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Due: Wednesday, Dec 7, 2022 Collaborators/References: Piazza

## 1. The SWAP test

The SWAP test tests whether two pure quantum states  $|\phi\rangle$  and  $|\psi\rangle$  are the same. Before the measurement in the first qubit is made, the circuit does the following (ignoring normalization):

$$|+\rangle |\psi\rangle |\phi\rangle \rightarrow |0\rangle |\psi\rangle |\phi\rangle + |1\rangle |\phi\rangle |\psi\rangle$$
$$\rightarrow |+\rangle |\psi\rangle |\phi\rangle + |-\rangle |\phi\rangle |\psi\rangle$$

(a) If  $|\phi\rangle = |\psi\rangle$ , then the state of the circuit before the measurement is

$$(|+\rangle + |-\rangle) |\psi\rangle |\psi\rangle = |0\rangle |\psi\rangle |\psi\rangle.$$

So the probability that we observe  $|0\rangle$  in the first wire is  $\boxed{1}$ .

(b) The state of the circuit before the measurement is

$$\frac{1}{2}\left|0\right\rangle \left(\left|\psi\right\rangle \left|\phi\right\rangle + \left|\phi\right\rangle \left|\psi\right\rangle\right) + \frac{1}{2}\left|1\right\rangle \left(\left|\psi\right\rangle \left|\phi\right\rangle - \left|\phi\right\rangle \left|\psi\right\rangle\right).$$

The probability that we observe  $|0\rangle$  in the first wire is

$$\frac{1}{4} \left( \left\langle \psi \right| \left\langle \phi \right| + \left\langle \phi \right| \left\langle \psi \right| \right) \left( \left| \psi \right\rangle \left| \phi \right\rangle + \left| \phi \right\rangle \left| \psi \right\rangle \right) = \frac{1}{4} \left( 1 + 1 \right) = \left| \frac{1}{2} \right|$$

(c) Suppose we apply the SWAP test with the inputs being two identical density matrices:

$$\rho_1 = \rho_2 = p |0\rangle \langle 0| + (1-p) |1\rangle \langle 1|$$
.

We can do this problem probabilistically. The initial states of the circuit and associated probabilities are:

$$Pr(|+\rangle |0\rangle |0\rangle) = p^{2}$$

$$Pr(|+\rangle |0\rangle |1\rangle) = p(1-p)$$

$$Pr(|+\rangle |1\rangle |0\rangle) = (1-p)p$$

$$Pr(|+\rangle |1\rangle |1\rangle) = (1-p)(1-p)$$

From the previous parts, the probability that we observe  $|0\rangle$  on the top wire is

$$p^{2} + (1-p)(1-p) + \frac{1}{2}[p(1-p) + (1-p)p] = \boxed{1-p+p^{2}}.$$

## 2. kl-qubit code

The generalization of the 9-qubit code to a *kl*-qubit code has the codewords:

$$|0\rangle_{L} = \frac{1}{2^{l/2}} (|\underbrace{000\dots0}_{k}\rangle + |\underbrace{111\dots1}_{k}\rangle)^{\otimes l}$$
$$|1\rangle_{L} = \frac{1}{2^{l/2}} (|\underbrace{000\dots0}_{k}\rangle - |\underbrace{111\dots1}_{k}\rangle)^{\otimes l}$$

Here k, l are odd numbers. In the known case where k = l = 3, we know that the code can correct 1 bit error and 1 phase error. To correct the bit error we need to measure 2 syndrome bits. To correct the phase error we also need to measure 2 syndrome bits.

## 3. Quantum Hamming code

Given a qubit  $|\psi\rangle = \alpha |0\rangle + \beta |1\rangle$ . The 7-qubit Hamming encoded version is

$$|\psi\rangle \rightarrow \alpha |0_L\rangle + \beta |1_L\rangle$$

where

$$|0_L\rangle = \frac{1}{\sqrt{8}}(|0000000\rangle + |1010101\rangle + |0110011\rangle + |1100110\rangle + |0001111\rangle + |1011010\rangle + |0111100\rangle + |1101001\rangle)$$

$$|1_L\rangle = \frac{1}{\sqrt{8}}(|1111111\rangle + |0101010\rangle + |1001100\rangle + |0011001\rangle + |1110000\rangle + |0100101\rangle + |111000011\rangle + |0010110\rangle)$$

as usual. Here, the parity check matrix is (following Nielsen and Chuang's convention):

$$H = \begin{pmatrix} 0 & 0 & 0 & 1 & 1 & 1 & 1 \\ 0 & 1 & 1 & 0 & 0 & 1 & 1 \\ 1 & 0 & 1 & 0 & 1 & 0 & 1 \end{pmatrix}.$$

It is easy to check that H all 7-bit strings in  $|0_L\rangle$  and  $|1_L\rangle$  are annihilated by H. This means we can do parity check on  $\alpha |0_L\rangle + \beta |1_L\rangle$  to find the  $\sigma_x$  error on qubit 3. This is because the corrupted state takes the form

$$|\psi'\rangle = \sigma_{z,6}(\alpha |0_L + e_3\rangle + \beta |1_L + e_3\rangle)$$

- (a) Suppose there is a  $\sigma_x$  error on qubit 3 and  $\sigma_z$  error on qubit 6. Since the 7-qubit Hamming code can correct one bit flip and one phase flip either on one or different qubits, the one-qubit state that we will obtain after error correction is simply  $|\psi'\rangle = |\psi\rangle = \alpha |0\rangle + \beta |1\rangle$ .
- (b) Suppose there is a  $\sigma_x$  error on qubit 3 and  $\sigma_y$  error on qubit 6.

## 4. CSS code with distance 2

(a) We are given a CSS code  $CSS(C_1, C_2)$  where

$$C_1 = \{0000, 0011, 0101, 1001, 0110, 1010, 1100, 1111\} \qquad C_2 = \{0000, 1111\}$$

The dimension of  $CSS(C_1, C_2)$  is  $|C_1|/|C_1| = 8/2 = 4$ . The  $CSS(C_1, C_2)$  is defined by

$$CSS(C_1, C_2) = \{|x + C_2| | x \in C_1\}.$$

By inspection, we see that

$$CSS(C_1, C_2) = \text{span}\{|0000 + C_2\rangle, |0011 + C_2\rangle, |0101 + C_2\rangle, |1001 + C_2\rangle\}.$$

(b) Suppose we apply a  $\sigma_x$  to qubits 1 and 2 of this code. Then the we have

$$CSS(C_1, C_2) \rightarrow \text{span}\{|1100 + C_2\rangle, |1111 + C_2\rangle, |1001 + C_2\rangle, |0101 + C_2\rangle\} = CSS(C_1, C_2).$$

So, this operation takes codewords to states in the code. Labeling the codewords as

$$|0000 + C_2\rangle$$
,  $|0011 + C_2\rangle$ ,  $|0101 + C_2\rangle$ ,  $|1001 + C_2\rangle = |a\rangle$ ,  $|b\rangle$ ,  $|c\rangle$ ,  $|d\rangle$ .

we see that applying  $\sigma_x$  on qubits 1 and 2 does the following:

$$|a\rangle \rightarrow |b\rangle \quad |b\rangle \rightarrow |a\rangle \quad |c\rangle \rightarrow |d\rangle \quad |d\rangle \rightarrow |c\rangle$$
.

(c) Since  $\sigma_z$  does nothing but change the phases on the basis states of the CSS code, the vector space spanned by the transformed states still remains the same. Because of this, this operation also takes codewords in  $CSS(C_1, C_2)$  to codewords in  $CSS(C_1, C_2)$ . Under this operation, the four basis codewords are mapped to themselves, with potentially a phase flip:

$$|a\rangle \rightarrow |a\rangle \quad |b\rangle \rightarrow |b\rangle \quad |c\rangle \rightarrow -|c\rangle \quad |d\rangle \rightarrow -|d\rangle$$

(d) Now suppose we apply H to all four qubits. To answer this question we can be explicit and write down what happens to all the basis states. But we can also obtain a solution as follows. The CSS code above is spanned by

$$\{|0000\rangle + |1111\rangle, |0011\rangle + |1100\rangle, |0101\rangle + |1100\rangle, |1001\rangle + |0110\rangle\}$$

After applying  $H^{\otimes 4}$ , the four codewords are transformed to

$$\{|++++\rangle + |----\rangle, |++--\rangle + |--++\rangle, |+-+-\rangle + |--++\rangle, |+--+\rangle\}$$

Is the span of these states the CSS code? To this end, we consider the components  $|0000\rangle$ ,  $|0011\rangle$ ,  $|0101\rangle$ ,  $|1001\rangle$ . It is not difficult to see that for each of these components, each of the resulting states also has  $|1111\rangle$ ,  $|1100\rangle$ ,  $|1010\rangle$ , and  $|0110\rangle$  respectively. Moreover, there are no components with an odd number of 1's, since the amplitudes for these components cancel in the expansion of the states from the +, – basis into the 0,1 basis. As a result, the operation  $H^{\otimes 4}$  also take the four codewords to states in the code. In particular,

$$|a\rangle \rightarrow |a\rangle + |b\rangle + |c\rangle + |d\rangle$$

$$|b\rangle \rightarrow |a\rangle + |b\rangle - |c\rangle - |d\rangle$$

$$|c\rangle \rightarrow |a\rangle - |b\rangle + |c\rangle - |d\rangle$$

$$|d\rangle \rightarrow |a\rangle - |b\rangle - |c\rangle + |d\rangle$$