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Note: Einsten's
 Noether's theorem for a scalar field in flat space
                                                                                                                                      implicit sum notation
  Def: The smooth one-pavemeter subgroup of transformations is used here.
  \widetilde{x}^{\mu} = \chi_{\varepsilon}^{\mu}(x) \widetilde{\phi}(\widetilde{x}) = \overline{F}_{\varepsilon}(\phi(x))
 is an infinitesimal symmetry of the action S[p]= \dn x & (p(x), \partial p(x), \tau)
  if \delta S = \frac{d}{d\epsilon|_{\epsilon=0}} \widetilde{S}[\widetilde{\phi}] = 0 for all \phi(x) \left(\widetilde{S}[\widetilde{\phi}] = \int_{X_{\epsilon}(u)} d^{n}x \mathcal{L}(\widetilde{\phi}(\widetilde{x}), \widetilde{\widetilde{S}}\widetilde{\phi}(\widetilde{x}), \widetilde{x})\right)
 Theorem: if \tilde{x}'' = X_{\varepsilon}''(x), \tilde{\phi}(\tilde{x}) = \tilde{F}_{\varepsilon}(\phi(x)) is an infinitesimal symmetry
  of S[d] = \ \dots d' \times \mathcal{L}(\phi(\times), \partial \phi(\times), \times) \quad \text{for all (nice) } \ \mathcal{U} \sum \mathcal{L}(\phi')
                                                                                                                                                 integration over
 then \partial_{\mu} \left[ \mathcal{L} S \times^{\mu} + \frac{\partial \mathcal{L}}{\partial \phi_{,\mu}} (S \phi - \phi_{,\nu} S \times^{\nu}) \right] = 0 when \phi satisfies the
                                                                                                                                                      U should be defined
 Euler-Logrenge eqs. \frac{\partial \mathcal{L}}{\partial \phi} - \partial_m \frac{\partial \mathcal{L}}{\partial \phi_{,m}} = 0. \left( S \times^m = \frac{1}{2} | X_{\varepsilon}^m(x), S \phi(\phi) = \frac{1}{2} | F_{\varepsilon}(\phi) \right)
 Note: in Minkowski space this is a continuity equation duit o
proof Very similar to the particle case. Let's set De=0 = dele=0
 By assumption SS = D_{\varepsilon=0} \widetilde{S}[\tilde{\phi}] = \frac{1}{4\varepsilon} \widetilde{S}[\tilde{\phi}] = 0.
                                                                                                                                  Note that
                                                                                                                                  Dero (A(E)B(E))
 \widetilde{S}[\widetilde{\phi}] = \int_{\mathcal{X}_{\varepsilon}(u)} J^{n}\widetilde{x} \, \mathcal{L}(\widetilde{\phi}(\widetilde{x}), \widetilde{\partial}\widetilde{\phi}(\widetilde{x}), \widetilde{x})
                                                                                                                               \left( \mathcal{D}_{\varepsilon=0} \mathcal{A}(\varepsilon) \right) \mathcal{B}(o) + \mathcal{A}(o) \left( \mathcal{D}_{\varepsilon=0} \mathcal{B}(\varepsilon) \right)
              = Judnx | det Je | L ( p (x), d p (x), x)
     where J_{\varepsilon} is the Jacobian metrix, (J_{\varepsilon})^{\alpha}v = \frac{\partial X_{\varepsilon}^{\alpha}}{\partial x^{\alpha}} = \frac{\partial \tilde{x}^{\alpha}}{\partial x^{\alpha}}
      Since everything is smooth, (J_e^{-1})^m v = \frac{\partial x^n}{\partial \tilde{x}^n}
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Some information that we need: Let A be an invertible matrix. · det(A) = det(A) tr(A' dA) (Jecobi's formule) $\cdot \frac{1}{4} A^{-1} = -A^{-1} \frac{1}{4} A^{-1}$ Now. · De= (Je) " = De= 3×" = 3, 8×" $D_{\varepsilon=0} \det J_{\varepsilon} = \det (J_{0}) \det (J_{0}') \det (J_{0}'') = \det (D_{\varepsilon=0} J_{\varepsilon}) = \partial_{u} \delta_{x}^{u}$ · De= | det Je| = | let Jo | De= det Je = 2 & & x " · De : 0 J' = - J' (De JE) J' = - De : Je · $D_{\varepsilon} = \widetilde{\phi}(\check{x}) = D_{\varepsilon} = F_{\varepsilon}(\phi(x)) = \delta\phi(\phi(x))$ $\frac{\partial \tilde{\phi}}{\partial \tilde{x}^{\mu}} = \frac{\partial \tilde{x}^{\mu}}{\partial \tilde{x}^{\mu}} = \frac{\partial$ $\Rightarrow D_{\varepsilon=0} \frac{\partial \vec{\phi}}{\partial \vec{x}^{m}} = \left(\frac{1}{64} \delta d\right) \partial_{\nu} \phi \left(\vec{J_{0}}\right)^{\nu}_{m} - \delta \vec{F_{0}} \left(\phi(x)\right) \partial_{\nu} \phi \partial_{m} \delta x^{\nu}$ $= \partial_{m} \delta \phi - \phi_{,\nu} \partial_{m} \delta x^{\nu}$ $o = SS = \int_{\mathcal{U}} d^{4}x \left\{ \left(D_{e^{-x}} \right| \det J_{\epsilon} \right) \mathcal{L} \left(\phi(x), \partial \phi(x), x \right) + |\det J_{0}| D_{\epsilon^{-x}} \mathcal{L} \left(\delta(\tilde{x}), \tilde{\partial} \tilde{b}(\tilde{x}), \tilde{x} \right) \right\}$ = \ 1 x \ 2 2 3 5 x 4 + 3 2 5 x 4 + 3 6 5 d + 3 6 (3 x 5 d - b, v 2 x 5 x) \ > use this to check if S5 =0 for the given trensformation. Here is where we use the EL egs. $\frac{\partial \mathcal{L}}{\partial t,n} \partial_n \delta \phi = \partial_n \left(\frac{\partial \mathcal{L}}{\partial t,n} \delta \phi \right) - \left(\partial_n \frac{\partial \mathcal{L}}{\partial \phi_n} \right) \delta \phi = \partial_n \left(\frac{\partial \mathcal{L}}{\partial \phi_n} \delta \phi \right) - \frac{\partial \mathcal{L}}{\partial \phi} \delta \phi$

$$\frac{\partial \mathcal{L}}{\partial f_{n}} f_{n} \partial_{n} sx^{*} = \partial_{n} \left(\frac{\partial \mathcal{L}}{\partial f_{n}} f_{n} sx^{*} \right) - \partial_{n} \left(\frac{\partial \mathcal{L}}{\partial f_{n}} f_{n} \right) sx^{*}$$

$$= \partial_{n} \left(\frac{\partial \mathcal{L}}{\partial f_{n}} f_{n} sx^{*} \right) - \left(\partial_{n} \frac{\partial \mathcal{L}}{\partial f_{n}} f_{n} sx^{*} - \frac{\partial \mathcal{L}}{\partial f_{n}} f_{n} sx^{*} \right)$$

$$= \partial_{n} \left(\frac{\partial \mathcal{L}}{\partial f_{n}} f_{n} sx^{*} \right) - \frac{\partial \mathcal{L}}{\partial f_{n}} f_{n} sx^{*} - \frac{\partial \mathcal{L}}{\partial f_{n}} f_{n} sx^{*} + \frac{$$