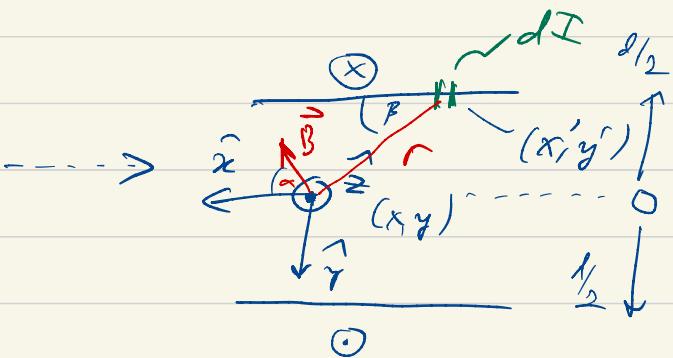
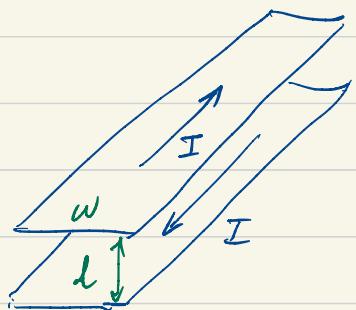


# 8.311 Electromagnetic Theory

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Part 5, due April 6, 2022

(7)



$$\text{Have } dI = \frac{I}{w} dw$$

$$dB \text{ (due to } dI) = \frac{\mu_0 dI}{2\pi r}$$

$$= \frac{\mu_0 dI}{2\pi} \left[ \frac{1}{(x-x')^2 + (y-y')^2} \right] \frac{dI}{w}$$

$$\therefore dB_x = dB \cos \alpha = dB \sin \beta$$

$$= \frac{\mu_0 dI}{2\pi} \frac{y - y'}{(x - x')^2 + (y - y')^2}$$

$$dR_y = -dB \sin \alpha$$

$$= -\frac{\mu_0 dI}{2\pi} \frac{x - x'}{(x - x')^2 + (y - y')^2}$$

2 p Inter

$$\Rightarrow B_x \Big|_{y=0} = 2 \int dB_x = 2 \int \frac{\mu_0}{2\pi} \frac{y - y'}{(x - x')^2 + (y - y')^2} \frac{I}{w} dx'$$

$$= 2 \left\{ \begin{array}{l} x' = \frac{-w}{2} \\ x' = \frac{w}{2} \end{array} \right. \frac{\mu_0}{2\pi} \frac{-d/2}{(x - x')^2 + (d/2)^2} \frac{I}{w} dx'$$

$$= \frac{\mu_0 I}{\pi w} \left\{ \tan^{-1} \frac{x - w/2}{d/2} - \tan^{-1} \frac{x + w/2}{d/2} \right\}$$

$B_y \Big|_{y=0} = 0$  due to symmetry.

$\int_0^L$

$$\boxed{\vec{B}(x, y, 0) = \frac{\mu_0 I}{\pi w} \left\{ \tan^{-1} \frac{x-w/2}{d/2} - \tan^{-1} \frac{x+w/2}{d/2} \right\}}$$

(2)

$d < w$  ... then

$$\overbrace{\hspace{10cm}}^{\textcircled{X}}$$

$$(i) \quad B_x \rightarrow \frac{\mu_0 I}{\pi w} \left\{ +\frac{\pi}{2} + \frac{\pi}{2} \right\} (?)$$

I'm getting confused about the signs here,  
but the answer is that

$$\boxed{B_x \rightarrow \frac{\mu_0 I}{w}}$$

in  $\leftarrow$  direction

(ii) Vector potential?

Here  $\vec{B} = +\frac{\mu_0 I}{w} \hat{x}$  uniform...

So can take  $\vec{A} = B_y \hat{z}$

Check  $\vec{\nabla} \times \vec{A} = \partial_y A_z \hat{z} = B \hat{z}$  ✓

So take

$$\boxed{\vec{A} = +\frac{\mu_0 I}{w} y \hat{z} + \text{constant}}$$

(iii) magnetic force per unit length on each ship...

Field on one ship due to the other is

$$B_F = \frac{1}{2} B = \frac{\mu_0 I}{2w}$$

Force per unit length--

$\sim l \cdot A$

$$\frac{F}{e} = \frac{1}{e} \int_{\text{strip}} j B_F d^3 r$$

$$= \frac{1}{l} B_F \int j d^3 r = I B_F = \boxed{\frac{\mu_0 I^2}{2w}}$$

(iv) magnetic energy per unit length--

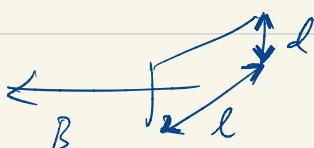
$$\frac{U}{e} = \frac{1}{e} \int_{\text{gap}} \frac{B^2}{2\mu_0} d^3 r = \boxed{\frac{B^2}{2\mu_0} \frac{1}{e} \cdot ldw = \frac{\mu_0 d I^2}{2w}}$$

W/W per

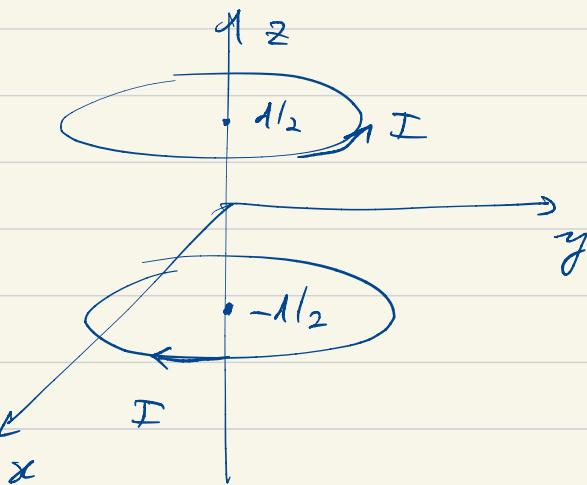
$$\frac{\mu_0 I}{w}$$

Self inductance--

$$\frac{L}{e} = \frac{\Phi}{I} \frac{1}{e} = \frac{1}{Il} \int_{\text{gap}} B \cdot d^2 \sigma = \frac{\mu \cdot I \cdot ld}{Il w} = \boxed{\frac{\mu d}{w}}$$



③ Calculate  $\vec{B}$  near center of  
anti - de Laval nozzle config



We know (from notes) that

$$\vec{B}(0, 0, z) = \frac{\mu_0 I}{2} \left\{ \frac{-R^2}{(R^2 + (d/2 + z)^2)^{3/2}} + \frac{R^2}{(R^2 + (d/2 - z)^2)^{3/2}} \right\}$$

Taylor expand this near  $z = 0$  we get

$$\vec{B}(0, 0, z) \approx \frac{4\pi d R^2 \mu_0 I z}{(d^2 + 4R^2)^{5/2}}$$

To find  $B_x, B_y$  we use rotational symmetry + the fact that  $\nabla \cdot \vec{B} = 0$

$$\rightarrow \partial_x B_x + \partial_y B_y = -\partial_z B_z$$

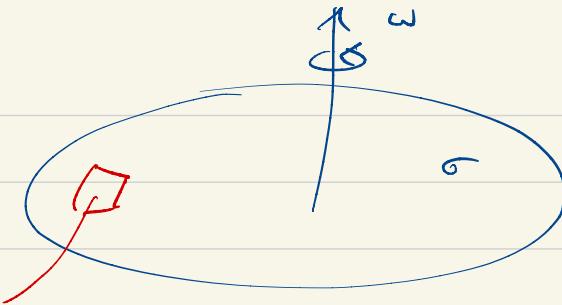
$$\Rightarrow \partial_x B_x = \partial_y B_y = -\frac{1}{2} \partial_z B_z$$

$$= -2y \frac{dR^2 \mu_0 I}{(d^2 + 4R^2)^{5/2}}$$

S.

$$\boxed{\vec{B} \Big|_{r \approx d} = \frac{48dR^2 \mu_0 I}{(d^2 + 4R^2)^{5/2}} \left( -\frac{x}{z} \hat{x} - \frac{y}{z} \hat{y} + \hat{z} \right)}$$

(4)

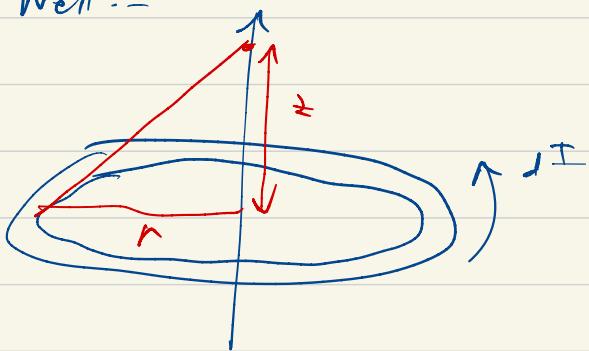


$$dr \ r \ d\theta$$

$$dI = \sigma r w dr$$

(i) induced  $\vec{B}$  field on axis ...

Well --



$$dB = \frac{\mu_0 dI}{2} \frac{r^2}{(r^2 + z^2)^{3/2}}$$

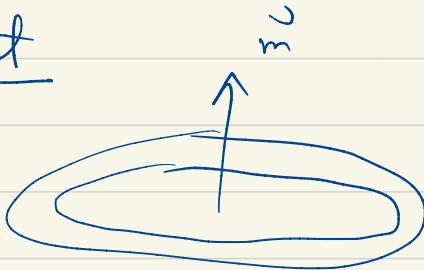
$$= \frac{\mu_0 \sigma w}{2} \frac{r^3}{(r^2 + z^2)^{3/2}} dr$$

$$\Rightarrow B = \int_{r=0}^{r=R} \frac{\mu_0 \sigma w}{2} \frac{r^3 dr}{(r^2 + z^2)^{3/2}} = (\text{...mathematica})$$

$$B = \mu_0 \sigma w \left\{ -|z| + \frac{R^2 + 2z^2}{2(R^2 + z^2)^{1/2}} \right\}$$

(ii) magnetic moment

$$m = \int dm$$



$$dm = A dI = \pi r^2 (\sigma r \omega) dr$$

$$= \pi \sigma \omega r^3 dr$$

$$\therefore [m] = \int_{r=0}^{r=R} \pi \sigma \omega r^3 dr = \left[ \frac{\pi \sigma \omega R^4}{4} \right]$$

Relate  $m$  to  $B$ ?

Compare  $B @ \infty$  to  $\tilde{B}$  due to  $m$  ...

$$B @ \infty \dots \rightarrow B = \mu_0 \sigma \omega / z \left\{ 1 + \frac{1 + R^2/z^2}{\sqrt{1 + R^2/z^2}} \right\}$$

expand  $R/z \ll 1 \dots$  to get

$$\boxed{\beta = \sigma \mu_0 w |z| \frac{1}{8} \frac{R^4}{|z|^4} = \frac{\mu_0 \sigma w R^4}{8|z|^3}}$$

Now field due to magnetic dipole moment -

$$\boxed{B} = \left( \frac{3 \vec{r} \cdot (\vec{m} \cdot \vec{z}) - m r^2}{r^5} \right) \frac{\mu_0}{4\pi}$$

and  $\vec{r} = |z| \vec{m}/m$

$$\approx \frac{3m |z|^2 - m |z|^2}{|z|^5} \frac{\mu_0}{4\pi}$$

$$= \frac{\mu_0}{4\pi} \frac{2m}{|z|^3} = \frac{\mu_0}{4\pi} \frac{2}{|z|^3} \cdot \frac{\pi \sigma w R^4}{4} = \boxed{\frac{\mu_0 \sigma w R^4}{8|z|^3}}$$

C