Due: 10/25/2021

## Kinetic Theory

- 1. Poisson Brackets:
- (a) Show that for observable  $\mathcal{O}(\mathbf{p}(\mu), \mathbf{q}(\mu))$ ,  $d\mathcal{O}/dt = \{\mathcal{O}, \mathcal{H}\}$ , along the time trajectory of any micro state  $\mu$ , where  $\mathcal{H}$  is the Hamiltonian.
- (b) If the ensemble average  $\langle \{\mathcal{O}, \mathcal{H}\} \rangle = 0$  for any observable  $\mathcal{O}(\mathbf{p}, \mathbf{q})$  in phase space, show that the ensemble density satisfies  $\{\mathcal{H}, \rho\} = 0$ .

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- **2.** Equilibrium density: Consider a gas of N particles of mass m, in an external potential  $U(\vec{q})$ . Assume that the one body density  $\rho_1$  ( $\vec{p}$ ,  $\vec{q}$ , t), satisfies the Boltzmann equation. For a stationary solution,  $\partial \rho_1/\partial t = 0$ , it is sufficient from Liouville's theorem for  $\rho_1$  to satisfy  $\rho_1 \propto \exp\left[-\beta\left(p^2/2m + U(\vec{q})\right)\right]$ . Prove that this condition is also necessary by using the H-theorem as follows.
- (a) Find  $\rho_1(\vec{p}, \vec{q})$  that minimizes  $H = N \int d^3\vec{p} d^3\vec{q} \rho_1(\vec{p}, \vec{q}) \ln \rho_1(\vec{p}, \vec{q})$ , subject to the constraint that the total energy  $E = \langle \mathcal{H} \rangle$  is constant. (Hint: Use the method of Lagrange multipliers to impose the constraint.)
- (b) For a mixture of two gases (particles of masses  $m_a$  and  $m_b$ ) find the distributions  $\rho_1^{(a)}$  and  $\rho_1^{(b)}$  that minimize  $H = H^{(a)} + H^{(b)}$  subject to the constraint of constant total energy. Hence show that the kinetic energy per particle can serve as an empirical temperature.

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**3.** (Optional) Evolving a canonical harmonic oscillator density: A dilute gas of non-interacting particles is in equilibrium in a harmonic potential, such that the density for each particle has the form

$$\rho_0(\vec{q}, \vec{p}) = \exp\left[-\beta \left(\frac{Kq^2}{2} + \frac{p^2}{2m}\right)\right] \left(\frac{\beta}{2\pi}\right)^3 \left(\frac{K}{m}\right)^{3/2}.$$

At time t = 0, and external force  $\vec{F}(t)$  is applied, changing the one particle Hamiltonian to  $H_0 - \vec{q} \cdot \vec{F}(t)$ .

(a) Write down the (Liouville) equation governing subsequent evolution of the one particle density.

- (b) Confirm that the density at later times satisfies,  $\rho(\vec{q}, \vec{p}, t) = \rho_0 (\vec{q} \langle \vec{q} \rangle_t, \vec{p} \langle \vec{p} \rangle_t)$ , and find the equations of motion for  $\langle \vec{q} \rangle_t$  and  $\langle \vec{p} \rangle_t$ .
- (c) Compute the entropy S(t) associated with the probability density  $\rho$ .
- (d) Would a similar time dependent shift of the density work in the case of the canonical weight associated with a general potential  $\mathcal{V}(\vec{q})$  (e.g.  $\mathcal{V}(\vec{q}) \propto q^4$ ) driven by an external force?

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**4.** Zeroth-order hydrodynamics: The hydrodynamic equations resulting from the conservation of particle number, momentum, and energy in collisions are (in a uniform box):

$$\begin{cases} \partial_t n + \partial_\alpha (nu_\alpha) = 0 \\ \partial_t u_\alpha + u_\beta \partial_\beta u_\alpha = -\frac{1}{mn} \partial_\beta P_{\alpha\beta} \\ \partial_t \varepsilon + u_\alpha \partial_\alpha \varepsilon = -\frac{1}{n} \partial_\alpha h_\alpha - \frac{1}{n} P_{\alpha\beta} u_{\alpha\beta} \end{cases},$$

where n is the local density,  $\vec{u} = \langle \vec{p}/m \rangle$ ,  $u_{\alpha\beta} = (\partial_{\alpha}u_{\beta} + \partial_{\beta}u_{\alpha})/2$ , and  $\varepsilon = \langle mc^2/2 \rangle$ , with  $\vec{c} = \vec{p}/m - \vec{u}$ .

(a) For the zeroth order density

$$f_1^0(\vec{p}, \vec{q}, t) = \frac{n(\vec{q}, t)}{(2\pi m k_B T(\vec{q}, t))^{3/2}} \exp \left[ -\frac{(\vec{p} - m \vec{u}(\vec{q}, t))^2}{2m k_B T(\vec{q}, t)} \right],$$

calculate the pressure tensor  $P_{\alpha\beta}^0 = mn \langle c_{\alpha}c_{\beta}\rangle^0$ , and the heat flux  $h_{\alpha}^0 = nm \langle c_{\alpha}c^2/2\rangle^0$ .

- (b) Obtain the zeroth order hydrodynamic equations governing the evolution of  $n(\vec{q}, t)$ ,  $\vec{u}(\vec{q}, t)$ , and  $T(\vec{q}, t)$ .
- (c) Show that the above equations imply  $D_t \ln (nT^{-3/2}) = 0$ , where  $D_t = \partial_t + u_\beta \partial_\beta$  is the material derivative along streamlines.
- (d) Write down the expression for the function  $H^0(t) = \int d^3\vec{q}d^3\vec{p}f_1^0(\vec{p},\vec{q},t) \ln f_1^0(\vec{p},\vec{q},t)$ , after performing the integrations over  $\vec{p}$ , in terms of  $n(\vec{q},t)$ ,  $\vec{u}(\vec{q},t)$ , and  $T(\vec{q},t)$ .
- (e) Using the hydrodynamic equations in (b) calculate  $dH^0/dt$ .
- (f) Discuss the implications of the result in (e) for approach to equilibrium.

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**5.** Diffusion: Consider a mixture of two gases (a) and (b), in a box of volume V.

(a) Write down the Boltzmann equations for the one particle densities  $f_a$ , and,  $f_b$ , in terms of the Liouville operators  $\mathcal{L}_{\alpha} \equiv [\partial_t + (\vec{p}_{\alpha}/m_{\alpha}) \cdot \nabla]$ , and collision operators

$$C_{\alpha,\beta} = -\int d^3\vec{p}_2 d^2\vec{b}_{\alpha\beta} |\vec{v}_1 - \vec{v}_2| [f_{\alpha}(\vec{p}_1, \vec{q}_1) f_{\beta}(\vec{p}_2, \vec{q}_1) - f_{\alpha}(\vec{p}_1', \vec{q}_1) f_{\beta}(\vec{p}_2', \vec{q}_1)],$$

where  $\alpha = a, b$  and  $\beta = a, b$ .

- (b) Assuming that the collision terms are much more dominant than the Liouville streams (dilute limit), write down a zeroth order solution to the Boltzmann equations.
- (c) Write down the hydrodynamic equations governing  $n_a(\vec{q},t)$  and  $n_b(\vec{q},t)$ .
- (d) Write down the one particle densities corresponding to a configuration in which  $n_a(\vec{q}) + n_b(\vec{q}) = n$  is uniform across a system at rest and at uniform temperature, i.e.  $\vec{u} = 0$  with n and T constant throughout. Does a non-uniform mixture, with spatially varying  $n_a(\vec{q})$  and  $n_b(\vec{q})$ , come to equilibrium in zeroth order hydrodynamics?
- (e) The first order solutions to the Boltzmann equation are given by

$$f_{\alpha}^{1}(\vec{q}, \vec{p}, t) = f_{\alpha}^{0} \left[ 1 - \tau_{\alpha} \mathcal{L}_{\alpha} [\ln f_{\alpha}^{0}] \right],$$

where  $\tau_{\alpha}$  is a characteristic time between collisions. Compute  $\vec{u}_{\alpha} = \langle \vec{p}_{\alpha}/m_{\alpha} \rangle$  at first order.

(f) Show that in first order hydrodynamics the densities relax by diffusion, and identify the diffusion constant.

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6. Viscosity: Consider a classical gas between two plates separated by a distance w. One plate at y = 0 is stationary, while the other at y = w moves with a constant velocity  $v_x = u$ . A zeroth order approximation to the one particle density is,

$$f_1^0(\vec{p}, \vec{q}) = \frac{n}{(2\pi m k_B T)^{3/2}} \exp\left[-\frac{1}{2m k_B T} \left((p_x - m\alpha y)^2 + p_y^2 + p_z^2\right)\right],$$

obtained from the *uniform* Maxwell–Boltzmann distribution by substituting the average value of the velocity at each point. ( $\alpha = u/w$  is the velocity gradient.)

(a) The above approximation does not satisfy the Boltzmann equation as the collision term vanishes, while  $df_1^0/dt \neq 0$ . Find a better approximation,  $f_1^1(\vec{p})$ , by linearizing the Boltzmann equation, in the single collision time approximation, to

$$\mathcal{L}\left[f_1^1\right] pprox \left[\frac{\partial}{\partial t} + \frac{\vec{p}}{m} \cdot \frac{\partial}{\partial \vec{q}}\right] f_1^0 pprox -\frac{f_1^1 - f_1^0}{\tau_{\times}},$$

where  $\tau_{\times}$  is a characteristic mean time between collisions.

- (b) Calculate the net transfer  $\Pi_{xy}$  of the x component of the momentum, of particles passing through a plane at y, per unit area and in unit time.
- (c) Note that the answer to (b) is independent of y, indicating a uniform transverse force  $F_x = -\Pi_{xy}$ , exerted by the gas on each plate. Find the coefficient of viscosity, defined by  $\eta = F_x/\alpha$ .

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- 7. Effusion: The probability distribution for speed c of particles of mass m in a gas at temperature T is proportional to  $c^2e^{-\frac{c^2}{2\sigma^2}}$ , with  $\sigma^2 = k_BT/m$ . Some particles are allowed to leak (effuse) out of a small hole with diameter much less than the mean free path.
- (a) Show that the probability distribution for speed of the escaping particles is proportional to  $c^3e^{-\frac{c^2}{2\sigma^2}}$ .
- (b) Find the average kinetic energy of the escaping particles.
- (c) What is the fraction of escaping particles with kinetic energy greater than  $\mathcal{E}$ ?

 $\dagger$  Reviewing the problems and solutions provided on the course web-page for preparation for  $Test\ 2$  should help you with the above problems.