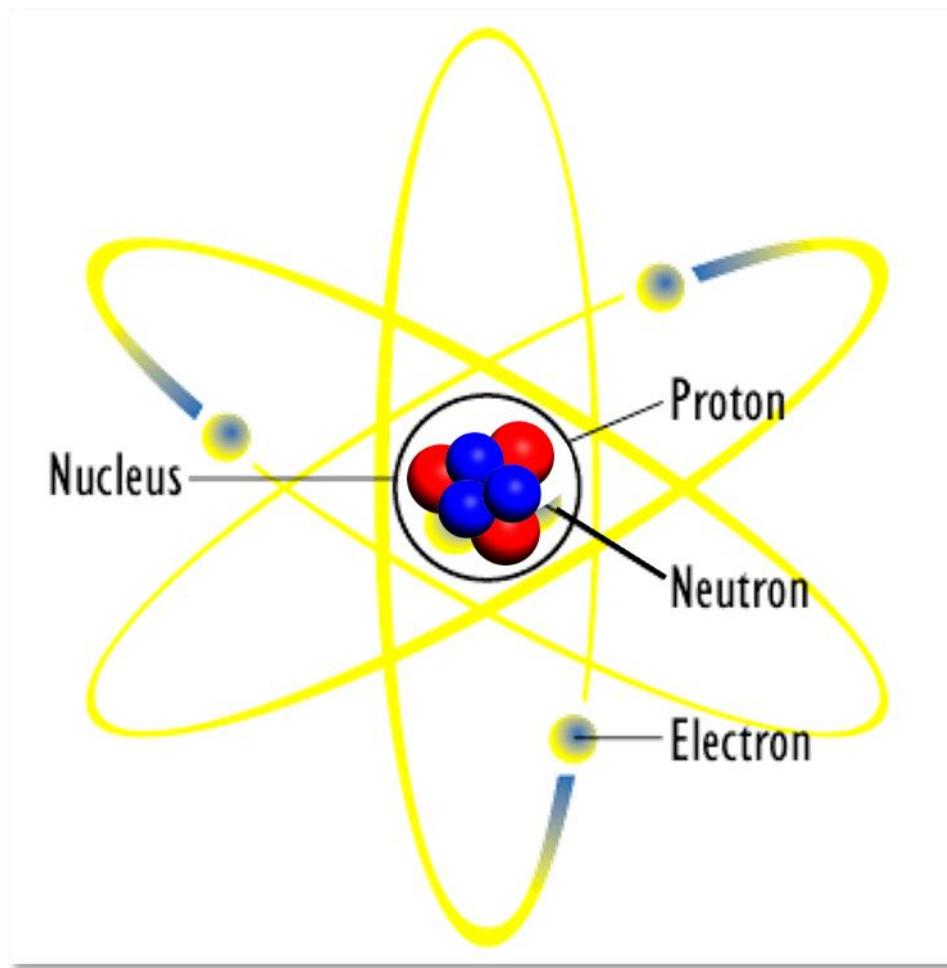
**90th
anniversary**

**Fermi surface
formation**

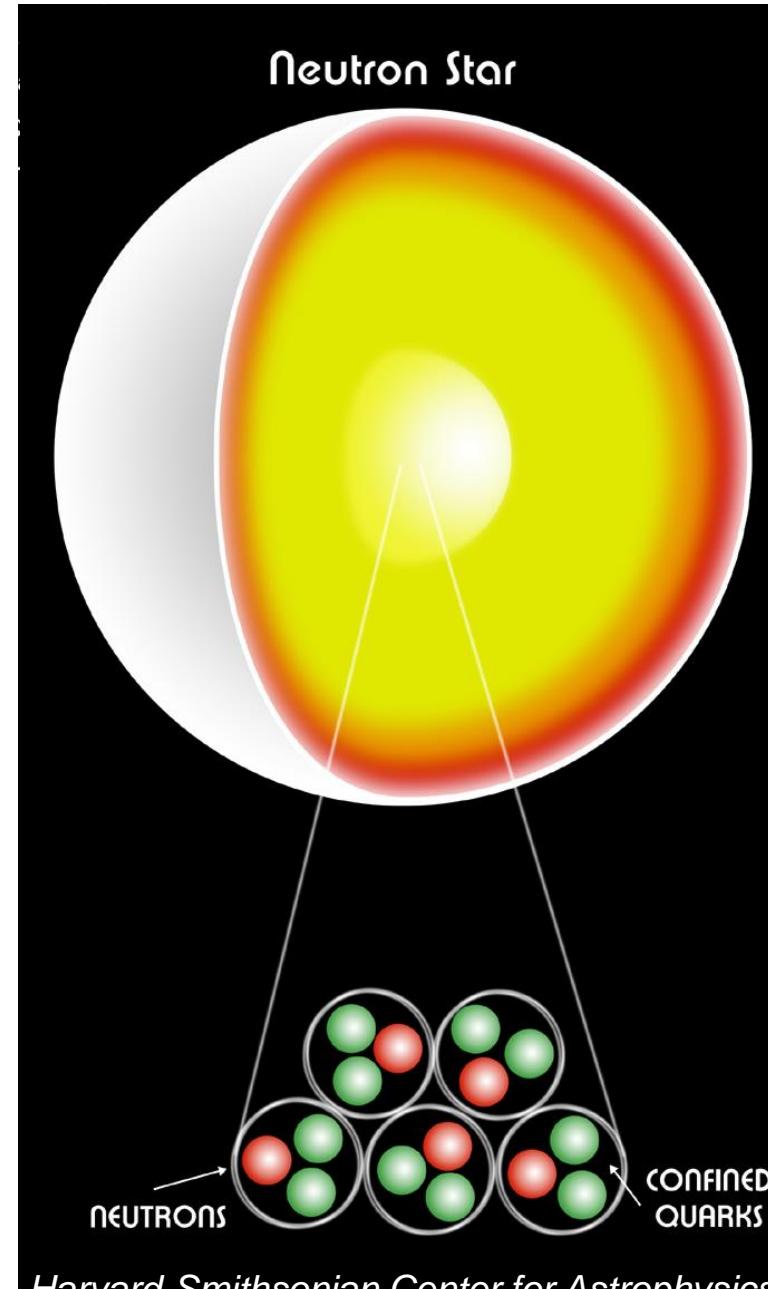
See also: Drake et al., PRA 2012, selectively probe the central portion of an inhomogeneous gas.

Z. Yan, P. Patel, B. Mukherjee, Z. Hadzibabic, T. Yefsah, J. Struck, MWZ, PRL 2017

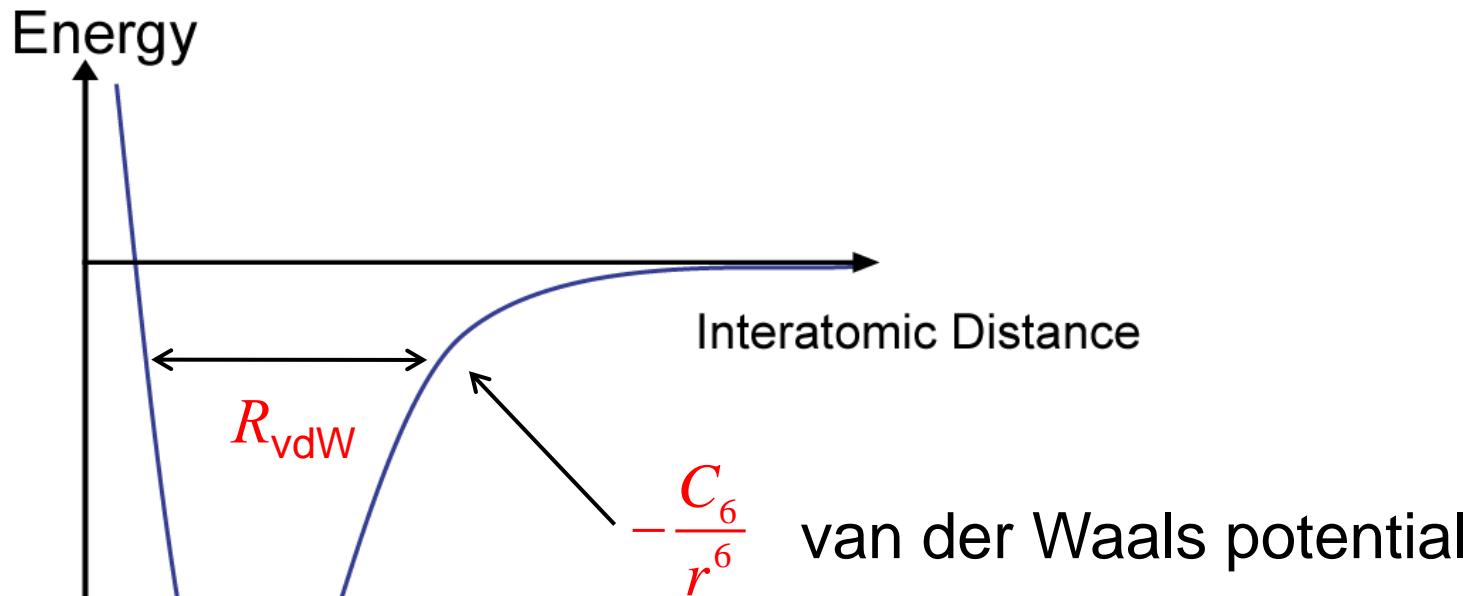
Fermions – The Building Blocks of Matter



Lithium-6



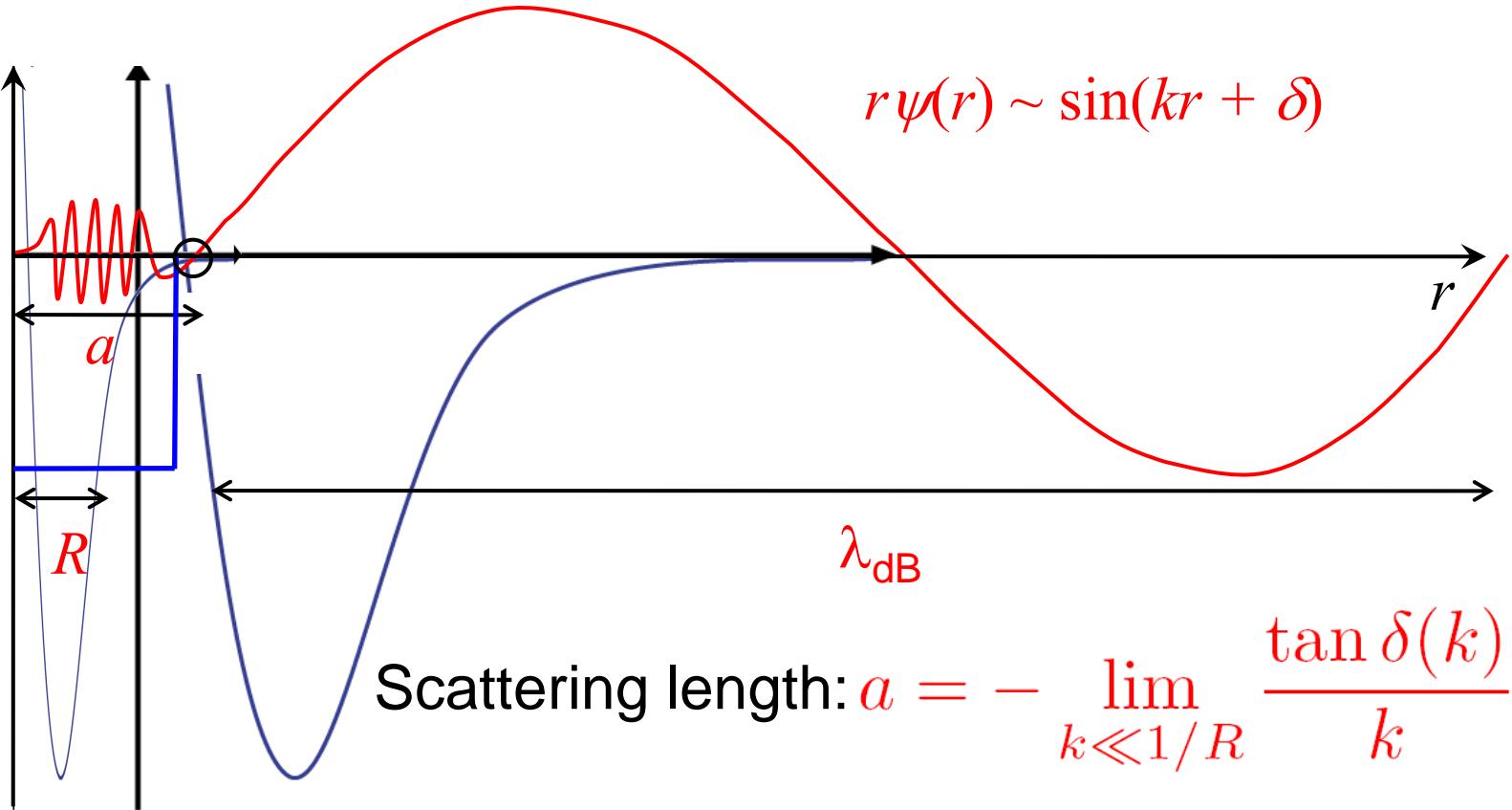
Interatomic interactions



$$R_{\text{vdW}} \approx \left(\frac{m C_6}{\hbar^2} \right)^{1/4}$$

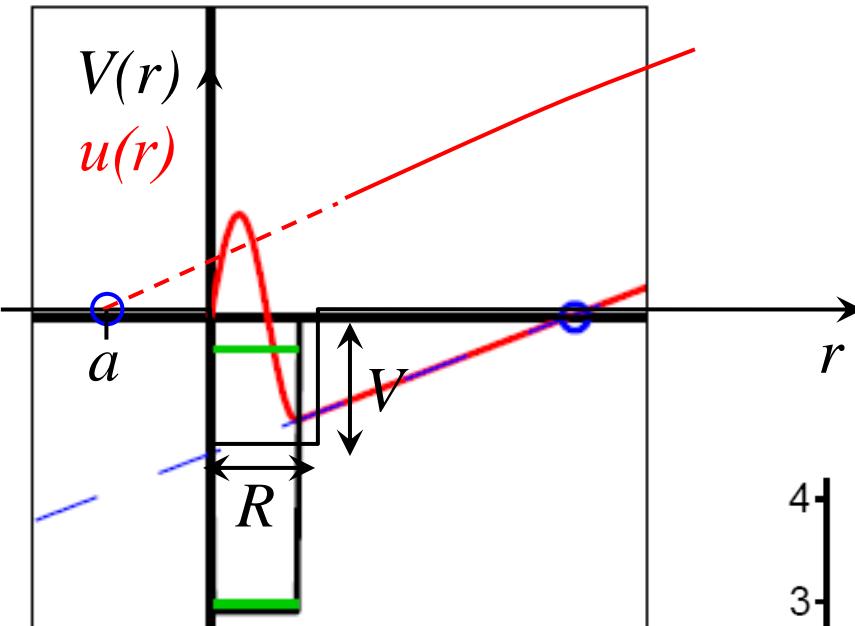
- For Alkali atoms: $R_{\text{vdW}} \sim 50\text{-}200 a_0$
- Ultracold collisions: $\lambda_{dB} \approx 1 \mu m \gg R$
→ atoms do not probe the details of the potential

Interatomic interactions

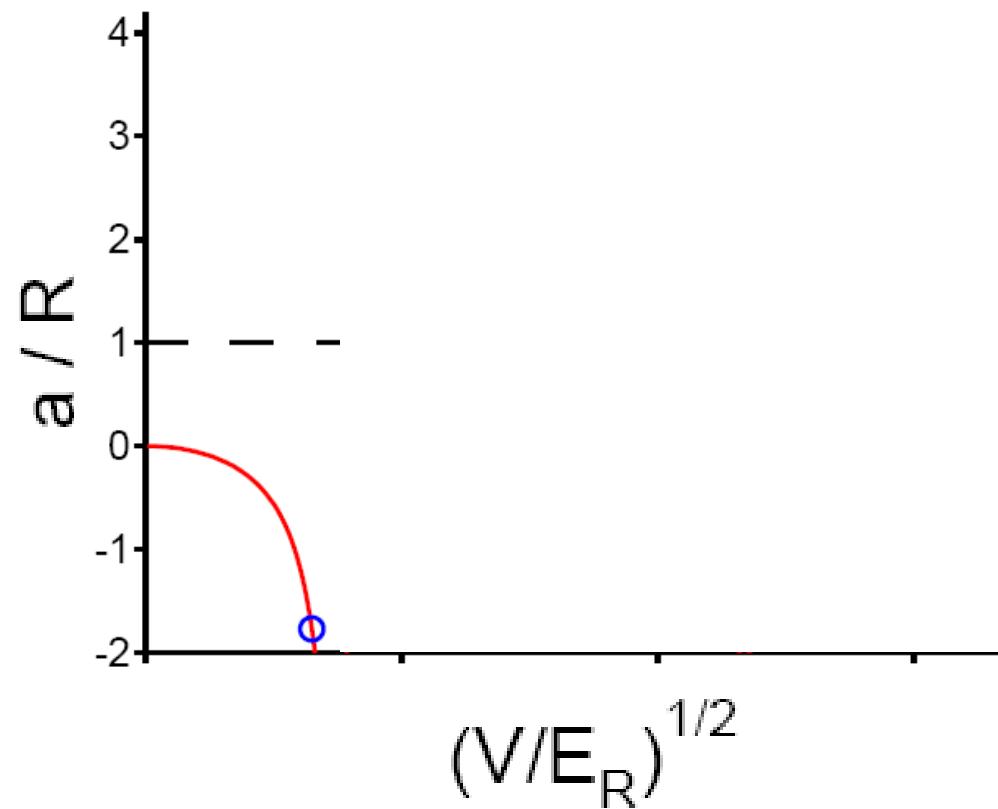


- For Alkali atoms: $R_{vdW} \sim 50-200 a_0$
- Ultracold collisions: $\lambda_{dB} \approx 1 \mu m \gg R$
- atoms do not probe the details of the potential

Scattering Resonances



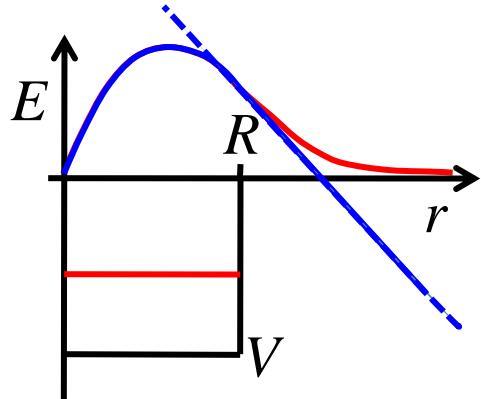
$$u(r) = \begin{cases} \sin(k(r - a)) & r > R \\ \text{(s-wave) scattering length} & \end{cases}$$



Tunable Interactions

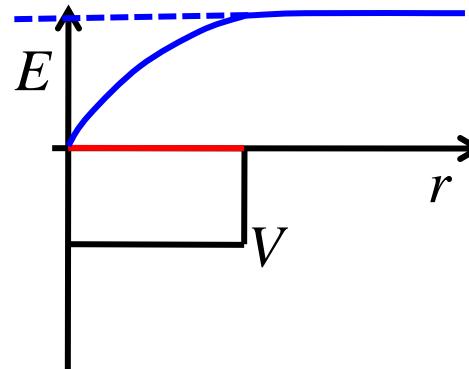
Vary interaction strength between spin up and spin down

Example: tunable square well (with $k_F R \ll 1$):



strong attraction
deep bound state

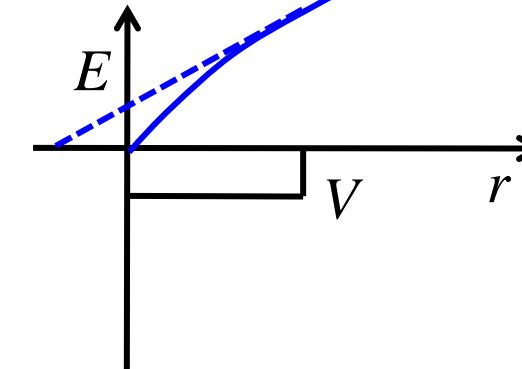
$$a > 0$$



Resonance
bound state appears

scattering length

$$a \rightarrow \pm\infty$$



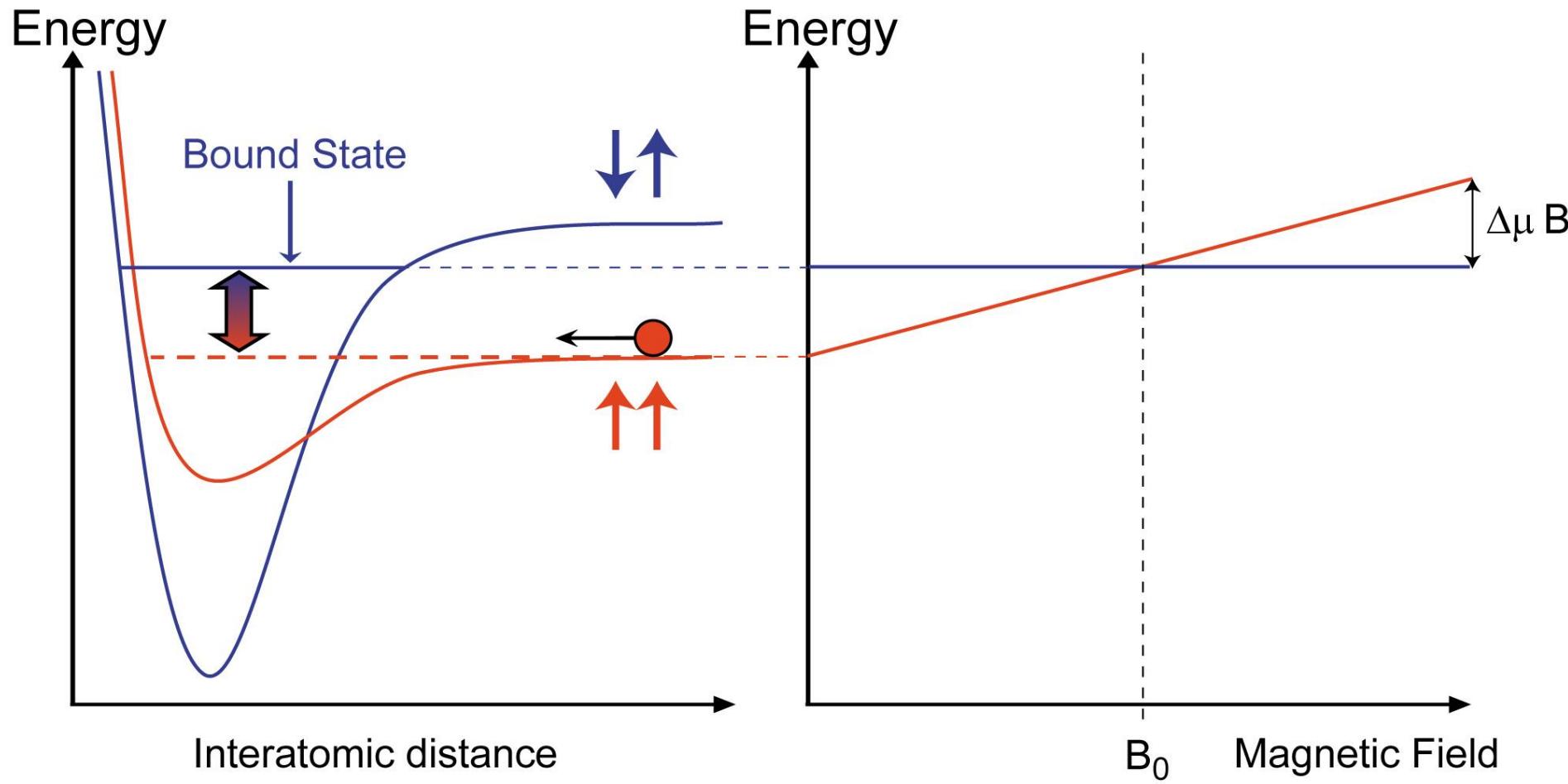
weak attraction
no bound state

$$a < 0$$

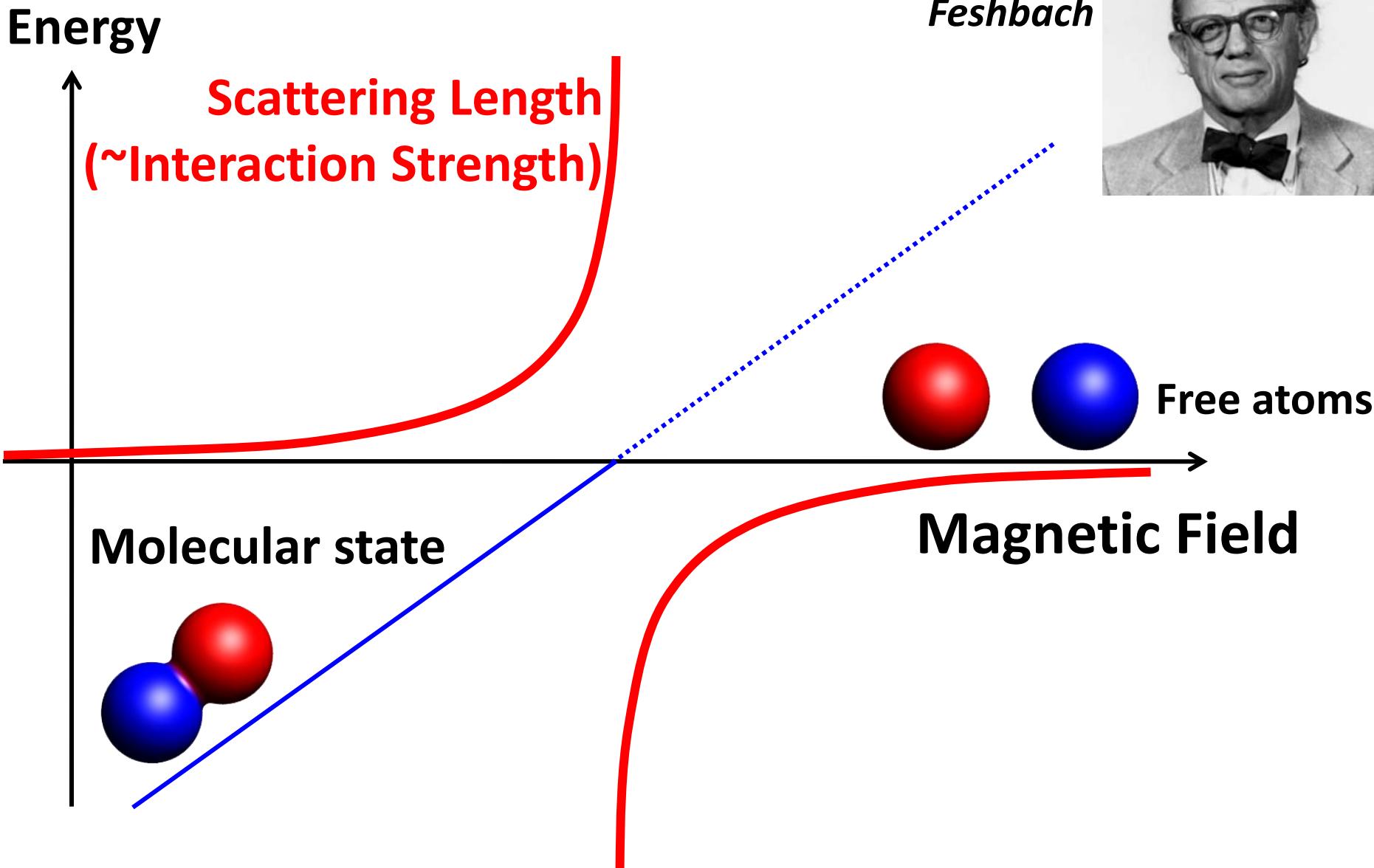
Important for Many-Body
Physics:

$$\frac{\text{Interparticle distance}}{\text{Scattering length}} = \frac{1}{k_F a}$$

Feshbach resonances: Tuning the interactions



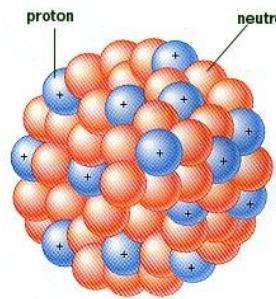
Feshbach resonances: Tuning the interactions



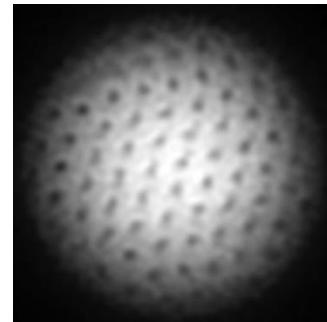
Strongly Interacting Fermi Systems

Length scales

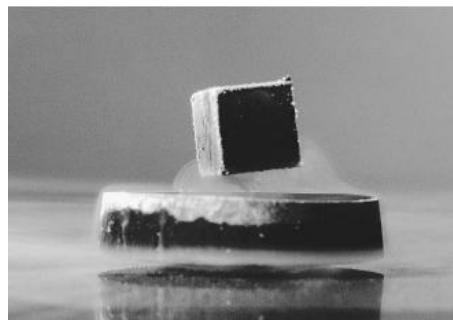
10^{-15} m



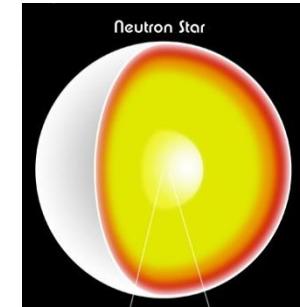
1 mm



1 m



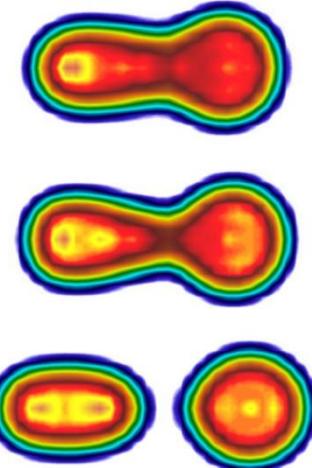
10^4 m



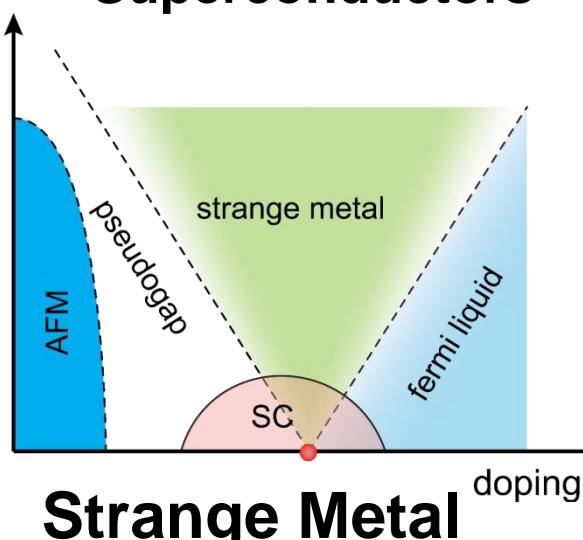
10^7 m



Nuclei



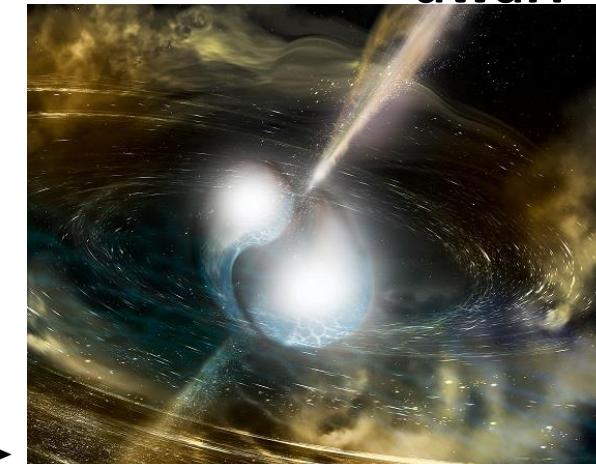
Ultracold
Gases



High- T_c
Superconductors

Nuclear Fission

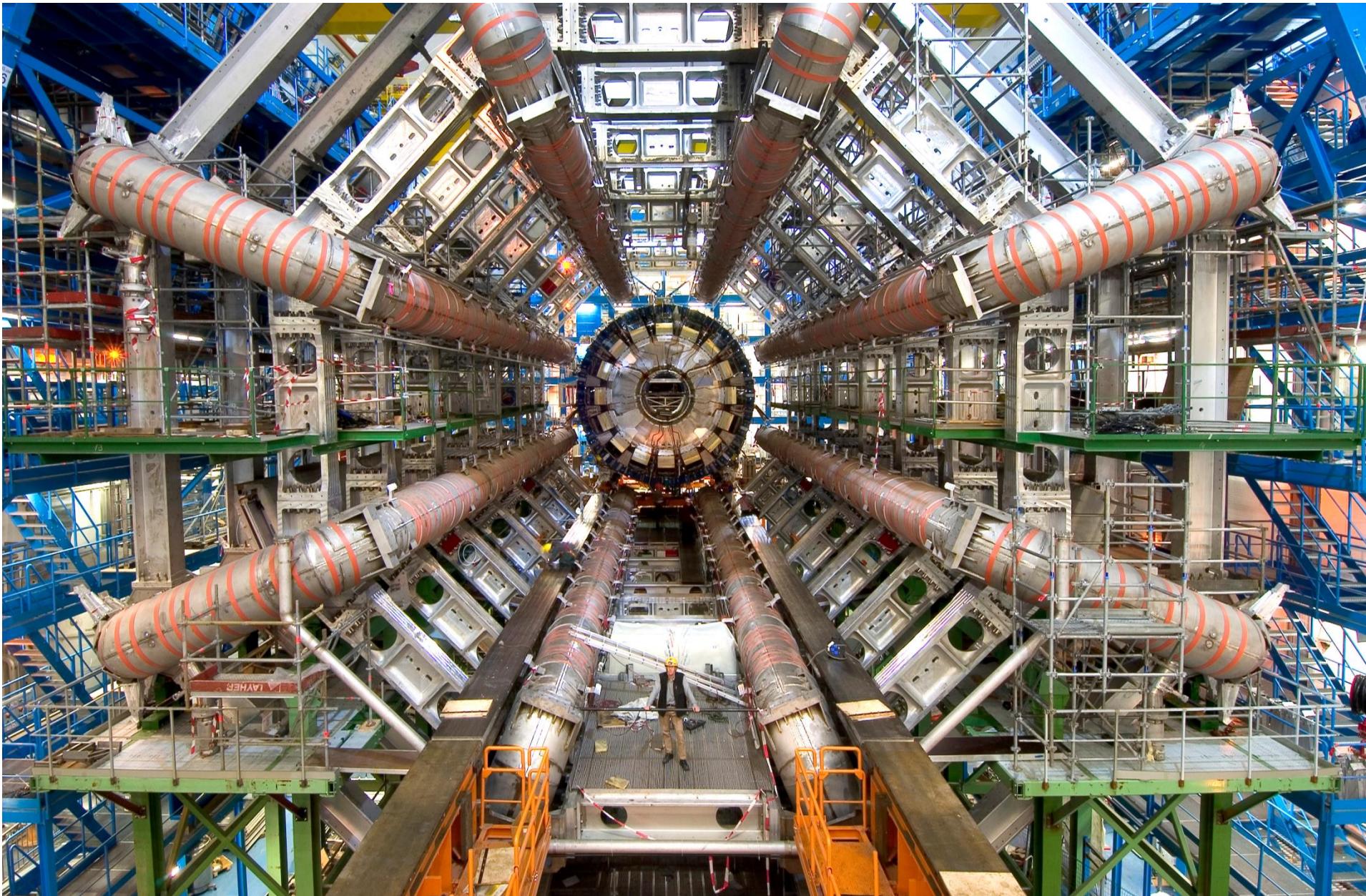
Neutron Star

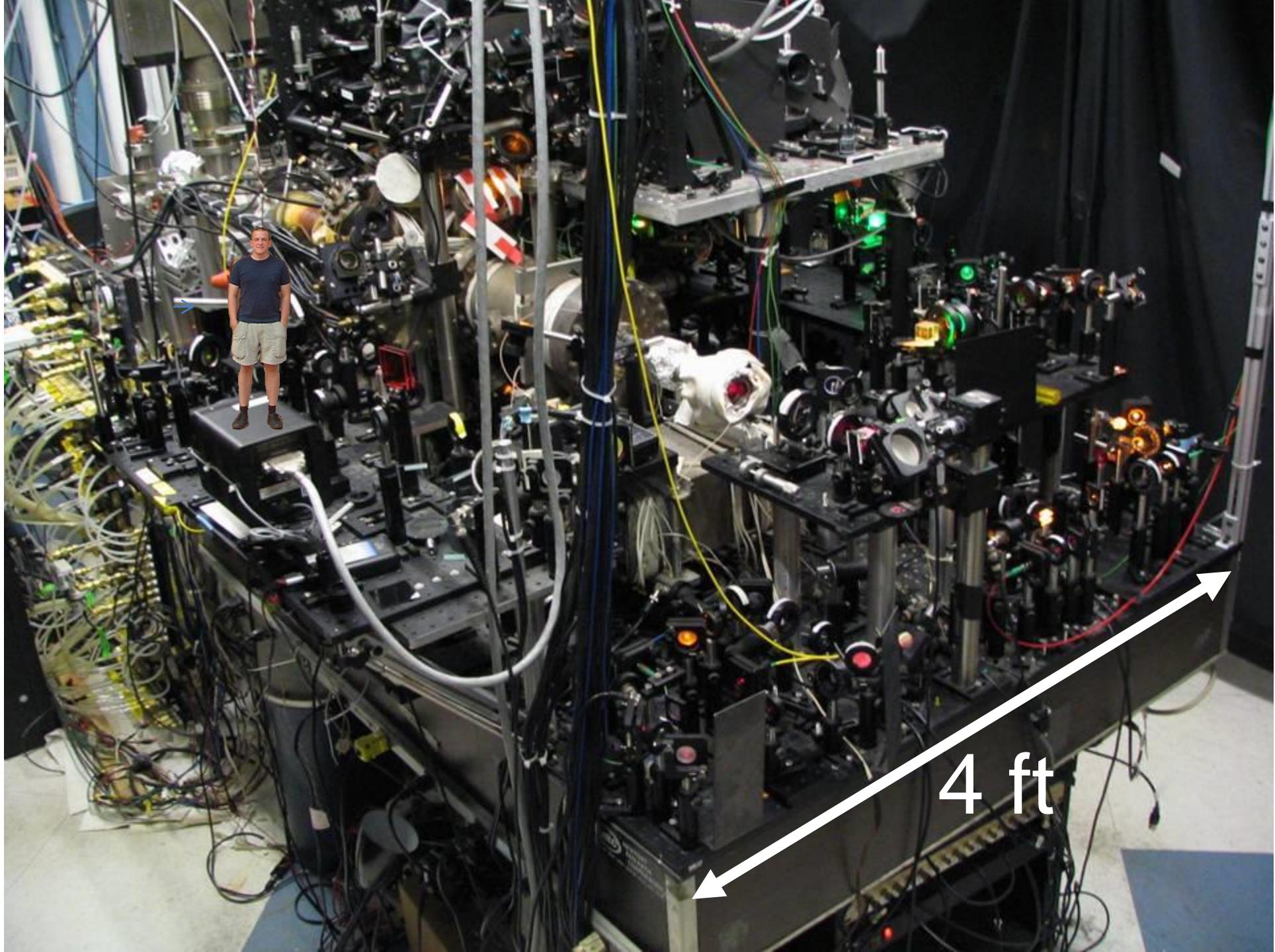


White
dwarf

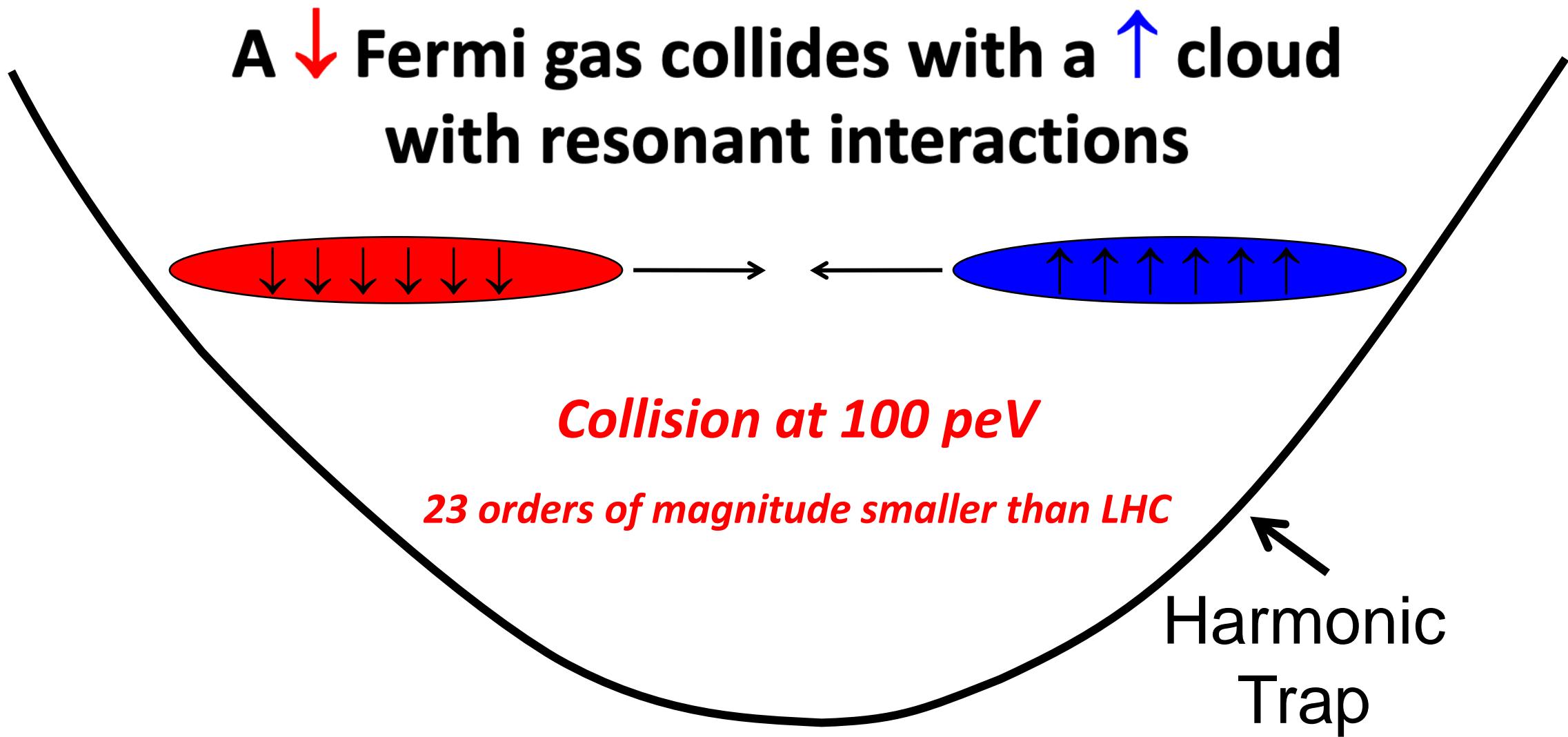
Neutron Star Merger

Large Hadron Collider (LHC)



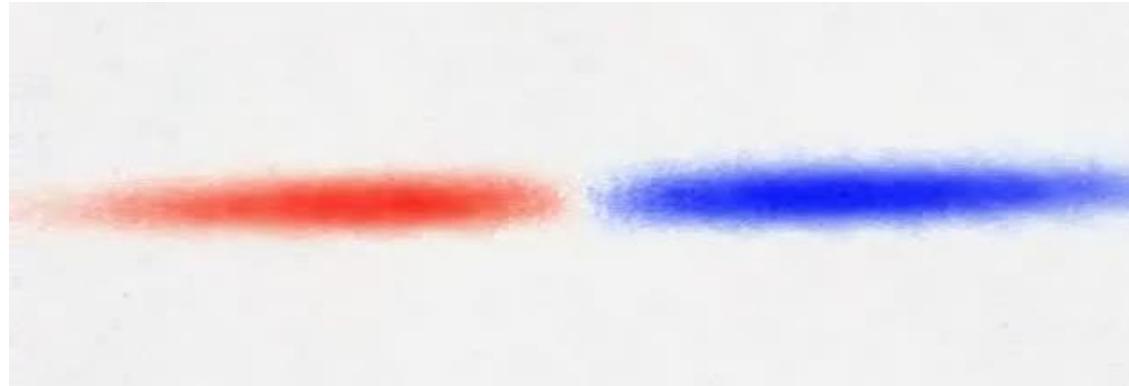


Little Fermi Collider (LFC)

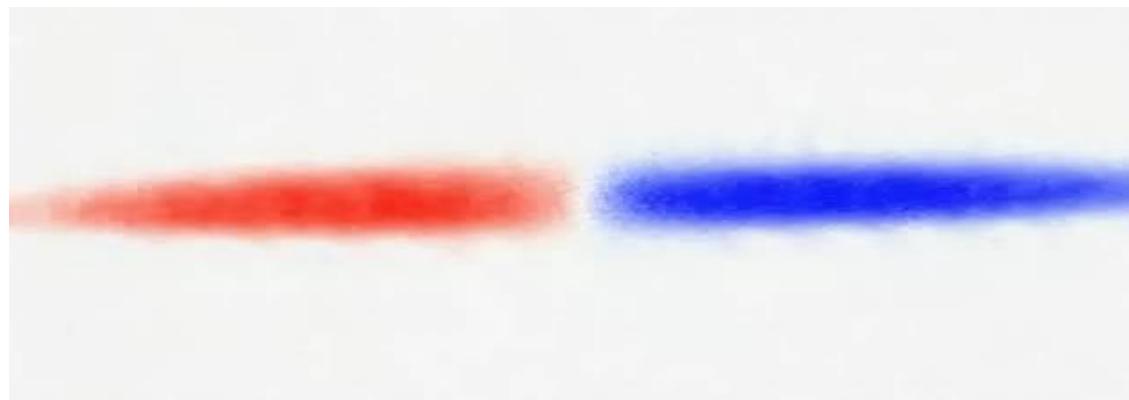


Little Fermi Collider (LFC)

Without Interactions

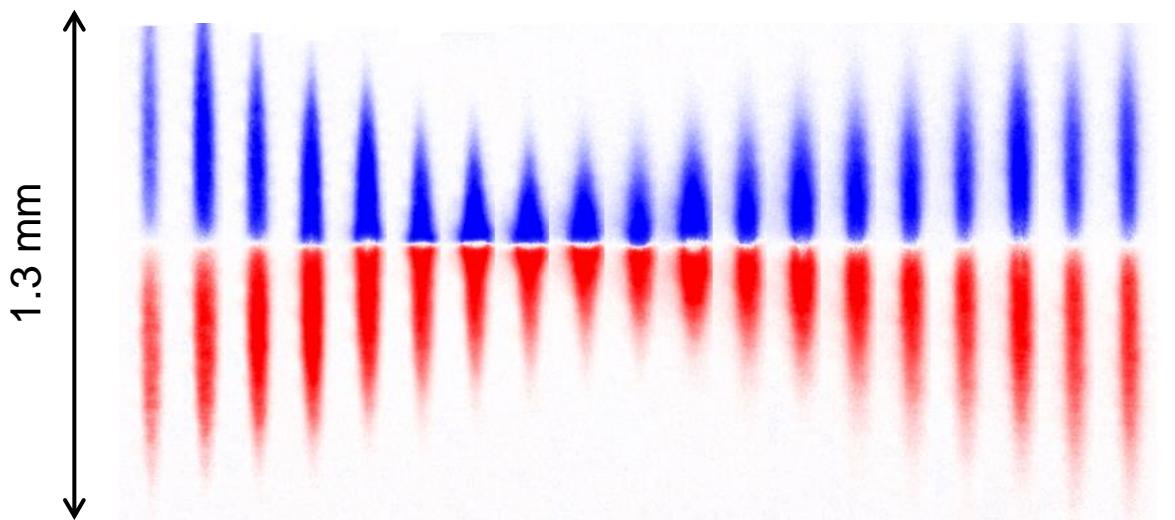


Resonant Interactions

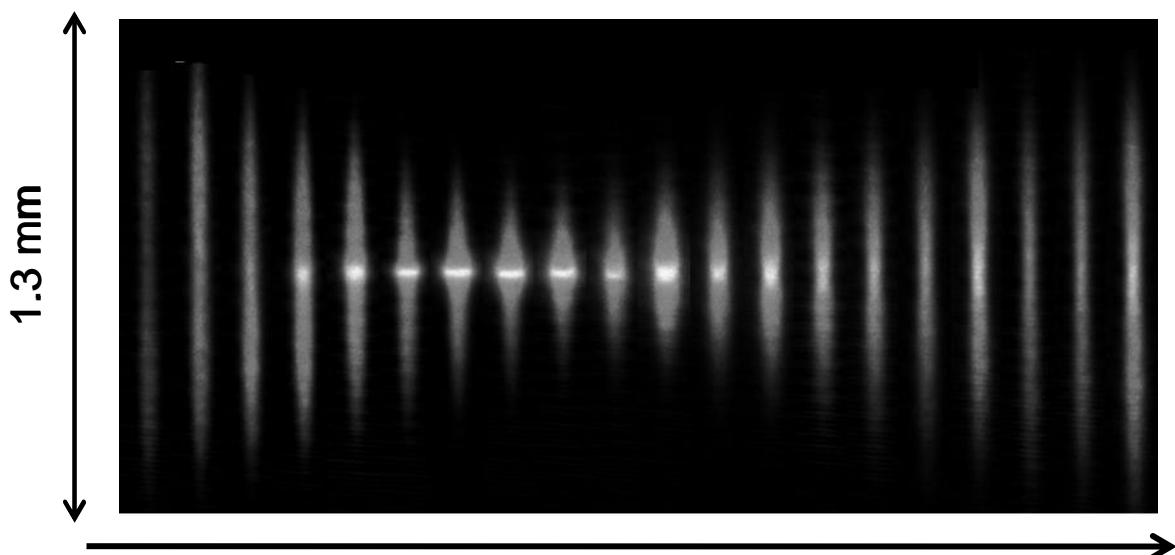


The bouncing gas

First collision

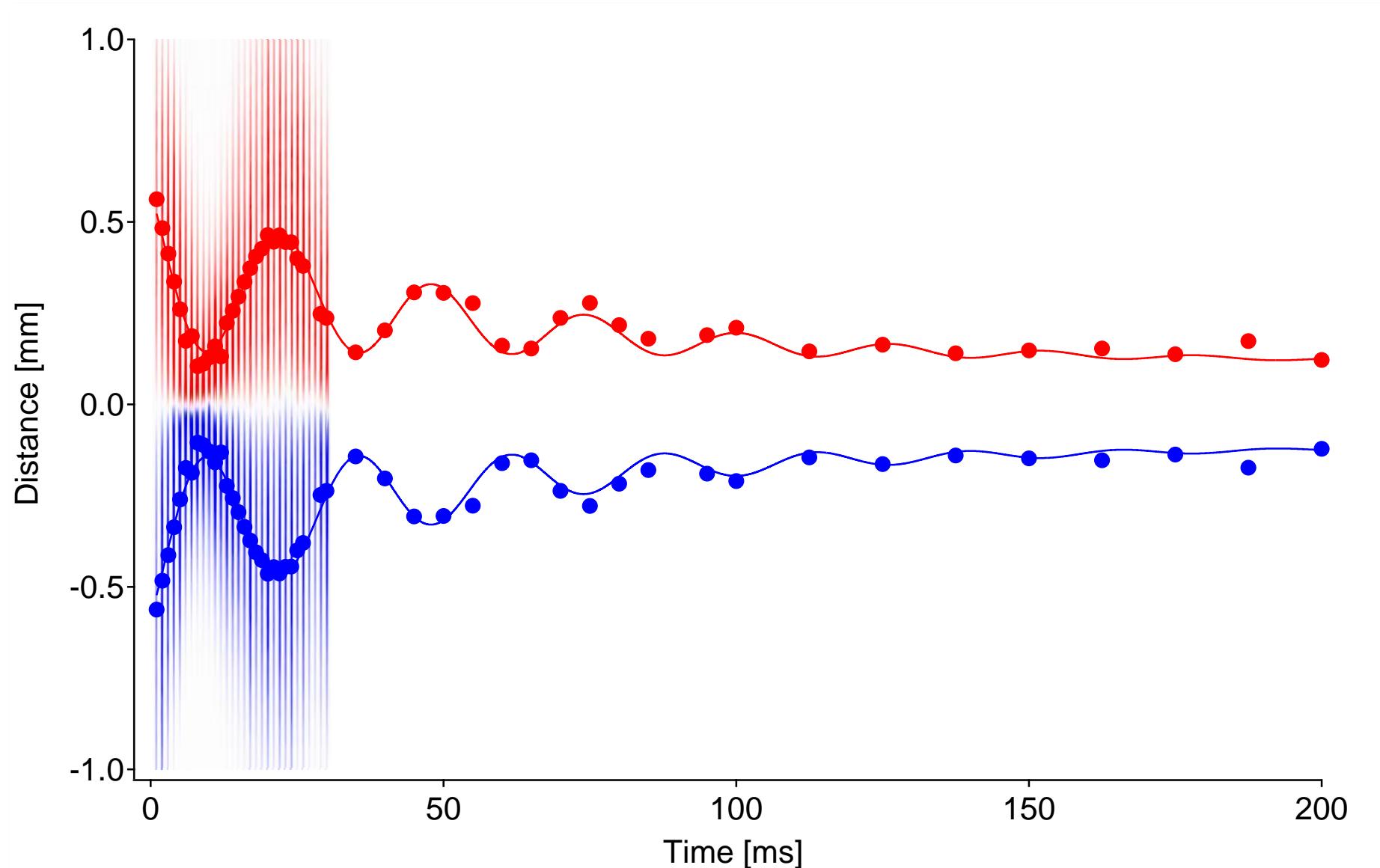


Difference density

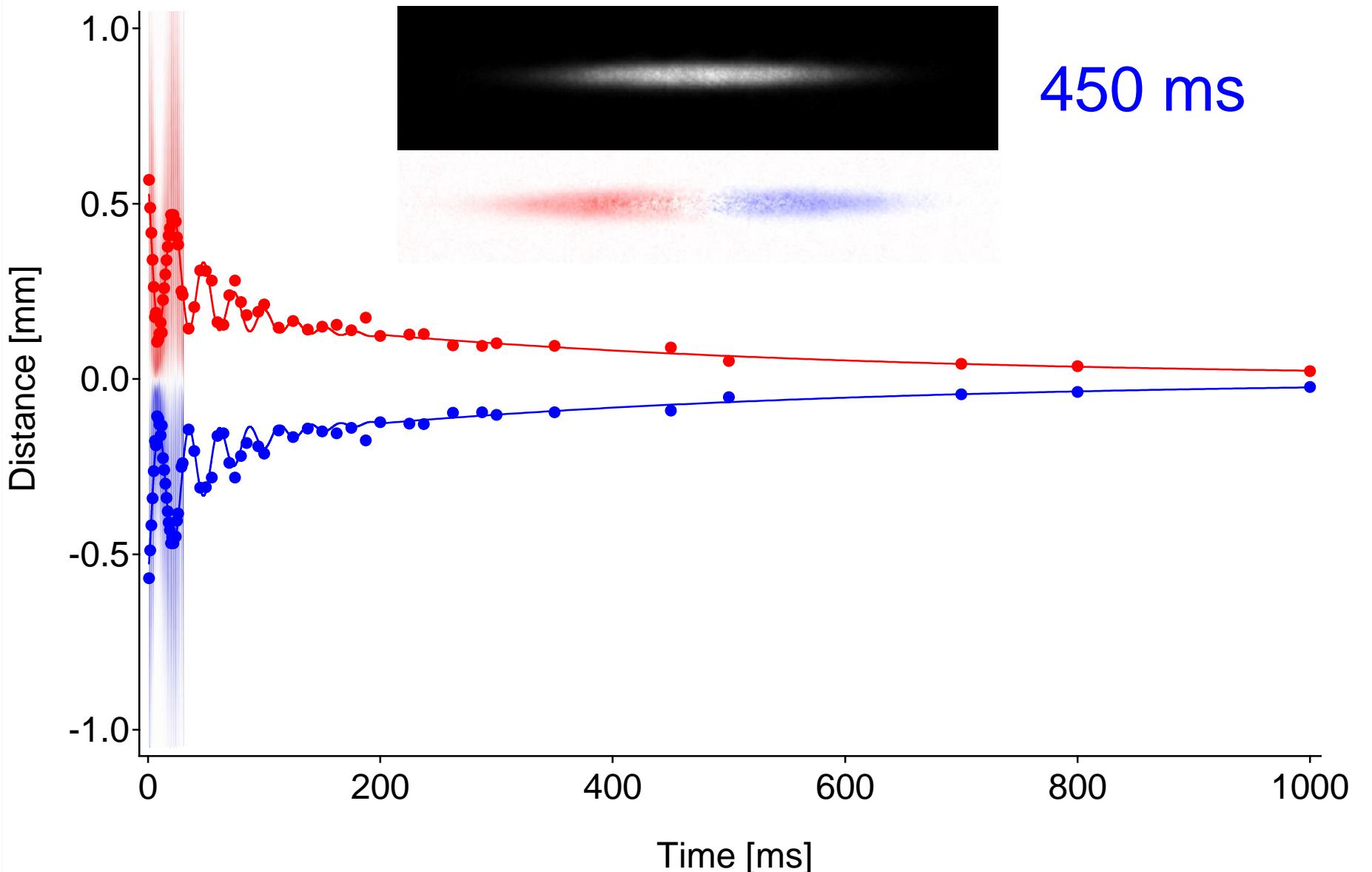


Total density

Later times



Much later times



Quantum limit of spin diffusion

Mean free path \sim Interparticle spacing d

Diffusion constant:

$$D \sim \text{mean free path} \times \text{average velocity}$$
$$\cancel{d} \quad \times \quad \frac{\hbar}{m \cancel{d}}$$

$$D \sim \frac{\hbar}{m} = \frac{\text{Planck's constant}}{\text{Particle mass}} = \frac{(0.1 \text{ mm})^2}{1 \text{ s}}$$

→ Quantum Limit of Diffusion

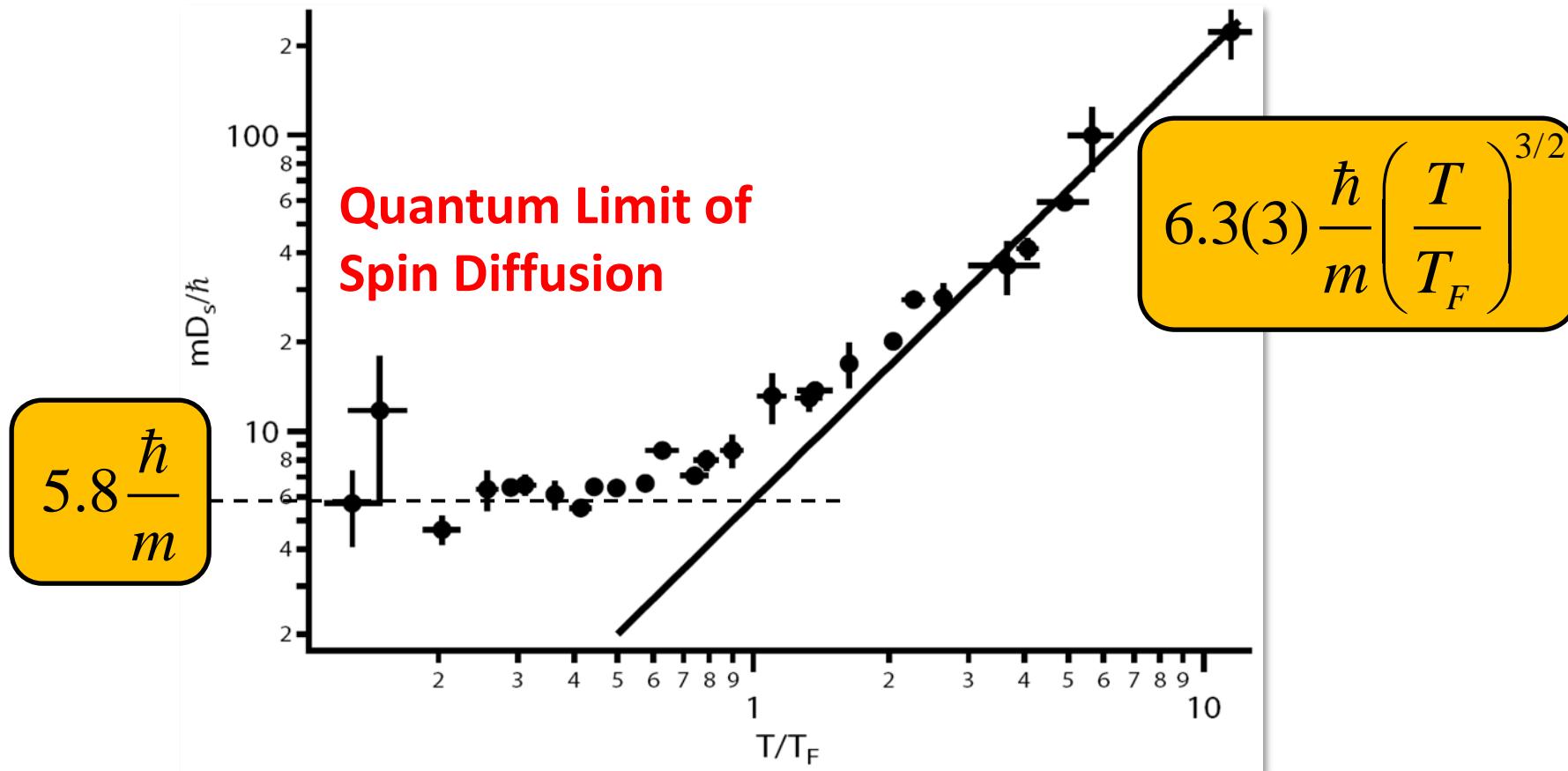
In a hot relativistic fluid (e.g. Quark-Gluon Plasma): $D \sim \frac{\hbar c^2}{T}$

$$mc^2 \rightarrow T$$

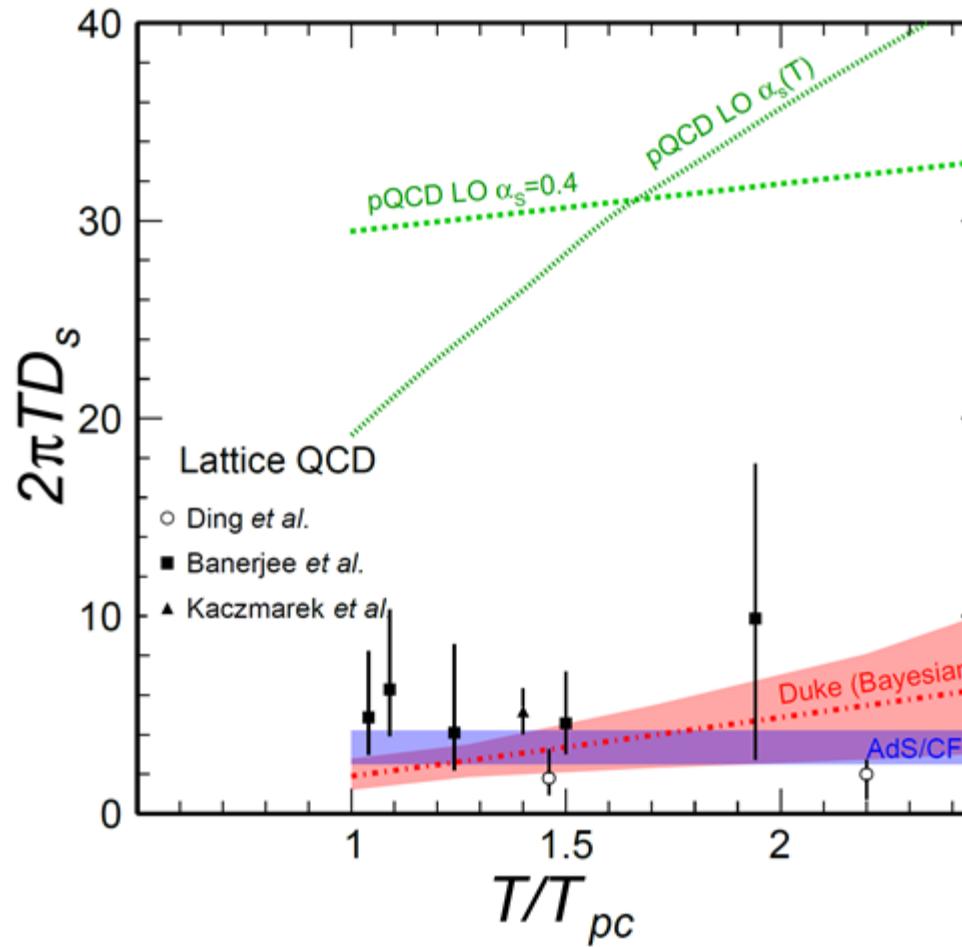
Spin Diffusion vs Temperature

Spin current = $-D \cdot$ Spin density gradient

Universal high-T behavior:



Diffusivity in the Quark-Gluon Plasma



**Can Fermi Gases become
superfluid?**

Superconductivity

Electrons are Fermions

Discovery of superconductivity 1911



Heike
Kamerlingh-Onnes
1911

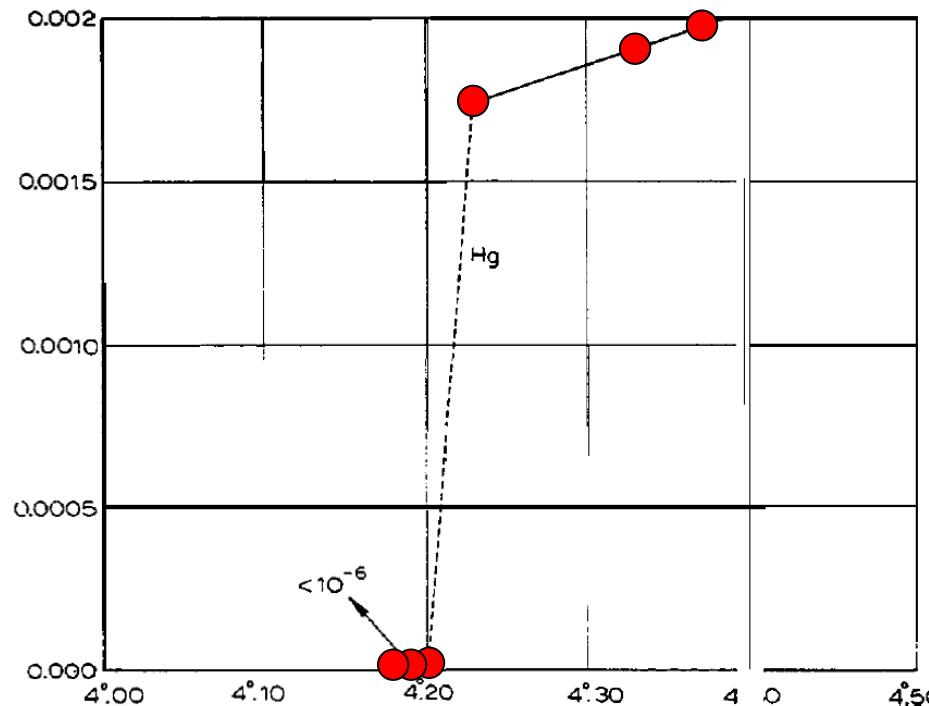


Fig. 17.

Fermionic Superfluidity

Condensation of Fermion Pairs

- Helium-3 (Lee, Osheroff, Richardson 1971)
- Superconductors: *Charged superfluids of electron pairs*
Frictionless flow \Leftrightarrow Resistance-less current



John Bardeen



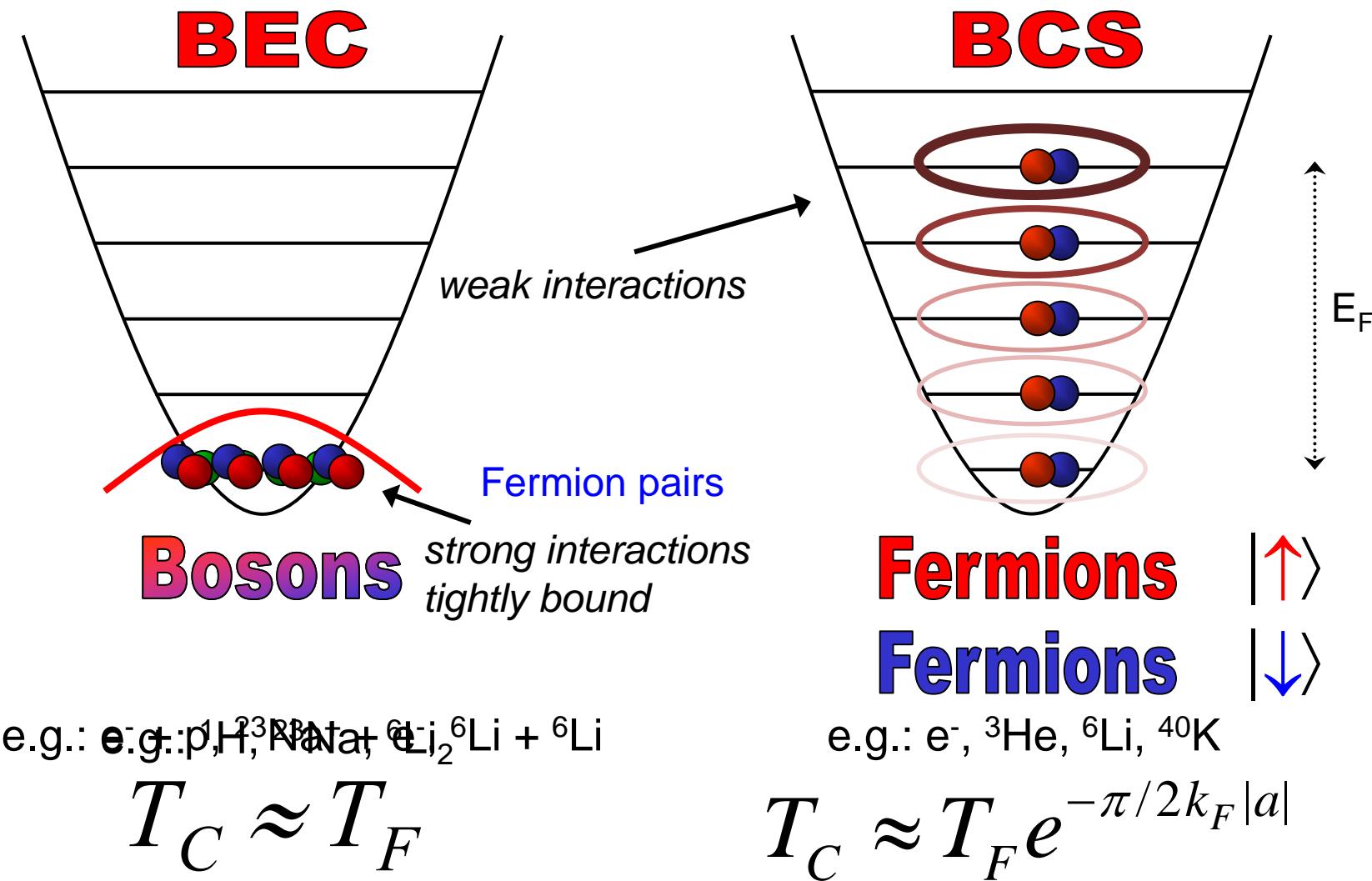
Leon N. Cooper



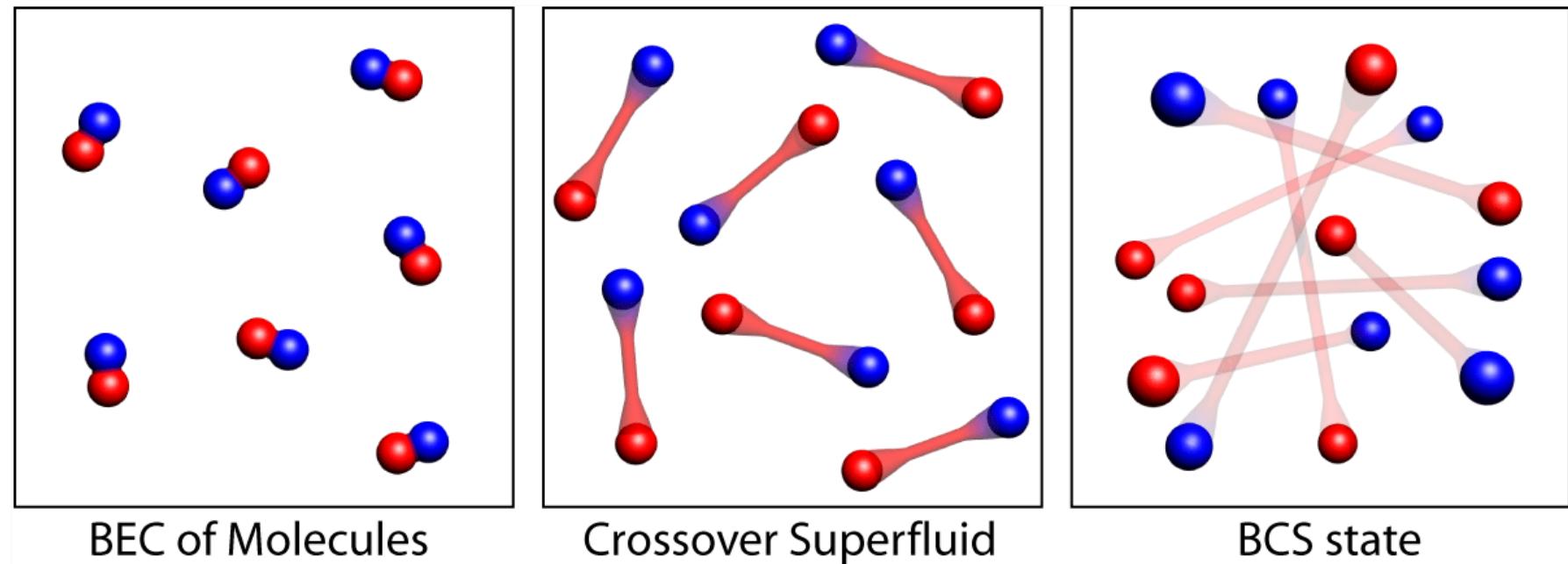
John R. Schrieffer

- Neutron stars
In the core: Quark superfluid

Bosons vs Fermions



From BEC to BCS



$\xleftarrow{\gg 1} \quad 1 \quad 0 \quad -1 \quad \xrightarrow{\ll -1}$

Weakly Interacting Bosons

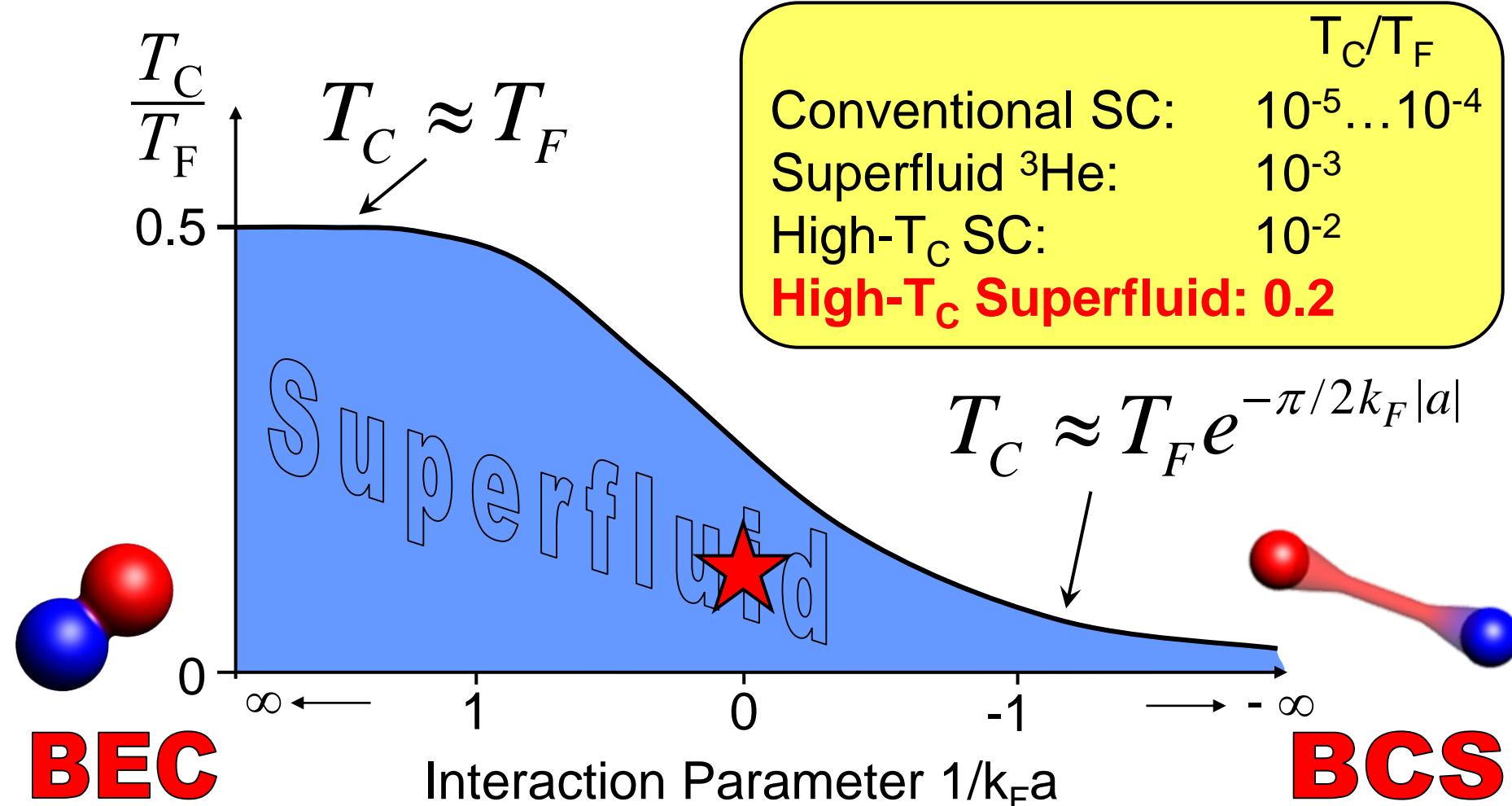
→ Strongly Interacting Bosons

→ Strongly Interacting Fermions

→ Weakly Interacting Fermions

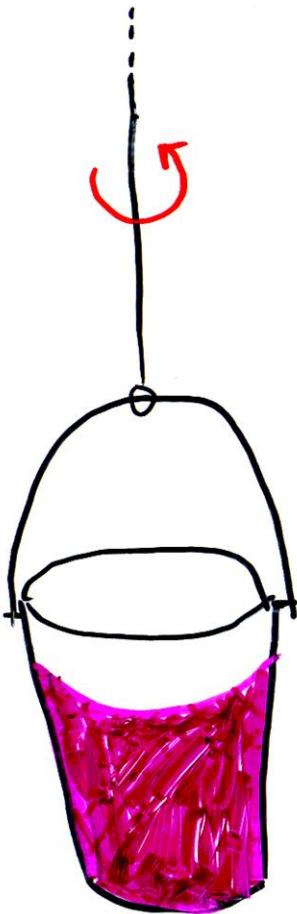
$$(k_F a)^{-1} = \frac{\text{Interparticle Distance}}{\text{Scattering Length}}$$

Critical Temperature for Fermionic Superfluidity



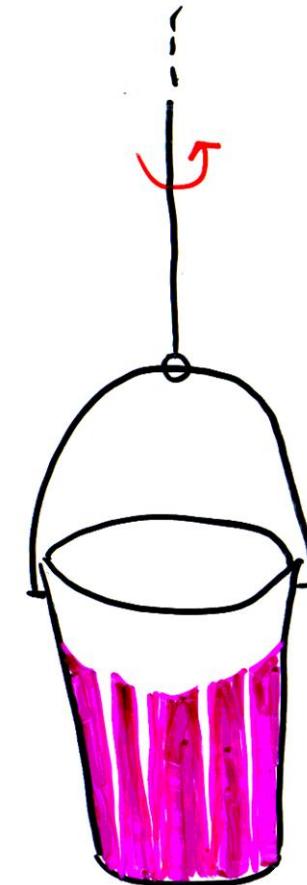
Scaled to the density of electrons in solids:
Superconductivity far above room temperature!

Rotating Fluids



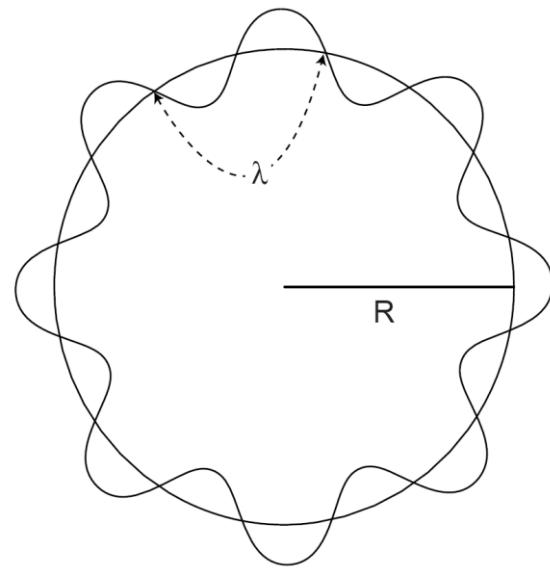
Normal

Fluid



Quantum

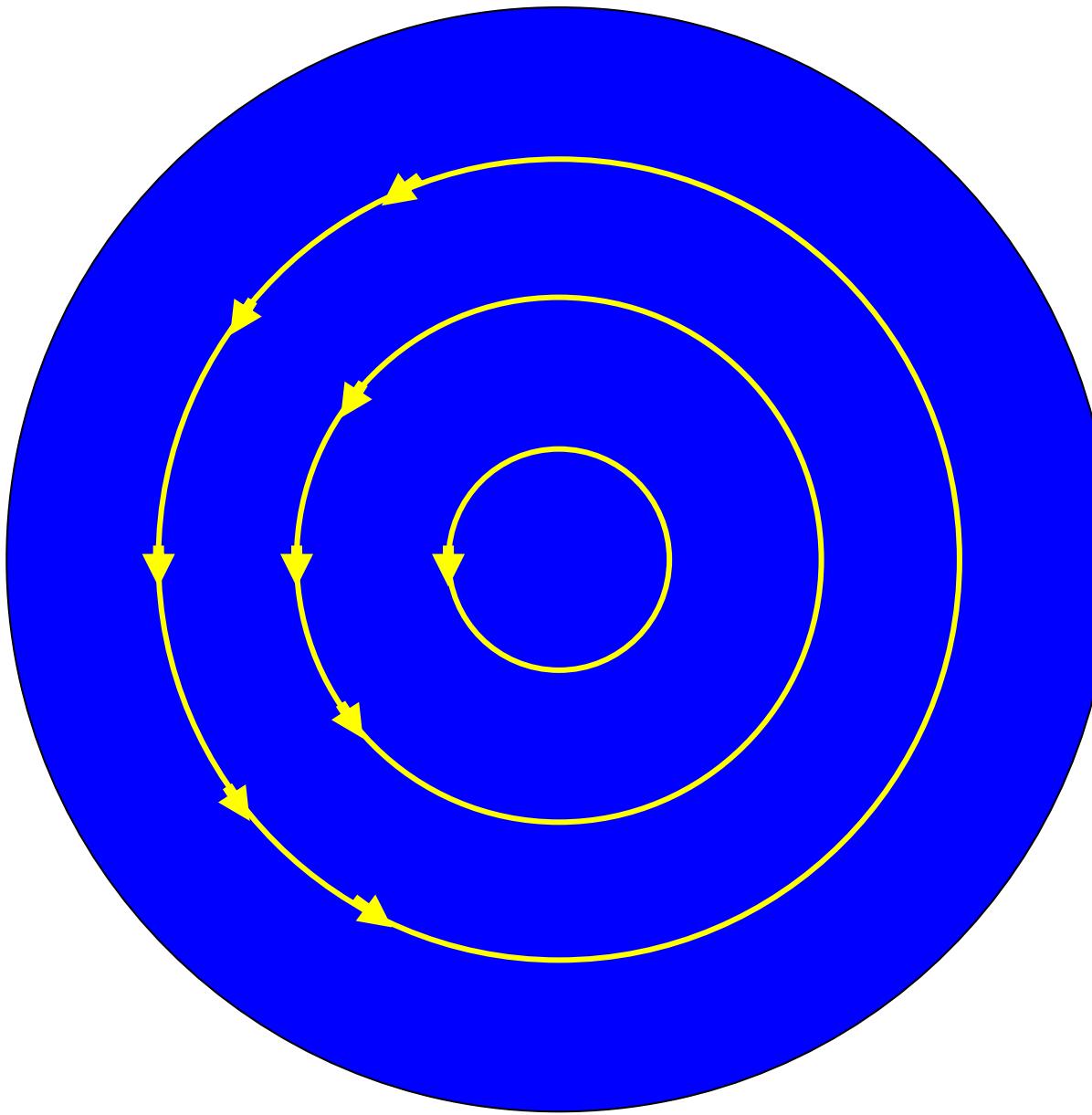
Rotating superfluid



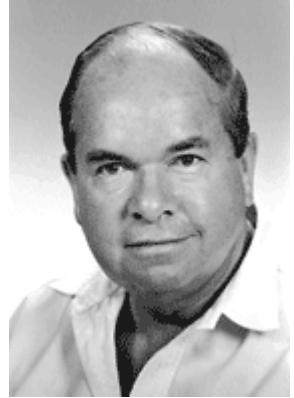
$$\oint p \, dx = \ell h$$



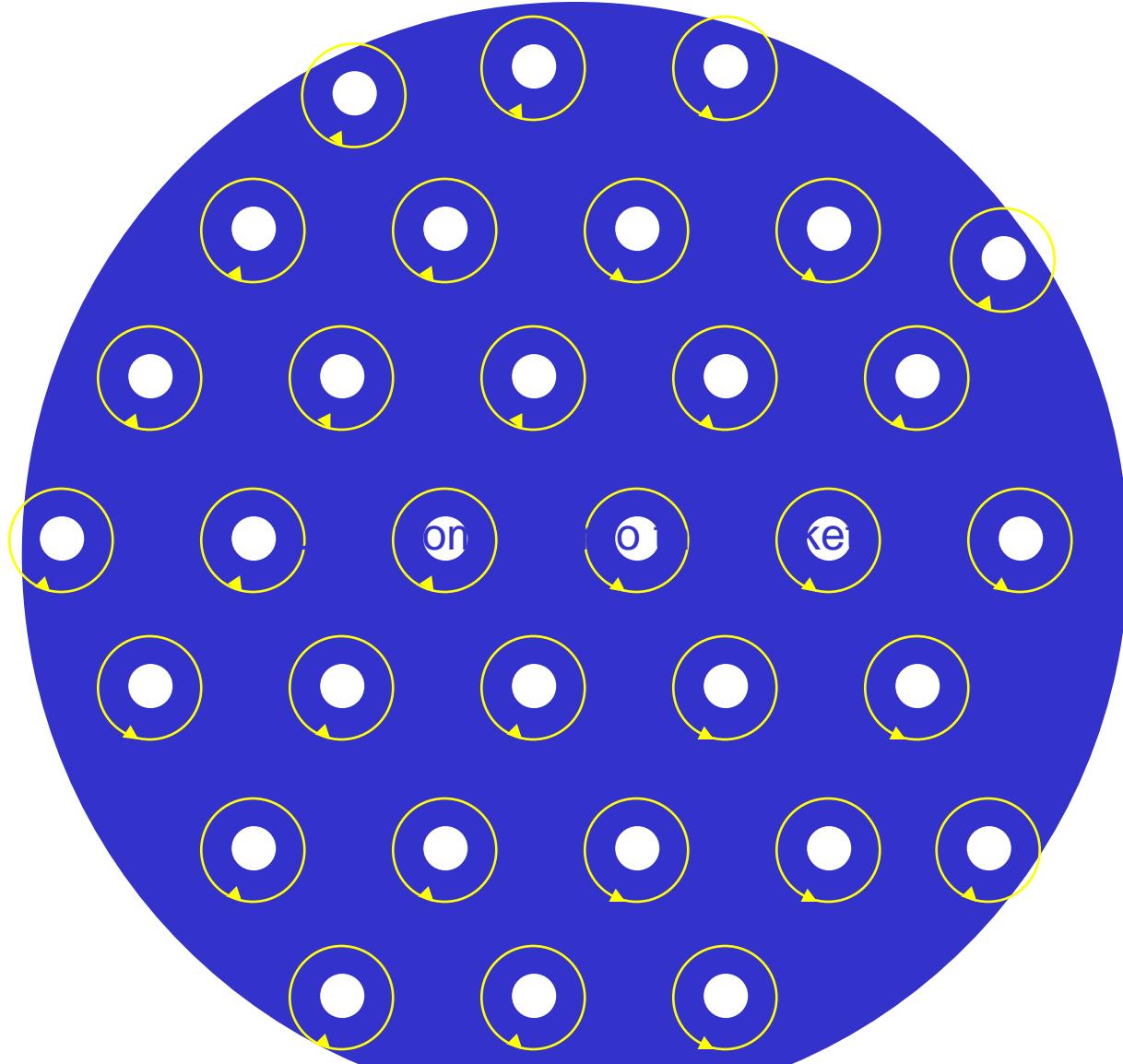
Does it rotate like this?



Like this!

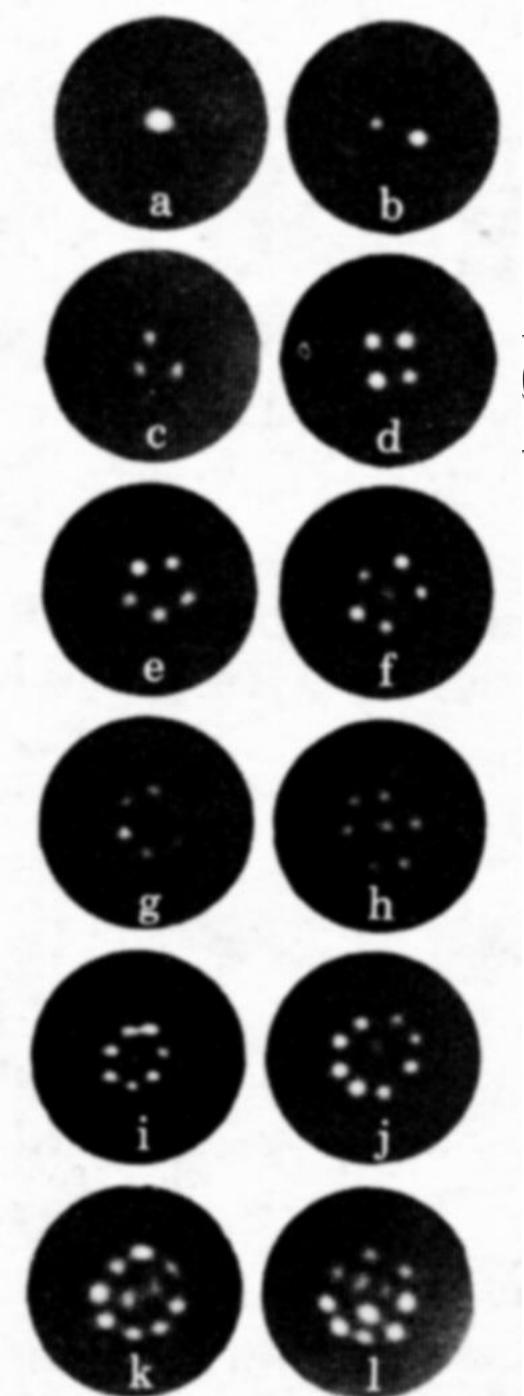


Aleksei A.
Abrikosov



Abrikosov lattice (triangular lattice)

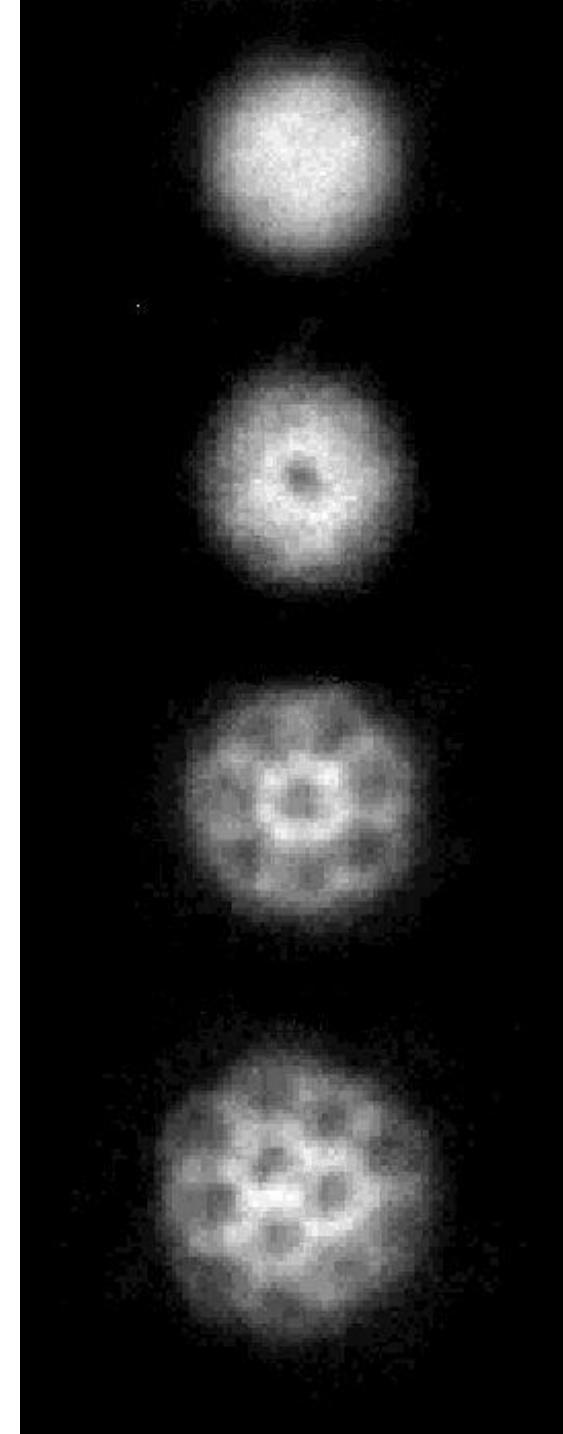
Vortex Arrays in Bosonic Gases / Fluids



Berkeley
(R.E. Packard, 1979)
Helium-4

ENS
(J. Dalibard, 2000)
Rubidium BEC

*Also: Phase engineering
of single vortices in BEC;
JILA (1999)*



THE DIRECT OBSERVATION OF INDIVIDUAL FLUX LINES
IN TYPE II SUPERCONDUCTORS

U. ESSMANN and H. TRÄUBLE

*Institut für Physik am Max-Planck-Institut für Metallforschung, Stuttgart and
Institut für theoretische und angewandte Physik der Technischen Hochschule Stuttgart*

Received 4 April 1967

Neutral superfluids under rotation

$$\vec{F} = 2m\vec{v} \times \vec{\omega}$$

Coriolis force in rotating frame



Superconductors in magnetic field

$$\vec{F} = q\vec{v} \times \vec{B}$$

Lorentz Force

U. Essmann and H. Träuble,
Physics Letters A, 24, 526 (1967)

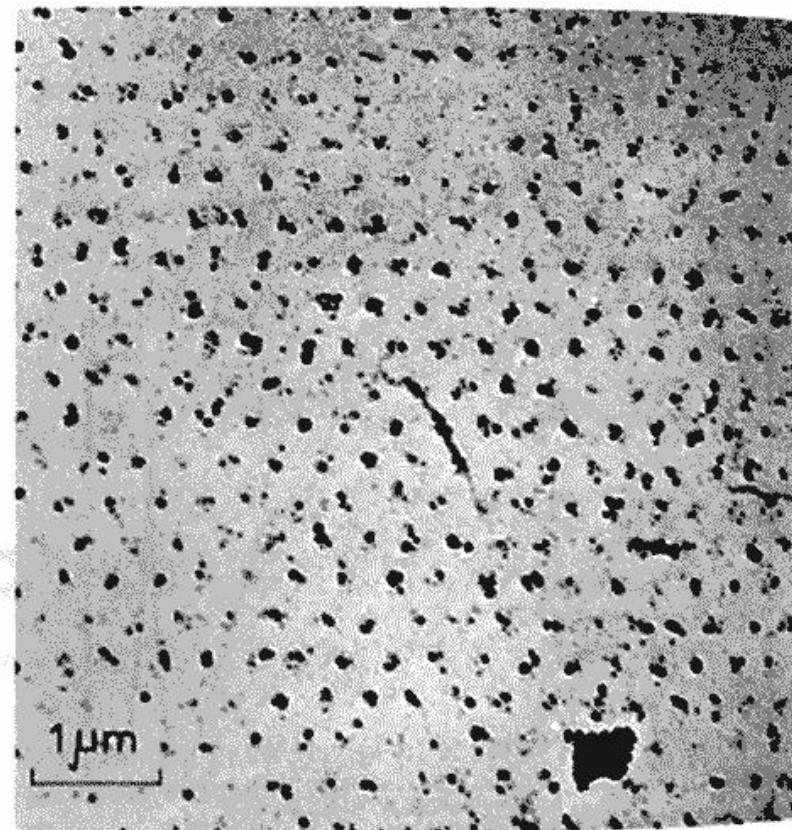
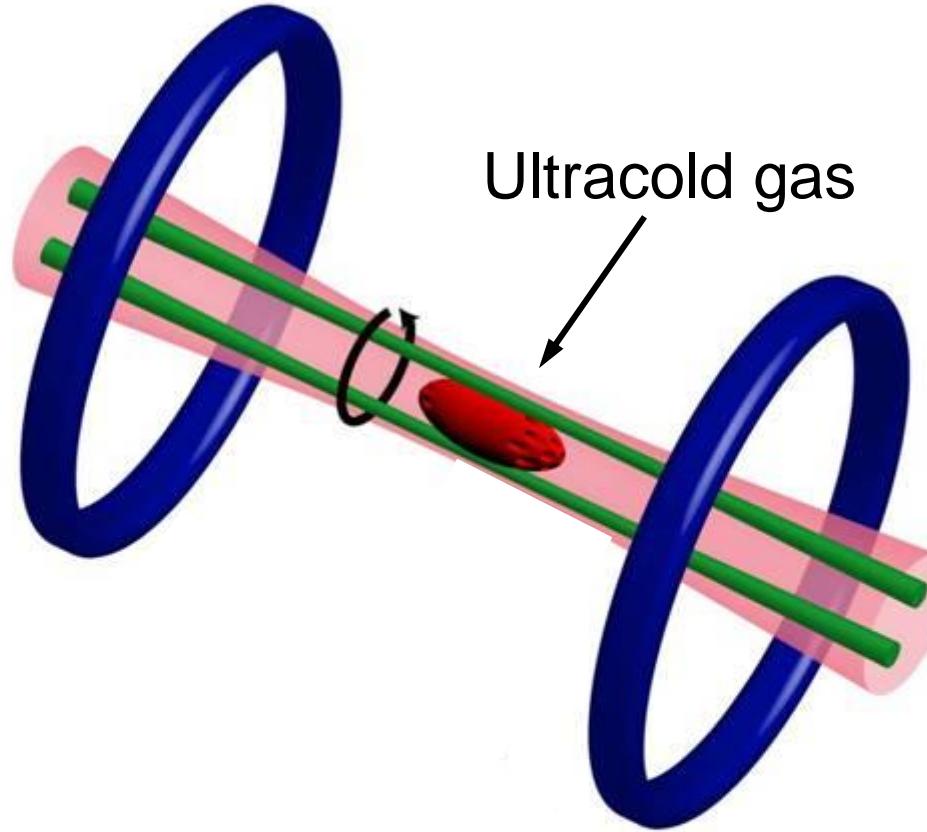
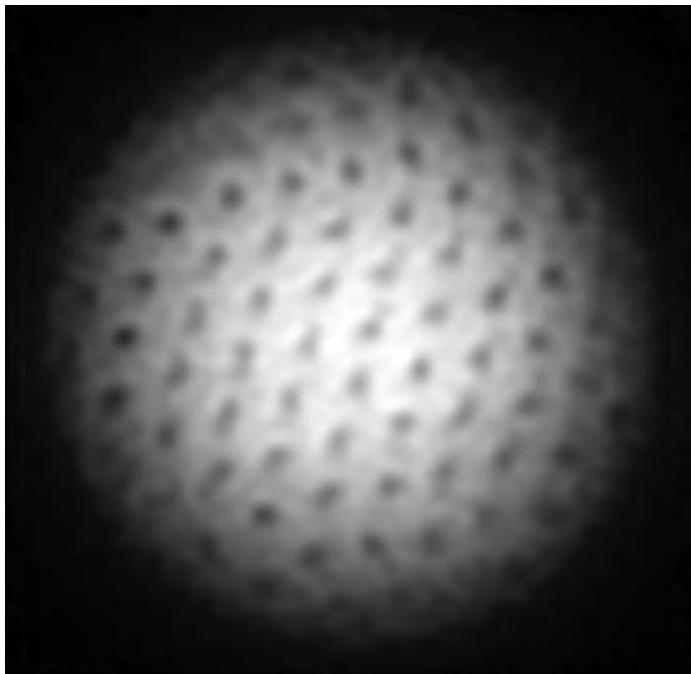
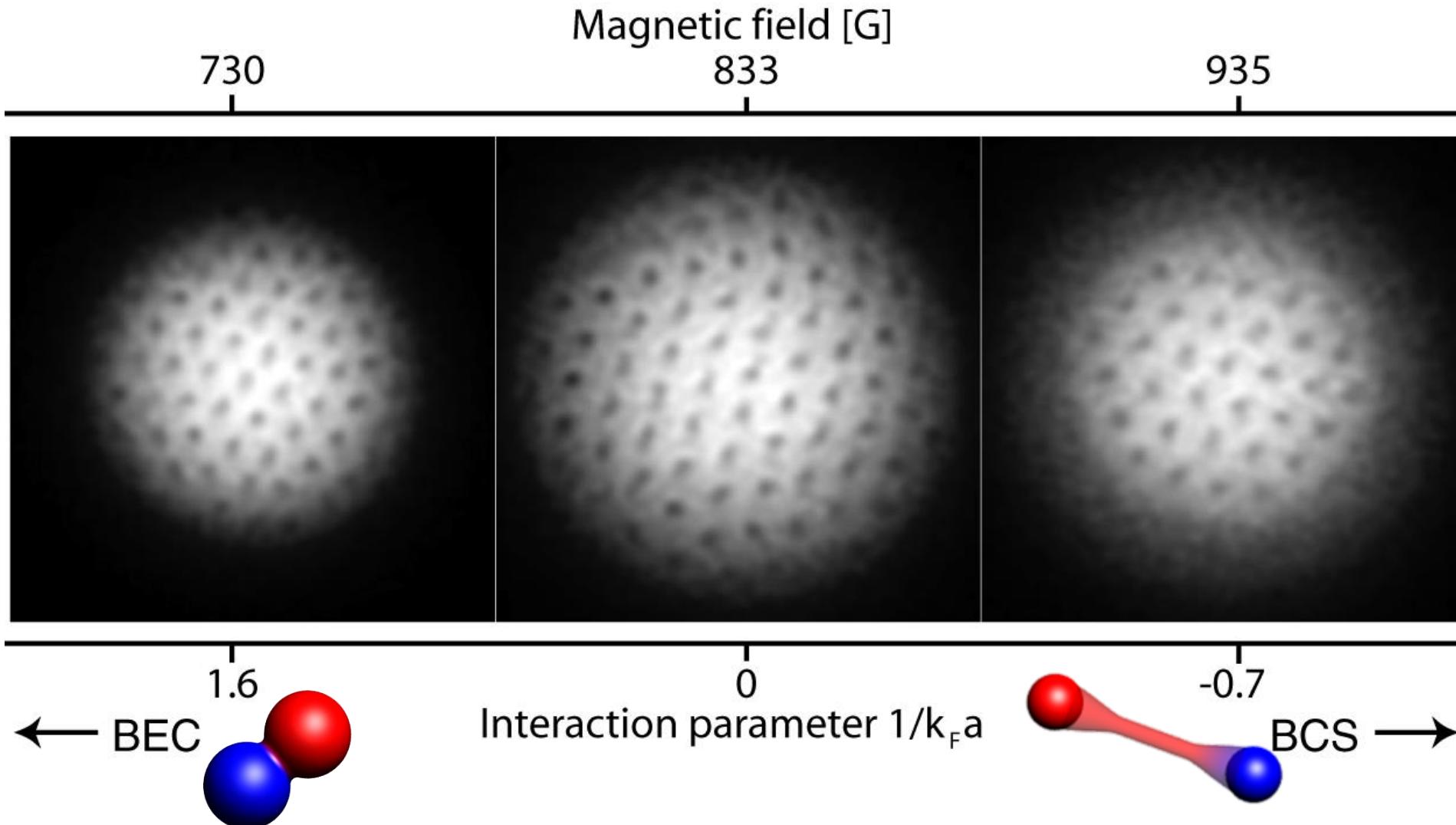


Fig. 1. "Perfect" triangular lattice of flux lines on the surface of a lead-4at%indium rod at 1.1°K. The black dots consist of small cobalt particles which have been stripped from the surface with a carbon replica.

Demonstration of superfluidity in a Fermi gas



Vortex lattices in the BEC-BCS crossover

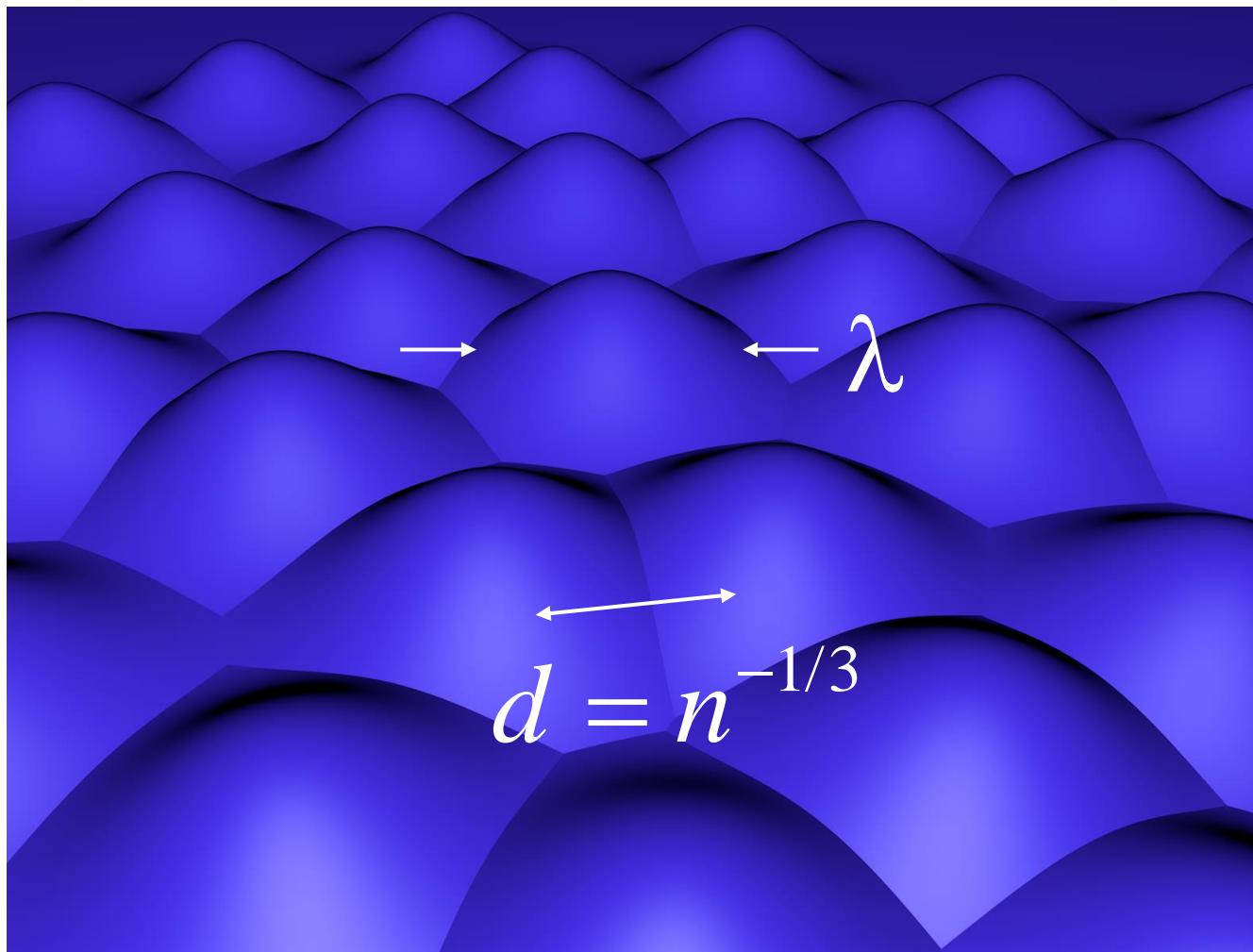


M.W. Zwierlein, J.R. Abo-Shaeer, A. Schirotzek, C.H. Schunck, W. Ketterle,
Nature 435, 1047-1051 (2005)

The Unitary Fermi Gas

Only two length scales:
Interparticle spacing $n^{-1/3}$
De Broglie wavelength λ_{dB}

Only two corresp. energy scales:
Fermi energy E_F
Temperature $k_B T$



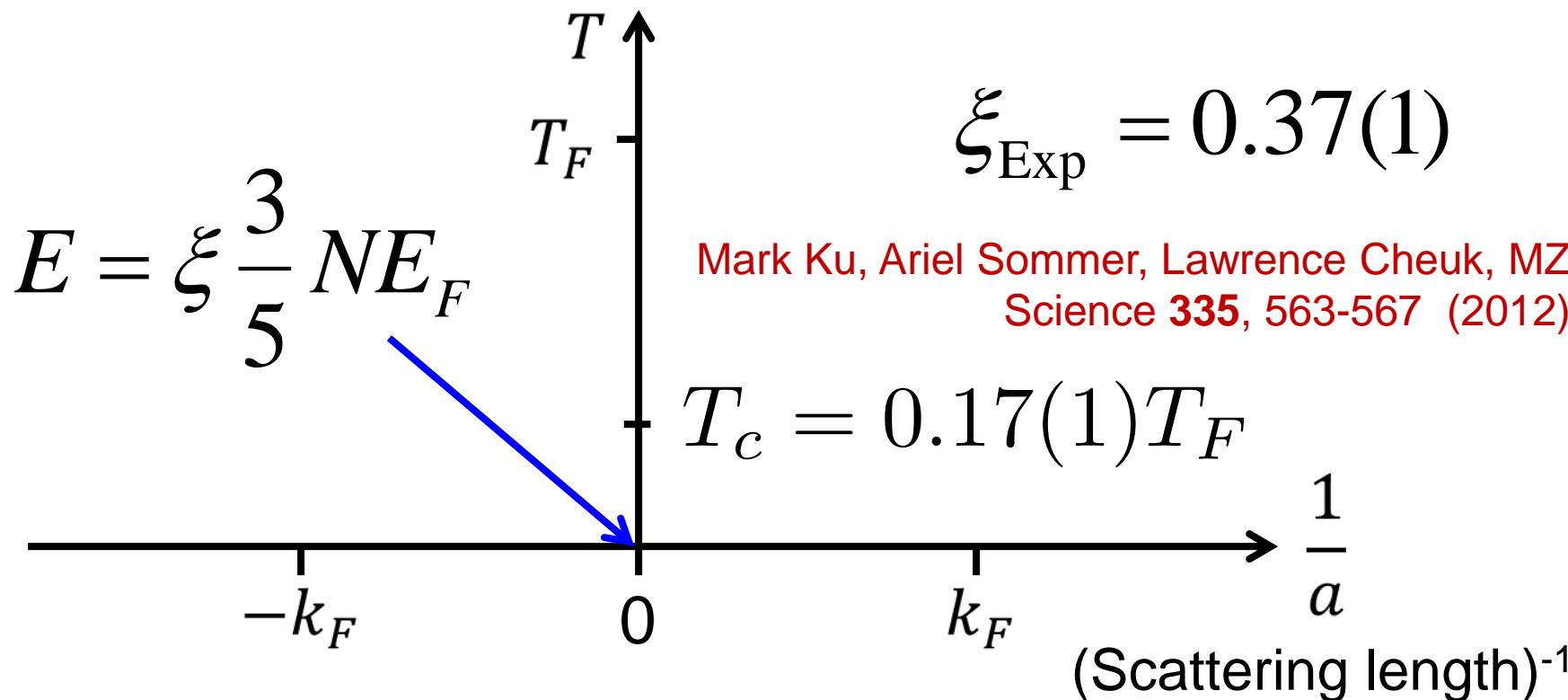
The Unitary Fermi Gas

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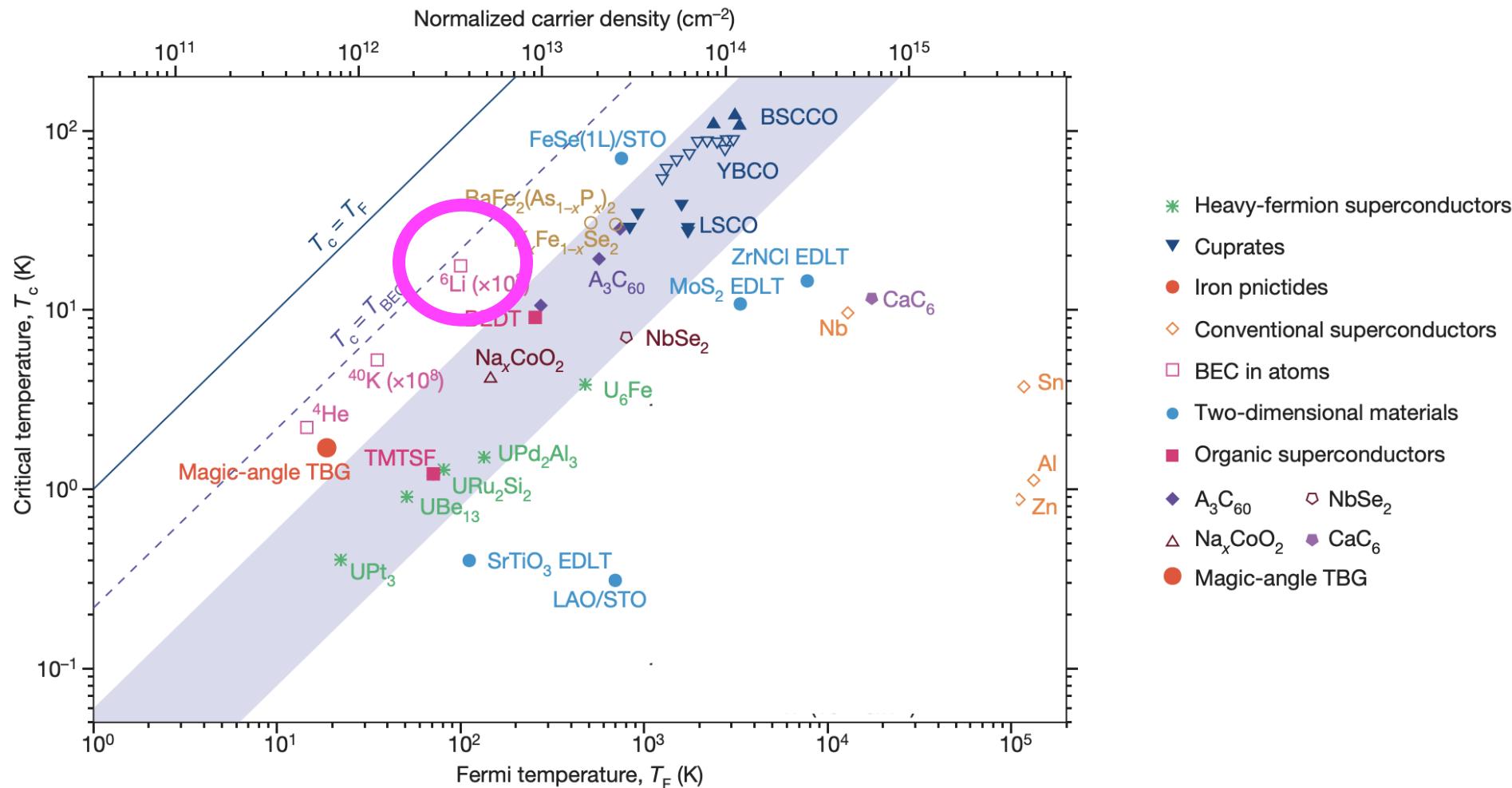
→ Universal equation of state

Duke/NC State, ENS, JILA, Innsbruck, Swinburne, MIT



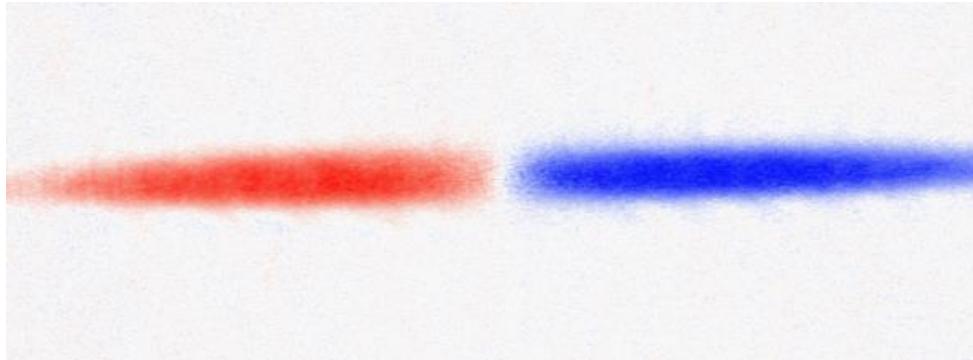
The Unitary Fermi Gas

Comparison to other dilute Fermi systems



Transport in Strongly Interacting Systems

e.g. Unitary Gas @ Feshbach Resonance



Mean free path \sim Interparticle spacing d

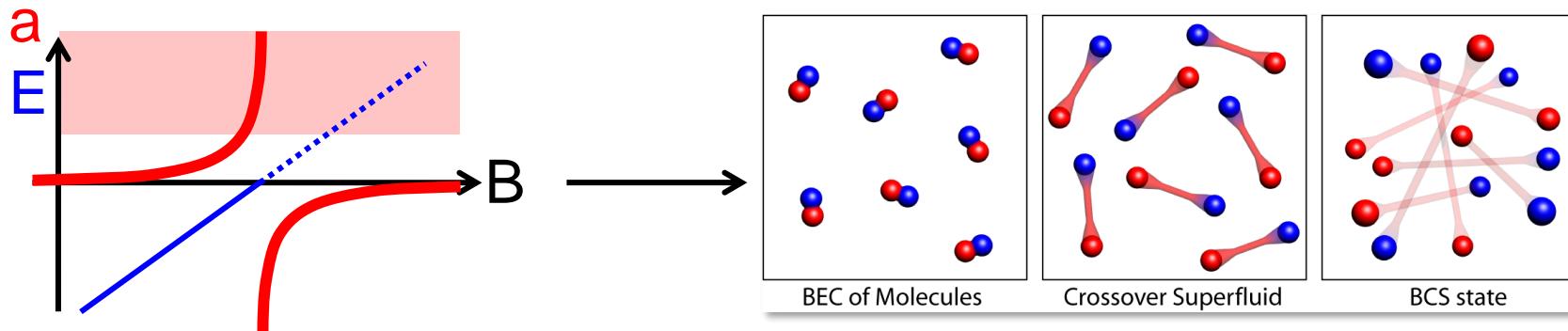
Average velocity $v \square \frac{\hbar}{md}$

Quantum Limit of Diffusion (charge, spin, momentum, thermal)

$$D \square v l \square \frac{\hbar}{m}$$

Sound in strongly interacting Fermi Gases

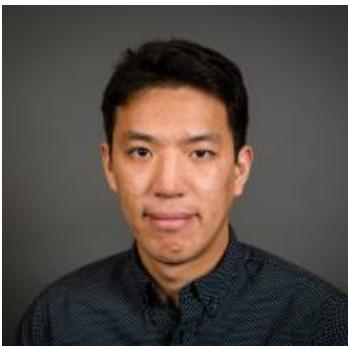
Feshbach Resonances → e.g. BEC-BCS Crossover



Quantum Gases → Quantum Fluids



Parth
Patel



Zhenjie
Yan



Biswaroop
Mukherjee

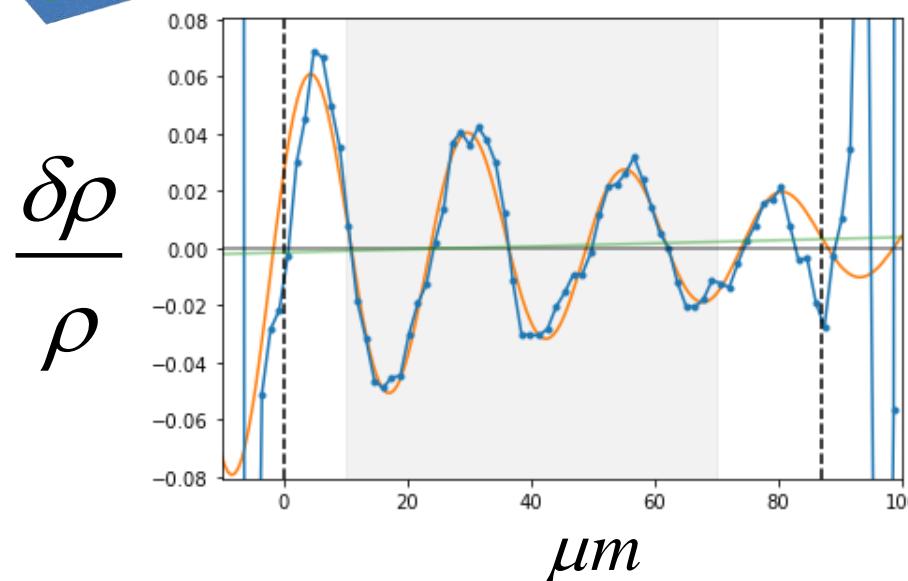
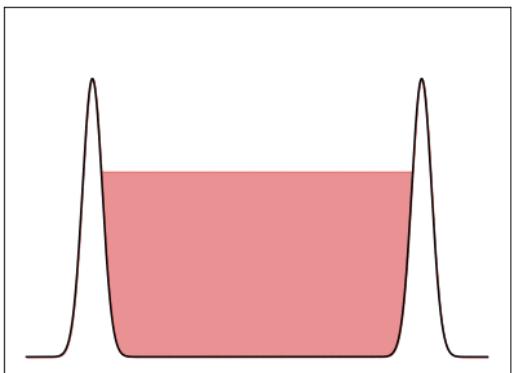
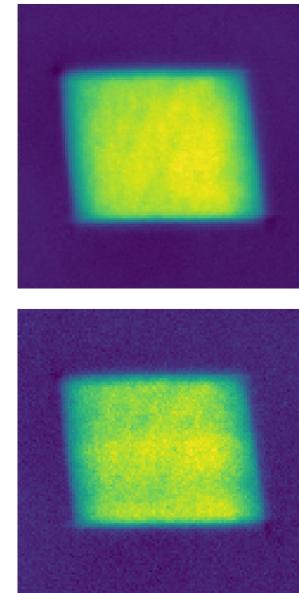
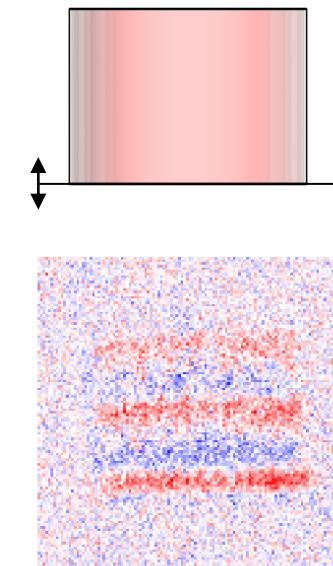
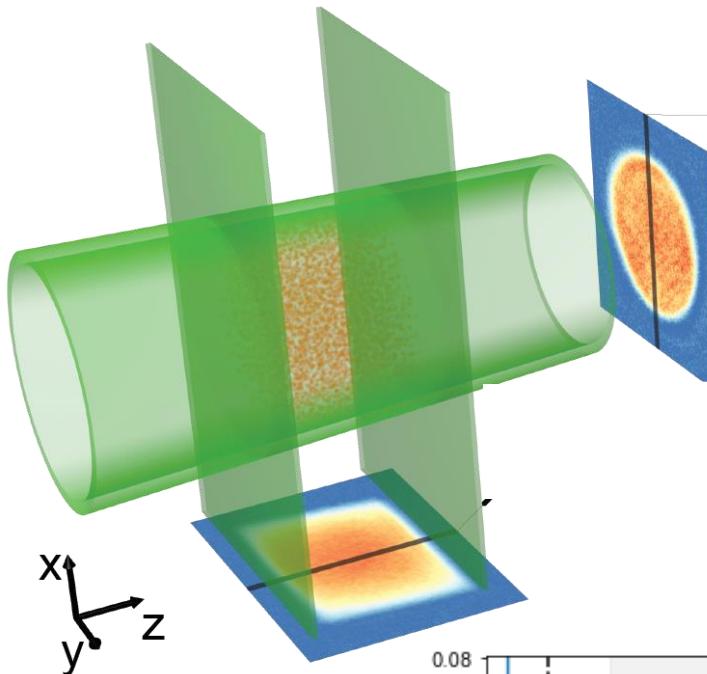


Dr. Richard
Fletcher



Dr. Julian
Struck

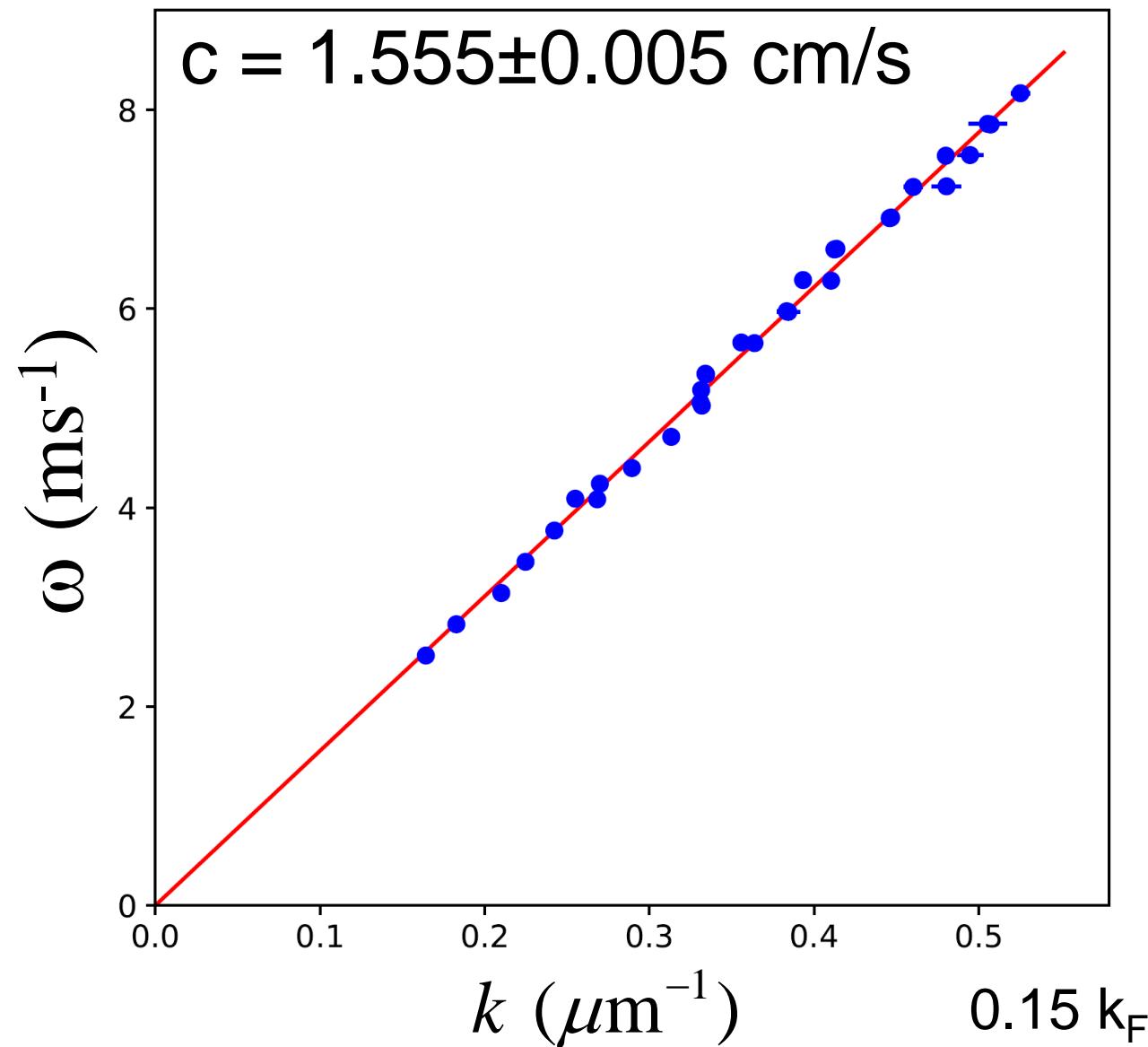
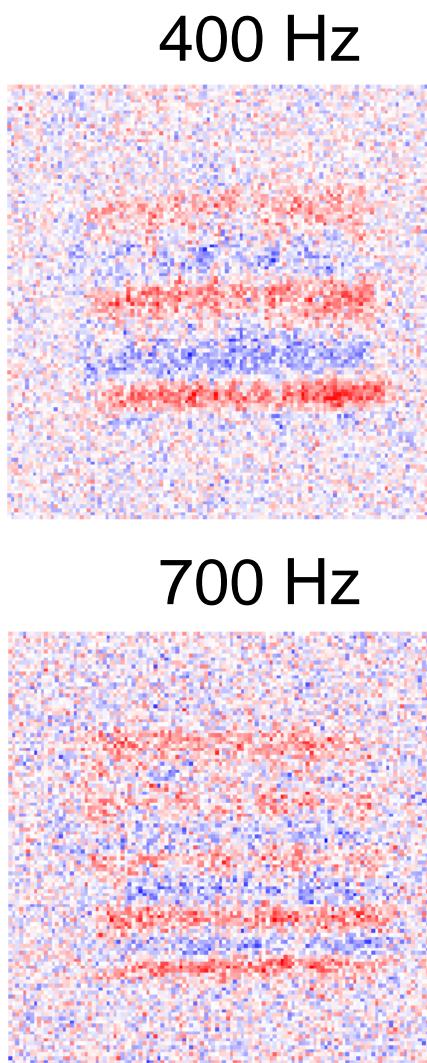
Creating Sound Waves in a Box



See also: 2D Bose gas, ENS (Dalibard),

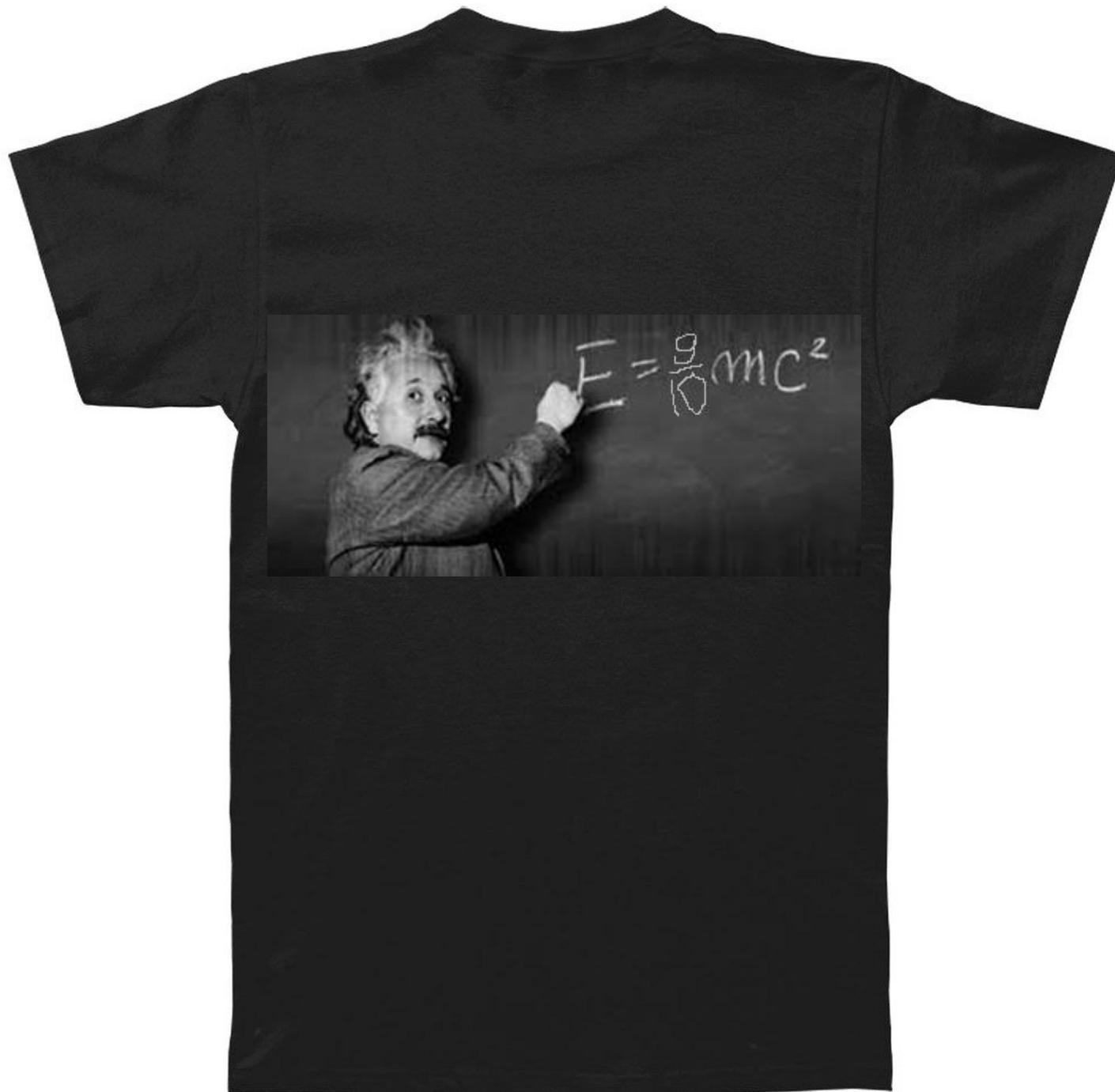
3D Bose gas, Cambridge (Hadzibabic), Fermi gas, NCState (Thomas), 2D Fermi Gas, Hamburg (Moritz)

Dispersion relation



$$E = \frac{9}{10} m c^2$$

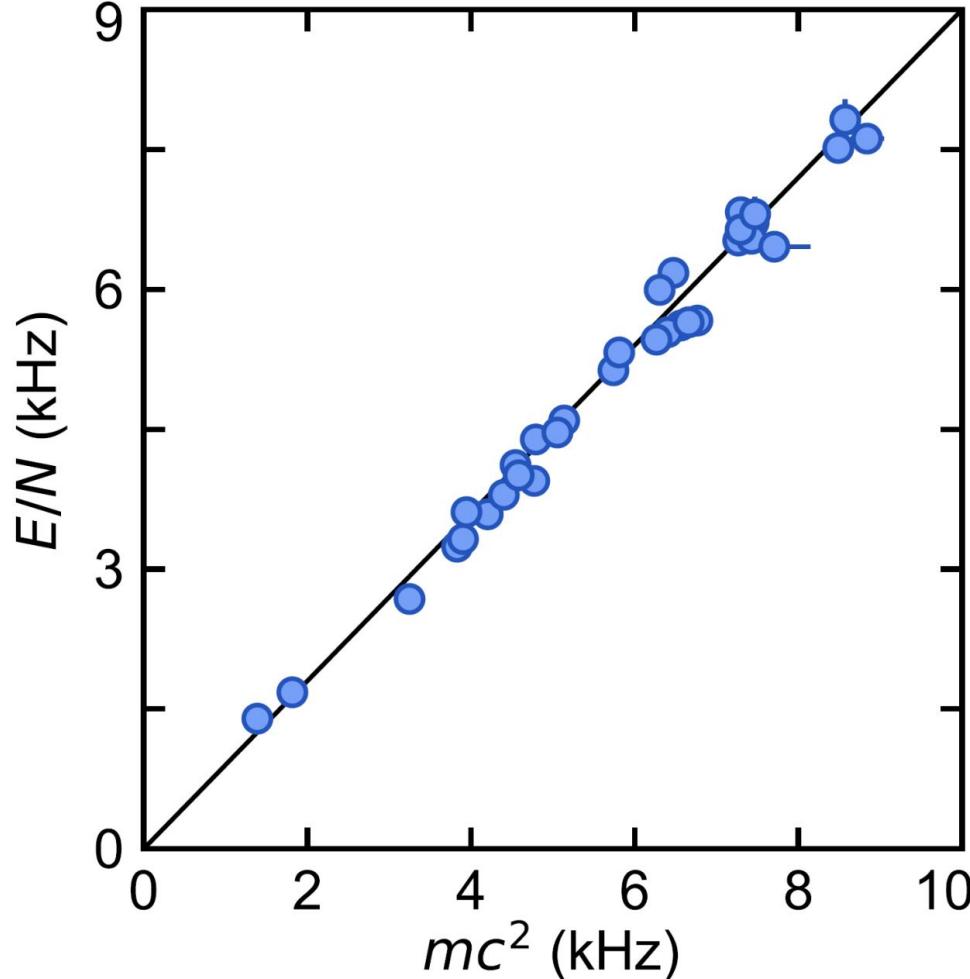
?



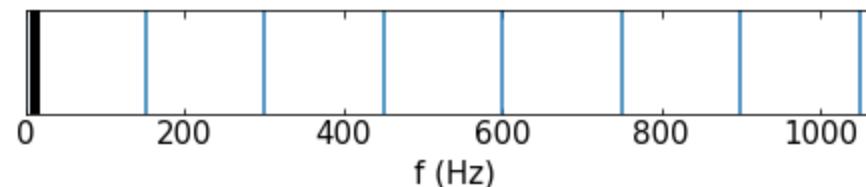
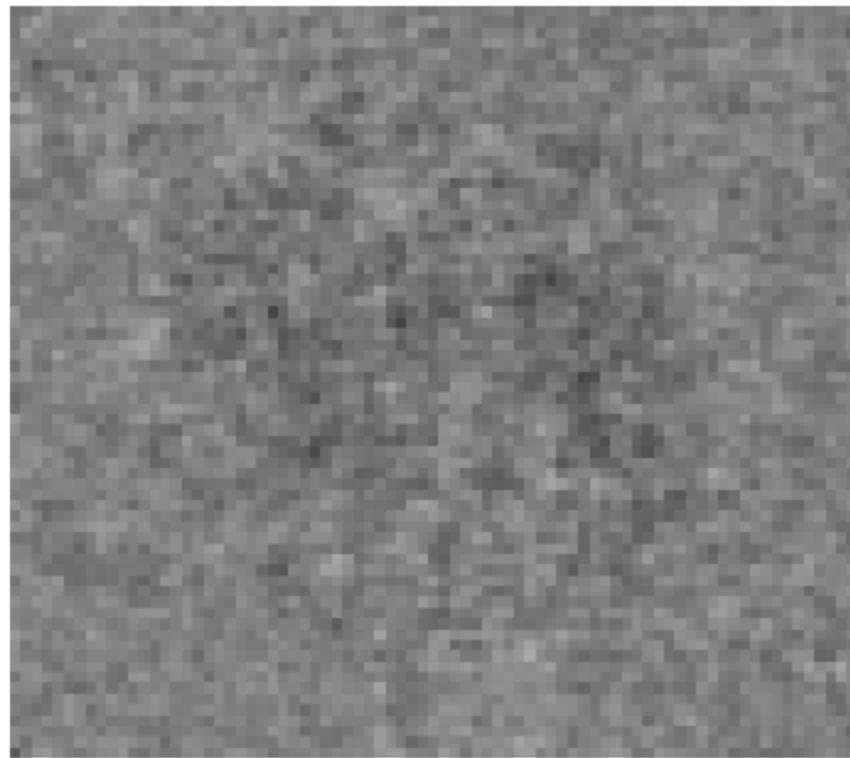
$$E = \frac{9}{10} m c^2$$

At unitarity:

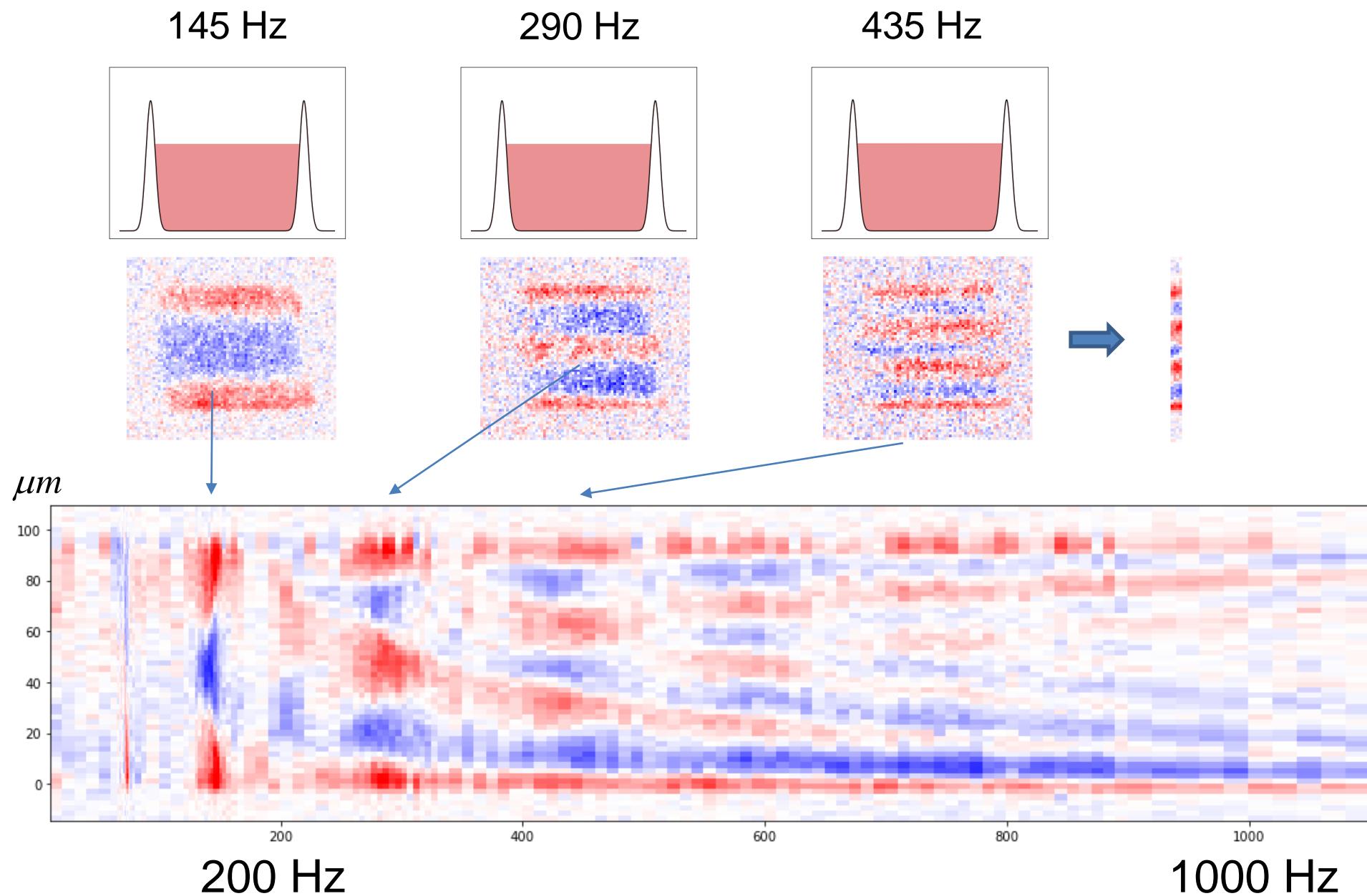
$$mc^2 = \left. \frac{\partial p}{\partial n} \right|_S = f(s) \left. \frac{\partial p_0}{\partial n} \right|_S = \frac{5}{3} \frac{p}{n} = \frac{10}{9} \frac{E}{N}$$



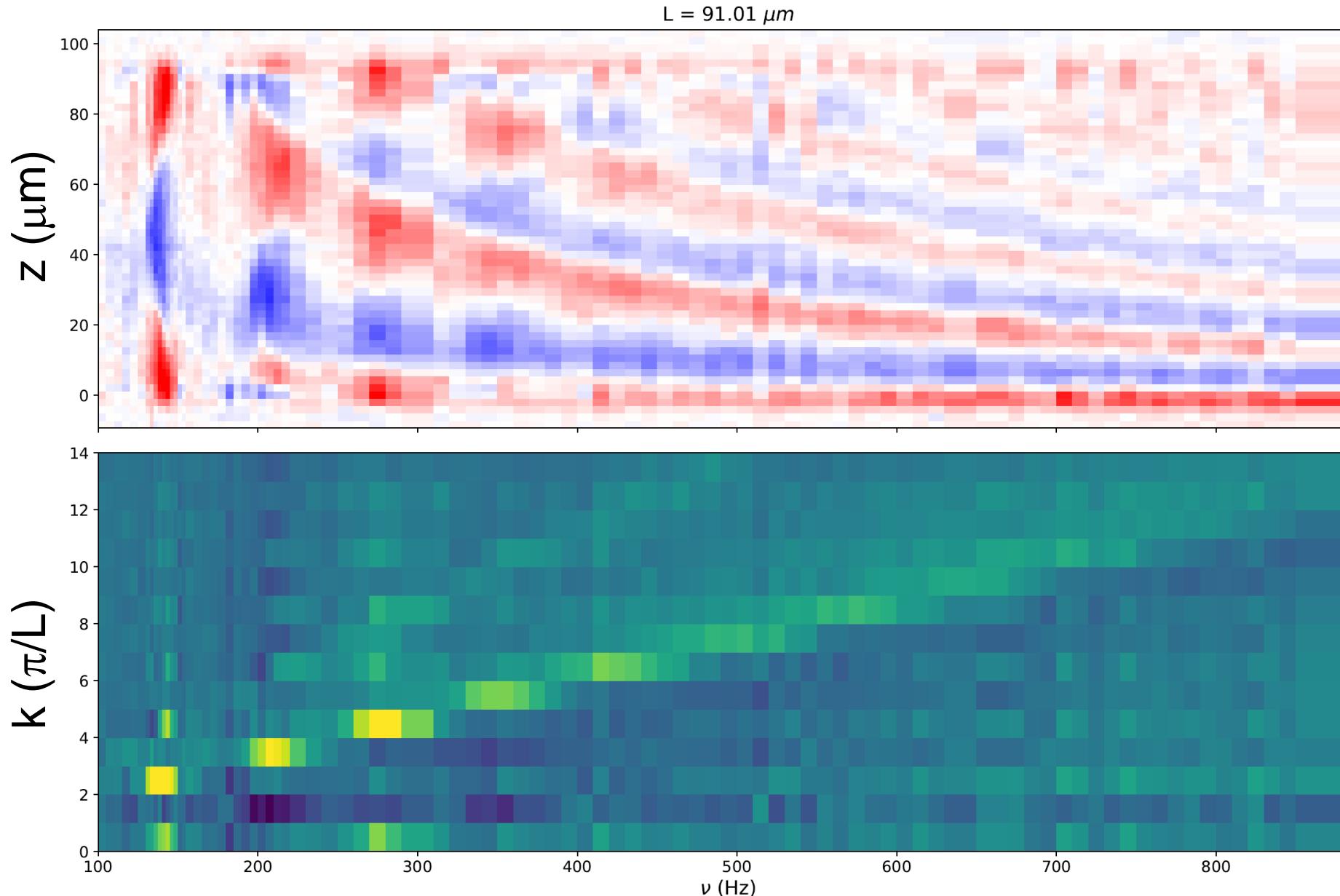
Resonant Modes



Resonant Modes



One-sided shaking: Even and odd

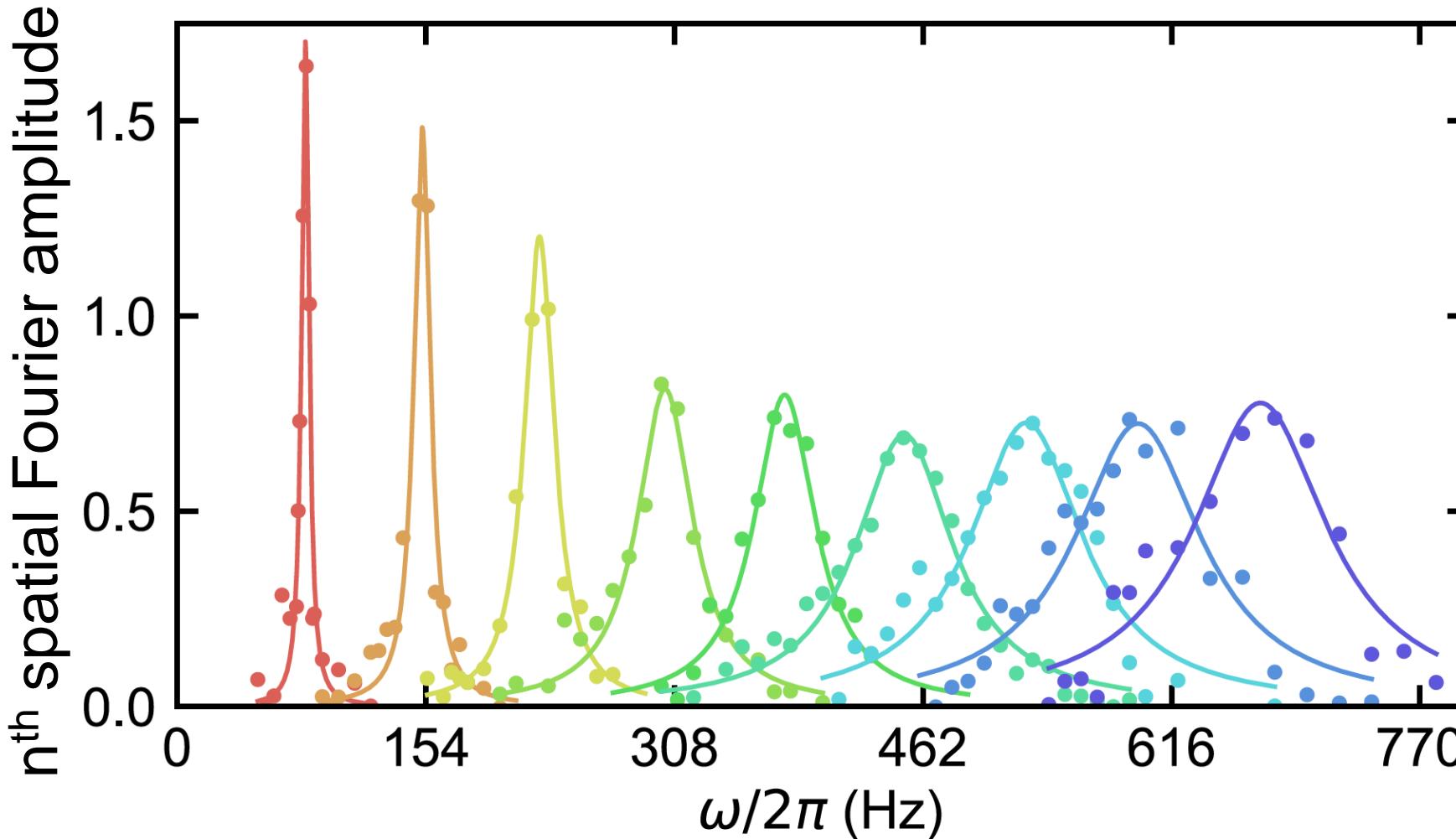


Unitary Hydrodynamics

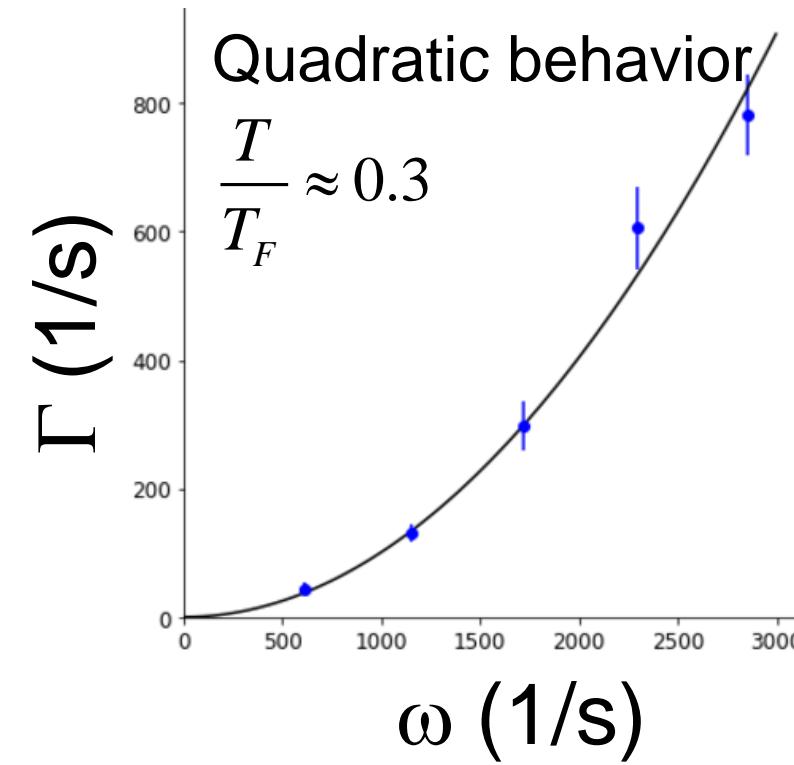
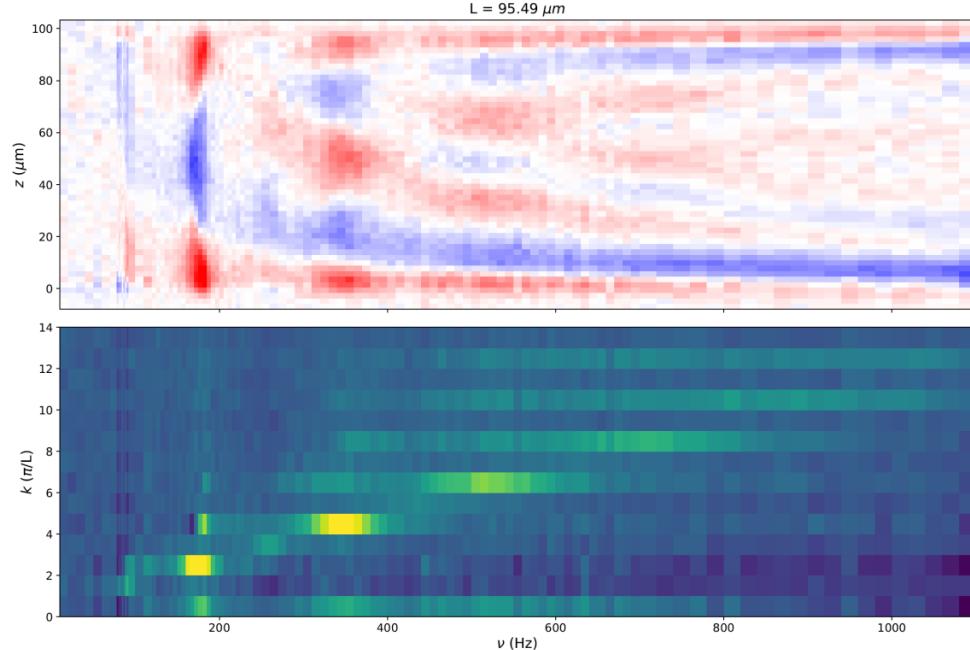
Mass	$\frac{\partial \rho}{\partial t} + \frac{\partial j}{\partial z} = 0$	Continuity equation
Momentum	$\frac{\partial j}{\partial t} + \frac{\partial p}{\partial z} = \frac{4}{3} \eta \frac{\partial^2 v}{\partial^2 z}$	Navier-Stokes equation viscosity damps fluid
	No bulk viscosity (ζ) due to scale invariance	
Heat	$\frac{\partial s}{\partial t} + \frac{\partial (vs)}{\partial z} = \frac{\kappa}{T} \frac{\partial^2 T}{\partial^2 z}$	thermal conduction increases entropy
Sound Diffusivity due to viscosity & thermal conductivity		$D = \frac{4}{3} \frac{\eta}{\rho} + \frac{4}{15} \frac{\kappa T}{P}$
Dispersion relation		$\omega^2 = c^2 k^2 + i \omega D k^2$
Damping rate:	$\Gamma = D k^2$	

Sound resonances

A direct measurement of the density response function



Sound Diffusivity from Sonogram Peaks



$$D = 2.1 \frac{\hbar}{m}$$

Quantum limited sound diffusion

Spin diffusion: Sommer et al., MIT 2011

Viscosity: Schaefer, Thomas (e.g. Science 2011)

Theory: Enss, Haussmann, Zwerger, Taylor, Randeria et al.

Viscosity of Superfluid Helium-4

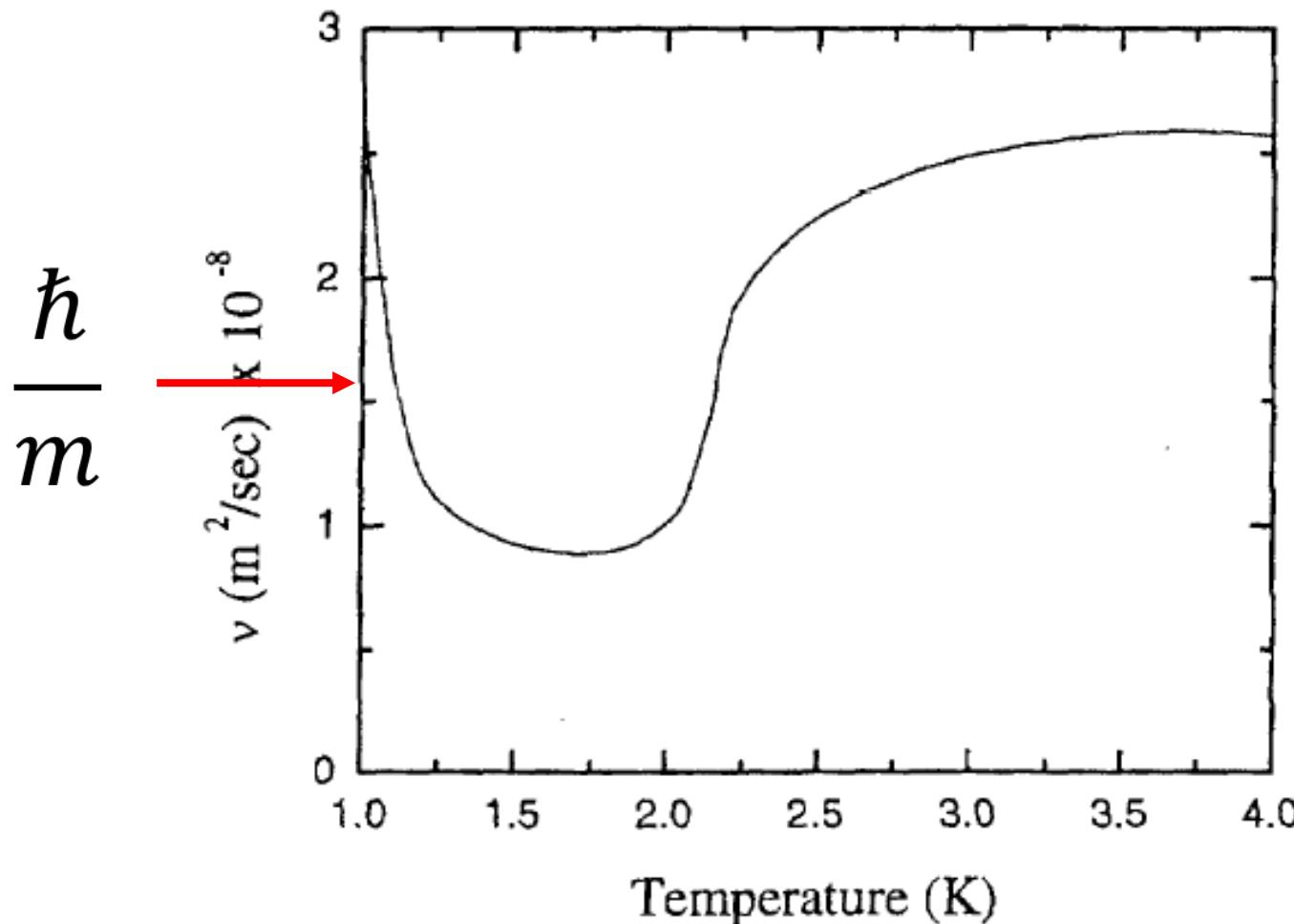
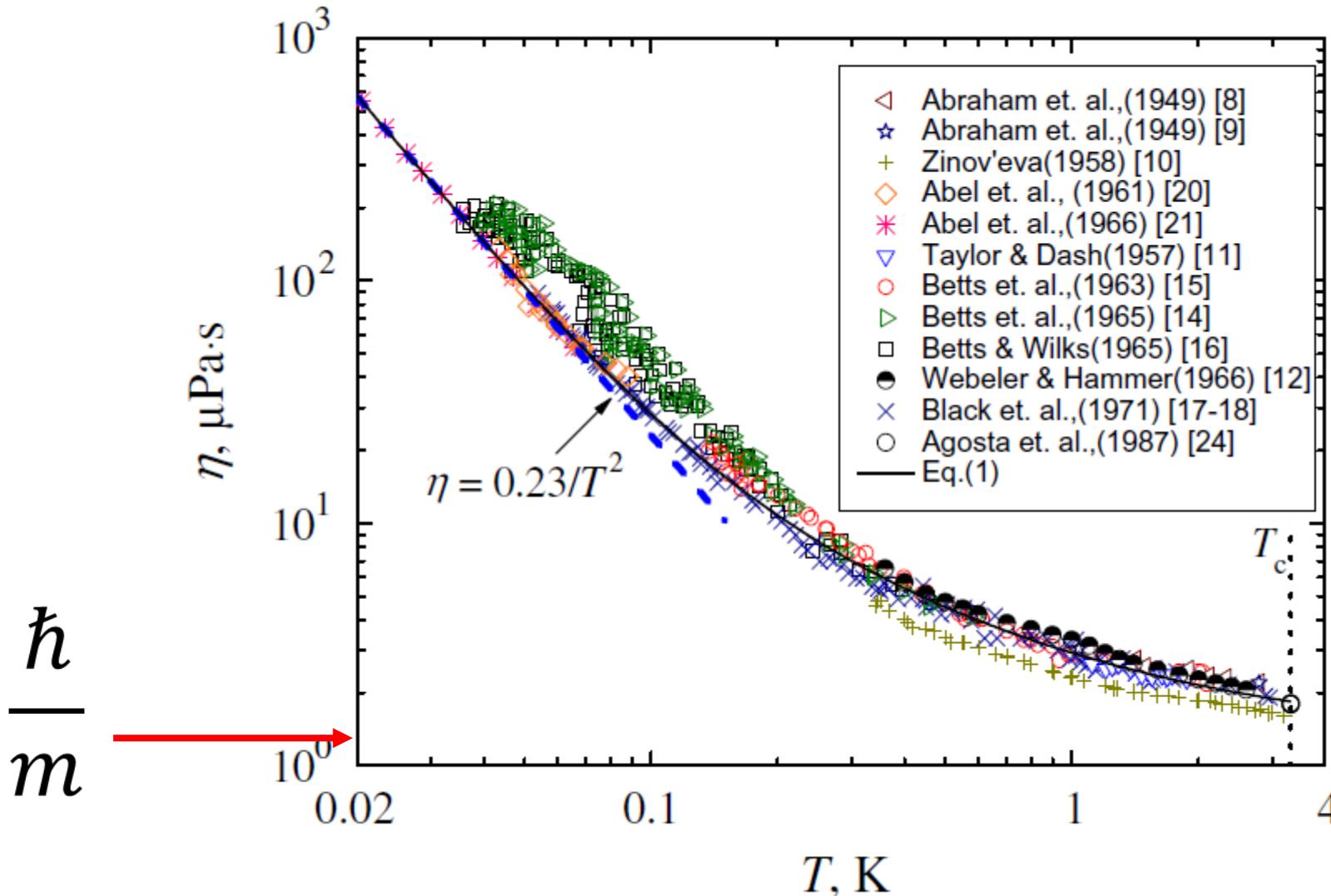


FIG. 12.3. The recommended values of the kinematic viscosity of liquid ${}^4\text{He}$, $\nu = \eta/\rho$, as a function of temperature at the saturated vapor pressure.

From: Donnelly, Barenghi, JPC Reference Data, 1998

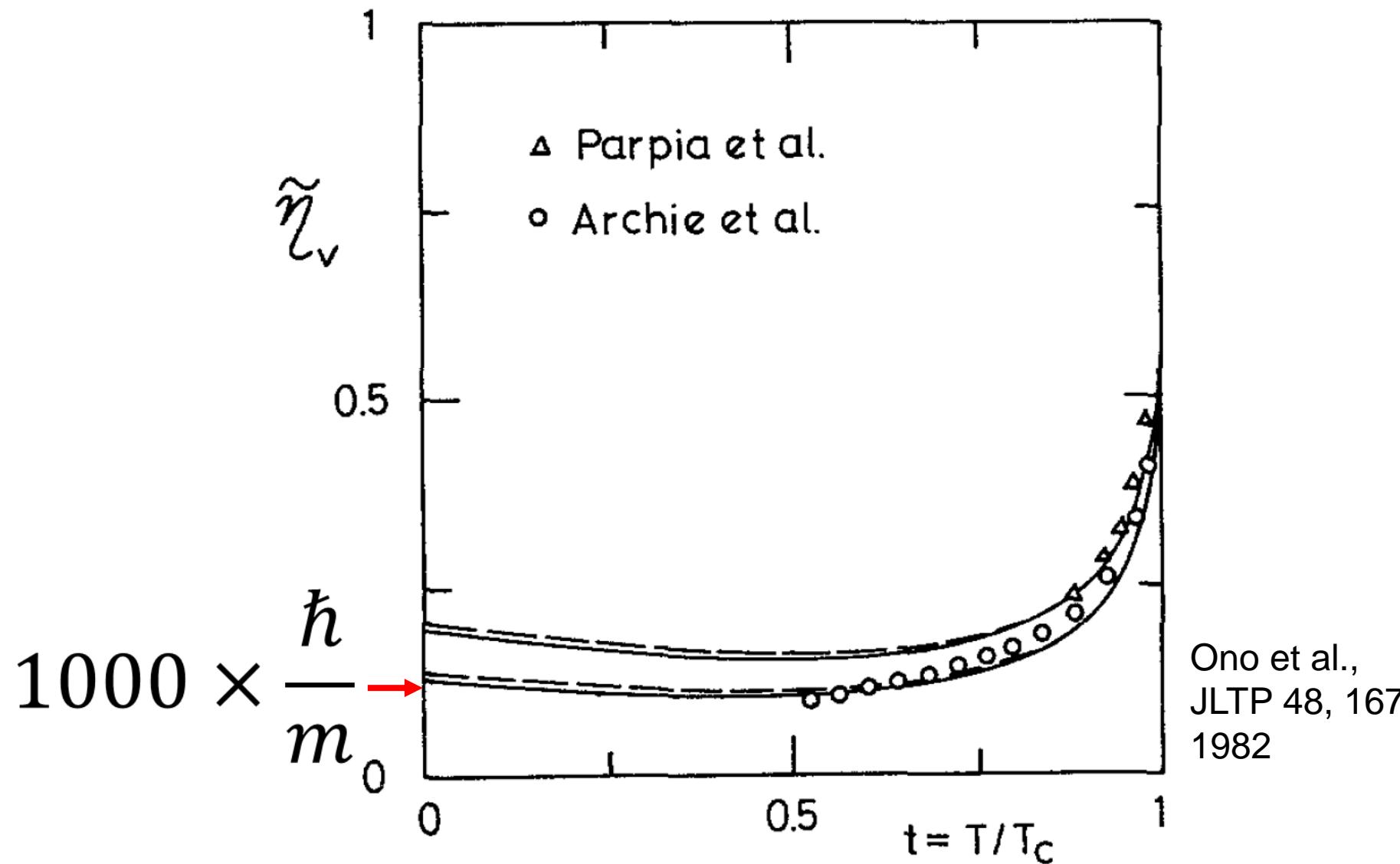
Viscosity of Superfluid Helium-3

Helium-3 above T_c : Fermi Liquid behavior

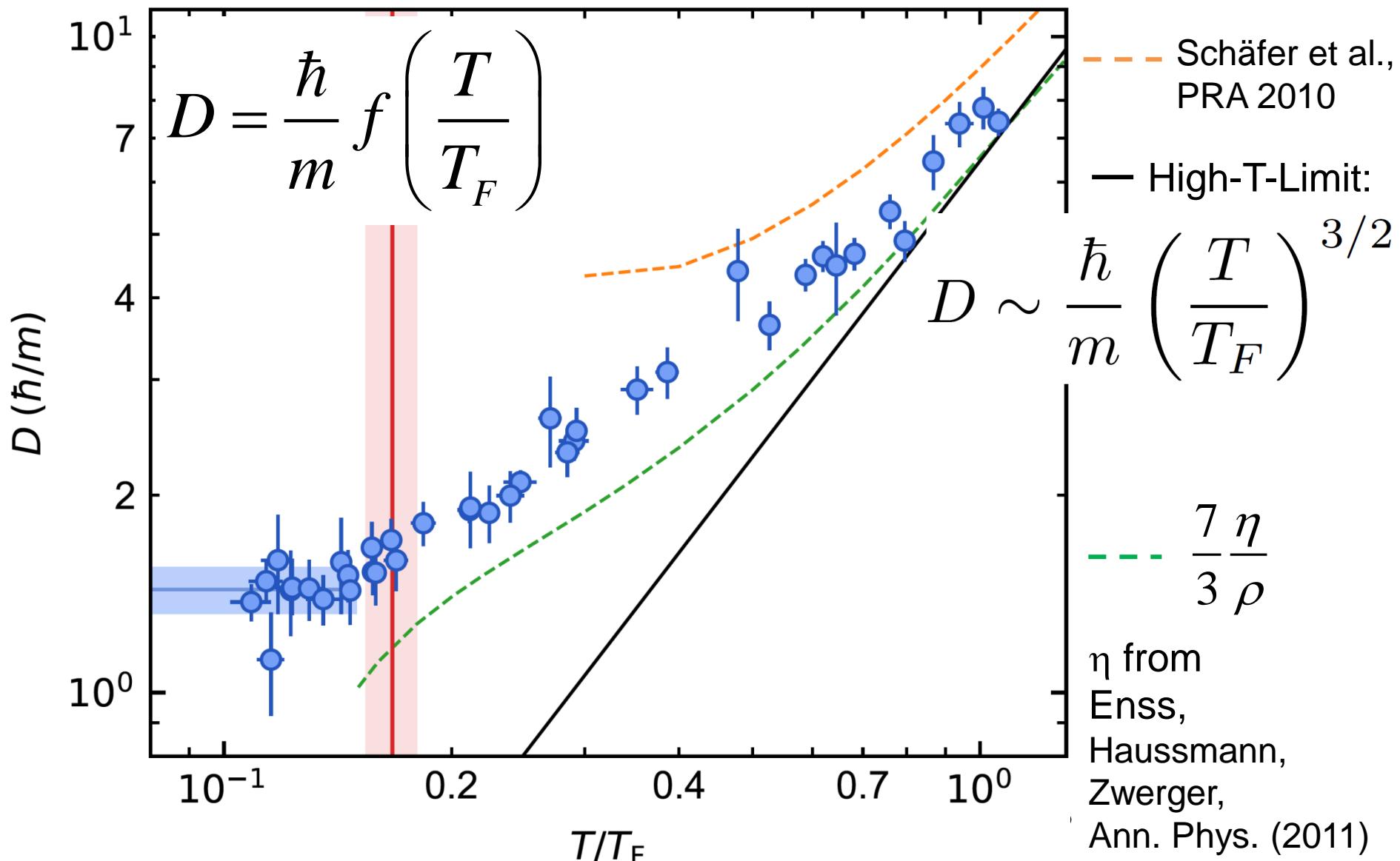


Viscosity of Superfluid Helium-3

Helium-3 below T_c : Sudden decrease, then constant



Quantum Limited Sound Diffusion



Comparison to state of the art theory

Quantum critical thermal transport in the unitary Fermi gas
Bernhard Frank, Wilhelm Zwerger, Tilman Enss
Phys. Rev. Research 2, 023301 (2020)

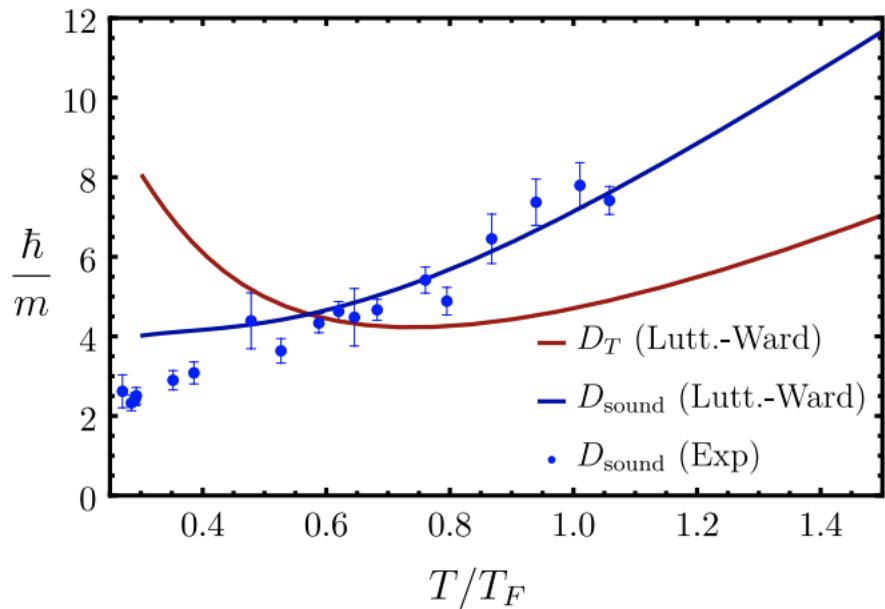


FIG. 1. Thermal diffusivity D_T (red) and sound diffusivity D_{sound} (blue) vs temperature T/T_F in the quantum critical regime of the unitary Fermi gas. Theoretical results from Luttinger-Ward calculations are shown in comparison with sound diffusion measurements [23].

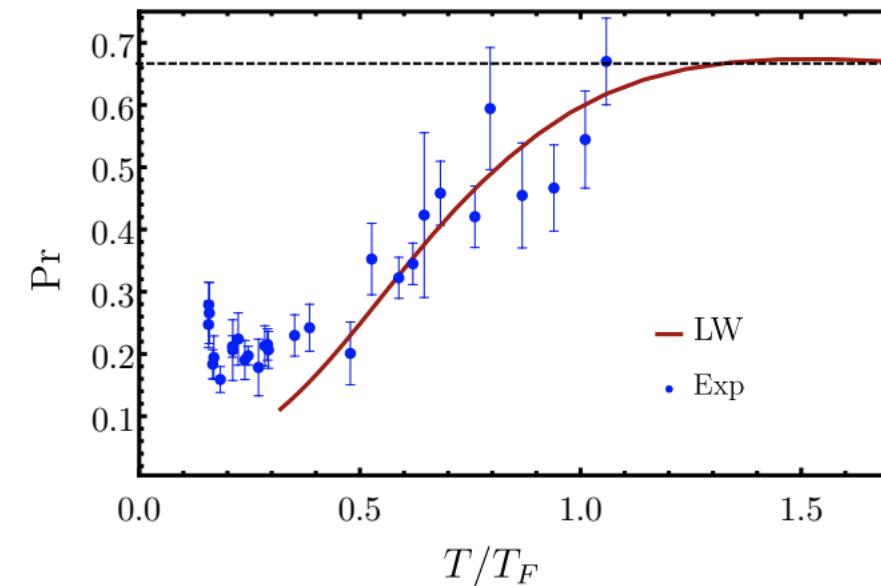


FIG. 8. Prandtl number $P_r = D_\eta/D_T$ vs temperature T/T_F from Luttinger-Ward calculations (red line) and sound attenuation measurements [23]; the dashed line marks the high-temperature limit $P_r = 2/3$.

Unitary Two-Fluid Hydrodynamics

For 1D flow, linearized:



Laszlo Tisza

$$\frac{\partial \rho}{\partial t} + \frac{\partial j}{\partial z} = 0$$

Continuity equation

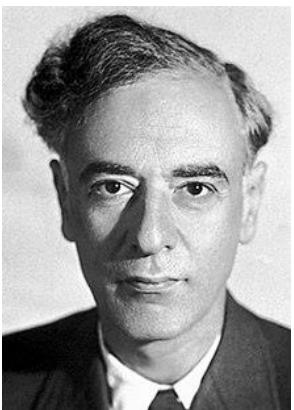
$$\frac{\partial j}{\partial t} + \frac{\partial p}{\partial z} = \frac{4}{3} \eta \frac{\partial^2 v_n}{\partial^2 z}$$

Navier-Stokes eq.
viscosity

damps normal fluid

$$\frac{\partial s}{\partial t} + \frac{\partial(v_n s)}{\partial z} = \frac{\kappa}{T} \frac{\partial^2 T}{\partial^2 z}$$

Entropy moving
with normal fluid,
thermal conduction
increases entropy



Lev Landau

Superfluid flow:

bulk viscosity

$$\frac{\partial v_s}{\partial t} + \frac{\partial \mu/m}{\partial z} = \rho_s \zeta_3 \frac{\partial^2(v_s - v_n)}{\partial^2 z}$$

normal-superfluid
interconversion

No other transport coefficients ($\zeta_1, \zeta_2, \zeta_4$)

due to scale invariance (Son, Stringari&Pitaevskii)

+Thermodynamics relating p, ρ, T, s, μ

First and Second Sound

Two fluids → Two sound modes

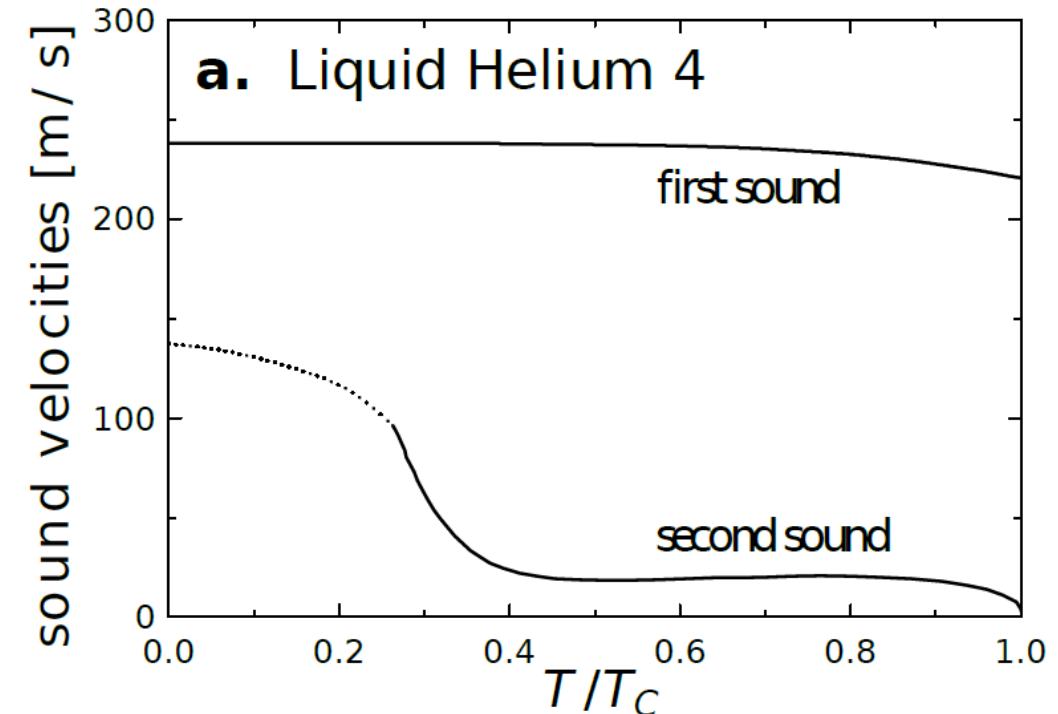
Review: *Stringari&Pitaevskii*
arXiv:1510.01306

For small expansivity: Density and Heat

$$c_{10}^2 = \left. \frac{\partial p}{\partial \rho} \right|_{S,N} = \frac{1}{\rho \kappa_s}$$

First sound, a density wave
normal and superfluid
oscillate in phase

$$c_{20}^2 = \left. \frac{\rho_s}{\rho_n} \frac{\partial T}{\partial \sigma} \right|_\rho \frac{\sigma^2}{\gamma} = \frac{\rho_s}{\rho_n} \frac{\sigma^2 T}{c_p}$$

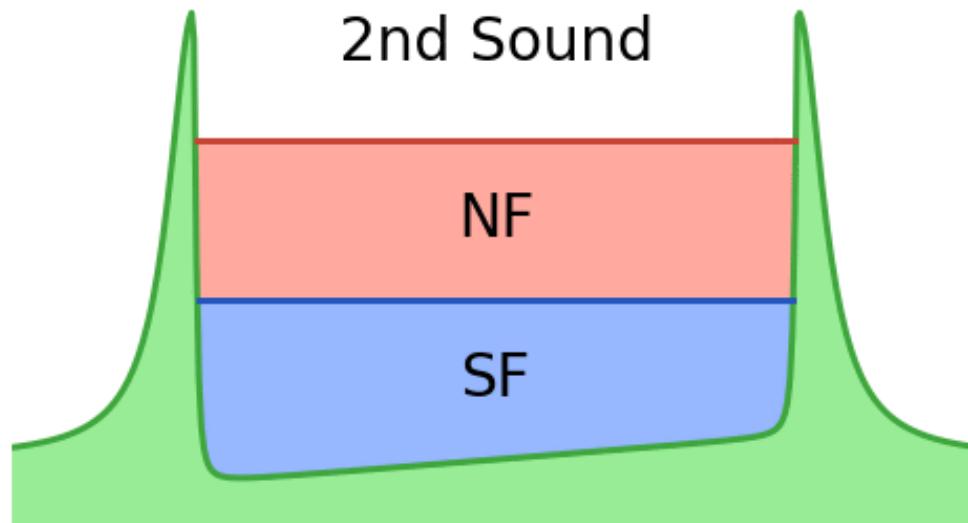


Second sound, a heat wave,
normal and superfluid oscillate out of phase

Second sound in a box

(Isentropic) expansivity at unitarity $-\frac{1}{\rho} \frac{\partial \rho}{\partial T} \Big|_S = -\frac{3}{2} \frac{1}{T}$ always non-zero
→ Can drive temperature wave using potential acting on density

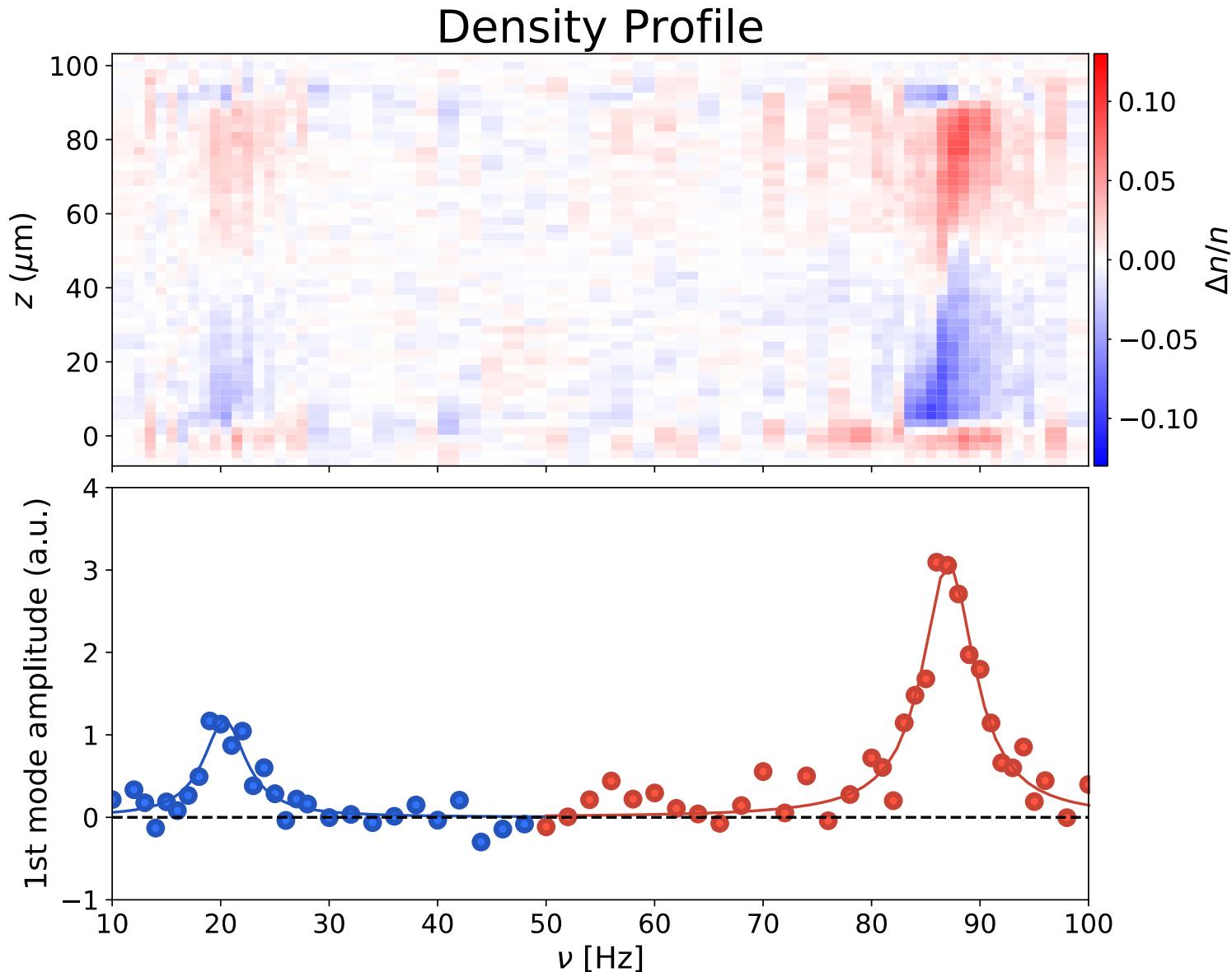
Apply oscillating gradient to the box



Observation of Second Sound in Quasi-1D geometry:
Sidorenkov et al., Grimm, Stringari&Piatevskii, Nature 2013

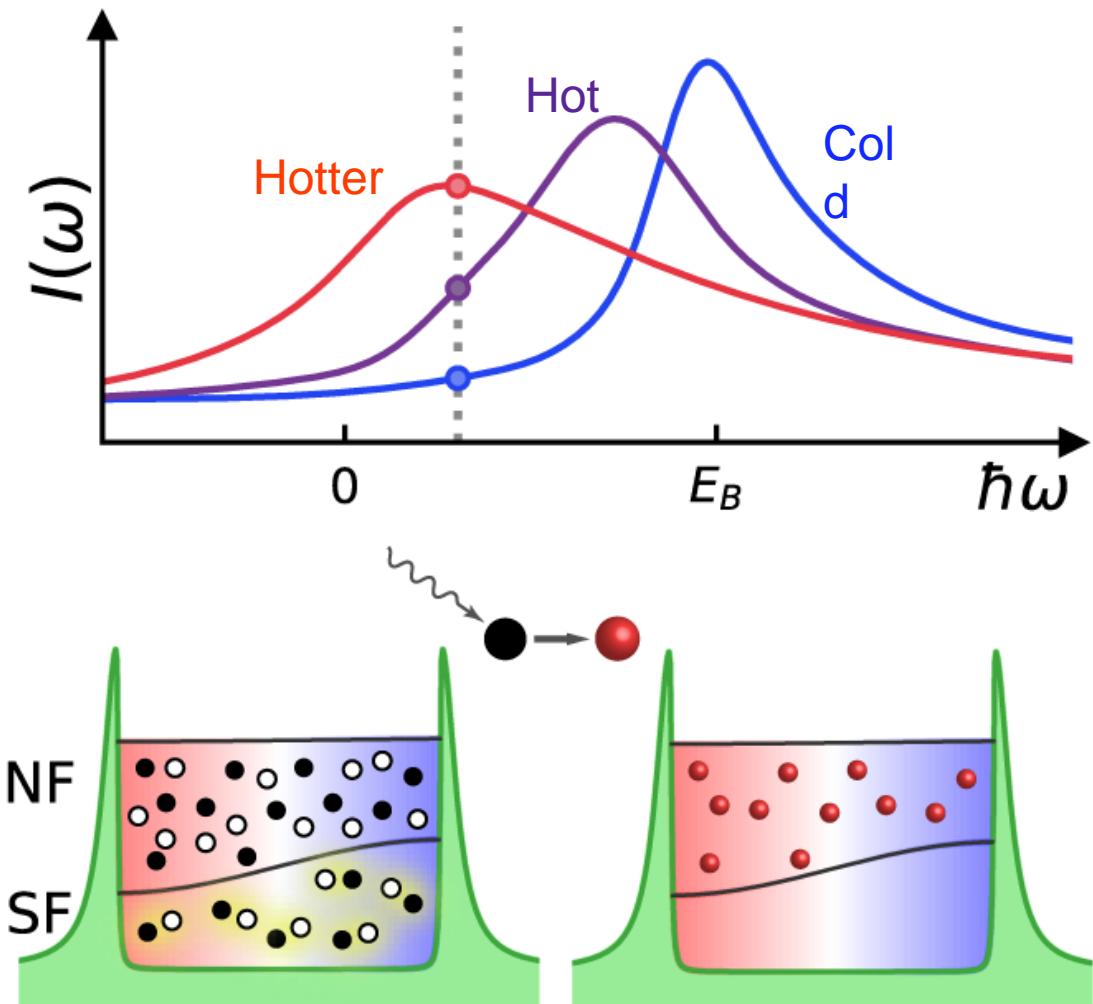
On bosons: weakly interacting situation: JILA (Debbie Jin 1996), MIT
Hydrodynamic gas: van der Straten

Second Sound seen in Density

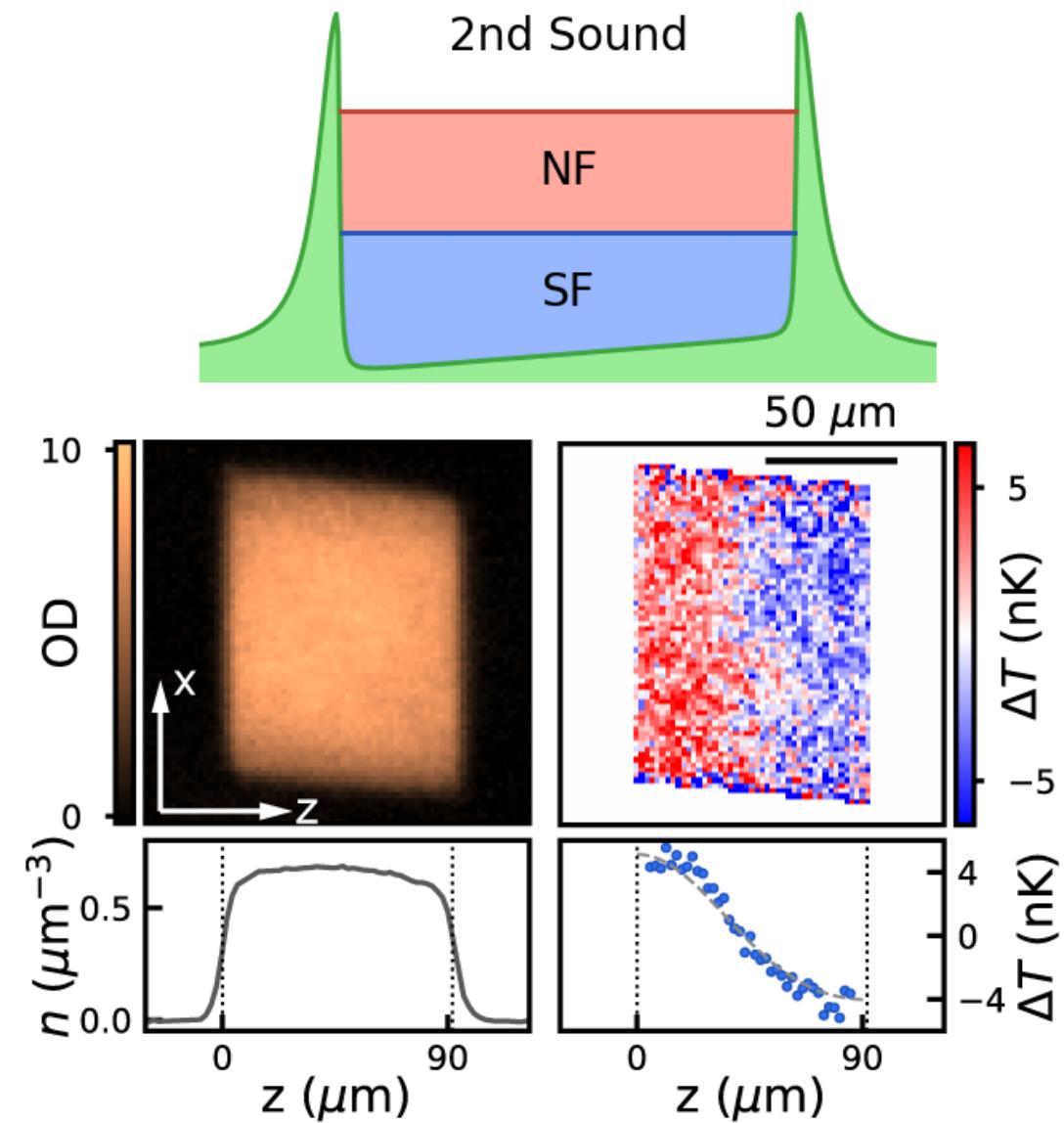


A local thermometer for heat transport

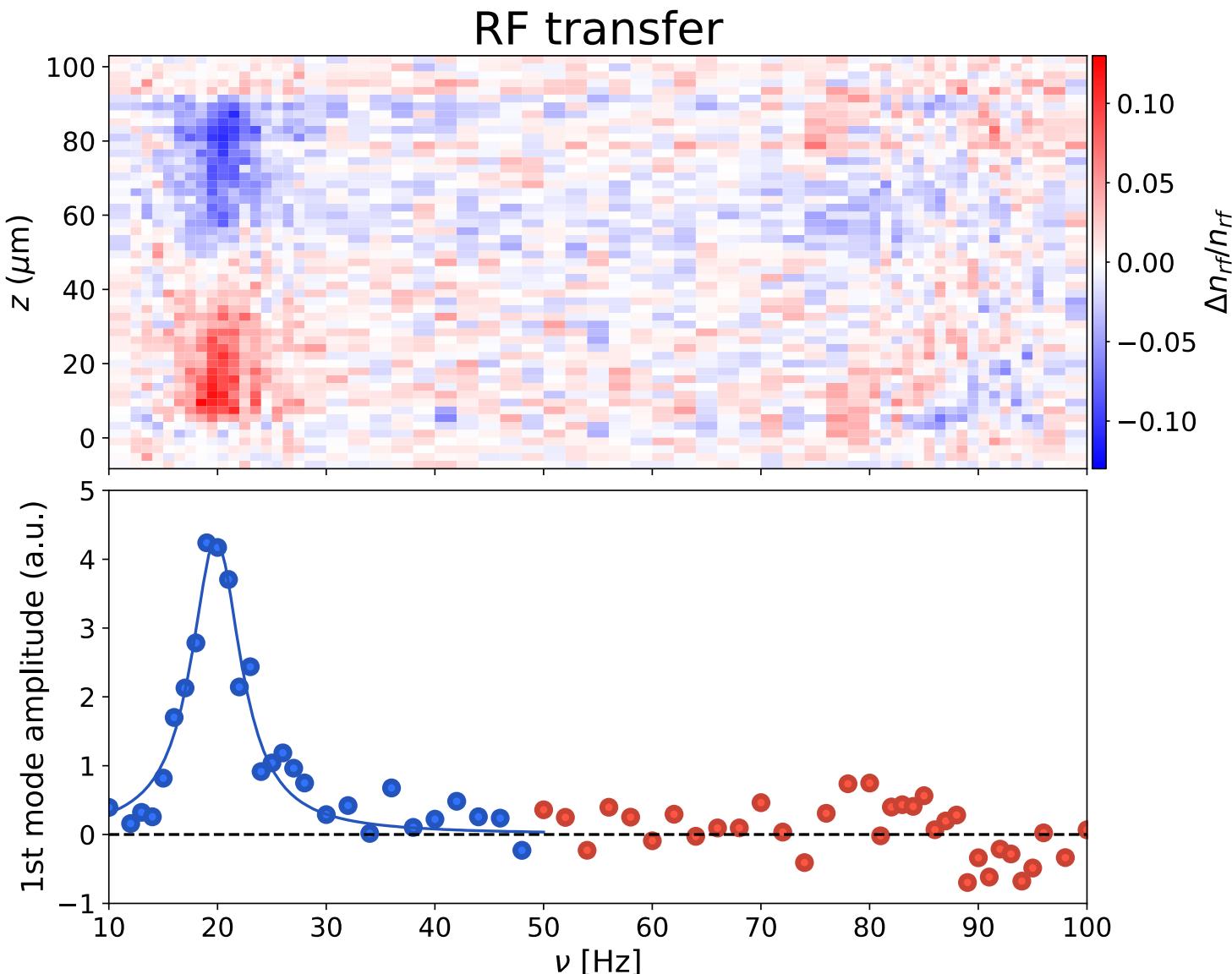
RF absorption spectra



2nd Sound Creation

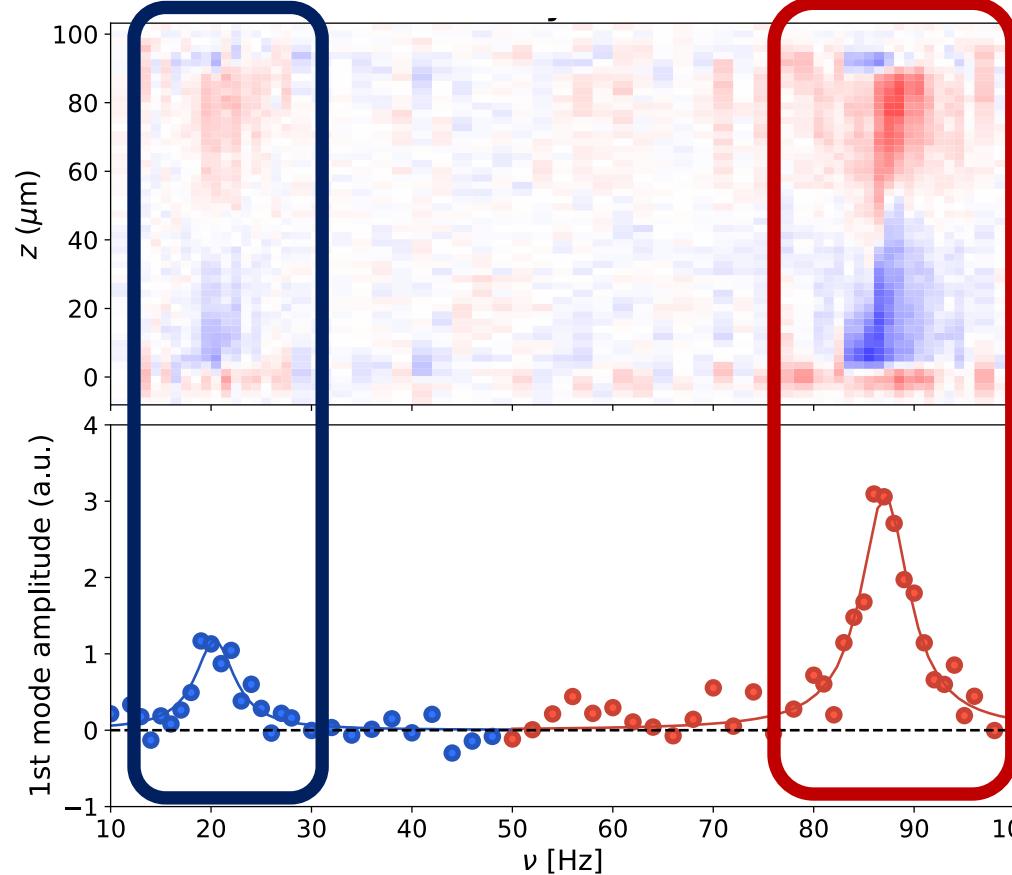


Second Sound with local thermometer

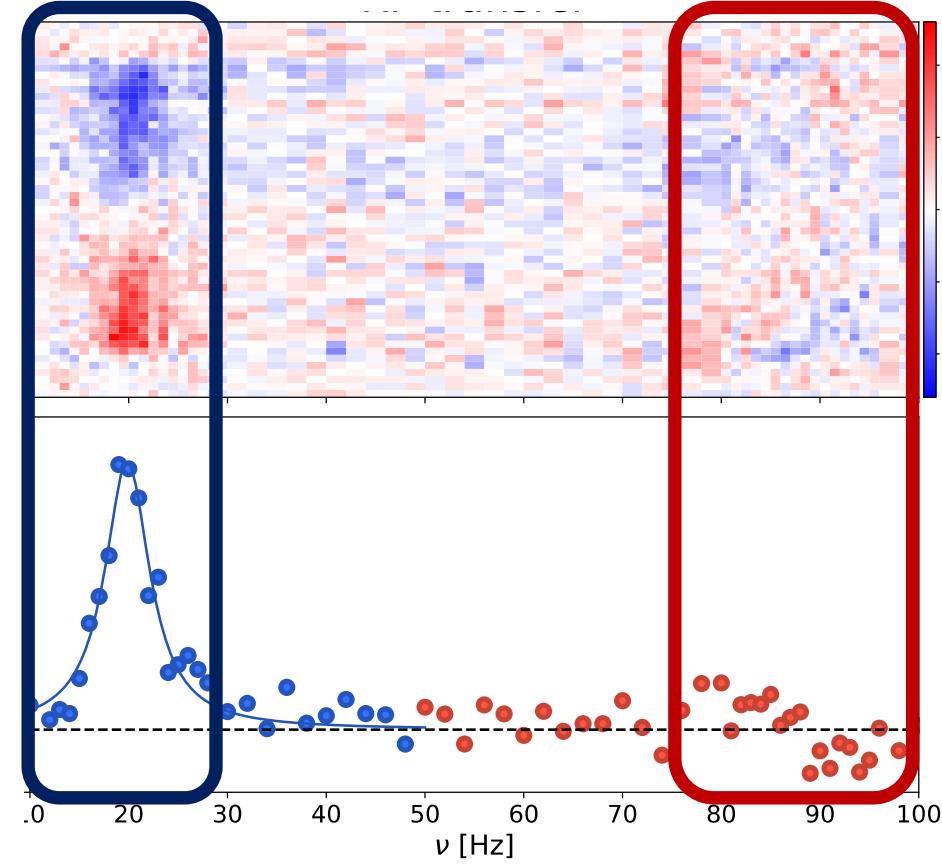


First and Second Sound

Density Probe

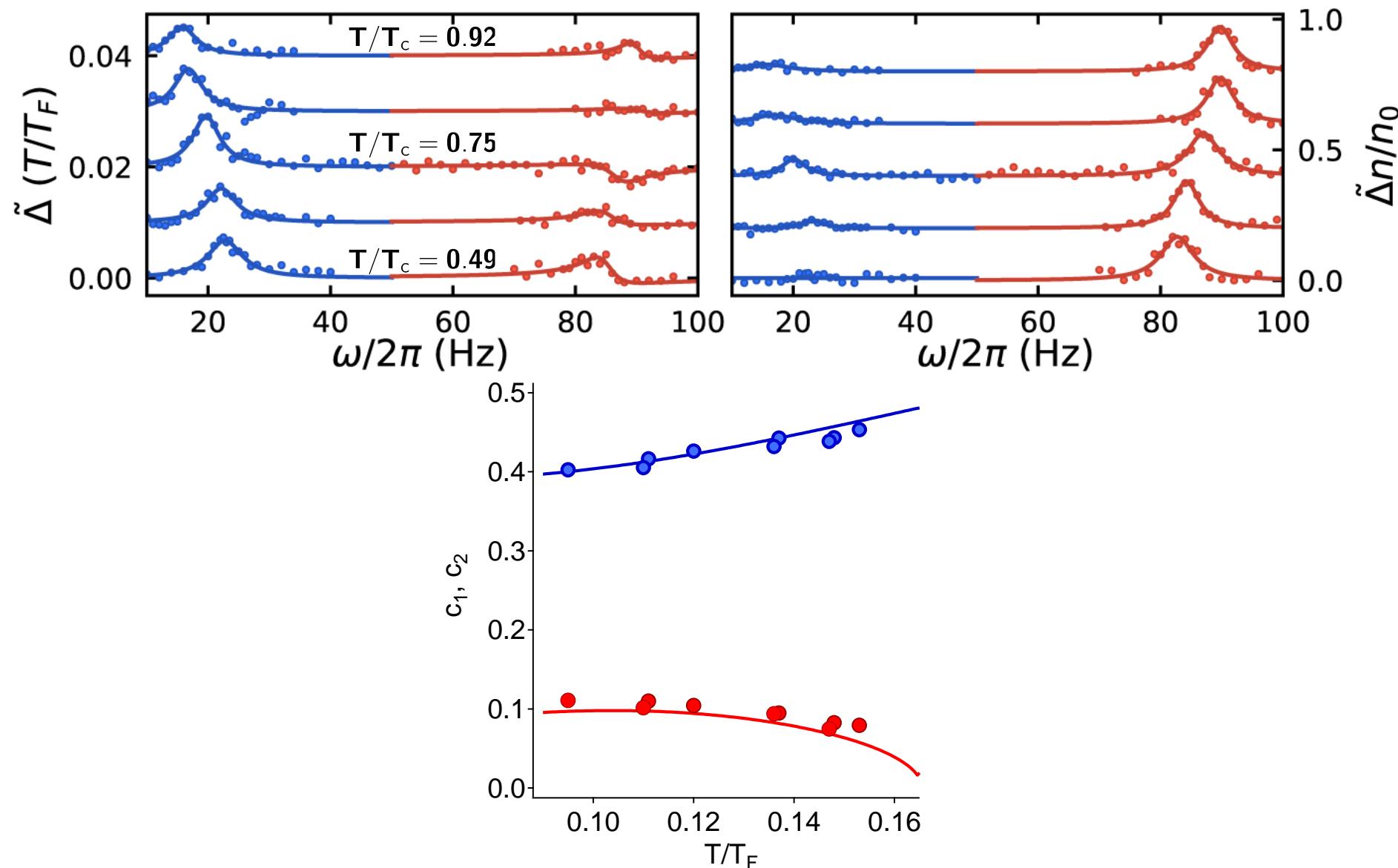


Temperature Probe
after RF transfer, n_{RF}

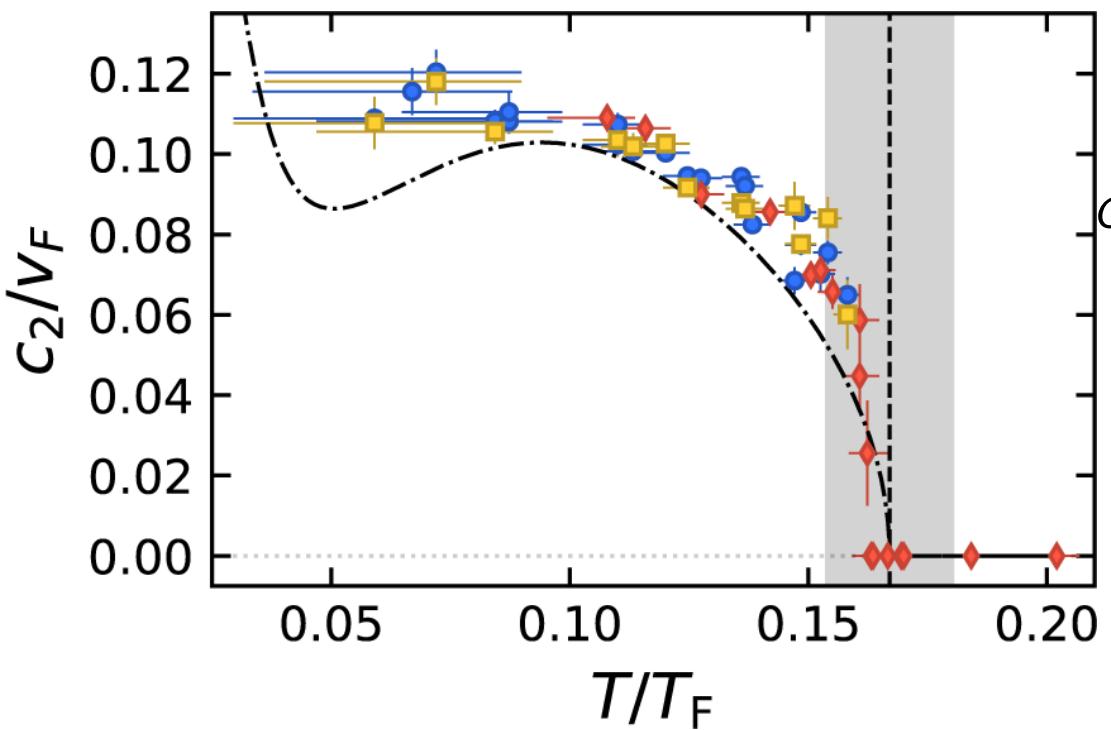


→ Speeds and Decay rates of First and Second sound

Speed of first and second sound

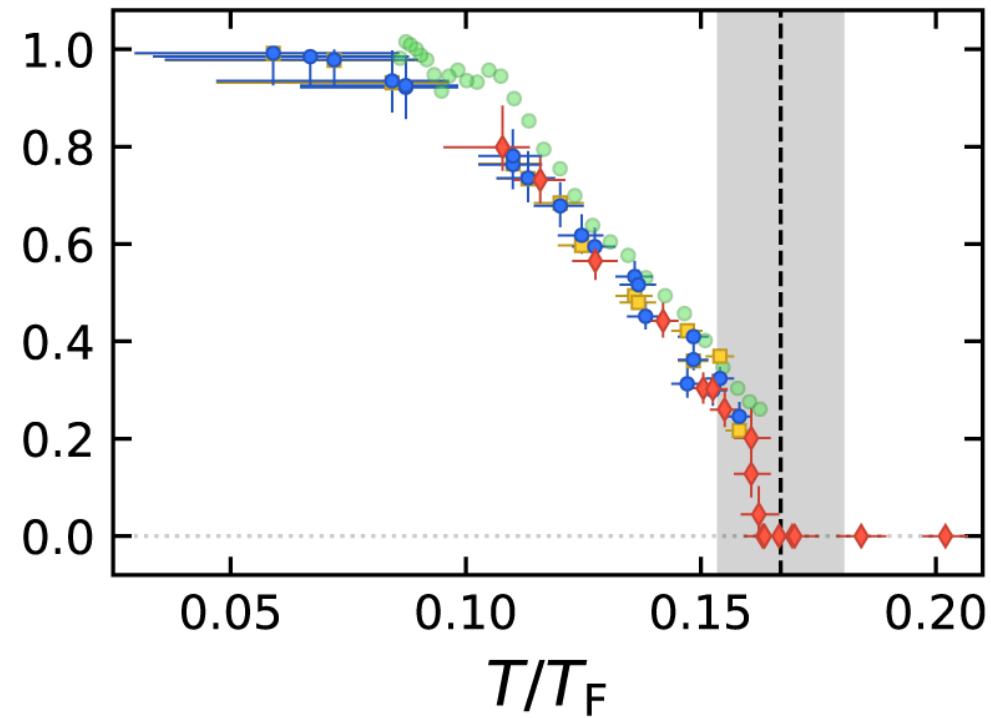


Superfluid fraction



$$C_2^2 = \frac{\rho_S}{\rho_N} \frac{TS^2}{C_p}$$

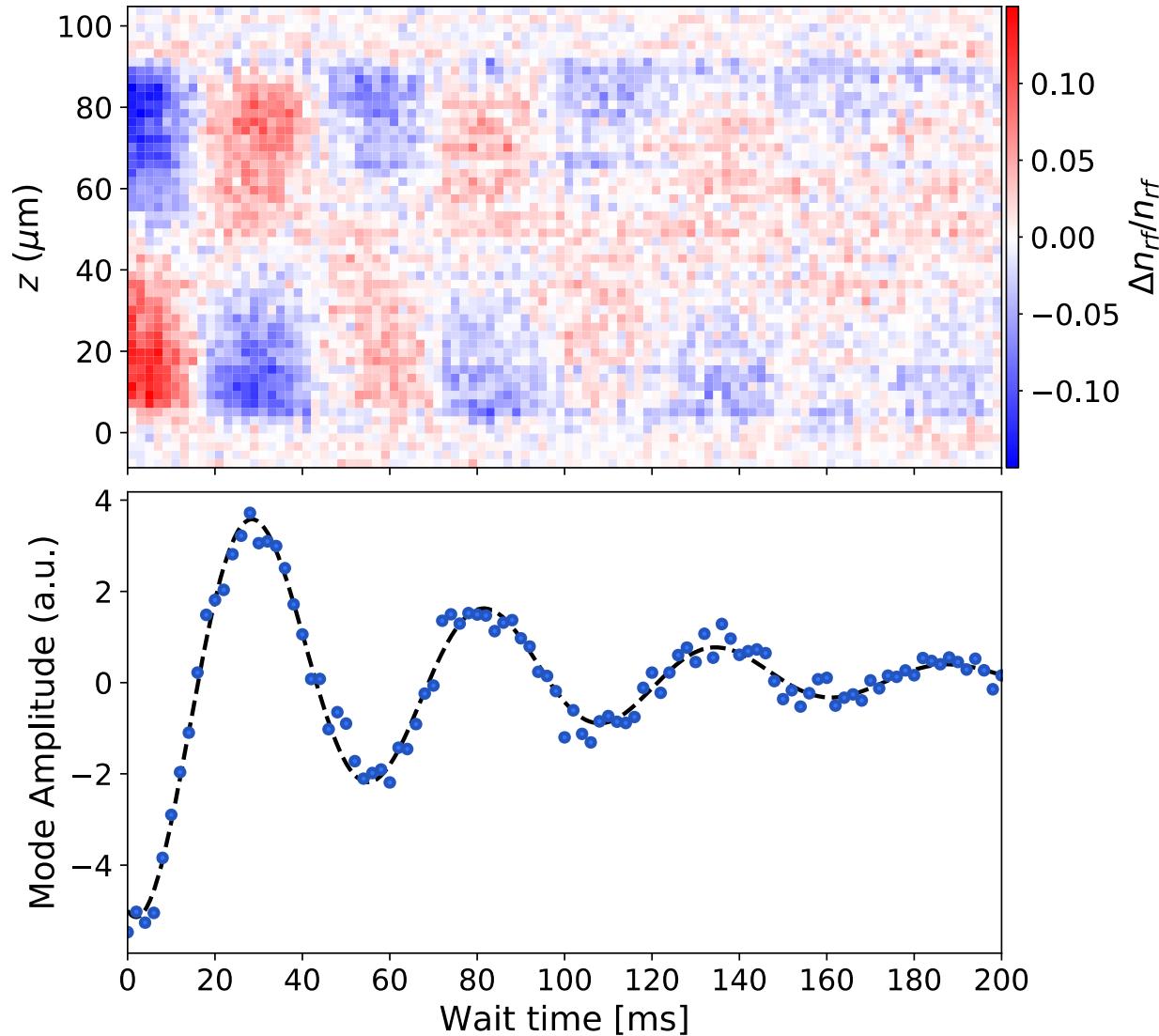
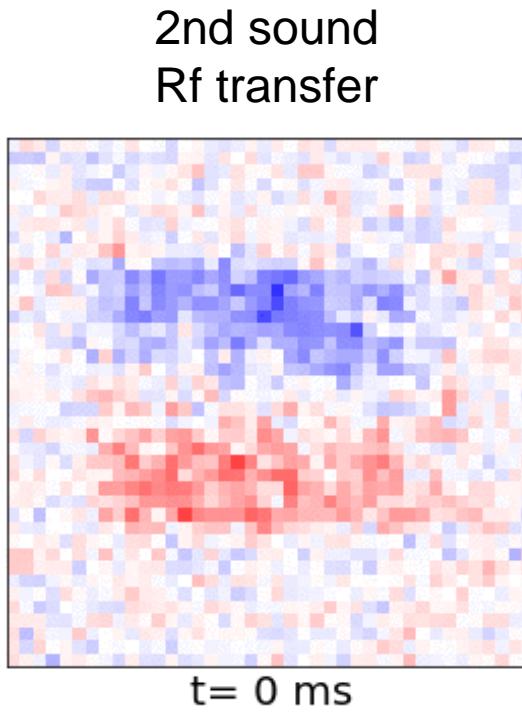
A blue arrow points from the equation $C_2^2 = \frac{\rho_S}{\rho_N} \frac{TS^2}{C_p}$ to the right, indicating the relationship between C_2^2 and the superfluid fraction ρ_S/ρ .



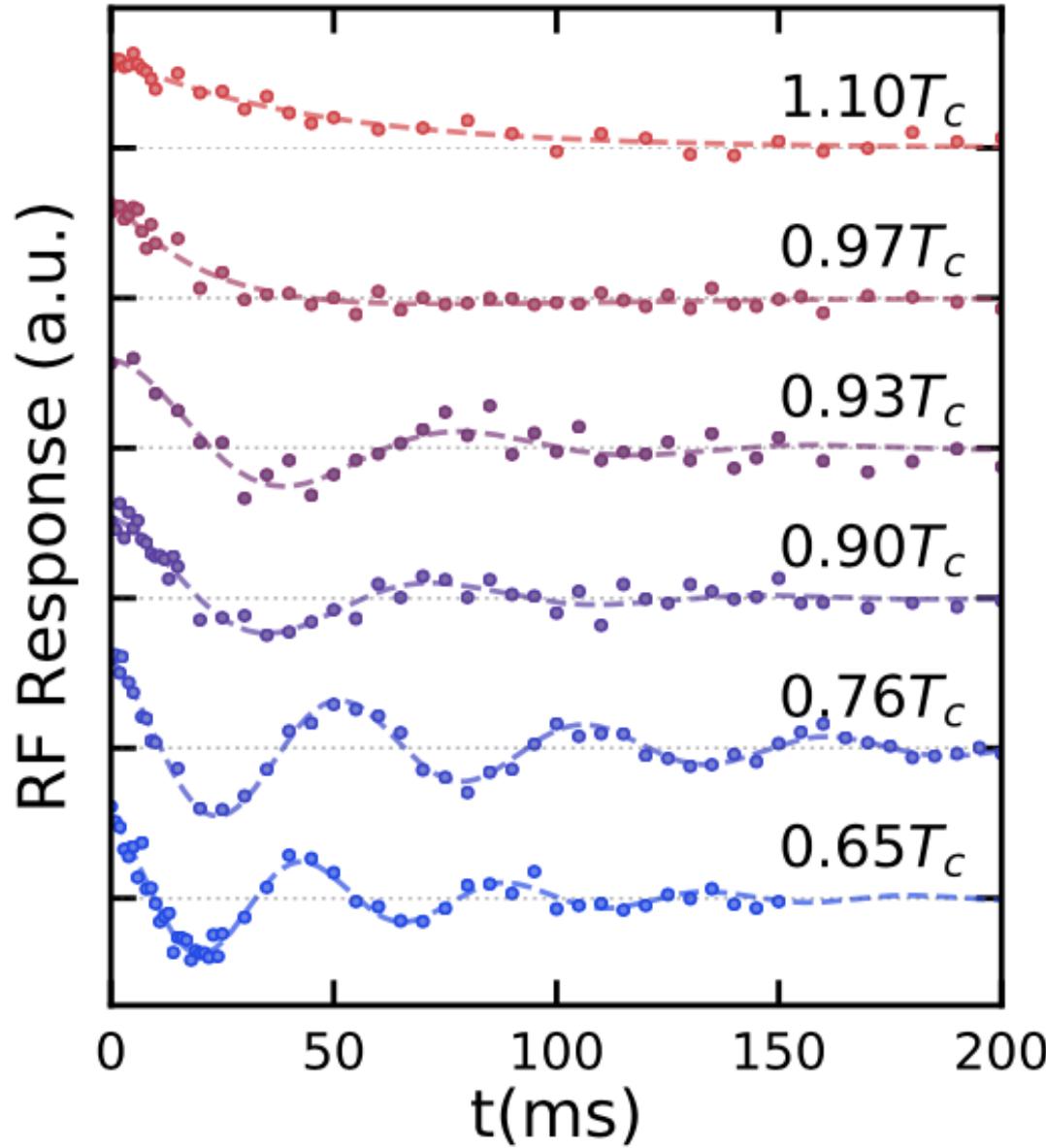
- Steady state response
- Free evolution, resonant shaking
- ◆ Free evolution, local heater

Damping of Second Sound

Temperature Probe

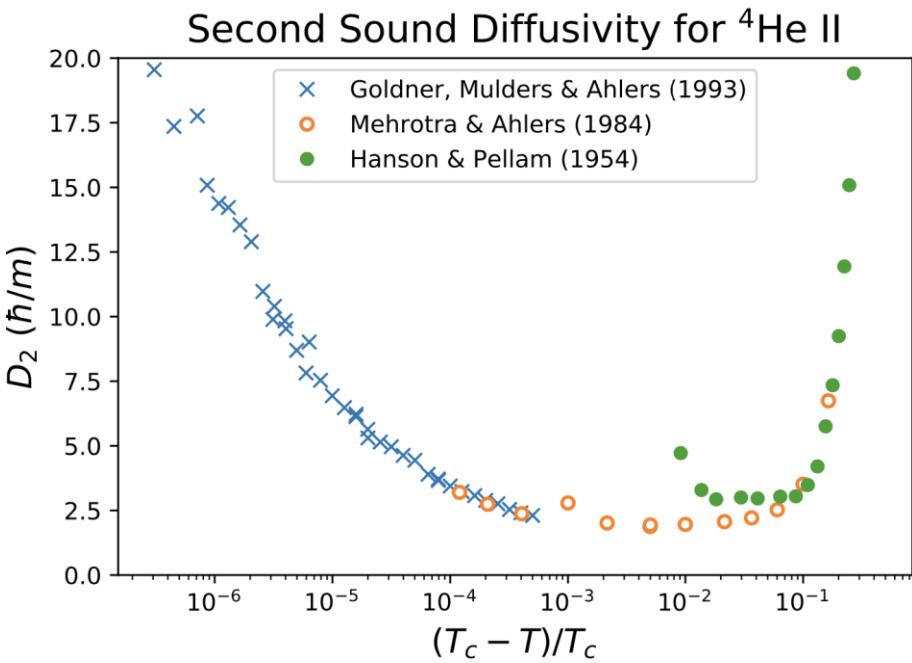
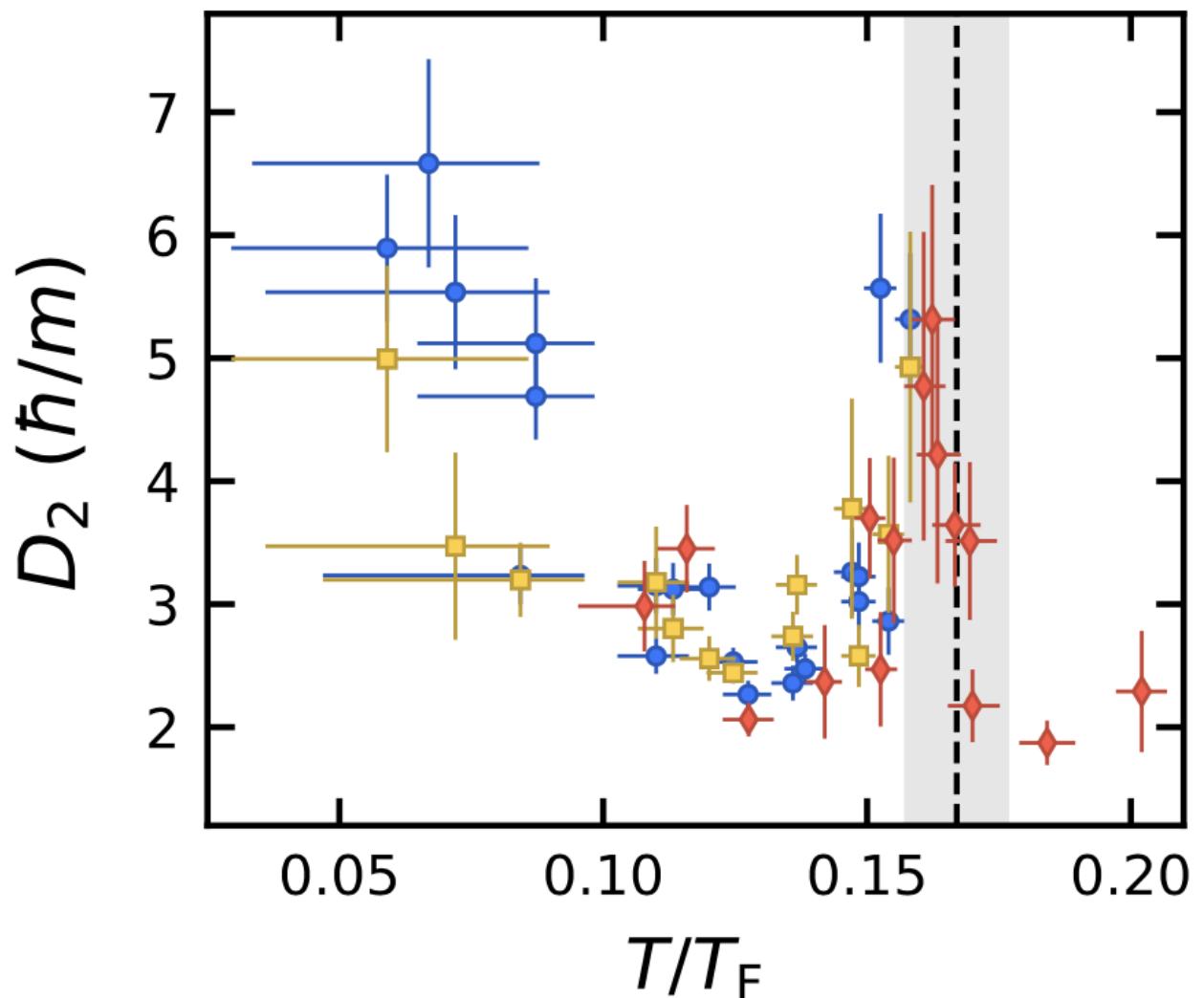


Transition from normal to superfluid



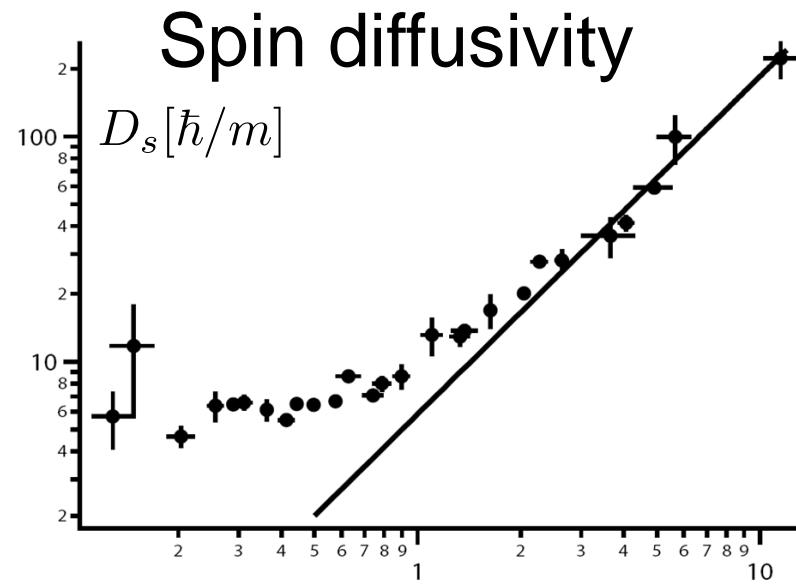
Diffusivity of Second Sound

$$D_2 = \Gamma_2/k^2$$

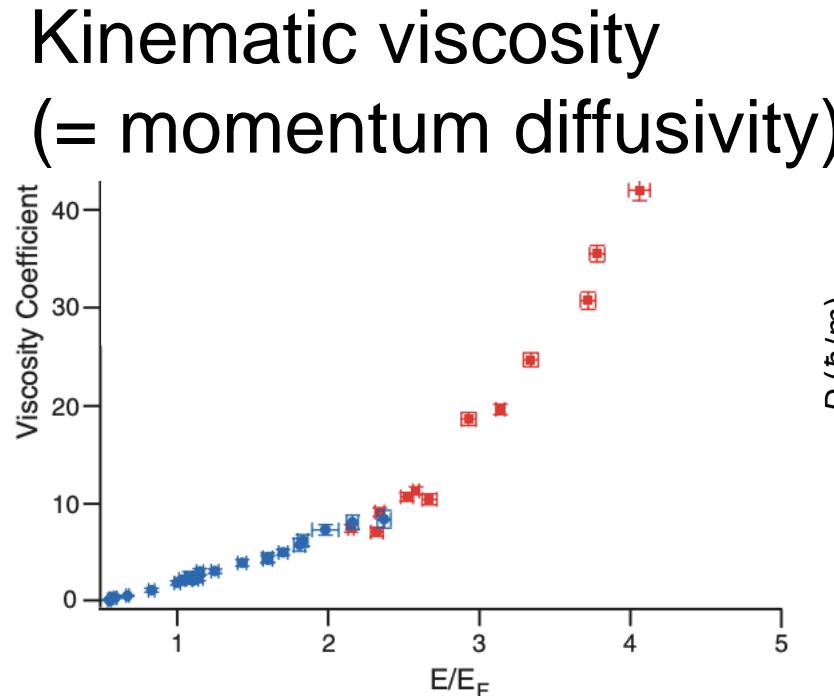


$$D_2 = \frac{\kappa}{c_p} + \frac{\rho_s}{\rho_n} \left(\xi_3 \rho + \frac{4}{3} \frac{\eta}{\rho} \right)$$

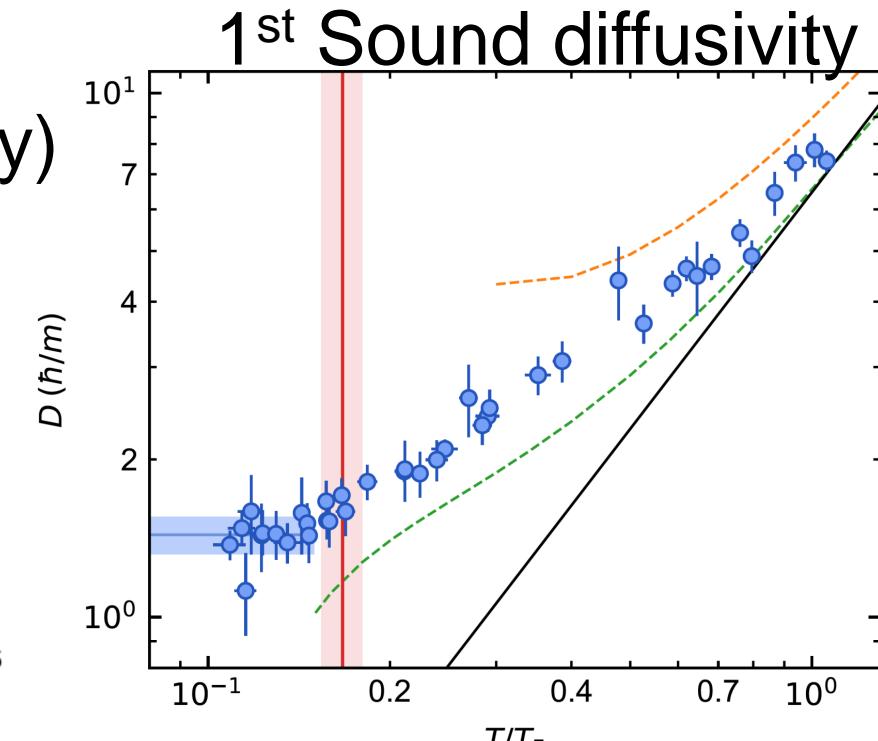
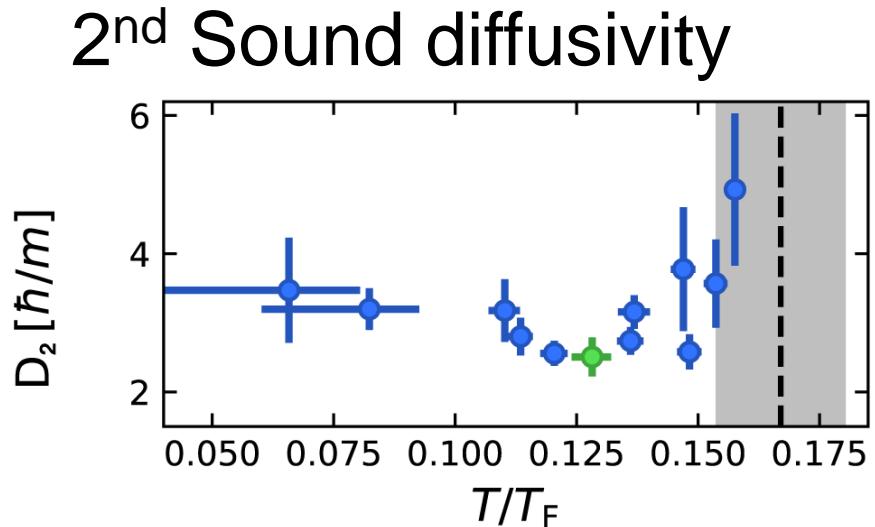
Diffusivities in the Unitary Gas



A.T. Sommer, M.J.H. Ku, G. Roati,
MZ, Nature 472, 201 (2011)



Cao et al., John Thomas group,
Science 331, 58 (2011)



P. Patel, Z. Yan, B. Mukherjee,
R. Fletcher, J. Struck, MZ,
Science 370, 1222 (2020)

Diffusivity (charge, spin, momentum, thermal)

$$D \propto v l \propto \frac{\hbar}{m}$$

Transport in Strongly Correlated Quantum Gases

Unifying themes of strongly interacting Fermi systems:

Loss of quasi-particle description: $\tau^{-1} \approx E_F / \hbar$

Quantum-Limited diffusivities:

$$D \approx \frac{\hbar}{m}$$

Quantum gas simulators are poised to further elucidate
Interplay between charge, spin, heat and Hall transport

- Universal Sound Diffusion in a Strongly Interacting Fermi Gas
Parth B. Patel, Zhenjie Yan, Biswaroop Mukherjee, Richard J. Fletcher, Julian Struck, MZ Science, 370, 1222 (2020)
- Spin Transport in a Mott Insulator of Ultracold Fermions
Matthew Nichols, Lawrence Cheuk, Melih Okan, Thomas Hartke, Enrique Mendez, T. Senthil, Ehsan Khatami, Hao Zhang, MZ, Science 363, 383 (2019)
- Doublon-Hole Correlations and Fluctuation Thermometry in a Fermi-Hubbard Gas
T. Hartke, B. Oreg, N. Jia, MZ, PRL 125, 113601 (2020)
- Bose polarons near quantum criticality
Zoe Z. Yan, Yiqi Ni, Carsten Robens, MZ, Science, 368, 190-194 (2020)
- Geometric squeezing into the lowest Landau level
Richard J. Fletcher, Airlia Shaffer, Cedric C. Wilson, Parth B. Patel, Zhenjie Yan, Valentin Crépel, Biswaroop Mukherjee, MZ, arXiv:1911.12347 (Science 2021, to be published)

Bosons and Fermions



Bosons and Fermions

BEC 1/Fermi3

Fermionic Superfluids

Biswaroop Mukherjee

Parth Patel

Zhenjie Yan

Dr. Richard Fletcher

Visiting Professors:

Zoran Hadzibabic

Chris Vale

Former Members:

Ariel Sommer

Mark Ku

Dr. Tarik Yefsah

(→ CNRS, ENS Paris)

Dr. Julian Struck

(→ ENS Paris)

Fermi 1

NaK Dipolar Molecules

Yiqi Ni

Alex Yu Chuang

Eric Wolf

Dr. Carsten Robens

Former Members:

Zoe Yan

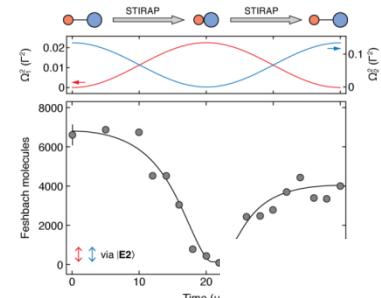
Jee Woo Park (→ SNU Korea)

Dr. Sebastian Will

(→ Prof., Columbia U.)

Cheng-Hsun Wu (→ Google X)

Dr. Huanqian Loh (→ Singap.)



Fermi 2

Fermi Gas Microscope

Thomas Hartke

Botond Oreg

Dr. Ningyuan Jia

Former Members:

Lawrence Cheuk

(→ Prof. @Princeton)

Melih Okan

Matthew Nichols

Dr. Hao Zhang

Katherine Lawrence

Vinay Ramasesh

Thomas Gersdorf

Dr. Thomas Lompe

(→ Hamburg)

Dr. Waseem Bakr

(→ Prof. @Princeton)

Fermi3

Rotating Superfluids

Airlia Shaffer-Moag

Cedric Wilson

Biswaroop Mukherjee

Parth Patel

Zhenjie Yan

Dr. Richard Fletcher

Valentin Crepel

