

Ciao Tutti!

Effective Field Theory approach for Dark Energy/ Modified Gravity



Instituut-Lorentz

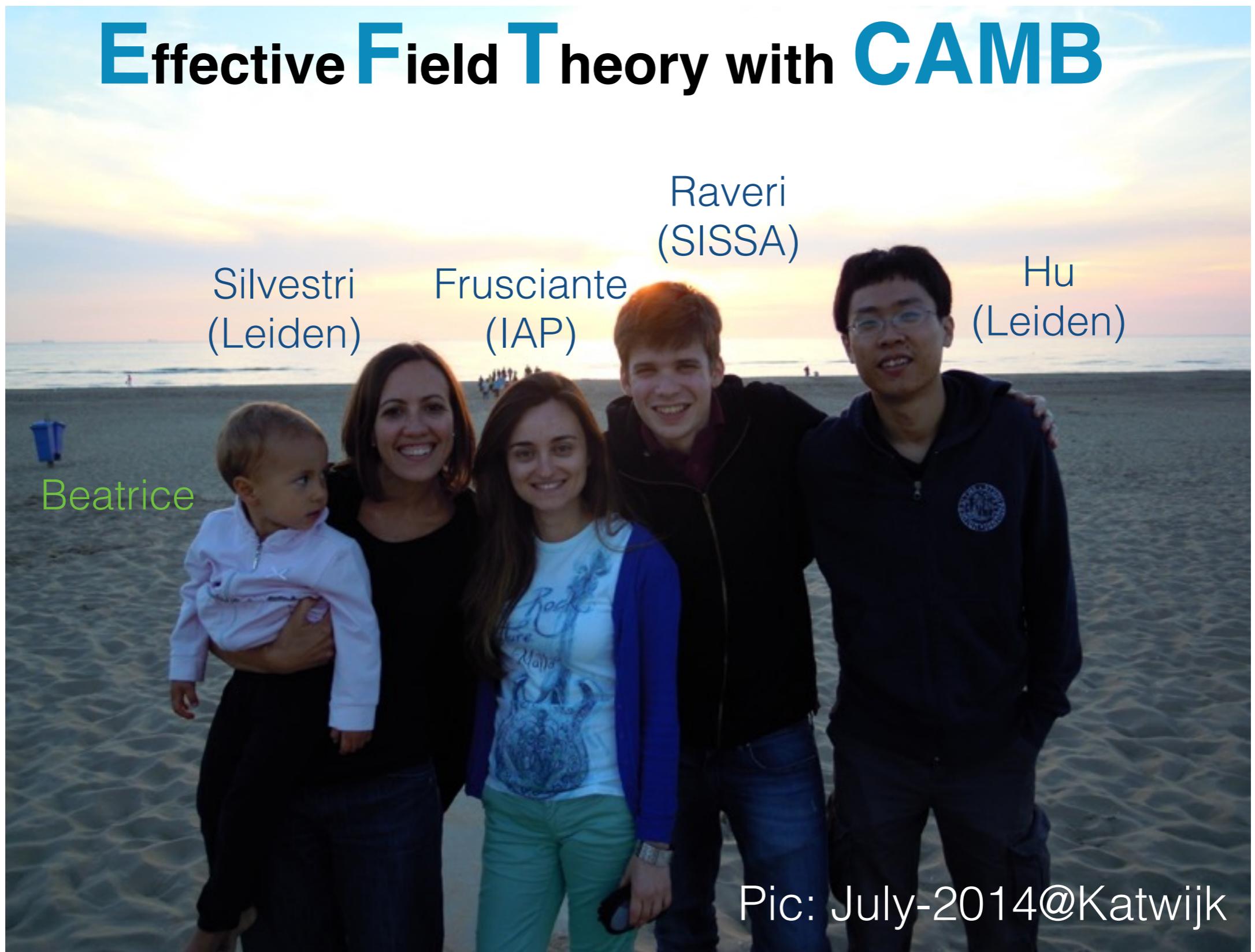
Bin HU

Lorentz Institute, Leiden University

Padova, May, 2015

EFTCAMB team

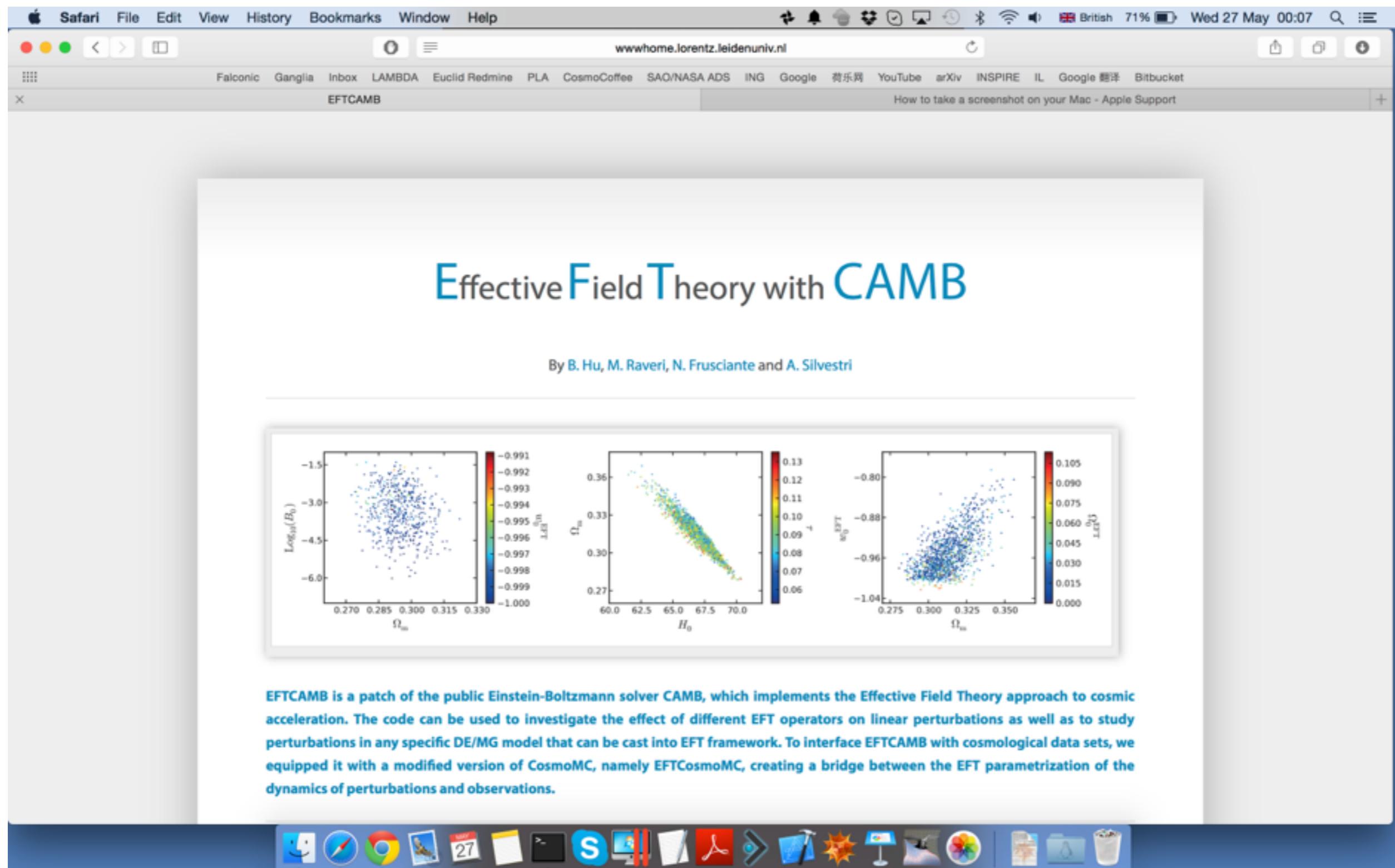
Effective Field Theory with CAMB

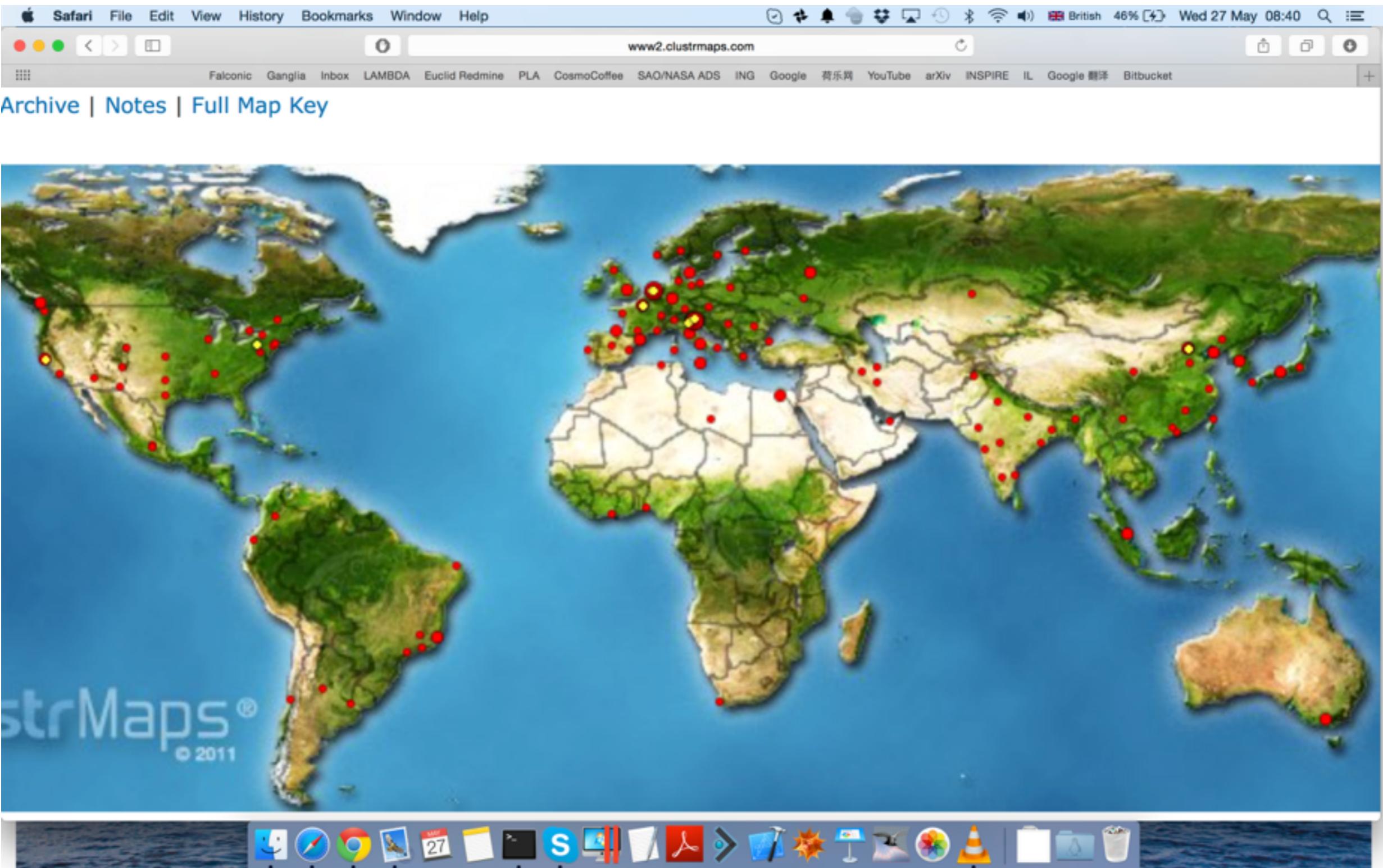


[Hu et.al. PRD89,103530(2014); PRD90,043513(2014); PRD91,063524(2015)]

<http://wwwhome.lorentz.leidenuniv.nl/~hu/codes/>

<http://wwwhome.lorentz.leidenuniv.nl/~hu/codes/>





clicked 1500+ times
but 300+ from italy and 300+ from Leiden

Outline

1. The test of gravity at linear perturbation level:
EFT for DE/MG
2. The structure of EFTCAMB
3. The results from EFTCosmoMC and Planck-2015
4. Conclusion

1. The test of gravity at linear perturbation level: EFT for DE/MG

- Effective Field Theory (EFT) approach provides generic parametrisation of the action of the scalar ***perturbation*** DE/MG with ***single*** scalar field.

$$S = \int d^4x \sqrt{-g} \left\{ \frac{m_0^2}{2} [1 + \Omega(\tau)] R + \Lambda(\tau) - a^2 c(\tau) \delta g^{00} + \frac{M_2^4(\tau)}{2} (a^2 \delta g^{00})^2 - \frac{\bar{M}_1^3(\tau)}{2} a^2 \delta g^{00} \delta K^\mu_\mu - \frac{\bar{M}_2^2(\tau)}{2} (\delta K^\mu_\mu)^2 \right. \\ \left. - \frac{\bar{M}_3^2(\tau)}{2} \delta K^\mu_\nu \delta K^\nu_\mu + \frac{a^2 \hat{M}^2(\tau)}{2} \delta g^{00} \delta R^{(3)} + m_2^2(\tau) (g^{\mu\nu} + n^\mu n^\nu) \partial_\mu (a^2 g^{00}) \partial_\nu (a^2 g^{00}) + \dots \right\} + S_m[\chi_i, g_{\mu\nu}],$$

[Bloomfield et. al. JCAP08(2013)010]

[Gubitosi et. al. JCAP 1302 (2013) 032]

* There are **7 independent** functions at linear level, EFT functions

- * Ω , Λ and c relate with background operators, only one are independent
- * EFT functions depend on **time** only

$$\boxed{\begin{aligned}\mathcal{H}^2 &= \frac{a^2}{3m_0^2(1+\Omega)}(\rho_m + 2c - \Lambda) - \mathcal{H}\frac{\dot{\Omega}}{1+\Omega}, \\ \dot{\mathcal{H}} &= -\frac{a^2}{6m_0^2(1+\Omega)}(\rho_m + 3P_m) - \frac{a^2(c+\Lambda)}{3m_0^2(1+\Omega)} - \frac{\ddot{\Omega}}{2(1+\Omega)},\end{aligned}}$$

$$\boxed{\begin{aligned}c &= -\frac{m_0^2\ddot{\Omega}}{2a^2} + \frac{m_0^2\mathcal{H}\dot{\Omega}}{a^2} + \frac{m_0^2(1+\Omega)}{a^2}(\mathcal{H}^2 - \dot{\mathcal{H}}) - \frac{1}{2}(\rho_m + P_m), \\ \Lambda &= -\frac{m_0^2\ddot{\Omega}}{a^2} - \frac{m_0^2\mathcal{H}\dot{\Omega}}{a^2} - \frac{m_0^2(1+\Omega)}{a^2}(\mathcal{H}^2 + 2\dot{\mathcal{H}}) - P_m.\end{aligned}}$$

1.1 The logic of construction of the action

1. Choose the time coordinate (clock), by asking

$$\delta\varphi(t, \vec{x}) \equiv \varphi(t, \vec{x}) - \bar{\varphi}(t) = 0$$

(breaking time translation
diffemorphism)

2. Build the block of the EFT by the operators which keep the unbroken 3D spatial Diffs

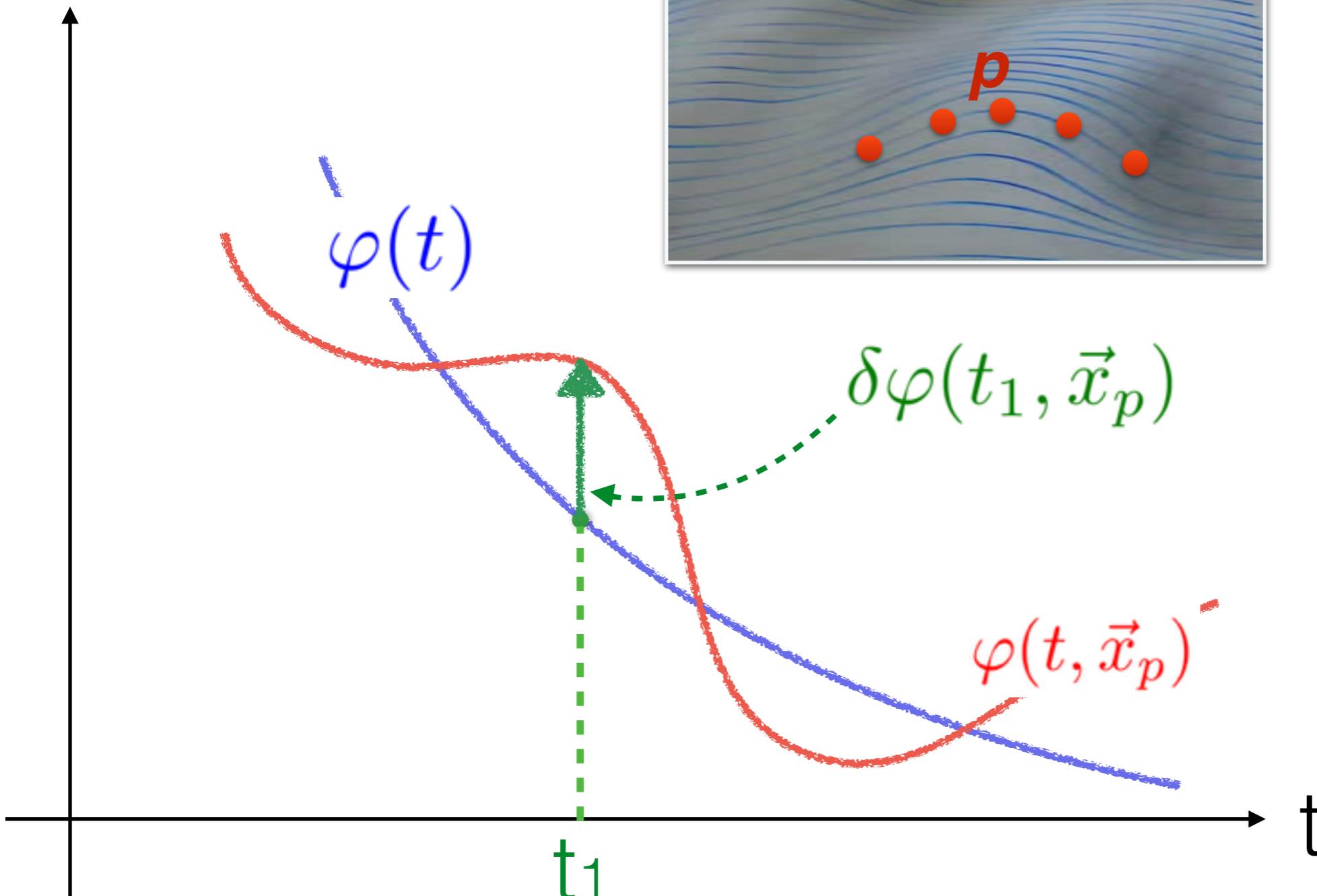
$$\delta g^{00}, \delta K_{\mu\nu}, \delta R_{\mu\nu\rho\sigma} \text{ (or } C_{\mu\nu\rho\sigma}), \delta R_{\mu\nu}, \text{ and } \delta R,$$

3. Multiply these operators by a only time dependent function

$$\begin{aligned}
S = & \int d^4x \sqrt{-g} \left\{ \frac{m_0^2}{2} [1 + \Omega(\tau)] R + \Lambda(\tau) - a^2 c(\tau) \delta g^{00} \right. \\
& + \frac{M_2^4(\tau)}{2} (a^2 \delta g^{00})^2 - \frac{\bar{M}_1^3(\tau)}{2} a^2 \delta g^{00} \delta K^\mu{}_\mu - \frac{\bar{M}_2^2(\tau)}{2} (\delta K^\mu{}_\mu)^2 \\
& - \frac{\bar{M}_3^2(\tau)}{2} \delta K^\mu{}_\nu \delta K^\nu{}_\mu + \frac{a^2 \hat{M}^2(\tau)}{2} \delta g^{00} \delta R^{(3)} \\
& \left. + m_2^2(\tau) (g^{\mu\nu} + n^\mu n^\nu) \partial_\mu (a^2 g^{00}) \partial_\nu (a^2 g^{00}) + \dots \right\} \\
& + S_m[\chi_i, g_{\mu\nu}], \tag{1}
\end{aligned}$$

1.2 Another point of view of EFT

Covariant approach



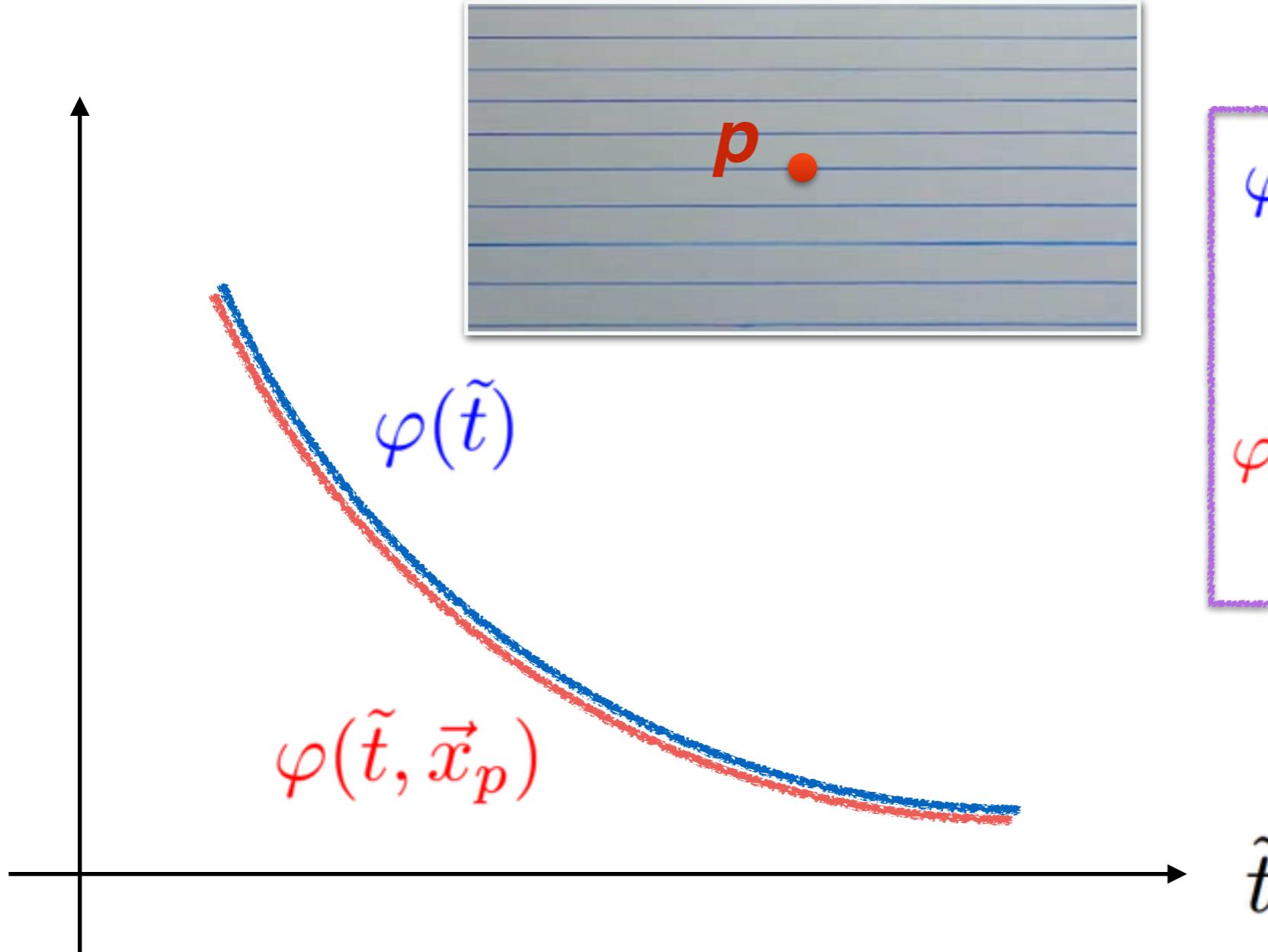
$\varphi(t)$: background
field config

$\varphi(t, \vec{x}_p)$: field config
at point ' p '

$\delta\varphi(t_1, \vec{x}_p)$: field fluct.
at point ' p '

Valid in ***ALL*** the gauge

EFT approach



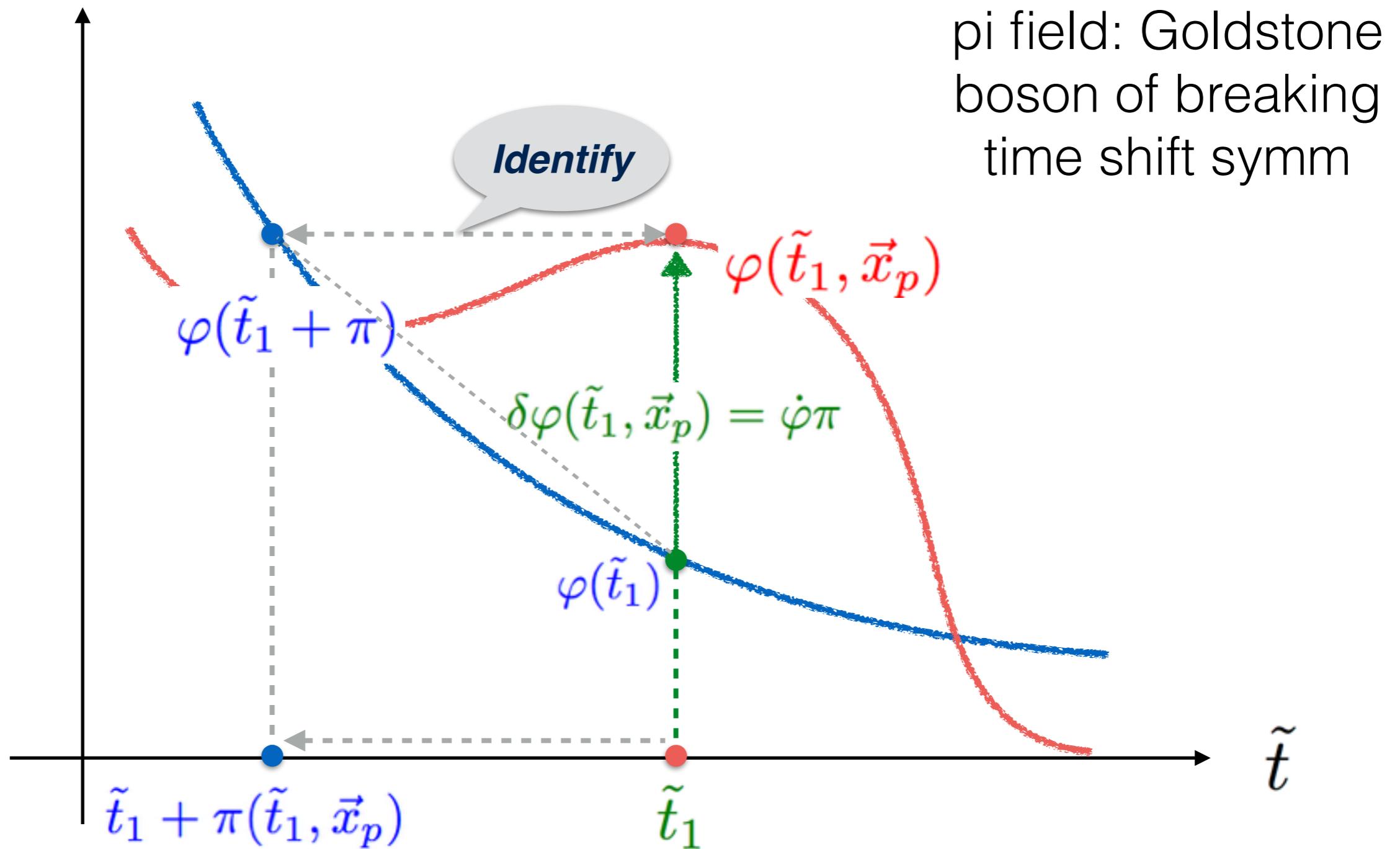
$\varphi(\tilde{t})$: background
field config

$\varphi(\tilde{t}, \vec{x}_p)$: field config
at point 'p'

Only Valid in the unitary gauge

$$\delta\varphi(t, \vec{x}) \equiv \varphi(t, \vec{x}) - \bar{\varphi}(t) = 0$$

EFT approach=> Covariant approach



Stuckburg trick: restore full covariance

1.3 Parametrizations

1. Full mapping

(From the covariant form)

e.g.

$$f(R) = -m^2 \frac{c_1(R/m^2)^n}{c_2(R/m^2)^n + 1},$$

[Hu,Sawicki PRD**76**, 064004 (2007)]

$$\Lambda = \frac{m_0^2}{2} [f - R f_R] \quad ; \quad c = 0 \quad ; \quad \Omega = f_R$$

(Work in progress with Rizzato et. al.)

2. Pure EFT parametrization

(Phenomenological param)

Constant models: $\Omega(a) = \Omega_0$;

Linear models: $\Omega(a) = \Omega_0 a$;

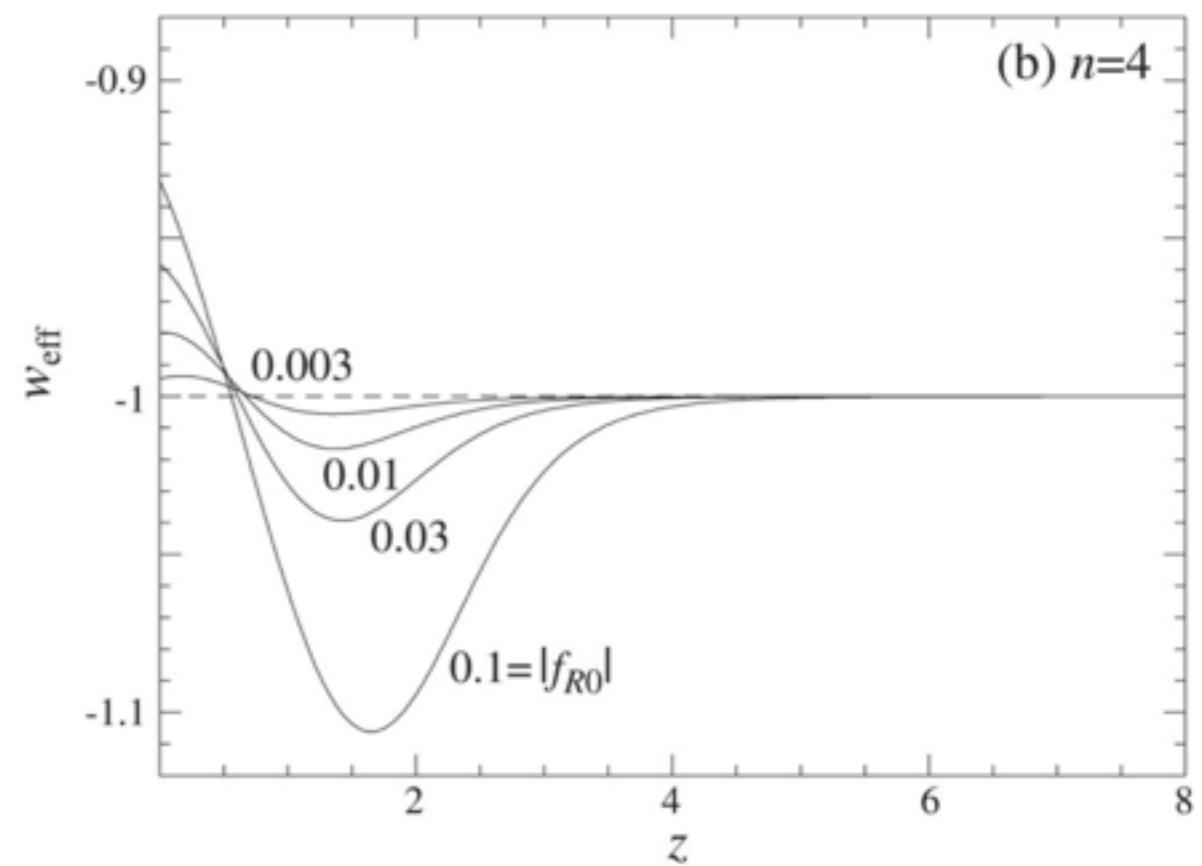
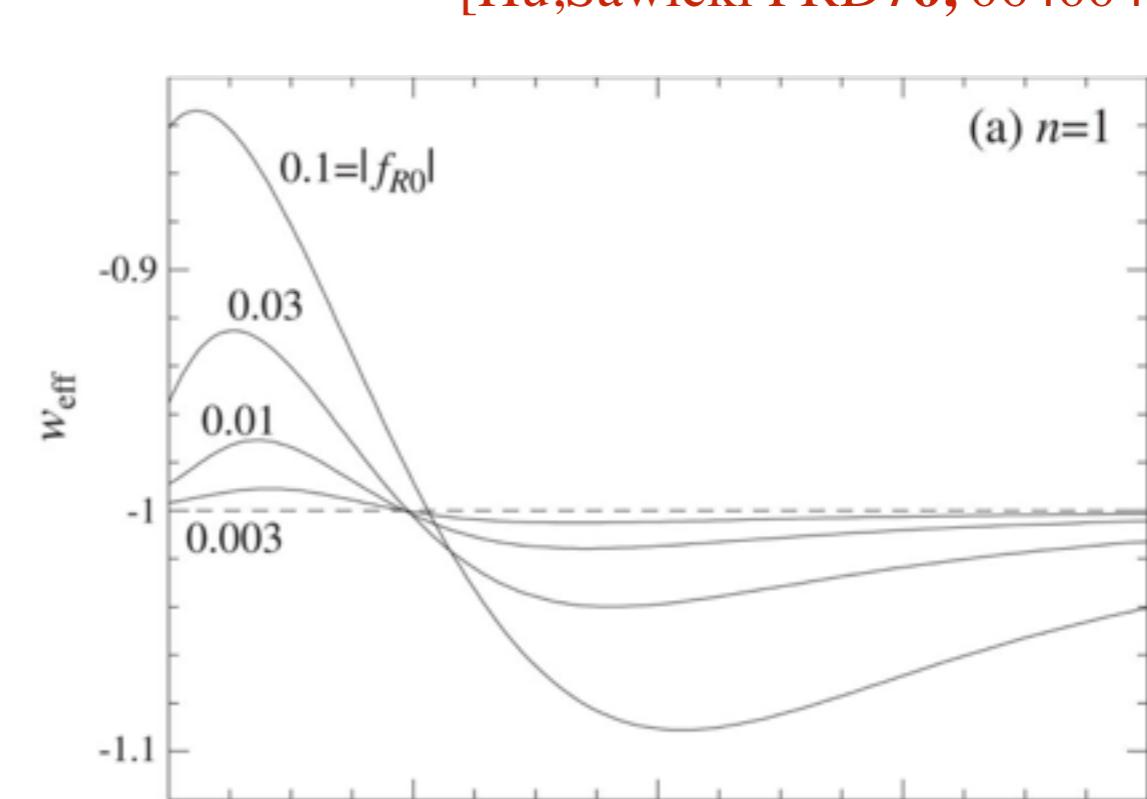
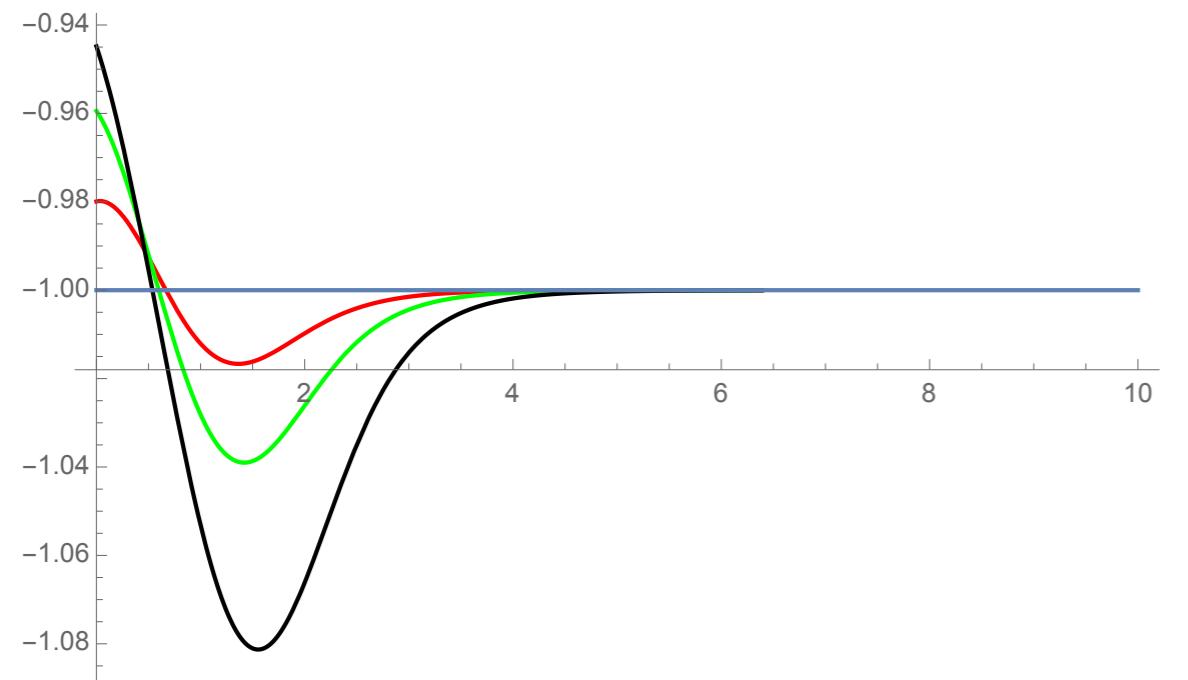
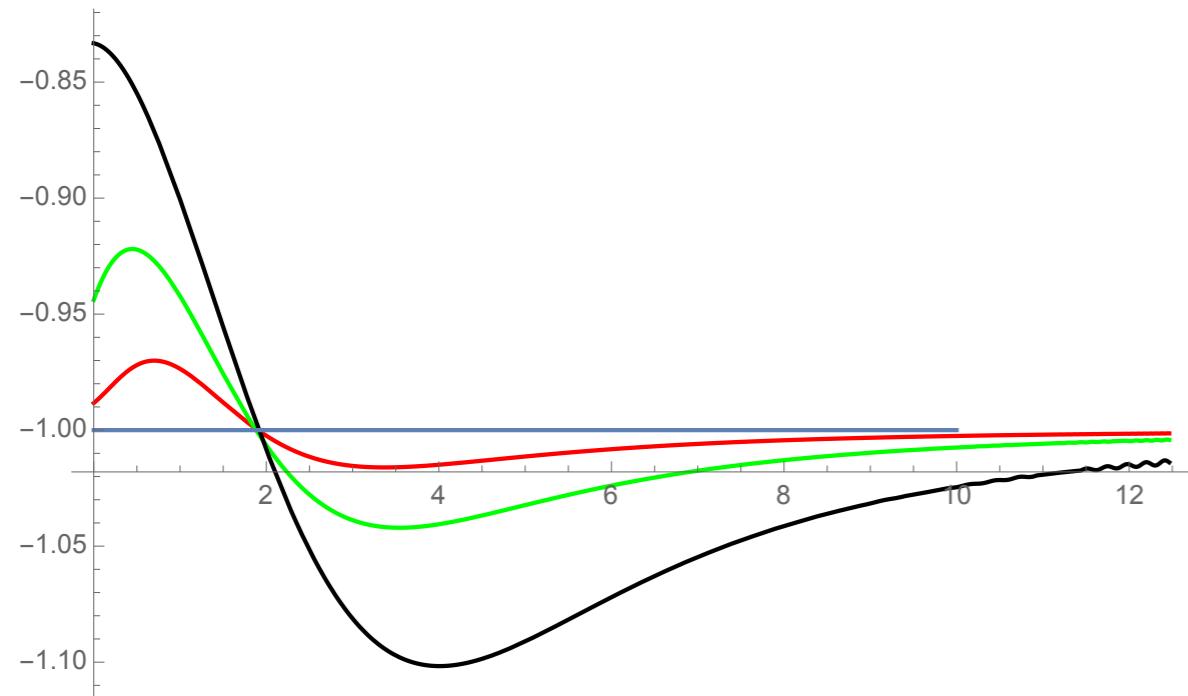
Power law models: $\Omega(a) = \Omega_0 a^s$;

Exponential models: $\Omega(a) = \exp(\Omega_0 a^s) - 1$.

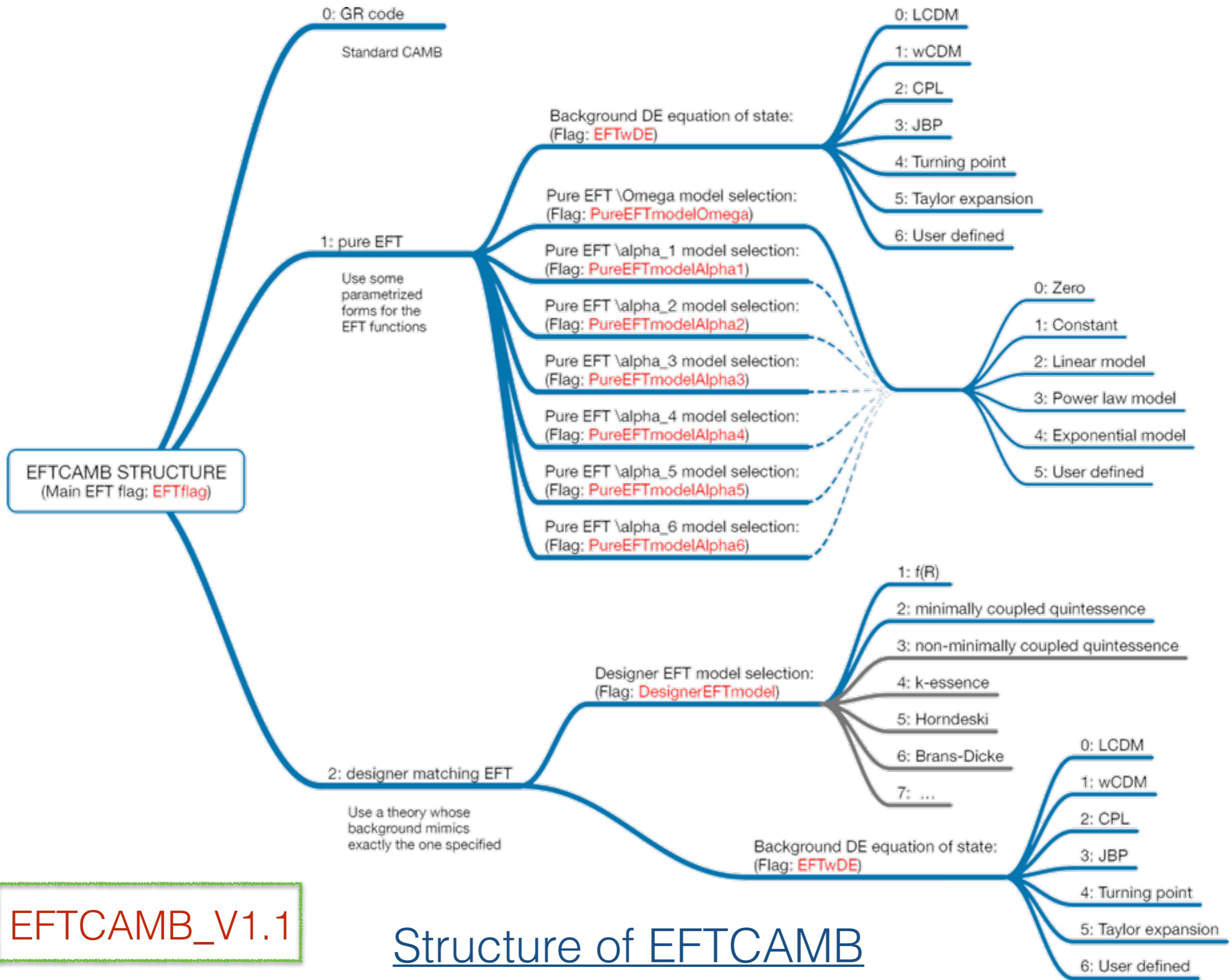
**Have to make sure
that your parametrisation
to be viable, e.g. ghost-free!**

(Plots from Matteo Rizzato)

[Hu,Sawicki PRD76, 064004 (2007)]



2. The structure of EFTCAMB



2.1 Background parametrization—EoS

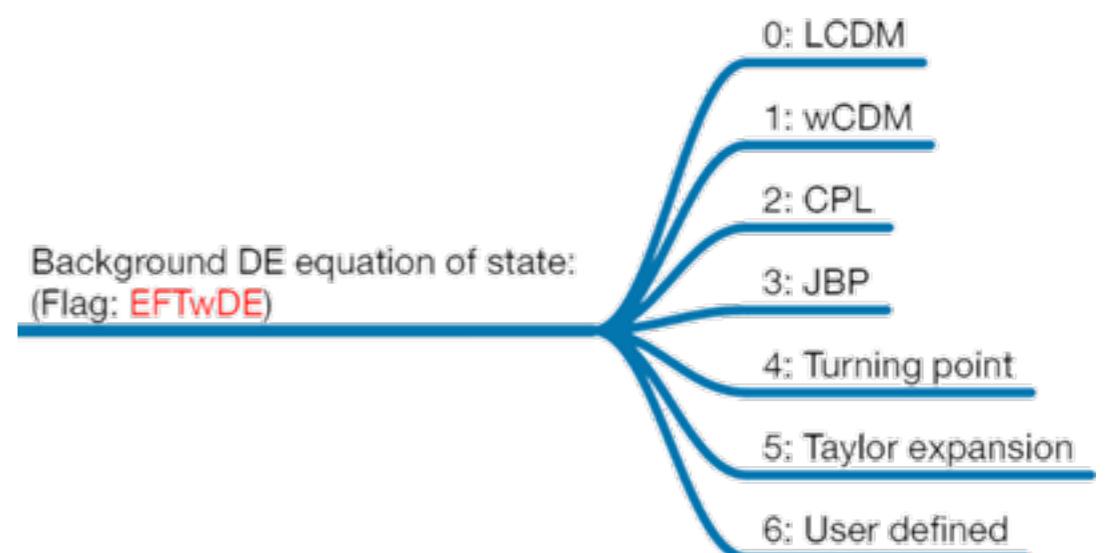
EFTCAMB provides 6 different kinds of parametrization of EoS
(Flag: **EFTwDE**), including:

LCDM ($w=-1$),

wCDM ($w=w_0$),

CPL ($w=w_0+w_a \cdot a$),

... ...



2.2.1 EFT parametrization: Pure EFT

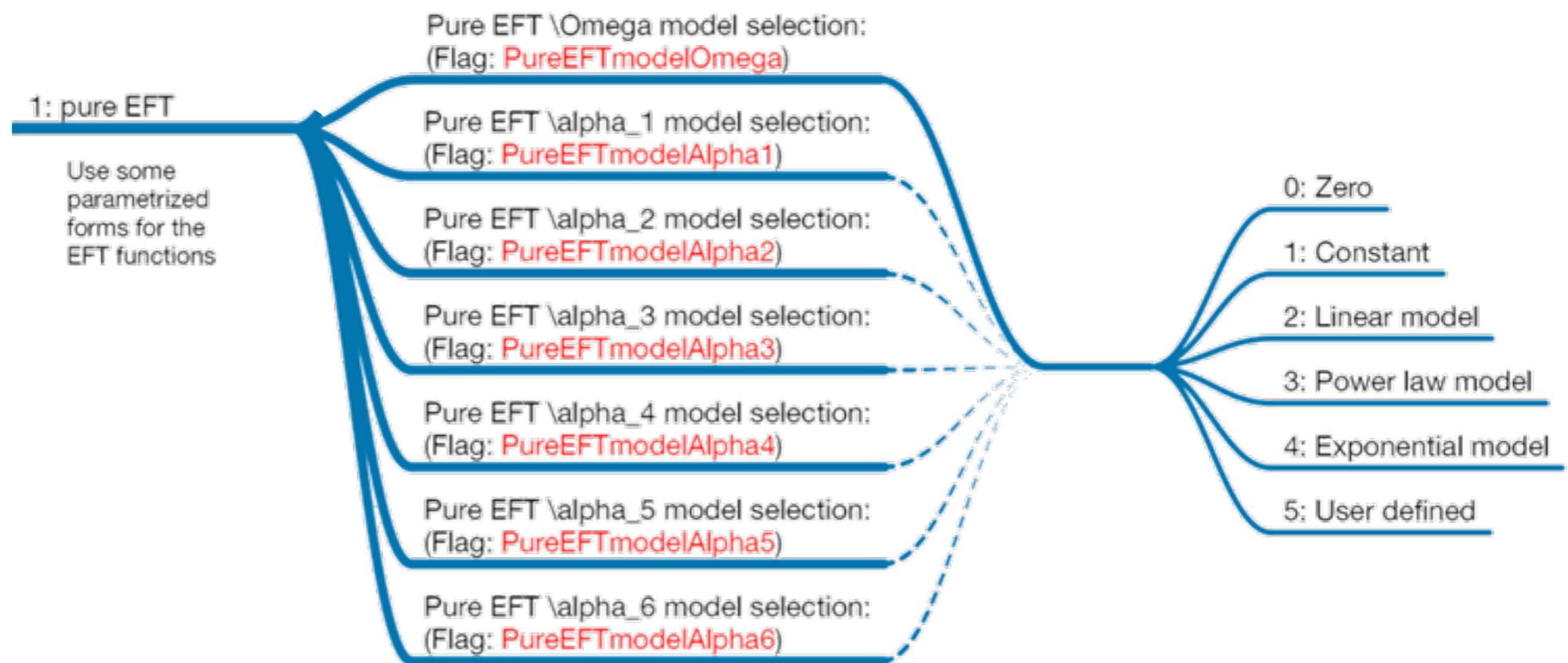
Phenomenological parametrization, e.g.

Constant models: $\Omega(a) = \Omega_0$;

Linear models: $\Omega(a) = \Omega_0 a$;

Power law models: $\Omega(a) = \Omega_0 a^s$;

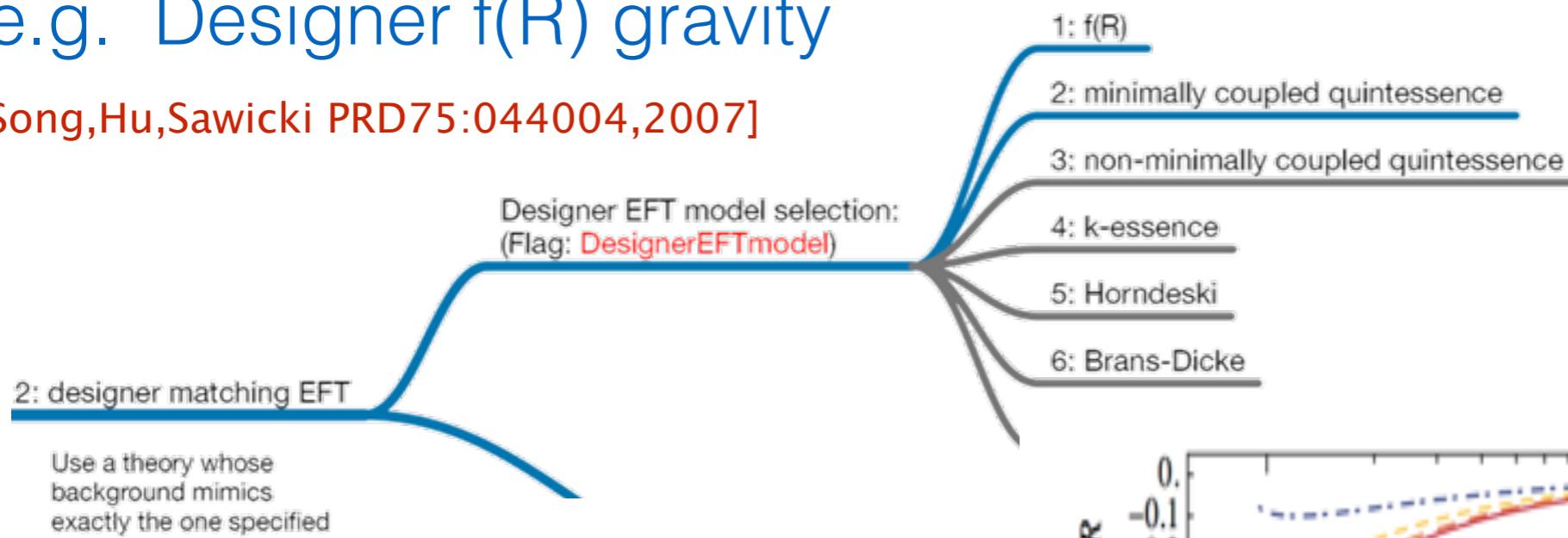
Exponential models: $\Omega(a) = \exp(\Omega_0 a^s) - 1$.



2.2.2 EFT parametrization: Full mapping—designer mapping

e.g. Designer f(R) gravity

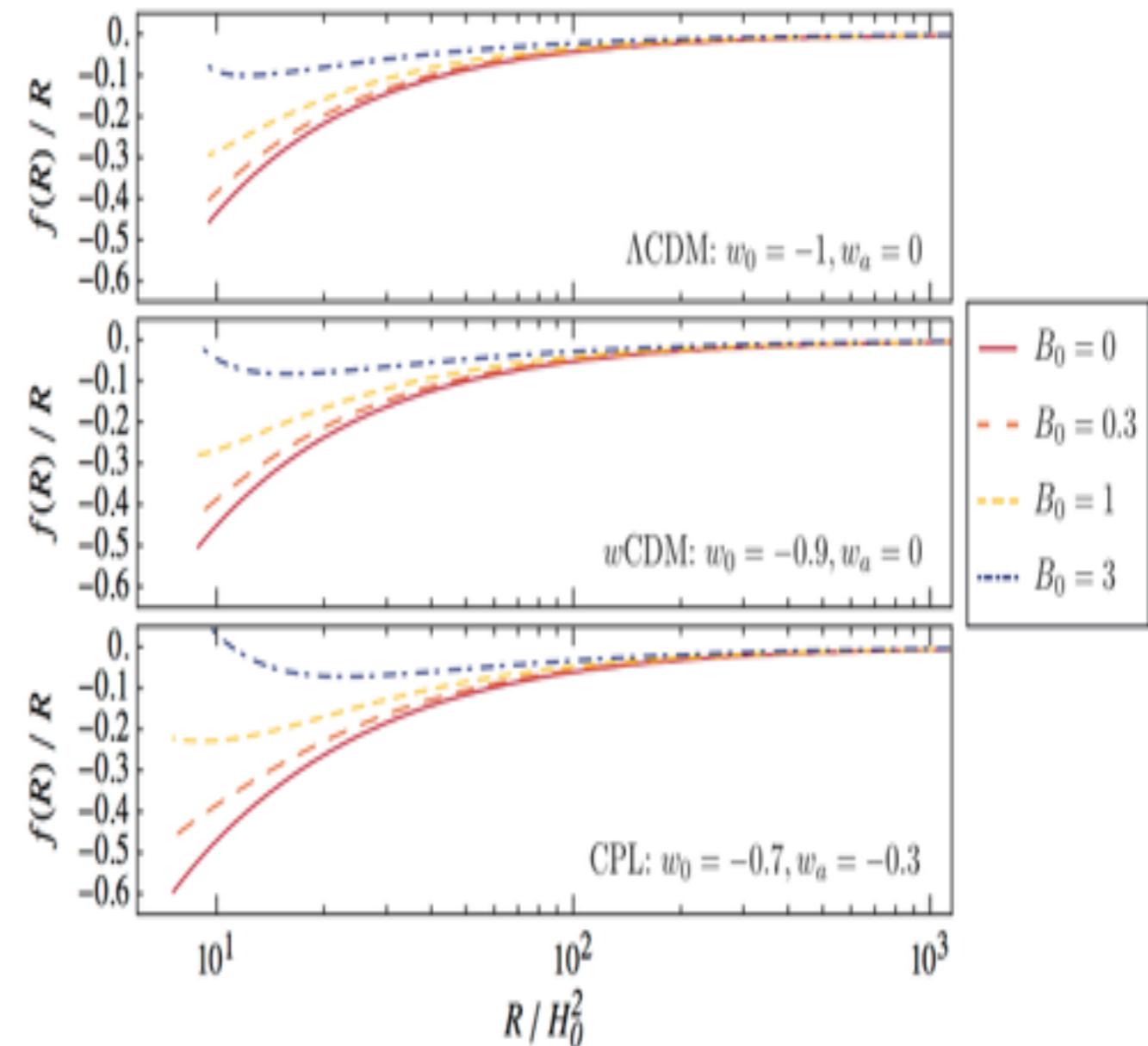
[Song,Hu,Sawicki PRD75:044004,2007]



$$f'' - \left(1 + \frac{H'}{H} + \frac{R''}{R'}\right) f' + \frac{R'}{6H^2} f = -\frac{R'}{3M_P^2 H^2} \rho_{\text{DE}},$$

$$B_0 \sim \frac{6f_{RR}}{(1+f_R)} H^2|_{a=1}$$

GR limit: $B_0 \rightarrow 0$,
effective mass $\rightarrow \text{Infty}$



2.2.3 EFT parametrization: Full mapping (coming soon v2.0)

A few examples work in progress:

1. Hu-Sawicki $f(R)$ model (with Rizzato et.al.)
2. Horava gravity (with Frusciante,Raveri,Silvestri,Vernieri)
3. K-mouflage

Advantage: the stable parameter regime are fully controlled by the user instead of the code per se!

2.3 Perturbation equations

We implement the pi field into the Einstein-Boltzmann solver

CAMB → **EFTCAMB**

Evolving the full **Einstein** equation, **Klein-Golden** equation (pi field), **fluid equation** (CDM,baryon, massive neutrino), **Boltzmann hierarchy** equation sets (CMB, massless neutrino)

- EFT: **DO NOT rely on QS approx!**

time-time Einstein equation:

$$k^2\eta = -\frac{a^2}{2m_0^2(1+\Omega)} [\delta\rho_m + \dot{\rho}_Q\pi + 2c(\dot{\pi} + \mathcal{H}\pi)] + \left(\mathcal{H} + \frac{\dot{\Omega}}{2(1+\Omega)}\right) k\mathcal{Z} + \frac{\dot{\Omega}}{2(1+\Omega)} [3(3\mathcal{H}^2 - \dot{\mathcal{H}})\pi + 3\mathcal{H}\dot{\pi} + k^2\pi]$$

momentum Einstein equation:

$$\frac{2}{3}k^2(\sigma_* - \mathcal{Z}) = \frac{a^2}{m_0^2(1+\Omega)} [(\rho_m + P_m)v_m + (\rho_Q + P_Q)k\pi] + k\frac{\dot{\Omega}}{(1+\Omega)} (\dot{\pi} + \mathcal{H}\pi),$$

space-space off-diagonal Einstein equation:

$$k\dot{\sigma}_* + 2k\mathcal{H}\sigma_* - k^2\eta = -\frac{a^2P\Pi_m}{m_0^2(1+\Omega)} - \frac{\dot{\Omega}}{(1+\Omega)} (k\sigma_* + k^2\pi),$$

space-space trace Einstein equation:

$$\begin{aligned} \ddot{h} = & -\frac{3a^2}{m_0^2(1+\Omega)} [\delta P_m + \dot{P}_Q\pi + (\rho_Q + P_Q)(\dot{\pi} + \mathcal{H}\pi)] - 2\left(\frac{\dot{\Omega}}{1+\Omega} + 2\mathcal{H}\right) k\mathcal{Z} + 2k^2\eta \\ & - 3\frac{\dot{\Omega}}{(1+\Omega)} \left[\ddot{\pi} + \left(\frac{\ddot{\Omega}}{\dot{\Omega}} + 3\mathcal{H}\right) \dot{\pi} + \left(\mathcal{H}\frac{\ddot{\Omega}}{\dot{\Omega}} + 5\mathcal{H}^2 + \dot{\mathcal{H}} + \frac{2}{3}k^2\right) \pi \right], \end{aligned}$$

- For Klein-Golden Eq. Of π field

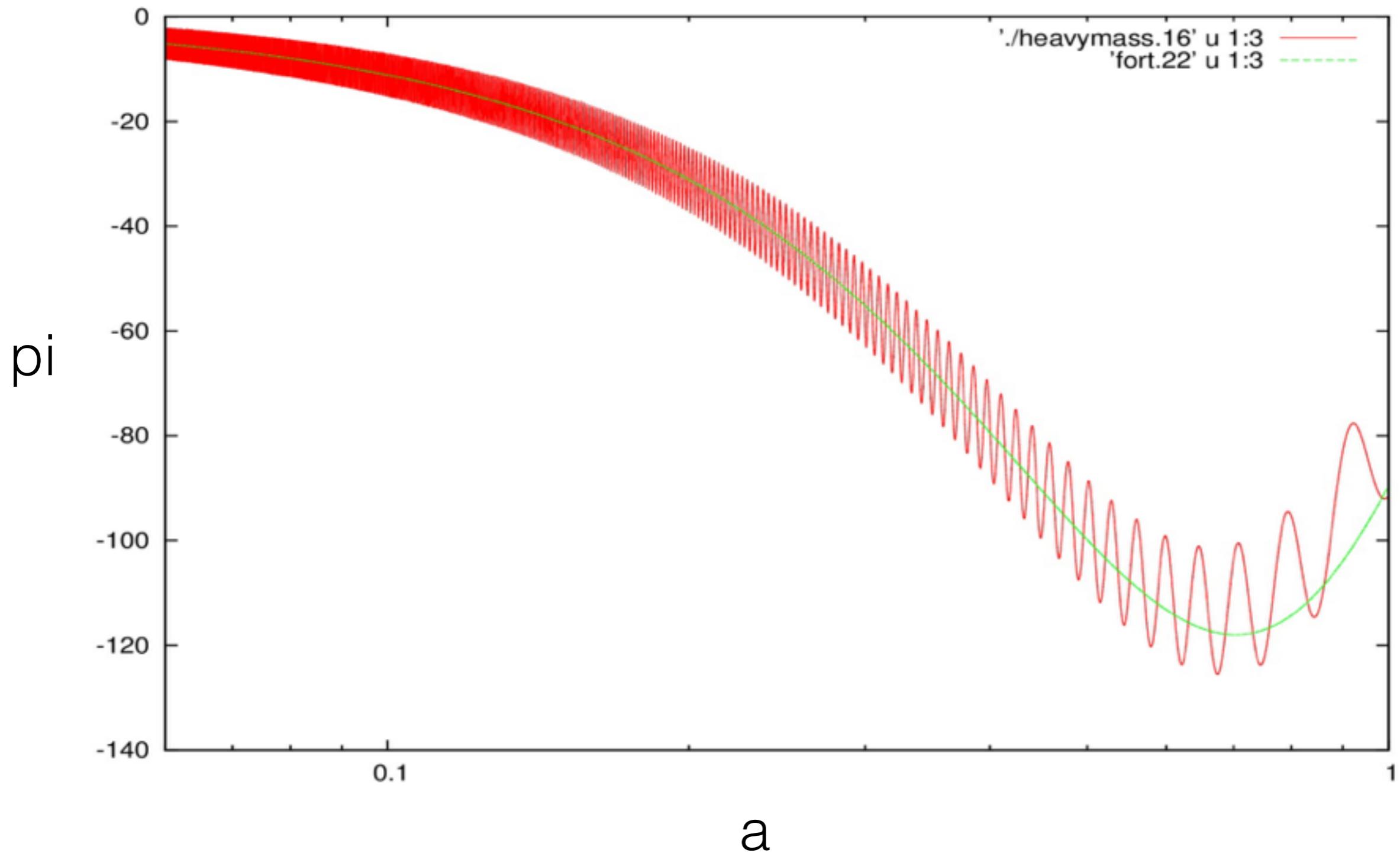
$$\begin{aligned}
 & \left(c + \frac{3m_0^2}{4a^2} \frac{\dot{\Omega}^2}{(1+\Omega)} \right) \ddot{\pi} + \left[\frac{3m_0^2}{4a^2} \frac{\dot{\Omega}}{(1+\Omega)} \left(\ddot{\Omega} + 4\mathcal{H}\dot{\Omega} + \frac{(\rho_Q + P_Q)a^2}{m_0^2} \right) + \dot{c} + 4\mathcal{H}c - \frac{\dot{\Omega}}{2(1+\Omega)}c \right] \dot{\pi} \\
 & + \left[\frac{3}{4} \frac{m_0^2}{a^2} \frac{\dot{\Omega}}{(1+\Omega)} \left(\frac{(3\dot{P}_Q - \dot{\rho}_Q + 3\mathcal{H}(\rho_Q + P_Q))a^2}{3m_0^2} + \mathcal{H}\ddot{\Omega} + 8\mathcal{H}^2\dot{\Omega} + 2(1+\Omega)(\ddot{\mathcal{H}} - 2\mathcal{H}^3) \right) \right. \\
 & \quad \left. - 2\dot{\mathcal{H}}c + \left(\dot{c} - \frac{\dot{\Omega}}{2(1+\Omega)}c \right) \mathcal{H} + 6\mathcal{H}^2c + \left(c + \frac{3m_0^2}{4a^2} \frac{\dot{\Omega}^2}{(1+\Omega)} \right) k^2 \right] \pi \\
 & + \left[c + \frac{3}{4} \frac{m_0^2}{a^2} \frac{\dot{\Omega}^2}{(1+\Omega)} \right] k\mathcal{Z} + \frac{1}{4} \frac{\dot{\Omega}}{(1+\Omega)} (3\delta P_m - \delta\rho_m) = 0,
 \end{aligned}$$

kinetic	friction	mass	sound speed	source
$A(\tau)$	$\ddot{\pi}$	$B(\tau)$	$\dot{\pi}$	$C(\tau)$
			π	$+ k^2 D(\tau) \pi + E(\tau) = 0$

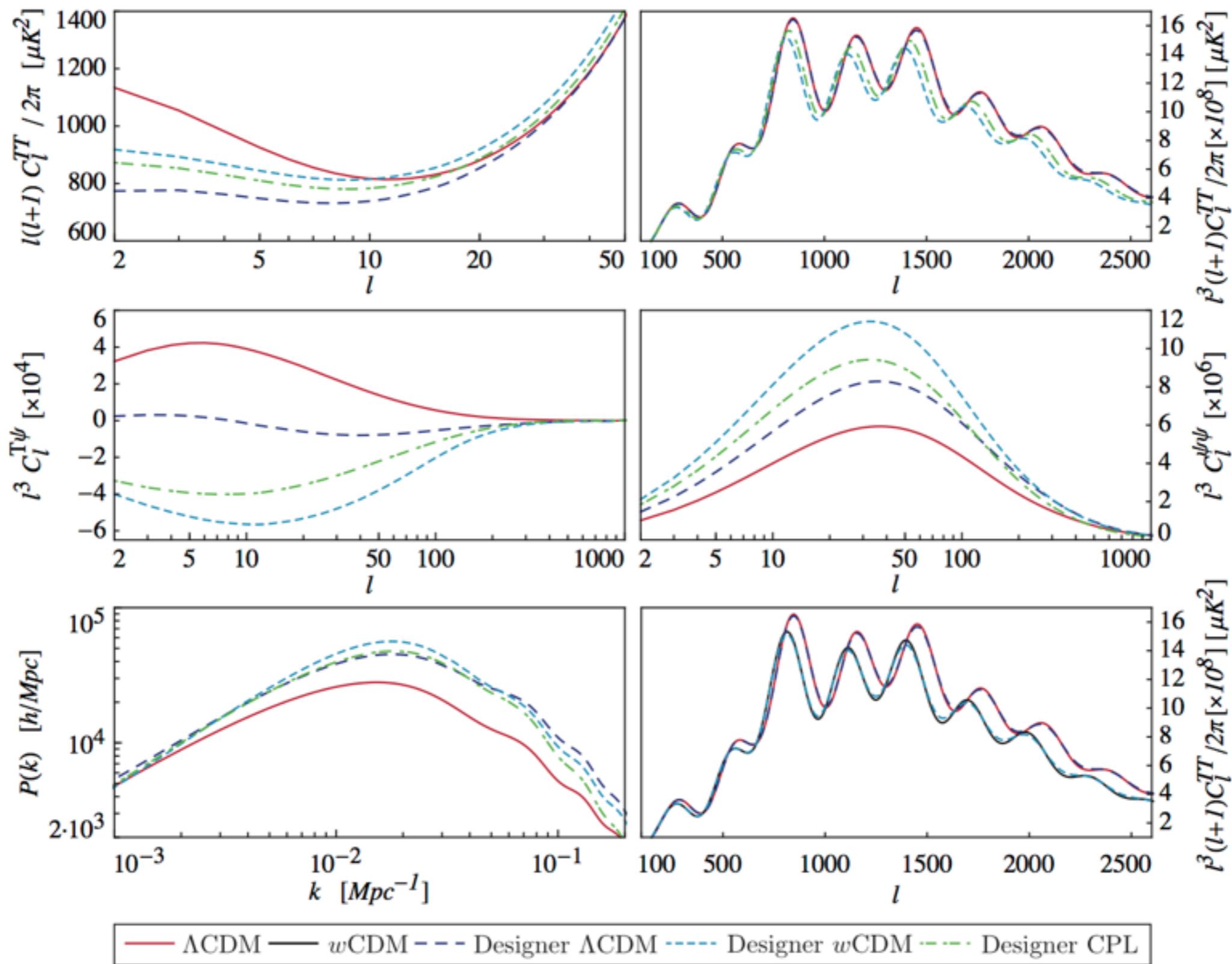
Have pass the viability condition:

1. Effective Newton constant does not change sign: $1+\Omega>0$
2. ghost instability: $A>0$
3. sound speed ≤ 1 : $D/A \leq 1$
4. mass square ≥ 0 : $C/A \geq 0$

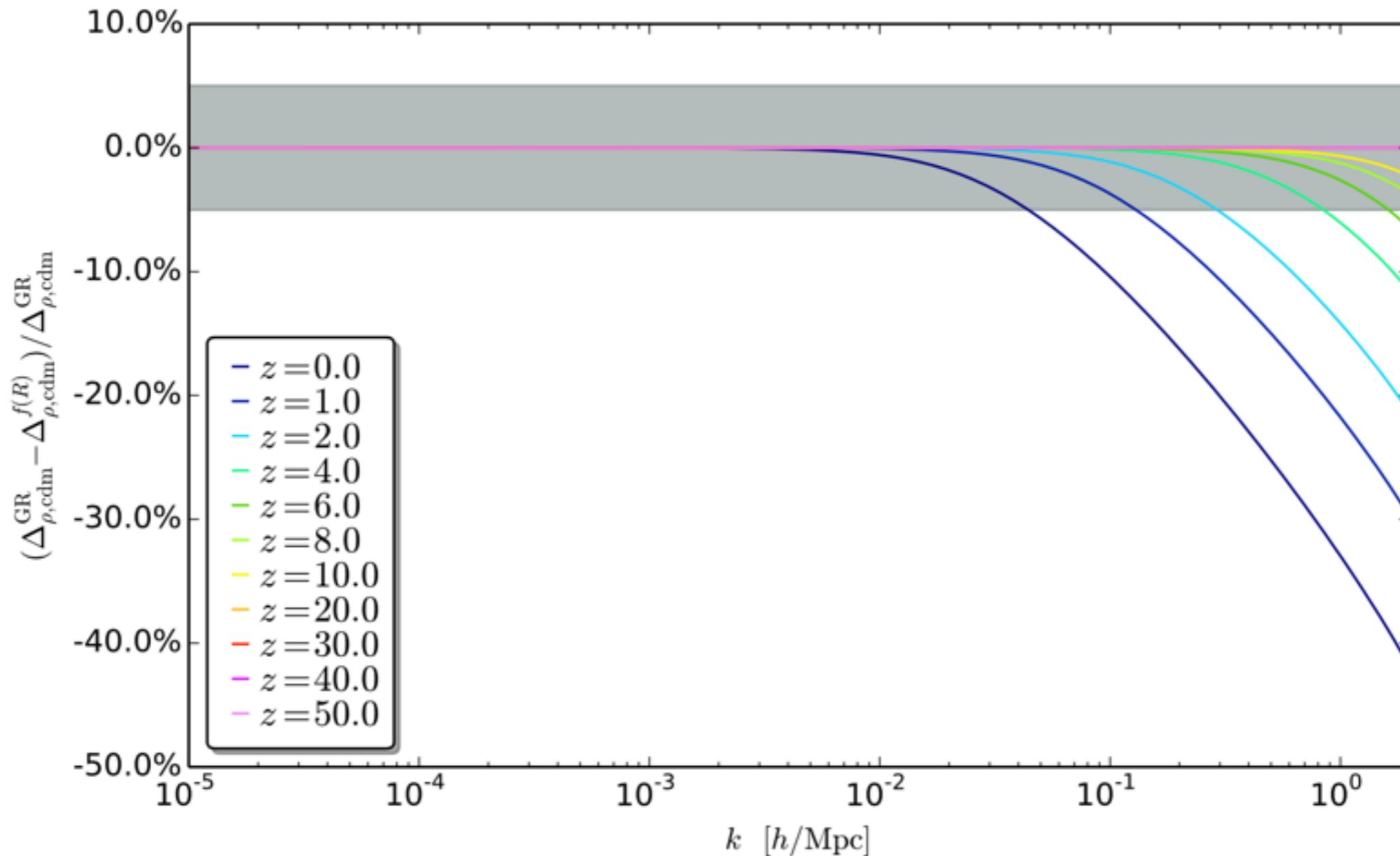
pi field solution: f(R) example



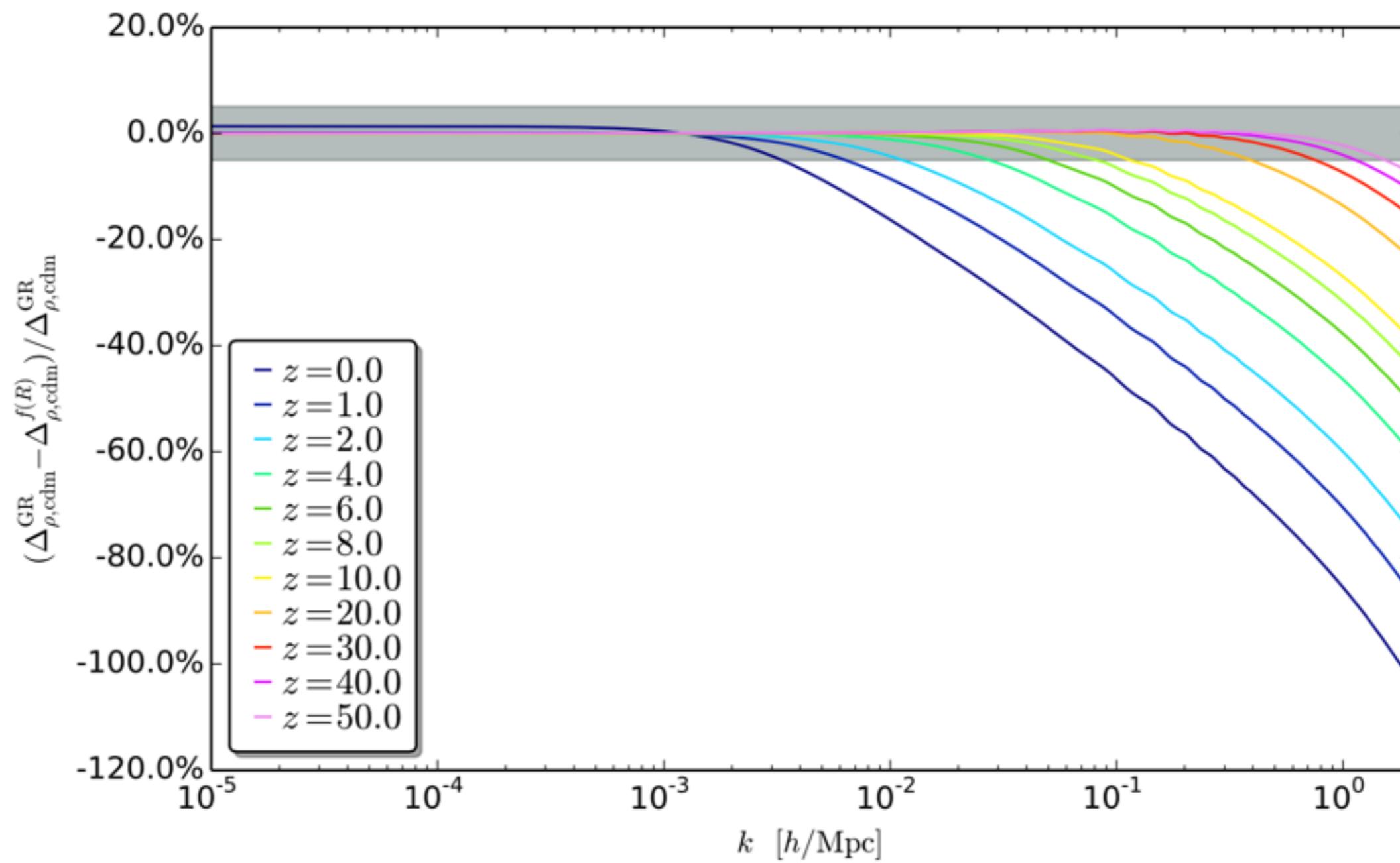
2.4 CMB spectra—example: f(R)



2.5 Transfer function of CDM



Designer $f(R)$ with LCDM background
B0=0.001

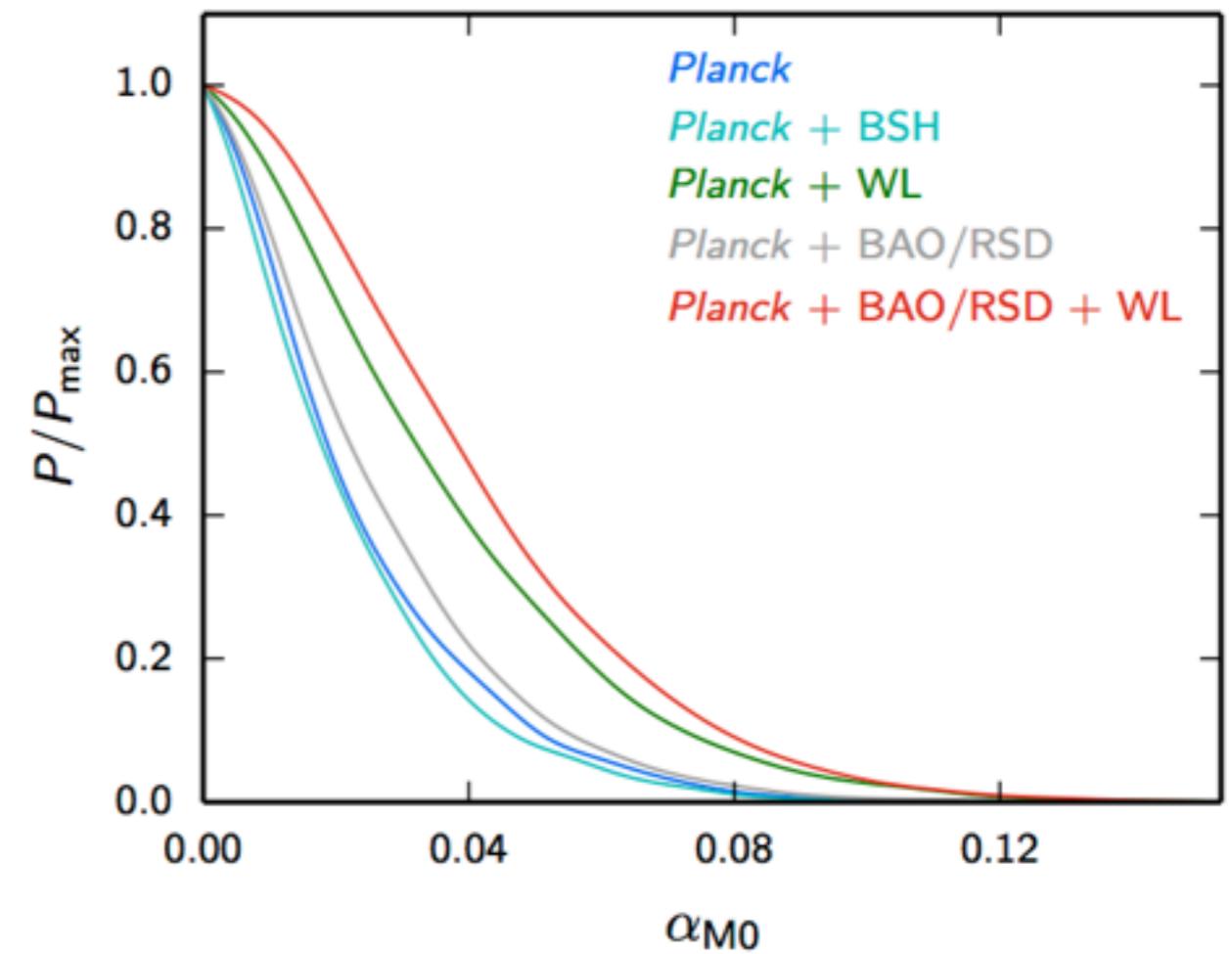
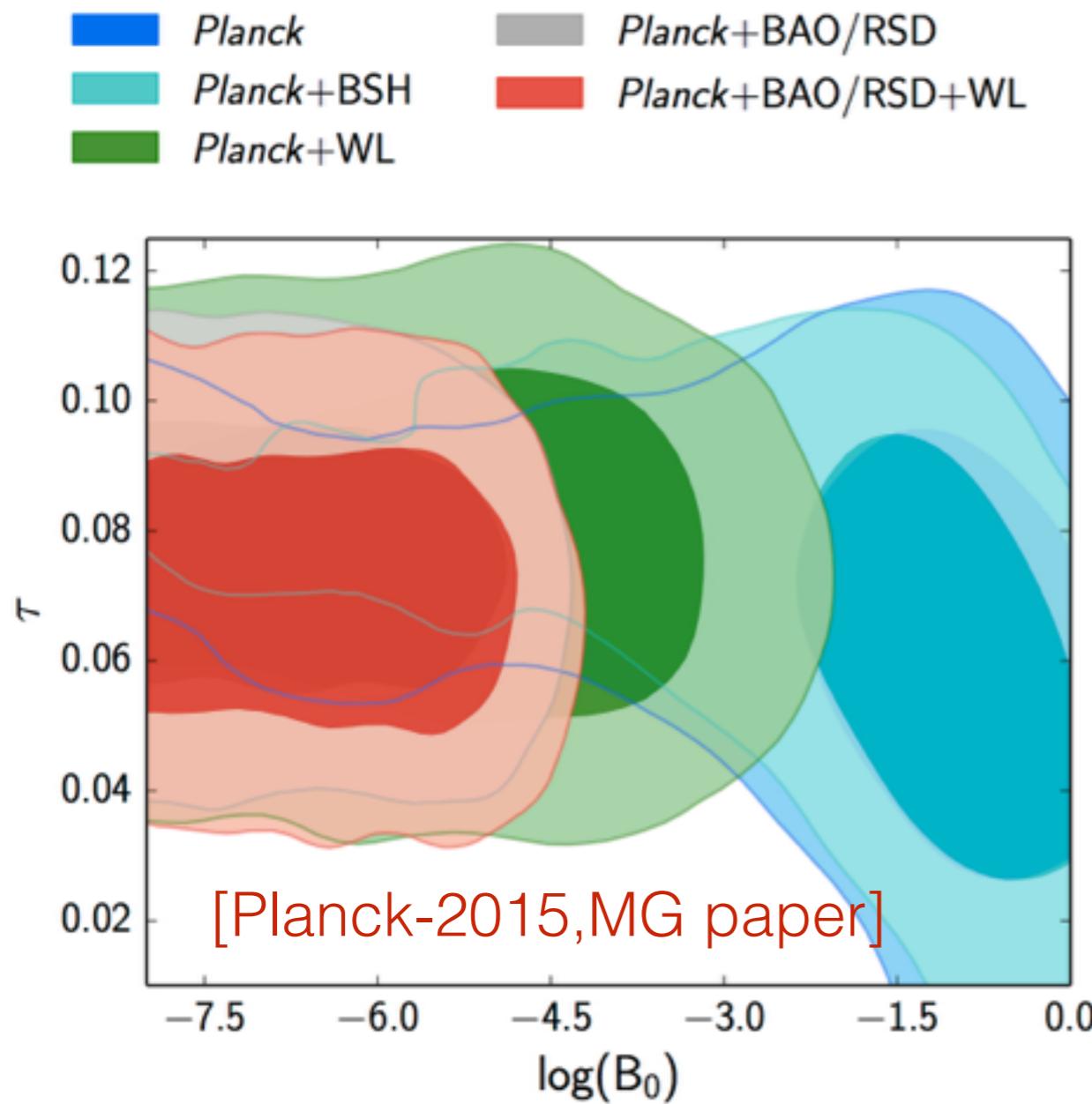


Designer $f(R)$ with wCDM background
 $B0=0.01$ and $w=-0.95$

3. Parameter estimation results from EFTCosmoMC and Planck-2015

CosmoMC → EFTCosmoMC

Designer $f(R)$

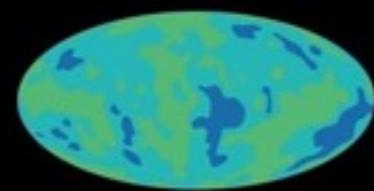


Linear EFT

[Thank Sabino push this
inside the collaboration!]

4. Conclusion

- EFTCAMB include most of viable **single field DE/MG model**
- For scalar field: full perturbative treatment, does not rely on quasistatic approx
- Support LCDM/wCDM/CPL background
- Automatical stability check for given parameterization
- Selected by Planck 2015 data release
- Selected by Theory Working Group of Euclid



KEEP
CALM
AND
TEST
GRAVITY

the EFTCAMB team

Thank you!