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Course 8.333 Stat Mech 1

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Test Final Exam

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9:00 am - 12:00 pm

(1)

# 1. FREE FERMIONS

(a) Joint probability, fermions  $\Rightarrow \gamma = -1$

$$P(\{n_\ell\}) = \frac{1}{Q_{-1}} \prod_{\ell} \exp[-\beta(\varepsilon(\ell) - \mu) n_\ell]$$

$$= \frac{\prod_{\ell} \exp[-\beta(\varepsilon(\ell) - \mu) n_\ell]}{\prod_{\ell} \exp[-\beta(\varepsilon(\ell) - \mu) n_\ell]}$$

$$= \prod_{\ell} \exp[-\beta(\varepsilon(\ell) - \mu) n_\ell] [1 + \exp(\beta\mu - \beta\varepsilon(\ell))]^{-1}$$

(b) Average occupation number

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$$\begin{aligned} \langle n_i \rangle_{-1} &= \frac{-\partial \ln \Omega_{-1}}{\partial (\beta \epsilon(i))} = \frac{1}{z^{-1} e^{\beta \epsilon(i)} + 1} \\ &= \frac{1}{e^{\beta \epsilon(i)} e^{-\beta \mu} + 1} = \frac{1}{e^{\beta (\epsilon(i) - \mu)} + 1} \end{aligned}$$

~~Answer is part (a); for a fixed  $\epsilon$~~

~~$$\begin{aligned} P(\{n_i\}) &= \frac{e^{-\beta (\epsilon(i) - \mu) n_i}}{1 + e^{-\beta (\epsilon(i) - \mu)}} \cdot \frac{e^{\beta (\epsilon(i) - \mu) n_i}}{e^{\beta (\epsilon(i) - \mu) n_i}} \\ &= \dots \end{aligned}$$~~

From here, can solve ...

$$\boxed{e^{\beta (\epsilon(i) - \mu)} = \frac{1}{\langle n_i \rangle_{-1}} - 1} \Rightarrow e^{-\beta (\epsilon(i) - \mu)} = \frac{\langle n_i \rangle_{-}}{1 - \langle n_i \rangle_{-}}$$

So

$$P(\{n_i\}_{-1}) = \prod_i \frac{e^{-\beta (\epsilon(i) - \mu) n_i}}{1 + e^{\beta (\mu - \epsilon(i))}} = \prod_i \frac{\left( \frac{\langle n_i \rangle_{-}}{1 - \langle n_i \rangle_{-}} \right)^{n_i}}{1 + \frac{\langle n_i \rangle_{-}}{1 - \langle n_i \rangle_{-}}}$$

(next page

$$\dots = \prod_l \frac{\left[ \langle n_l \rangle_- / (1 - \langle n_l \rangle_-) \right]^{n_l}}{1 + \frac{\langle n_l \rangle_-}{1 - \langle n_l \rangle_-}}$$

$$= \prod_l \frac{\left( \langle n_l \rangle_- / (1 - \langle n_l \rangle_-) \right)^{n_l}}{\frac{1}{1 - \langle n_l \rangle_-}}$$

$$= \prod_l (1 - \langle n_l \rangle_-) \cdot \left[ \frac{\langle n_l \rangle_-}{1 - \langle n_l \rangle_-} \right]^{n_l}$$

or

$$= \prod_l \langle n_l \rangle_-^{n_l} (1 - \langle n_l \rangle_-)^{+1 - n_l}$$

(c) Entropy for dist? Max entropy?

Entropy is

$$S = -k_B \sum_{n=1}^M p_n \ln p_n$$

Max entropy is obtained when  $p_n = 1/M$

$$\rightarrow S_{\max} = k_B \ln M$$

(d) Entropy for  $P(\{n_e\})$  , zero temp limit

~~$S = -k_B \sum_e \{ \langle n_e \rangle \ln \langle n_e \rangle + (1 - \langle n_e \rangle) \ln (1 - \langle n_e \rangle) \}$~~

Occupation numbers of different one-particle states are independent, so the corresponding entropies are additive ...

$$S = -k_B \sum_e \left\{ \langle n_e \rangle \ln \langle n_e \rangle + (1 - \langle n_e \rangle) \overset{\ln}{\downarrow} \ln (1 - \langle n_e \rangle) \right\}$$

This is just  $-k_B \sum_e p_e \ln p_e + (1 - p_e) \ln (1 - p_e)$

since  $n_e = 0$  or  $1$  ...

• in the zero-temperature limit,

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all occupation numbers are either 0 (excited)  
or 1 (ground)

So if  $n_l = 1$ , then

$$S = -k_B \sum_l 0 = 0$$

if  $n_l = 0$  then

$$S = -k_B \sum_l 0 = 0$$

So in either case contribution to entropy is  
zero,

$\Rightarrow$  system at  $T=0$  has  $\boxed{S=0}$ .

(e) Variance in total # of particles  
at  $T=0$  behavior?

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( ) For a fixed  $l$  have to calculate  $\langle e^{ikn_l} \rangle$

$$\langle e^{ikn_l} \rangle = \frac{e^{-\beta(\epsilon(l)-\mu)}}{1 + e^{-\beta(\epsilon(l)-\mu)}} \cdot e^{ik} + \frac{1}{1 + e^{-\beta(\epsilon(l)-\mu)}}$$

$$n_l = 0$$

$$= \frac{1 + e^{-\beta(\epsilon(l)-\mu) + ik}}{1 + e^{-\beta(\epsilon(l)-\mu)}}$$

→ now get cumulants...

$$\ln \langle e^{ikn_l} \rangle = (ik) \frac{1}{1 + e^{\beta(\epsilon(l)-\mu)}} - \frac{k^2}{2} \left[ \frac{e^{\beta(\epsilon(l)-\mu)}}{(1 + e^{\beta(\epsilon(l)-\mu)})^2} \right]$$

+ ...

$$\underline{\text{So}} \quad \langle n_l^2 \rangle_c = \frac{e^{\beta(\epsilon(l)-\mu)}}{(1 + e^{\beta(\epsilon(l)-\mu)})^2}$$

$$\rightarrow \langle N^2 \rangle_c = \sum_l \langle n_l^2 \rangle_c = \sum_l \frac{e^{\beta(\epsilon(l)-\mu)}}{(1 + e^{\beta(\epsilon(l)-\mu)})^2}$$

Zero temp  $\rightarrow \beta \rightarrow \infty$



$$\text{and } \varepsilon(\phi) < \mu = \varepsilon_F$$

So,

$$\langle N^2 \rangle_c \sim \sum_x \frac{e^{-\beta X}}{(1 + e^{-\beta X})^2} \rightarrow \boxed{0}$$

So all fluctuations in total particle number vanishes at zero temperature.



## 2. NON DEGENERATE QUANTUM GAS

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$$n = N/V > 2, \quad \mathcal{H} = \sum_i \mathbf{p}_i^2 / 2m, \quad \text{classical 3D}$$

$$n \lambda^3 / g \ll 1 \Rightarrow \text{non degenerate}$$

$$\lambda = \frac{h}{\sqrt{2\pi m k_B T}}$$

(a) Density

$$z = e^{\beta \mu}$$

$$n_\gamma \equiv \frac{N_\gamma}{V} = \frac{g}{\lambda^3} f_{3/2}^\gamma(z)$$

Pressure

$$P_\gamma = \frac{1}{\beta} \frac{g}{\lambda^3} f_{5/2}^\gamma(z)$$

(6) Find  $P(n, T)$

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To do this, follow procedure from textbook...

• First, we use the identity

$$f_n^\eta(z) \approx \sum_{\alpha=1}^{\infty} \eta^{\alpha+1} \frac{z^\alpha}{\alpha^n} \quad @ \text{ nondegenerate regime}$$

$$\Rightarrow \begin{cases} \frac{n\eta\lambda^3}{g} = f_{3/2}^\eta(z) = z + \eta \frac{z^2}{2^{3/2}} + \frac{z^3}{2^{3/2}} + \eta \frac{z^4}{4^{3/2}} \\ \beta P \eta \lambda^3 / g = f_{5/2}^\eta(z) \approx z + \eta \frac{z^2}{2^{5/2}} + \frac{z^3}{3^{5/2}} + \eta \frac{z^4}{4^{5/2}} \end{cases}$$

→ get  $z$  in terms of  $n\lambda^3 \dots$

$$z = \left( \frac{n\eta\lambda^3}{g} \right) - \frac{\eta}{2^{3/2}} \left( \frac{n\eta\lambda^3}{g} \right)^2 + \left( \frac{1}{4} - \frac{1}{3^{3/2}} \right) \left( \frac{n\eta\lambda^3}{g} \right)^3$$

Put this back into  $P_\eta$  to find

$$\begin{aligned} \frac{\beta P \eta \lambda^3}{g} &= \left( \frac{n\eta\lambda^3}{g} \right) - \frac{\eta}{2^{3/2}} \left( \frac{n\eta\lambda^3}{g} \right)^2 + \left( \frac{1}{4} - \frac{1}{3^{3/2}} \right) \left( \frac{n\eta\lambda^3}{g} \right)^3 \\ &+ \frac{\eta}{2^{5/2}} \left( \frac{n\eta\lambda^3}{g} \right)^2 - \frac{1}{8} \left( \frac{n\eta\lambda^3}{g} \right)^3 + \frac{1}{3^{5/2}} \left( \frac{n\eta\lambda^3}{g} \right)^3 + \dots \end{aligned}$$

Ultimately we'll have

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$$P_{\eta} = n_{\eta} k_B T \left[ 1 - \underbrace{\frac{\eta}{2^{5/2}} \left( \frac{n_{\eta} \lambda^3}{g} \right)}_1 + \left( \frac{1}{6} - \frac{2}{3^{5/2}} \right) \left( \frac{n_{\eta} \lambda^3}{g} \right)^2 + \dots \right]$$

2nd virial  
coeff.

(c) Expression for  $\epsilon_{\eta} = E/v$  up to  $n^2$

Well... we use the nice fact that

$$\epsilon_{\eta} \equiv \frac{E_{\eta}}{v} = \frac{3}{2} P_{\eta} \quad \text{to get}$$

$$\epsilon_{\eta} = \frac{3}{2} n_{\eta} k_B T \left[ \frac{3}{2} - \frac{3\eta}{2^{7/2}} \left( \frac{n_{\eta} \lambda^3}{g} \right) + \left( \frac{3}{16} - \frac{3}{3^{5/2}} \right) \left( \frac{n_{\eta} \lambda^3}{g} \right)^2 + \dots \right]$$

(d) Heat capacity  $C_{V/N}$

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From (c) we have

$$E = V \varepsilon_{\eta} = \left( \frac{N}{n_{\eta}} \right) \varepsilon_{\eta} = N k_B T \left[ \frac{3}{2} - \frac{3\eta}{2^{7/2}} \left( \frac{n \lambda^3}{g} \right) + \dots \right]$$

$$\frac{d}{dT} \quad \frac{C_V}{N} = \frac{1}{N} \frac{dE}{dT} = \text{?}$$

$$\lambda = \frac{h}{\sqrt{2\pi m k_B T}} = \left( \frac{h}{\sqrt{2\pi m k_B}} \right) T^{-1/2} = \Omega T^{-1/2}$$

$$\Rightarrow E = N k_B T \left( \frac{3}{2} - \frac{3\eta}{2^{7/2}} \frac{n}{g} \Omega^3 T^{-3/2} + \dots \right)$$

$$\frac{1}{N} \frac{dE}{dT} = \frac{d}{dT} \left[ k_B T \left( \frac{3}{2} - \frac{3\eta}{2^{7/2}} \frac{n}{g} \Omega^3 T^{-3/2} + \dots \right) \right]$$

$$= \frac{d}{dT} \left( \frac{3}{2} k_B T \right) + \frac{k_B}{2} \frac{3\eta}{2^{7/2}} \frac{n}{g} \Omega^3 T^{-3/2} + \dots$$

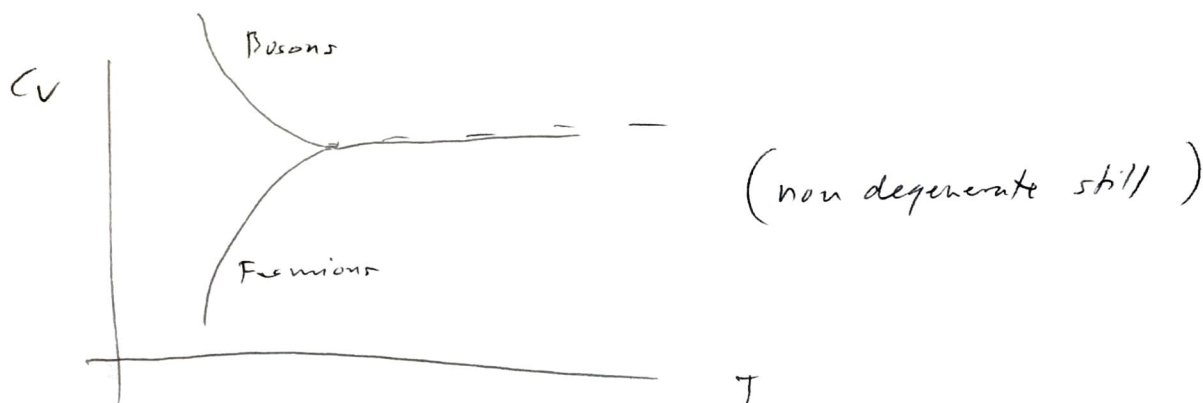
$$C_{V/N} = \frac{3}{2} k_B + \frac{k_B}{2} \frac{3\eta}{2^{7/2}} \frac{n}{g} \underbrace{\left( \Omega^3 T^{-3/2} \right)}_{\lambda^3} \quad \text{First QM correction}$$

Since the first order correction  $\propto \eta$

we have that

$\left\{ \begin{array}{l} C_V \text{ increases for bosons} \\ C_V \text{ decreases for fermions} \end{array} \right.$

$$\frac{C_V}{N} \sim \frac{3}{2} k_B \left( 1 + \frac{\eta}{2^{3/2}} \left( \frac{n \lambda^3}{g} \right) + \dots \right)$$



(e) Calculate

$$\alpha_P = \frac{1}{V} \frac{\partial V}{\partial T} \bigg|_{P, N}$$

To do this ... use  $V = N/n$  to get

$$\begin{aligned} \alpha_P &= \left( \frac{n}{N} \right) \frac{\partial (N/n)}{\partial T} = n \frac{\partial (1/n)}{\partial T} = n \frac{\partial (1/n)}{\partial n} \frac{\partial n}{\partial T} \\ &= n \cdot (-1) n^{-2} \frac{\partial n}{\partial T} = -\frac{1}{n} \frac{\partial n}{\partial T} \end{aligned}$$

And now, we have

$$n = \frac{g}{\lambda^3} f_{3/2}^\eta(z) \quad \text{so, } \dots \text{ next page}$$

Need to find an expression for

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$$n(P, T) \dots$$

To do this we'll need to find

$z$  in terms of  $P$ .

$$z = \frac{\beta P \lambda^3}{g} - \eta \frac{z^2}{z^{5/2}} = \frac{z^3}{z^{5/2}}$$

$$= \left( \frac{\beta P \lambda^3}{g} \right) - \frac{\eta}{z^{5/2}} \left( \frac{\beta P \lambda^3}{g} \right)^2 - \dots$$

Subbing this into

$$\frac{n \lambda^3}{g} = z + \eta \frac{z^2}{z^{3/2}} + \dots$$

$$= \left( \frac{\beta P \lambda^3}{g} \right) + \frac{\eta}{4\sqrt{2}} \left( \frac{\beta P \lambda^3}{g} \right)^2$$

$\downarrow$   
 $2^{5/2}$

~~$\left( \frac{\beta P \lambda^3}{g} \right)^3$~~   
 $3^{5/2}$

$$+ \left( \frac{1}{2^{5/2}} - \frac{1}{g} \right) \left( \frac{\beta P \lambda^3}{g} \right)^3 + \dots$$

$$\frac{50}{\quad} \quad n = \frac{P}{k_B T} \left[ 1 + \frac{\eta}{2^{5/2}} \left( \frac{\beta P \lambda^2}{g} \right)^2 + \left( \frac{1}{3^{5/2}} - \frac{1}{8} \right) \left( \frac{\beta P \lambda^2}{g} \right)^2 + \dots \right] \quad (14)$$

$$= \frac{P}{k_B T} \left[ 1 + \frac{\eta}{2^{5/2}} \left( \frac{\beta P \lambda^2}{g} \right) + \dots \right]$$

↑  
to first order correction only

Now

$$\alpha_P = \frac{1}{n} \frac{\partial n}{\partial T} \equiv \quad (\text{mathematica...})$$

~~$\alpha_P = \frac{1}{n} \frac{\partial n}{\partial T} \equiv \frac{20 g k_B T^{3/2}}{8 g k_B T^{5/2} + \sqrt{2} \Omega^3 P \eta}$~~       let  $\lambda = \frac{h}{\sqrt{2\pi m k_B T}} = \Omega T^{-1/2}$

then

$$\alpha_P = \frac{1}{n} \frac{\partial n}{\partial T}$$

↗ I don't wanna simplify this...

$$\alpha_P \approx \frac{7}{2T} - \frac{20 g k_B T^{3/2}}{8 g k_B T^{5/2} + \sqrt{2} \Omega^3 P \eta}$$

(f) what is  $c_p/N$ ?

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assuming that  $c_p/c_v = \gamma = 5/3$

Then

$$\frac{c_p}{N} = \frac{5}{3} \frac{c_v}{N} = \frac{5}{2} k_B + \frac{5}{2} k_B \frac{\eta}{2^{7/2}} \frac{\eta}{g} \underbrace{\Omega^3 T^{-3/2}}_{\lambda^3}$$

$$= \frac{5}{2} k_B + \frac{5}{2} k_B \frac{\eta}{2^{7/2}} \left( \frac{\eta \lambda^3}{g} \right) + \dots$$

$$= \frac{5}{2} k_B \left[ 1 + \frac{\eta}{2^{7/2}} \left( \frac{\eta \lambda^3}{g} \right) + \dots \right]$$



first QM correction.



# 3. TRAP SITES IN A BOSE GAS

(a) Use standard Bose statistics

$$\langle n_+ \rangle = \frac{1}{e^{\beta(\epsilon_+ - \mu)} - 1} = \frac{1}{e^{\beta(\frac{\hbar^2 k_+^2}{2m} - \mu)} - 1}$$

$$\langle n_+ \rangle = \frac{1}{e^{-\beta(\mu + \epsilon_w)} - 1}$$

trapped  $\rightarrow$  not free  
so energy is  $\epsilon = -w$

(b) For  $\langle n_g \rangle$  use standard Bose gas result ...

$$\langle n_g \rangle = \frac{g}{\lambda^3} f_{3/2}^+(z)$$

$$\langle n_g \rangle = \frac{g (2\pi m k_B T)^{3/2}}{h^3} f_{3/2}^+(e^{\beta\mu})$$

? Where does the density  $\rho$  come in ???

(c) Find  $n_c(T)$  for the onset of BEC

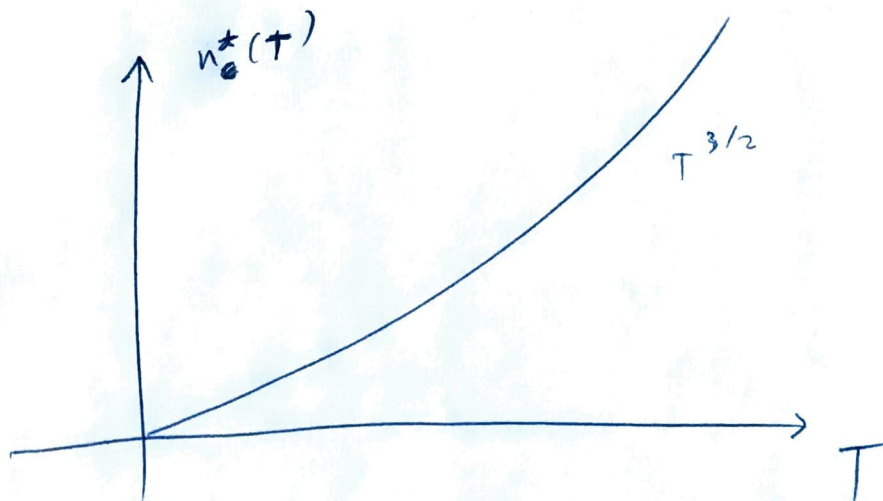
→ BEC occurs when  $\mu = -\omega$  state accommodates a macroscopic fraction of particles...

→ ~~the BEC occurs when  $\mu = -\omega$  state accommodates a macroscopic fraction of particles...~~

$$n_c^*(T) = \frac{g}{\lambda^3} f_{3/2}^+ \left( e^{-\beta \underbrace{\omega}_{\text{"}\mu\text{"}}} \right) \sim \omega$$

For small  $\omega$  ... same as  $\mu \rightarrow 0$

$$\underbrace{f_{3/2}^+}_{\sim 1} \left( \underbrace{e^{-\beta \omega}}_{\sim 1} \right) \sim \zeta_{3/2} \quad \text{so} \quad n_c^*(T) \sim T^{3/2}$$



(d) Total energy density in the BEC

The total energy density in the BEC

comes ~~entirely~~ from the continuum states  
+ particles in trapped states

① use usual results for energy density  
of an ideal Bose gas; but setting  $\mu$  to

(continuum) its limiting value of  $-\mu$   
(First contribution)

$$\epsilon_+ = \frac{3}{2} P_+ = \frac{3}{2\beta} \frac{g}{\lambda^3} f_{5/2}^+ (z).$$

$$\epsilon_+ = \frac{3}{2\beta} \frac{g}{\lambda^3} f_{5/2}^+ (e^{-\beta\mu})$$

multiply this  
by density of

trap  $> \rho$   
basically how  
many  
trap sites  
there are

② Second contribution from trapped states

$$E_{\text{trap}} = \cancel{\text{trap}} - \rho \omega \langle n_+ \rangle = -\rho \omega (n - \langle n_g \rangle)$$

$$= -\rho \omega \left( n - \frac{g}{\lambda^3} f_{3/2}^+ (e^{-\beta\mu}) \right)$$

↑  
density not  
in trap continuum

So, total energy density is

$$\epsilon_{\text{tot}} = -\rho \omega n + \frac{g}{\lambda^3} \left[ \frac{3}{2\beta} f_{5/2}^+ (e^{-\beta\mu}) + \rho \omega f_{3/2}^+ (e^{-\beta\mu}) \right]$$