

## Some potentially useful information

- Euler-Lagrange equations for generalized coordinates  $q_j$

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{q}_j} - \frac{\partial L}{\partial q_j} = \sum_{\alpha} \lambda_{\alpha} \frac{\partial f_{\alpha}}{\partial q_j}, \quad \text{or} \quad \frac{d}{dt} \frac{\partial L}{\partial \dot{q}_j} - \frac{\partial L}{\partial q_j} = \sum_{\beta} \lambda_{\beta} \frac{\partial g_{\beta}}{\partial \dot{q}_j}$$

constraints: holonomic  $f_{\alpha}(q, t) = 0$  or semiholonomic  $g_{\beta} = \sum_j a_{\beta j}(q, t) \dot{q}_j + a_{\beta t}(q, t) = 0$

- Generalized forces:  $d/dt(\partial L/\partial \dot{q}_j) - \partial L/\partial q_j = R_j$

Friction forces:  $\vec{f}_i = -h(v_i)\vec{v}_i/v_i$ ,  $\vec{v}_i = \dot{\vec{r}}_i$  gives  $R_j = -\partial \mathcal{F}/\partial \dot{q}_j$ ,  $\mathcal{F} = \sum_i \int_0^{v_i} dv'_i h(v'_i)$

- Hamilton's equations for canonical variables  $(q_j, p_j)$ :  $\dot{q}_j = \frac{\partial H}{\partial p_j}$ ,  $\dot{p}_j = -\frac{\partial H}{\partial q_j}$

- Hamiltonian for a Lagrangian quadratic in velocities

$$L = L_0(q, t) + \dot{\vec{q}}^T \cdot \vec{a} + \frac{1}{2} \dot{\vec{q}}^T \cdot \hat{T} \cdot \dot{\vec{q}} \Rightarrow H = \frac{1}{2} (\vec{p} - \vec{a})^T \cdot \hat{T}^{-1} \cdot (\vec{p} - \vec{a}) - L_0(q, t)$$

- The Moment of Inertia Tensor and its relations:

$$I_{ab} = \int dV \rho(\vec{r}) [\vec{r}^2 \delta_{ab} - r_a r_b] \quad \text{or} \quad I^{ab} = \sum_i m_i [\delta^{ab} \vec{r}_i^2 - r_i^a r_i^b]$$

$$I_{ab}^{(Q)} = M(\delta_{ab} \vec{R}^2 - R_a R_b) + I_{ab}^{(\text{CM})}, \quad \hat{I}' = \hat{U} \hat{I} \hat{U}^T$$

- Euler's Equations:

$$\begin{aligned} I_1 \dot{\omega}_1 - (I_2 - I_3) \omega_2 \omega_3 &= \tau_1 \\ I_2 \dot{\omega}_2 - (I_3 - I_1) \omega_3 \omega_1 &= \tau_2 \\ I_3 \dot{\omega}_3 - (I_1 - I_2) \omega_1 \omega_2 &= \tau_3 \end{aligned}$$

- Vibrations:  $L = \frac{1}{2} \dot{\vec{\eta}}^T \cdot \hat{T} \cdot \dot{\vec{\eta}} - \frac{1}{2} \vec{\eta}^T \cdot \hat{V} \cdot \vec{\eta}$  has Normal modes  $\vec{\eta}^{(k)} = \vec{a}^{(k)} \exp(-i\omega^{(k)}t)$

$$\det(\hat{V} - \omega^2 \hat{T}) = 0, \quad (\hat{V} - [\omega^{(k)}]^2 \hat{T}) \cdot \vec{a}^{(k)} = 0, \quad \vec{\eta} = \text{Re} \sum_k C_k \vec{\eta}^{(k)}$$

- Generating functions for Canonical Transformations:  $K = H + \partial F_i / \partial t$  and

$$F_1(q, Q, t): \quad p_i = \frac{\partial F_1}{\partial q_i}, \quad P_i = -\frac{\partial F_1}{\partial Q_i}, \quad F_2(q, P, t): \quad p_i = \frac{\partial F_2}{\partial q_i}, \quad Q_i = \frac{\partial F_2}{\partial P_i}$$

- Poisson Brackets:  $[u, v]_{q,p} = \sum_j \left[ \frac{\partial u}{\partial q_j} \frac{\partial v}{\partial p_j} - \frac{\partial u}{\partial p_j} \frac{\partial v}{\partial q_j} \right], \quad \frac{du}{dt} = [u, H] + \frac{\partial u}{\partial t}$

- Relations for Hamilton's Principle function,  $S = S(q_1, \dots, q_n; \alpha_1, \dots, \alpha_n, t)$

$$K = 0, \quad P_i = \alpha_i, \quad Q_i = \beta_i = \frac{\partial S}{\partial \alpha_i}, \quad p_i = \frac{\partial S}{\partial q_i}$$

- Relations for Hamilton's Characteristic function,  $W = W(q_1, \dots, q_n; \alpha_1, \dots, \alpha_n)$

$$K = H = \alpha_1, \quad P_i = \alpha_i, \quad \beta_1 + t = \frac{\partial W}{\partial \alpha_1}, \quad \beta_{i>1} = \frac{\partial W}{\partial \alpha_i}, \quad p_i = \frac{\partial W}{\partial q_i}$$

- Action Angle Variables:  $J = \oint p dq, \quad w = \frac{\partial W(q, J)}{\partial J}, \quad \dot{w} = \frac{\partial H(J)}{\partial J} = \nu(J)$