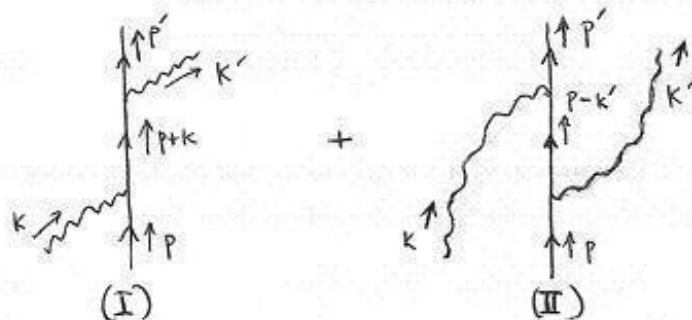


X-rays + gamma rays on electrons.

$e^- \gamma \rightarrow e^- \gamma$ (Compton scattering)

There two diagrams at lowest order



Let $\epsilon_\mu(k) + \epsilon_\mu(k')$
be the photon
polarization vectors

$$\begin{aligned}
 & \text{(I)} \quad \bar{u}(p') (-ie\gamma^\mu) \epsilon_\mu^*(k') \frac{i(\not{p} + \not{k} + m)}{(p+k)^2 - m^2 + i\epsilon} (-ie\gamma^\nu) \epsilon_\nu(k) u(p) + \text{(II)} \quad \bar{u}(p') (-ie\gamma^\mu) \epsilon_\mu(k) \frac{i(\not{p} - \not{k}' + m)}{(p-k')^2 - m^2 + i\epsilon} (-ie\gamma^\nu) \epsilon_\nu^*(k') u(p) \\
 & = -ie^2 \epsilon_\mu^*(k') \epsilon_\mu(k) \bar{u}(p') \left[\frac{\gamma^\mu (\not{p} + \not{k} + m) \gamma^\nu}{(p+k)^2 - m^2 + i\epsilon} + \frac{\gamma^\nu (\not{p} - \not{k}' + m) \gamma^\mu}{(p-k')^2 - m^2 + i\epsilon} \right] u(p)
 \end{aligned}$$

We can simplify this a little. Note that $k^2 = k'^2 = m_\gamma^2 = 0$.

$$\begin{aligned}
 p^2 = p'^2 = m^2. \quad \text{So} \quad (p+k)^2 &= p^2 + 2p \cdot k + k^2 = m^2 + 2p \cdot k \\
 (p-k')^2 &= p^2 - 2p \cdot k' + k'^2 = -2p \cdot k' + m^2
 \end{aligned}$$

$$\text{Also} \quad (\not{p} + m) \gamma^\mu u(p) = \gamma^\mu (-\not{p} + m) u(p) + \underbrace{2p^\mu}_{p_\mu \{ \gamma^\mu, \gamma^\nu \}} u(p)$$

Since $\not{p} u(p) = m u(p)$, $(-\not{p} + m) u(p) = 0$. Therefore

$$(\not{p} + m) \gamma^\mu u(p) = 2p^\mu u(p)$$

So the amplitude is now

$$i\mathcal{M} = -ie^2 \xi_\mu^*(k') \xi_\nu(k) \bar{u}(p') \left[\frac{\gamma^\mu \not{k} \gamma^\nu + 2\gamma^\mu \not{p}^\nu}{2p \cdot k} + \frac{\gamma^\nu \not{k}' \gamma^\mu - 2\gamma^\nu \not{p}^\mu}{2p \cdot k'} \right] u(p)$$

Photon polarizations + Ward identity

The interaction for QED has the general form

$$e \int d^4x j^\mu A_\mu \quad \text{where } j^\mu \text{ is the conserved electric charge current}$$

Let us consider a process with an outgoing photon with momentum k , polarization ϵ .

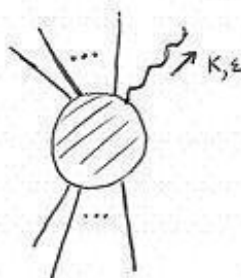
The amplitude can be written as

$$i\mathcal{M} = i\mathcal{M}^\mu(k) \xi_\mu^*(k)$$

where

$$\mathcal{M}^\mu(k) \propto \int d^4x e^{ik \cdot x} \langle f | j^\mu(x) | i \rangle$$

\uparrow final state, without photon under consideration \uparrow incoming state



Since $\partial_\mu j^\mu(x) = 0$,

$$k_\mu \mathcal{M}^\mu(k) \propto \int d^4x e^{ik \cdot x} k^\mu \langle f | j^\mu(x) | i \rangle = -i \int d^4x \partial_\mu (e^{ik \cdot x}) \langle f | j^\mu(x) | i \rangle$$

$$= i \int d^4x \, e^{ik \cdot x} \langle f | \partial_\mu j^\mu(k) | i \rangle = 0$$

This is a specific example of the Ward identity in QED.

When you replace the polarization vector for any external photon (can be incoming or outgoing) by the momentum k_μ , then the amplitude vanishes.

$$k_\mu M^\mu(k) = 0.$$

More on this later.