

## Tutorial n° 1.

### 2. Linear algebra

$$A. 1. (A^\dagger)^\dagger = \overline{(\overline{A}^\dagger)}^\dagger = \overline{\overline{(A^\dagger)^\dagger}} = \overline{\overline{A}} = A$$

$$2. (AB)^\dagger = \overline{(AB)^\dagger} = \overline{(\overline{B}^\dagger A^\dagger)} = \overline{B^\dagger} \overline{A^\dagger} = B^\dagger A^\dagger.$$

identique pour  $(Ax)^\dagger = x^\dagger A^\dagger$ .

$$3. \langle A^\dagger u, v \rangle = (A^\dagger u)^\dagger v = u^\dagger A^{\dagger\dagger} v = u^\dagger Av = \langle u, Av \rangle.$$

B. 1. Si  $A$  est hermitienne,  $A A^\dagger = A A - A^\dagger A$  donc  $A$  normale.  
 Si  $A$  est unitaire,  $A A^\dagger = A A^{-1} = 1 = A^{-1} A = A^\dagger A$  donc  $A$  normale.

$$2. (UV)^\dagger = V^\dagger U^\dagger = V^{-1} U^{-1} = (UV)^{-1} \text{ donc } UV \text{ est unitaire.}$$

$$3. (G+H)^\dagger = \overline{G+H}^\dagger = (\overline{G}+\overline{H})^\dagger = \overline{G}^\dagger + \overline{H}^\dagger = G^\dagger + H^\dagger = G+H$$

donc  $G+H$  est hermitienne

$$4. (v v^\dagger)^2 = \underbrace{v v^\dagger}_{\text{car } v \text{ est unitaire}} v v^\dagger = \langle v, v \rangle v v^\dagger = \|v\|^2 v v^\dagger = v v^\dagger$$

$(v v^\dagger)^\dagger = v^\dagger v^\dagger = v v^\dagger$  donc  $v v^\dagger$  est bien une matrice de projection.

Soit  $\lambda \in \mathbb{C}$  et  $u$  un vecteur.

$$\text{Gm a: } P^2 u = P u.$$

$$\begin{aligned} P u &= \lambda u \implies P(P u) = \underbrace{P(\lambda u)}_{\lambda^2 u} = \lambda u \\ &\implies \lambda^2 = \lambda \end{aligned}$$

D'où  $\lambda = 0$  ou  $\lambda = 1$ .

### 3. Quantum random access code.

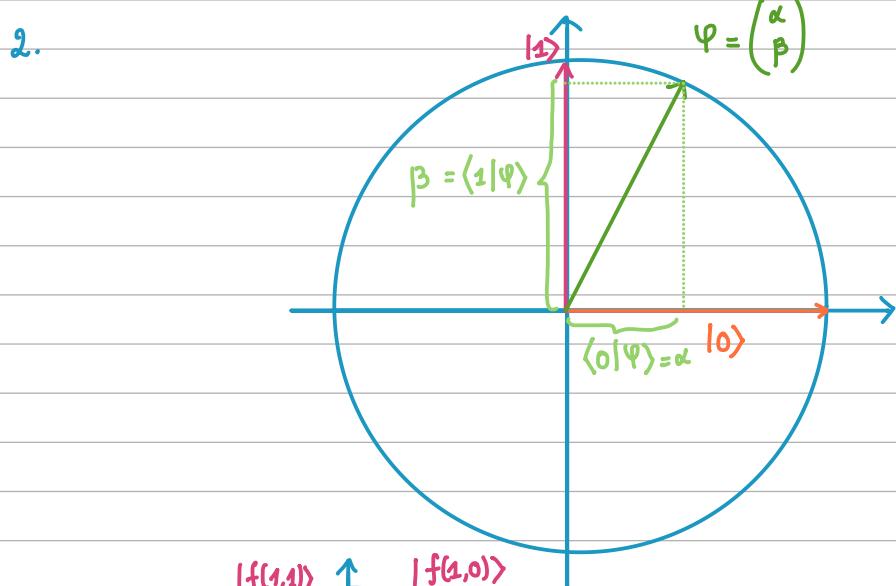
1. Gm a  $f: \{0,1\}^2 \rightarrow \{0,1\}$  donc on a nécessairement une collision.  
 Sans perdre en généralité, supposons  $f(0,x) = f(1,x)$ .

D'où, pour obtenir le 1<sup>er</sup> bit, on a :

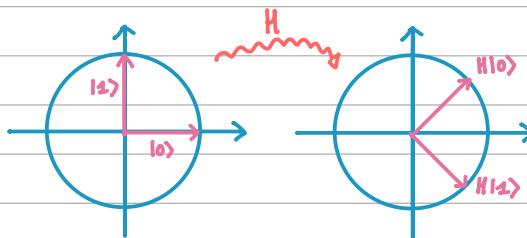
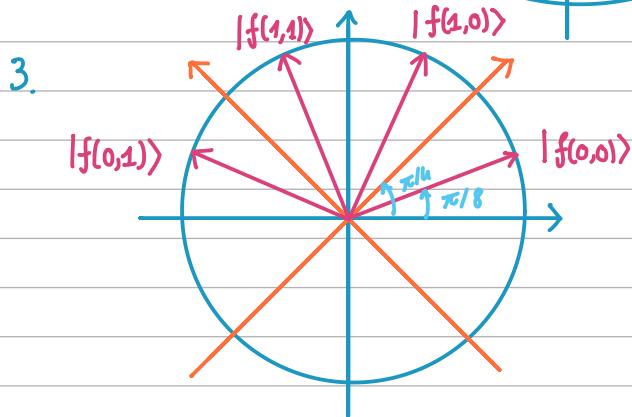
$$f(0,x) \underset{q_0}{\sim} 0 \quad \text{et} \quad f(1,x) \underset{q_1}{\sim} 1.$$

Par conséquent,  $|q_0 + q_1| = 1$  et  $|q_0| \geq p_1, |q_1| \geq p_2$ ,

d'où  $p \leq \frac{1}{2}$ .



Appliquer une matrice unitaire correspond à une rotation ou une symétrie de ce cercle unitaire (et du vecteur  $|\psi\rangle$  dedans).



C'est une symétrie sur l'axe  $x$  puis une rotation de  $45^\circ$ .

$$U_1 = \mathbb{1} \quad \text{et} \quad U_2 = R_{\pi/4}.$$

La probabilité de succès est de  $\cos^2(\pi/8) \approx 0,86$ .

#### 4. Tensor products

1.

$$A \otimes B = \begin{pmatrix} 0 & e^{2i\pi/3} & 0 & e^{i\pi/3} \\ e^{-2i\pi/3} & 0 & -1 & 0 \\ 0 & -1 & 0 & e^{-i\pi/3} \\ e^{i\pi/3} & 0 & e^{i\pi/3} & 0 \end{pmatrix}; \quad B \otimes A = \begin{pmatrix} 0 & 0 & e^{2i\pi/3} & e^{i\pi/3} \\ 0 & 0 & -1 & e^{-i\pi/3} \\ e^{-2i\pi/3} & -1 & 0 & 0 \\ e^{-i\pi/3} & 0 & e^{i\pi/3} & 0 \end{pmatrix}$$

$$2.$$

a)  $A \otimes (\lambda B) = \begin{pmatrix} a_{11}(\lambda B) & a_{12}(\lambda B) & \dots \\ a_{12}(\lambda B) & a_{22}(\lambda B) & \dots \\ \vdots & \vdots & \ddots \end{pmatrix} = \begin{pmatrix} (\bar{a}_{11}\lambda)B & (\bar{a}_{12}\lambda)B & \dots \\ (\bar{a}_{12}\lambda)B & (\bar{a}_{22}\lambda)B & \dots \\ \vdots & \vdots & \ddots \end{pmatrix} = \lambda A \otimes B$

d)  $(A \otimes B)^+ = \begin{pmatrix} a_{11}B & a_{21}B & \dots \\ a_{12}B & a_{22}B & \dots \\ \vdots & \vdots & \ddots \end{pmatrix}^+ = \begin{pmatrix} \bar{a}_{11}B^+ & \bar{a}_{12}B^+ & \dots \\ \bar{a}_{21}B^+ & \bar{a}_{22}B^+ & \dots \\ \vdots & \vdots & \ddots \end{pmatrix}$

$$= A^+ \otimes B^+$$

## Tutorial 2

### II. Warm-up calculations.

#### 1) Measurements and probabilities

$$P(|\Psi\rangle \text{ is measured as } |b\rangle) = |\langle b|\Psi\rangle|^2$$

Q1.  $\| |\Psi\rangle \| = 1$  Thus  $|\Psi\rangle$  is a state

measuring  $|\Psi\rangle$  leads to 0 with probability 1/2

measuring  $|\Psi\rangle$  leads to 1 with probability 1/2

Q2.  $\| |\Psi\rangle \| = 1$  Thus  $|\Psi\rangle$  is a state

measuring  $|\Psi\rangle$  leads to 0 with probability 1/2

measuring  $|\Psi\rangle$  leads to 1 with probability 1/2

Q3.  $\| |\Psi\rangle \| = 1$  Thus  $|\Psi\rangle$  is a state

measuring  $|\Psi\rangle$  leads to 0 with probability 1

measuring  $|\Psi\rangle$  leads to 1 with probability 0

Q4.  $\| |\Psi\rangle \| = 1$  Thus  $|\Psi\rangle$  is a state

measuring  $|\Psi\rangle$  leads to 0 with probability 1/2

measuring  $|\Psi\rangle$  leads to 1 with probability 1/2

Q5.  $\| |\Psi\rangle \| = 1$  Thus  $|\Psi\rangle$  is a state

measuring  $|\Psi\rangle$  leads to 00 with probability 1/2

measuring  $|\Psi\rangle$  leads to 01 with probability 0

measuring  $|\Psi\rangle$  leads to 10 with probability 0

measuring  $|\Psi\rangle$  leads to 11 with probability 1/2

Q6.  $\| |\Psi\rangle \| = 1$  Thus  $|\Psi\rangle$  is a state

measuring  $|\Psi\rangle$  leads to 00 with probability 1/2

measuring  $|\Psi\rangle$  leads to 01 with probability 0

measuring  $|\Psi\rangle$  leads to 10 with probability 0

measuring  $|\Psi\rangle$  leads to 11 with probability 1/2

Q7.  $\| |\Psi\rangle \| = 1$  Thus  $|\Psi\rangle$  is a state

measuring  $|\Psi\rangle$  leads to 00 with probability 1/4

measuring  $|\Psi\rangle$  leads to 01 with probability 1/2

measuring  $|\Psi\rangle$  leads to 10 with probability 0

measuring  $|\Psi\rangle$  leads to 11 with probability 1/4

Q8.  $\| |\Psi\rangle \| = 1$  Thus  $|\Psi\rangle$  is a state  
 measuring  $|\Psi\rangle$  leads to  $00$  with probability  $1/4$   
 measuring  $|\Psi\rangle$  leads to  $01$  with probability  $1/2$   
 measuring  $|\Psi\rangle$  leads to  $10$  with probability  $1/8$   
 measuring  $|\Psi\rangle$  leads to  $11$  with probability  $1/8$

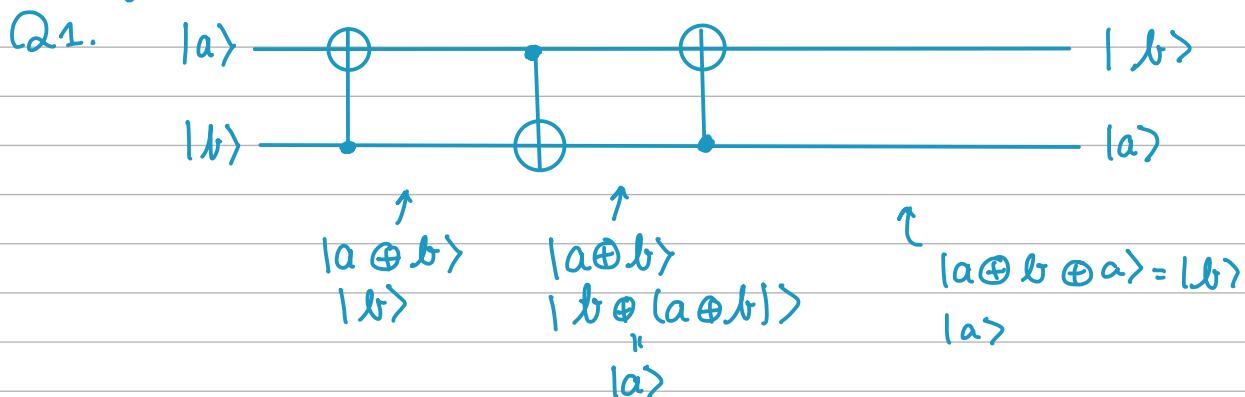
## 2) Partial measurements.

$$\begin{aligned}\sqrt{2} \langle 1 | \Psi \rangle &= \langle 1 | 0 \rangle \cdot \langle 1 | \Psi \rangle + \langle 1 | 1 \rangle \langle 1 | \Psi \rangle \\ &= \langle 1 | \Psi \rangle \\ &= \end{aligned}$$

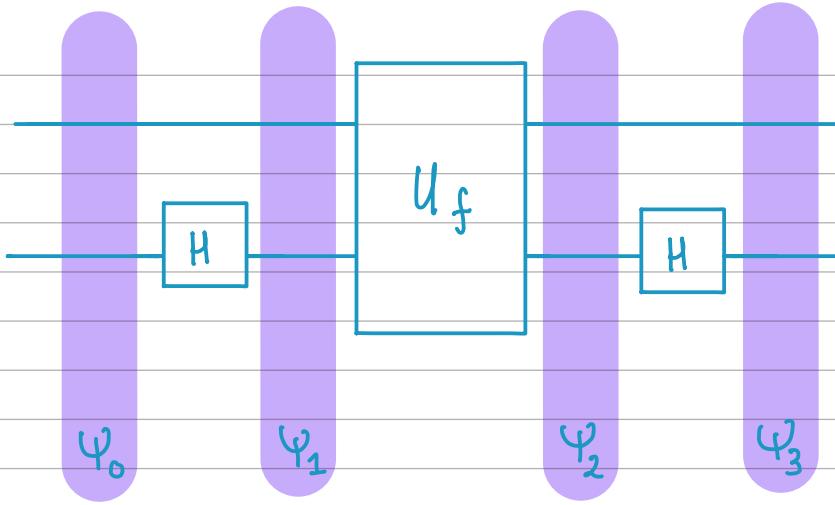
$$|\Psi'\rangle = H \otimes I |\Psi\rangle = \frac{1}{\sqrt{2}} (|+\rangle |\Psi\rangle + |- \rangle |\Psi\rangle) = \frac{1}{2} (|0\rangle (|\Psi\rangle + |\Psi\rangle) + |1\rangle (|\Psi\rangle - |\Psi\rangle))$$

$$\begin{aligned}\text{We have that } P(\text{measuring } |1\rangle) &= \| |\Psi\rangle - |\Psi'\rangle \|_2^2 \times \frac{1}{4} \\ &= (\langle \Psi | - \langle \Psi' |)(|\Psi\rangle - |\Psi'\rangle) \times \frac{1}{4} \\ &= \frac{1}{4} (\langle \Psi | \Psi \rangle + \langle \Psi' | \Psi \rangle - \langle \Psi | \Psi' \rangle - \langle \Psi' | \Psi' \rangle) \\ &= \frac{1}{4} (2 - \langle \Psi | \Psi \rangle - \langle \Psi' | \Psi' \rangle) \\ &= \frac{1}{2} (1 - \text{Re}(\langle \Psi | \Psi' \rangle)).\end{aligned}$$

## 3) Gates



Q2.



$$|\Psi_0\rangle = |a\rangle |b\rangle$$

$$|\Psi_1\rangle = (\mathbb{1} \otimes H) |a\rangle |b\rangle = \frac{1}{\sqrt{2}} |a\rangle (|0\rangle + (-1)^b |1\rangle)$$

$$\begin{aligned} |\Psi_2\rangle &= \frac{1}{\sqrt{2}} U_f |a\rangle |0\rangle + \frac{(-1)^b}{\sqrt{2}} U_f |a\rangle |1\rangle \\ &= \frac{1}{\sqrt{2}} |a\rangle |0\rangle |f(a)\rangle + \frac{(-1)^b}{\sqrt{2}} |a\rangle |1\rangle |f(a)\rangle \\ &= \frac{1}{\sqrt{2}} |a\rangle (|f(a)\rangle + (-1)^b |\overline{f(a)}\rangle) \end{aligned}$$

$$|\Psi_3\rangle = (\mathbb{1} \otimes H) |\Psi_2\rangle = \frac{1}{2} |a\rangle \left( (|0\rangle + (-1)^{\frac{f(a)}{2}} |0\rangle) + (-1)^b (|0\rangle + (-1)^{\frac{f(a)}{2}} |\overline{f(a)}\rangle) \right)$$

Par étude des cas  $b=0$  et  $b=1$ , on a bien que

$$|\Psi_3\rangle = |a\rangle (-1)^{\frac{f(a)+b}{2}} |b\rangle.$$

### III. Superdense coding.

Q1. Two bits

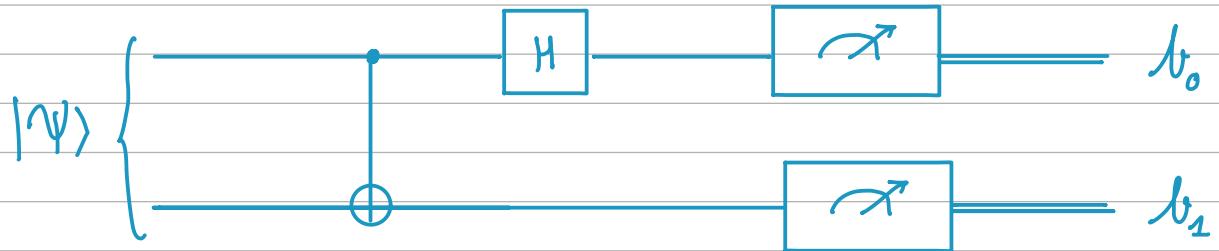
$$\begin{array}{ll} \text{Q2. } 00 \text{ mso } & |00\rangle + |11\rangle \rightarrow |00\rangle + |11\rangle \rightarrow |00\rangle + |11\rangle \\ 01 \text{ mso } & |00\rangle + |11\rangle \rightarrow |10\rangle + |01\rangle \rightarrow |10\rangle + |01\rangle \\ 10 \text{ mso } & |00\rangle + |11\rangle \rightarrow |00\rangle + |11\rangle \rightarrow |00\rangle - |11\rangle \\ 11 \text{ mso } & |00\rangle + |11\rangle \rightarrow |10\rangle + |01\rangle \rightarrow -|10\rangle + |01\rangle \end{array}$$

Tous ces états sont orthogonaux; ils forment une base

$$\mathcal{B} = \left\{ \frac{|00\rangle + |11\rangle}{\sqrt{2}}, \frac{|10\rangle + |01\rangle}{\sqrt{2}}, \frac{|00\rangle - |11\rangle}{\sqrt{2}}, \frac{|01\rangle - |10\rangle}{\sqrt{2}} \right\}.$$

Il suffit d'appliquer une matrice de passage de  $\mathcal{B}$

dans  $\mathcal{B}_{\text{base computationnelle}} = \{|00\rangle, |01\rangle, |10\rangle, |11\rangle\}$ .



## Tutorial #3.

### 2. Some properties of circuits

#### 2.1. Do circuits commute?

Q1.  $A = \frac{1}{\sqrt{2}} \otimes A'$  and  $B = B' \otimes \frac{1}{\sqrt{2}}$  commute:  $AB = BA = B' \otimes A$

Q2.  $A = \underbrace{\begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}}_X \otimes \frac{1}{\sqrt{2}}$  and  $B = H \otimes \frac{1}{\sqrt{2}}$  do not commute:  $AB \neq BA$ .

#### 2.2. Are circuits ambiguous?

Q1. This is exactly 2.1/Q1.

Q2. Let  $|\Phi\rangle = \alpha|100\rangle + \beta|110\rangle + \gamma|101\rangle + \delta|111\rangle$ .

When measuring the first qubit and then the second, we obtain that:

$$\Pr[1^{\text{st}} \text{ qubit is measured as } 0] = |\alpha|^2 + |\gamma|^2$$

$$\begin{aligned} \Pr[\text{measuring } 00] &= \Pr[1^{\text{st}} \text{ qubit } \sim 0] \times \Pr[2^{\text{nd}} \text{ qubit } \sim 0 | 1^{\text{st}} \text{ qubit } \sim 0]. \\ &= (|\alpha|^2 + |\gamma|^2) \times \frac{|\alpha|^2}{|\alpha|^2 + |\gamma|^2} = |\alpha|^2 \end{aligned}$$

We can do the same for all the other cases.

When measuring the 2<sup>nd</sup> qubit and then the 1<sup>st</sup> qubit, we have:

$$\Pr[2^{\text{nd}} \text{ qubit } \sim 0] = |\alpha|^2 + |\beta|^2$$

$$\begin{aligned} \text{and } \Pr[\text{measuring } 00] &= \Pr[2^{\text{nd}} \text{ qubit } \sim 0] \times \Pr[1^{\text{st}} \text{ qubit } \sim 0 | 2^{\text{nd}} \text{ qubit } \sim 0] \\ &= (|\alpha|^2 + |\beta|^2) \times \left( \frac{|\alpha|^2}{|\alpha|^2 + |\beta|^2} \right) \\ &= |\alpha|^2. \end{aligned}$$

Q3. Let  $|\Phi\rangle = \alpha |\Psi_0\rangle \otimes |0\rangle + \beta |\Psi_1\rangle \otimes |1\rangle$ .

We have that:

$$\Pr[|\Phi\rangle \xrightarrow{\text{act}} 0] = |\alpha|^2 + \Pr[|\Phi\rangle \xrightarrow{\text{act}} 1] = |\beta|^2$$

$$(\mathbf{U} \otimes \mathbf{I})|\Phi\rangle = \alpha (\mathbf{U}|\Psi_0\rangle) \otimes |0\rangle + \beta (\mathbf{U}|\Psi_1\rangle) \otimes |1\rangle.$$

$$\Pr[(\mathbf{U} \otimes \mathbf{I})|\Phi\rangle \xrightarrow{\text{act}} 0] = |\alpha|^2 + \Pr[(\mathbf{U} \otimes \mathbf{I})|\Phi\rangle \xrightarrow{\text{act}} 1] = |\beta|^2$$

### 3. The CHSH Game.

Q1. We have  $A: \{0,1\} \rightarrow \{0,1\}$  and  $B: \{0,1\} \rightarrow \{0,1\}$  two deterministic functions. With  $A(x) = x$  and  $B(y) = y$ , we have a probability of success of  $3/4$ .

We know that the probability of success for any deterministic strategy is in  $\{0, \frac{1}{4}, \frac{1}{2}, \frac{3}{4}, 1\}$ . Suppose that we have a strategy with a success rate  $> \frac{3}{4}$ , i.e. = 1, that is,

$$\forall x, y, \quad A(x) \oplus B(y) = x \wedge y.$$

This means  $A(0) \oplus B(0) = A(1) \oplus B(0) = A(0) \oplus B(1) = 0$

By disjunction, we have no case such that this is true.

We can conclude that  $\frac{3}{4}$  is the optimal success rate for a deterministic strategy.

Q2. (a)

$$\max_{n \in \mathbb{N}} \max_{\lambda \text{ a random on } [n]} \max_{(A_k, B_k)_{k \in [n]}} \sum_{k=0}^n \Pr[\lambda = k] \times \text{Success Rate}(A_k, B_k)$$

(b)  $\text{Success Rate}(\text{Shared Randomness}) \geq \max_{A, B} \text{Success Rate}(A, B) = 3/4$ .

We have that

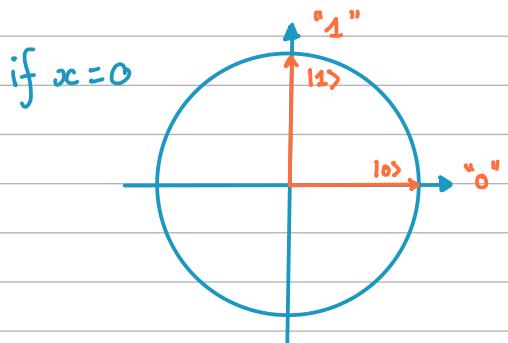
Fix  $n \in \mathbb{N}$ ,  $\lambda$  a random variable over  $[n]$  and  $(A_k, B_k)_{k \in [n]}$ .

$$\sum_{k=0}^n \Pr[\lambda = k] \times \text{Success Rate}(A_k, B_k) \leq \sum_{k=0}^n \Pr[\lambda = k] \cdot \frac{3}{4} = \frac{3}{4} \underbrace{\sum_{k=0}^n \Pr[\lambda = k]}_1 = \frac{3}{4}$$

Thus, Success Rate (Probabilistic Strategy).

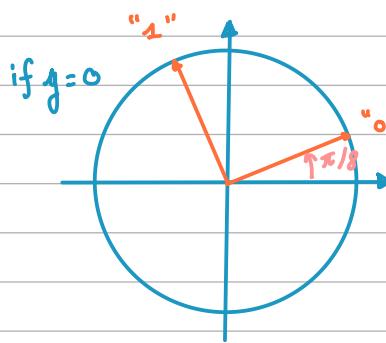
Q3. We make the following measurements :

Alice

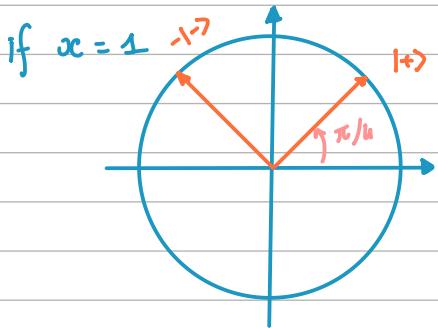


if  $x=0$

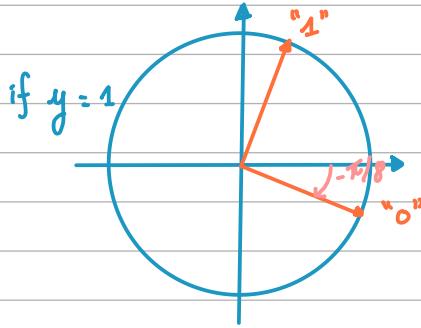
Bob



$$\text{Success Rate} = \cos^2 \frac{\pi}{8} \approx 0.85$$



if  $x=1$



if  $y=1$

# 1D $m \circ 5$

## II From order-finding to factoring

Q1

$x \in [1, l-1]$  invertible  $\Leftrightarrow \exists k \in [1, l-1], xk = 1$

$$\text{Bézout} \Leftrightarrow \gcd(x, l) = 1$$

$$\Leftrightarrow x \in \{1, \dots, l-1\}$$

Q2.  $\mathbb{Z}_l^{\times} = [1, l-1]$  and thus  $|\mathbb{Z}_l^{\times}| = l-1$ .

$$Q2. |\mathbb{Z}_N^{\times}| = |\mathbb{Z}_p^{\times} \times \mathbb{Z}_q^{\times}| = |\mathbb{Z}_p^{\times}| \times |\mathbb{Z}_q^{\times}| = (p-1)(q-1)$$

Q3. Take  $x \in \mathbb{U}([0, N])$ .

$$\begin{aligned} \Pr[x \in \mathbb{Z}_N^{\times}] &= \frac{|\mathbb{Z}_N^{\times}|}{|\mathbb{Z}_N|} = \frac{p-1}{p} \frac{q-1}{q} \\ &= 1 - \frac{1}{q} - \frac{1}{p} + \frac{1}{pq} \sim 1 - 2^{-\frac{1}{2}}. \end{aligned}$$

Q4. Soit  $k$  l'ordre de  $x$ .

$$(x)^{2k} = ((-1)^k x^k)^2 = (-1)^2 = 1.$$

On a donc  $|\{x \in \mathbb{Z}_N^{\times} \mid w(x) \text{ impair}\}| \leq |\{x \in \mathbb{Z}_N^{\times} \mid w(x) \text{ pair}\}|$ .

d'où  $\Pr[x \text{ a un ordre impair}] \geq \frac{1}{2}$ .

$$\begin{aligned} Q5. x^N - 1 &= (x^{N/2} - 1)(x^{N/2} + 1) = N \\ &\quad \overbrace{\hspace{10em}}^{\neq 1, N} \quad \overbrace{\hspace{10em}}^{\neq 1, N} \end{aligned}$$

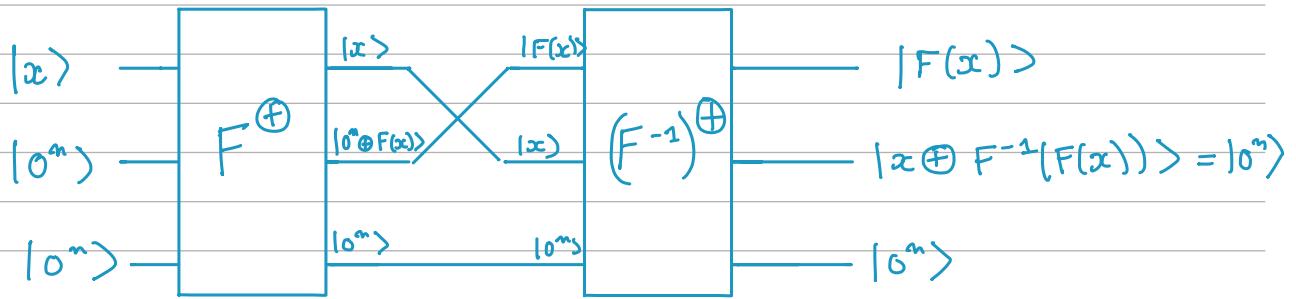
d'où on a un diviseur non-trivial de  $N$ .

Q6. Run Shor while obtaining even order on random  $x$ , return  $\gcd(x, N)$  or repeat.

Expected Run Time  $O\left(\frac{1}{1-\varepsilon} \times 2\right)$  appelle à Shor  
probabiliste pour que  $x^r \pm 1 \neq 1, N$   
↳ proba even-order

### III Modular exponentiation.

Q1.



Q2.  $F(x) := ax \bmod N$        $F^{-1}(x) := a^{-1}x \bmod N$

Q3. Exponentiation rapide.

Si impair

$$| y \leftarrow y \times y \times x \rangle$$

Si pair

$$| y \leftarrow y \times y \rangle$$

## II. Amplitude amplification

$$\text{Q1. } \Pr[\text{output} \in f^{-1}(1)] = p \text{ thus } \Pr[(\text{output}[k] \in f^{-1}(1)) \cap (\forall i < k, \text{output}[i] \notin f^{-1}(1))] \\ = p \cdot (1-p)^{k-1}$$

We have a geometric law  $\mathcal{G}(p)$  so  $E[\# \text{ steps}] = 1/p$ .

Q2. We write  $A$  in a universal set of gates, like  $\{\text{NOT}, H, \text{CNOT}\}$ , and we compute

$$A = U_1 \dots U_m \quad \text{and} \quad A^{-1} = A^\dagger = U_m^\dagger \dots U_1^\dagger.$$

$$\text{Q3. } |U\rangle = \sum_{i=1}^m \alpha_i |i\rangle = \sum_{i \in f^{-1}(0)} \alpha_i |i\rangle + \sum_{i \in f^{-1}(1)} \alpha_i |i\rangle$$

$$\text{Define } |G\rangle = \frac{\sum_{i \in f^{-1}(1)} \alpha_i |i\rangle}{\left( \sum_{i \in f^{-1}(1)} |\alpha_i|^2 \right)^{1/2}} = \frac{1}{\sqrt{\sum_{i \in f^{-1}(1)} |\alpha_i|^2}} \sum_{i \in f^{-1}(1)} \alpha_i |i\rangle$$

and

$$|B\rangle = \frac{1}{\sqrt{\sum_{i \in f^{-1}(0)} |\alpha_i|^2}} \sum_{i \in f^{-1}(0)} \alpha_i |i\rangle$$

$$\text{Q4. } Z_f |G\rangle = \sum_{i \in f^{-1}(1)} \frac{\alpha_i}{\sqrt{\sum_{i \in f^{-1}(1)} |\alpha_i|^2}} Z_f |i\rangle = \sum_{i \in f^{-1}(1)} \frac{\alpha_i}{\sqrt{\sum_{i \in f^{-1}(1)} |\alpha_i|^2}} (-1)^{f(i)} |i\rangle = - \sum_{i \in f^{-1}(1)} \frac{\alpha_i}{\sqrt{\sum_{i \in f^{-1}(1)} |\alpha_i|^2}} Z_f |i\rangle = -|G\rangle$$

Similarly,  $Z_f |B\rangle = |B\rangle$ .

This means  $Z_f$  is a "reflection through  $|B\rangle$ ".

$$\text{We have } (A R A^{-1} |U\rangle) = (A R A^{-1} A |0^m\rangle) \\ = A R |0^m\rangle \underbrace{\quad}_{= 1} \\ = A (2 |0^m\rangle \langle 0^m|0^m\rangle - |0^m\rangle) \\ = A |0^m\rangle \\ = |U\rangle.$$

Let  $|V\rangle$  be an orthogonal state to  $|U\rangle$ .

$A^{-1}|V\rangle$  is orthogonal to  $A^{-1}|U\rangle = |0^m\rangle$

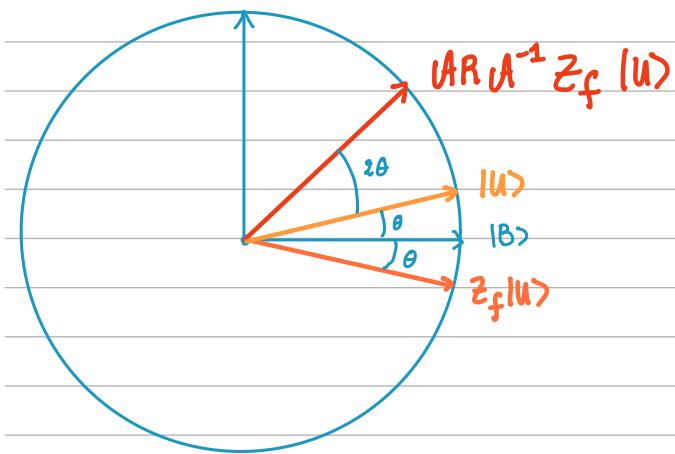
thus  $U^{-1}|V\rangle = \sum_{i \neq 0} \beta_i |i\rangle$  for some  $\beta_i \in \mathbb{C}$

then  $R U^{-1}|V\rangle = \sum_{i \neq 0} \beta^i R|i\rangle = - \sum_{i \neq 0} \beta_i |i\rangle = -|V\rangle$

and thus  $U R U^{-1}|V\rangle = -|V\rangle$ .

$U R U^{-1}$  is the "reflection through  $|U\rangle$ ".

Q6)



Q5. After  $k$  calls, we have an angle  $(2k+1)\theta$ .

We want  $(2k+1)\theta \approx \pi/2$  thus

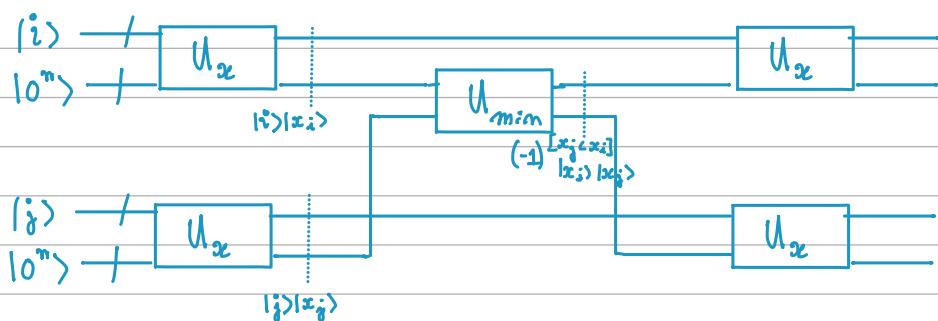
$$\begin{aligned} k &\approx \pi/4\theta \\ &\approx \pi/4 \sin^{-1}(1/\sqrt{p}) \\ \text{when } p &\text{ small} \quad \approx \pi/4 \sqrt{p} \end{aligned}$$

So, when  $p$  is small, we have  $\mathcal{O}(\sqrt{p})$  calls instead of  $\mathcal{O}(1/p)$ .

Q6 Grover's algorithm is when  $U = H^{\otimes n}$ , then  $\Pr[\text{output } e \in f^{-1}(1)] = \#f^{-1}(1)/2^n$ .

### III List-min

$$U_{\min} |x_i\rangle |x_j\rangle = (-1)^{[x_i < x_j]} |x_i\rangle |x_j\rangle$$



a) The worst case is  $x_1 > \dots > x_N$ .

The number of calls to Grover's algorithm is  $\mathcal{O}\left(\sum_{k=1}^N \sqrt{N/k}\right) = \mathcal{O}(N)$

$$\text{as } \sum_{k=1}^n \frac{1}{\sqrt{k}} = \mathcal{O}(n)$$

by comparing with an integral.

b) Let  $j$  be of rank  $n$ . Then  $\Pr[j \text{ chosen by the algorithm}] = 1/n$ .

By induction on the size of the list, we have that:

• for  $N = n = 1$ , we have  $\Pr[1 \text{ chosen}] = 1$ .

•  $\Pr[j \text{ chosen in the } N+1\text{-list}]$

$$= \frac{1}{N+1} + \sum_{i=1}^{N+1} \Pr[j \text{ is chosen after } i \mid i \text{ is chosen}] \times \frac{1}{N+1}$$

$$= \frac{1}{N+1} + \frac{1}{N+1} \sum_{i=1}^{N+1} P(j, N-i)$$

$$= \frac{1}{N+1} + \frac{1}{N+1} \sum_{i=1}^{N+1} \frac{1}{n} = \frac{1}{N+1} \left( 1 + \frac{N+1-n}{n} \right) = \frac{1}{n}$$

c)  $\mathbb{E}[\# \text{ Querres}] \leq C \sum_{n=1}^N P(n, N) \sqrt{N/n} = C \sqrt{N} \sum_{n=1}^N \frac{1}{n \sqrt{n}} \leq G(\sqrt{N})$ .

d) So, by Markov,  $\Pr[Q \geq t \mathbb{E}[Q]] \leq 1/t$

thus after  $3G(\sqrt{N})$  calls,  $\Pr[\text{success}] \geq \frac{2}{3}$ .

## IV Block ball and sphere

# 1D n° 8

## II. Parity check matrix

1) " $\Rightarrow$ " As  $H$  is a parity check then  $\ker(H^T) = C$  and thus  $\text{rank}(H) = \text{rank}(H^T)$   
 $= n - \dim(\ker H)$   
 $= n - k$

We have  $(\text{im } G) = C = (\ker H^T)$  thus  $GH^T = 0$ .

" $\Leftarrow$ " If  $GH^T = 0$  and  $y \in C = (\text{im } G)$  then  $y = xG$  for some  $x \in \mathbb{F}_2^k$ .

Then,  $y^T H^T = x^T G^T H^T = 0$  thus  $y \in \ker H^T$

and so  $C \subseteq \ker H^T$  and  $\dim(\ker H^T) = n - (n - k) = k = \dim C$ .

We can conclude  $C = \ker H^T$  and thus  $H$  is a parity check.

2) We have  $G = \text{span} (11011, 10101)$  and

we can use  $H^T = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \\ 0 & 1 & 1 \\ 1 & 0 & 1 \end{pmatrix}$  and it works!

Note: We can note that  $H = \begin{pmatrix} 1 & 1 & 1 \\ 1 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix}$  and we can always do that  
 (and to  $G$  too), thus we can use Gauss's pivot algorithm to  
 get  $H$ .

## III Computations around Shor's code.

$$1) |\bar{0}\rangle = \frac{1}{2\sqrt{2}} (|0\rangle^{\otimes 3} + |1\rangle^{\otimes 3})^{\otimes 3} \quad |\bar{1}\rangle = \frac{1}{2\sqrt{2}} (|0\rangle^{\otimes 3} - |1\rangle^{\otimes 3})^{\otimes 3}$$

2) (a) If  $E = \mathbb{1}^{\otimes 9}$  then the whole circuit is equivalent to  $\mathbb{1}^{\otimes 9}$ .

(b) If  $E = X \otimes \mathbb{1}^{\otimes 8}$  then the  $X$  error is propagated and corrected.

(c) If  $E = Z \otimes \mathbb{1}^{\otimes 8}$  then the  $Z$  error is propagated and corrected.

(d) If  $E = X_7 \otimes I^{\otimes 8}$  then the  $X_7$  error is propagated and corrected.

3) We have that  $X_7 \propto Y$  and thus we can correct any error on the first qubit (as  $X, Y, Z, I$  generates all unitary matrices).

#### IV Stabilizer codes

1) We consider  $(|000\rangle, |111\rangle)$  which forms a basis of  $\mathcal{C}_S$ .

2) " $\Rightarrow$ "  $H \subseteq S$

" $\Leftarrow$ "  $g = h_1 \dots h_k, H^{\dagger} h_g \in H,$

$$\begin{aligned} h_1 \dots h_k |\Psi\rangle &= h_1 \dots h_{k-1} |\Psi\rangle \\ &\vdots \\ &= h_1 |\Psi\rangle \\ &= |\Psi\rangle. \end{aligned}$$

3) If  $-I \in S$ , then for any  $u \in \mathcal{C}_S$ ,  $-I u = -u = u$  thus  $u = 0$ .

4) Rearranging the elements, assume  $g_1 g_2 \neq g_2 g_1$  then  $g_1 g_2 = -g_2 g_1$

$\forall u \in \mathcal{C}_{\{g_1, g_2\}}$   $g_1 g_2 u = g_2 g_1 u = u$

and thus  $-u = u$  and thus  $u = 0$ .

Lemma: Take  $a, b \in \{I, X, Y, Z\}$ , then either

$$[a, b] := ab - ba = 0 \quad \text{or} \quad \{a, b\} := ab + ba = 0$$

proof.  $XY = -YX$ ,  $XZ = -ZX$ ,  $YZ = -ZY$ , etc.

Note Let  $S$  be a stabilizer subgroup.

Define  $N(S) := \{x \in G_n \mid \forall y \in S, [x, y] = 0\}$ .

Then, for  $E \subseteq G_n$ ,

$E$  can be corrected iff  $\forall x, y \in E, x^T y \in N(S) \setminus S$ .

And logical operators  $\cong N(S) / S$ .

5) Show:  $|0\rangle \propto (|0\rangle^{\otimes 3} + |1\rangle^{\otimes 3})^{\otimes 3}$  and  $|1\rangle \propto (|0\rangle^{\otimes 3} - |1\rangle^{\otimes 3})^{\otimes 3}$

If  $\mathcal{C}_S$  is a  $[n, k, d]$ -stabilizer code, then  $S$  is generated by  $n-k$  stabilizers.   
↑  
 $\#$  logical qubits       $\#$  physical logic      We want to find stabilizers.

We can use  $Z_1 Z_2, Z_2 Z_3, Z_4 Z_5, Z_5 Z_6, Z_7 Z_8, Z_8 Z_9, X_1 \dots X_6, X_6 \dots X_9$

1 1 1

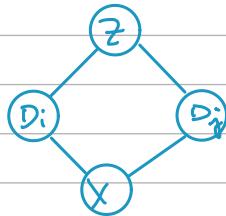
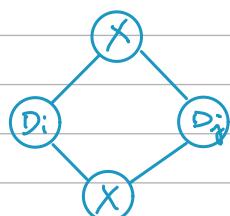
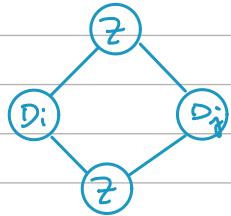
QEC Zoo

Philippe Faist FUB.

### TD n° 9

## 2. The Surface Code.

Q1. From the drawing, we have that



commutes.

Q2a. For any  $S, T \in G_n$ , we have  $[S, T] := ST - TS = 0$   
or  $\{S, T\} := ST + TS = 0$ .

Q2b.  $B^T A |\psi\rangle = B^T A \hat{S} |\psi\rangle = -\hat{S} B^T A |\psi\rangle$

thus  $B^T A |\psi\rangle$  is an eigenvector of  $\hat{S}$  with eigenvalue  $-1$ .

However,  $\mathcal{G}$  is a "proper" subspace of  $\hat{\mathcal{S}}$  with eigenvalue  $1$ .

By the Spectral Theorem, we have that the two proper spaces associated with two different eigenvalues are orthogonal.

Thus,  $B^T A |\psi\rangle \perp \mathcal{G}$ .

Q2c. Take  $|\phi\rangle, |\psi\rangle \in \mathcal{G}$  with  $|\phi\rangle \perp |\psi\rangle$ .

Either  $B^T A = 1$  and then  $\langle \phi | B^T A | \psi \rangle = \langle \phi | \psi \rangle = 0$

Either  $B^T A$  is a stabilizer and then  $\langle \psi | B^T A | \psi \rangle = \langle \psi | \psi \rangle = 0$ .

Q3a. Write  $|\psi\rangle = |d_1\rangle \dots |d_5\rangle$  and  $|d_1\rangle = \alpha|0\rangle + \beta|1\rangle$   
 $|d_4\rangle = \gamma|0\rangle + \delta|1\rangle$ .

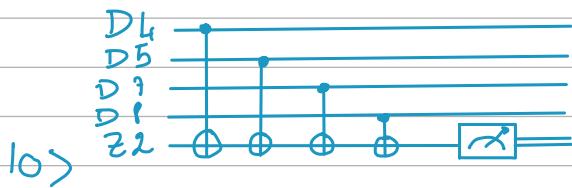
$$\hat{S}^{z_1} |\psi\rangle = |\psi\rangle \Leftrightarrow (\hat{z}|d_1\rangle)(\hat{z}|d_4\rangle) = |d_1\rangle |d_4\rangle$$

$$\Leftrightarrow (\alpha|0\rangle - \beta|1\rangle)(\gamma|0\rangle - \delta|1\rangle) \\ = (\alpha|0\rangle + \beta|1\rangle)(\gamma|0\rangle + \delta|1\rangle)$$

$$\Rightarrow \alpha\delta = 0 \text{ and } \beta\gamma = 0$$

$$\Rightarrow \delta^{z_1} = 0.$$

Q3b.



Q3c.

## II Entropies for counting

o) Chain rule:

$$\begin{aligned}
 H(X|Y) &= - \sum_{x,y} p_{x,y} \log \left( p_{x,y} / p_y \right) \\
 &= - \sum_{x,y} p_{x,y} \log p_{x,y} + \sum_y \underbrace{\left( \sum_x p_{x,y} \right)}_{p_y} \log p_y \\
 &= H(X, Y) - H(Y)
 \end{aligned}$$

$$\begin{aligned}
 1) \quad \sum_{j=1}^N H(X_j | X_1, \dots, X_{j-1}) &= \sum_{j=1}^N (H(X_1, \dots, X_j) - H(X_1, \dots, X_{j-1})) \\
 &\text{"telescopic" sum} \\
 &= H(X_1, \dots, X_N) - H(\ )
 \end{aligned}$$

$$= H(X_1, \dots, X_N)$$

$$\begin{aligned}
 2) \quad \sum_{j=1}^M H(X_{S_j}) &= \sum_{j=1}^M H((X_l)_{l \in S_j}) \\
 &= \sum_{j=1}^N \sum_{l \in S_j} H(X_l | X_{\{i \in S_j \mid i < l\}}) \\
 &\geq \sum_{j=1}^N \sum_{l \in S_j} H(X_l | X_1, \dots, X_{l-1}) \\
 &\geq k \sum_{l=1}^n H(X_l | X_1, \dots, X_{l-1}) \\
 &= k H(X_1, \dots, X_N).
 \end{aligned}$$

3) Consider  $(X_1, X_2, X_3) \in A$  chosen uniformly over  $A$ .

Then  $H(X_1, X_2, X_3) \leq H(X_1, X_2) + H(X_2, X_3) + H(X_1, X_3)$

$$\text{and } H(X_1, X_2, X_3) = -\sum_{a \in A} p_a \log p_a = \log |A|$$

thus

$$\begin{aligned}\log |A| &\leq \frac{1}{2}(H(A_{12}) + H(A_{13}) + H(A_{23})) \\ &\leq \frac{1}{2}(\log |A_{12}| + \log |A_{13}| + \log |A_{23}|) \\ &\leq \frac{1}{2} \times 3 \log n \leq \log n^{3/2}\end{aligned}$$

$$\text{and finally } |A| \leq n^{3/2}.$$

### III From divergences to entropies.

$$D(p \parallel q) = \mathbb{E}_{x \sim p} [\log(p_x / q_x)] = \sum_{x \in \mathcal{X}} p_x \log(p_x / q_x).$$

$$\begin{aligned}1) \quad -D(p_{xy} \parallel P_x \otimes P_y) &= -\sum_{\substack{x \in \mathcal{X} \\ y \in \mathcal{Y}}} p_{xy} \log(p_{xy} / p_x p_y) \\ &= H(X|Y)\end{aligned}$$

$$\begin{aligned}D(p_{xy} \parallel P_x \otimes P_y) &= -\sum_{\substack{x \in \mathcal{X} \\ y \in \mathcal{Y}}} p_{xy} \log(p_{xy} / p_x p_y) \\ &= I(X:Y).\end{aligned}$$

$$2) D(p \parallel \pi) = \text{Tr}(p(\log p - \log \pi))$$

Cannal quantique       $\mathcal{H}$  un espace de Hilbert (" $\mathcal{H} = \mathbb{C}^n$ ")

$\mathcal{L}(\mathcal{H})$

Un cannal quantique  $\Phi: \mathcal{L}(\mathcal{H}) \rightarrow \mathcal{L}(\mathcal{H}')$  linéaire  
qui preserve les états (et qui est donc CPTP)

$\Leftrightarrow \Phi$  trace preserving     $\Leftrightarrow \forall x \in \mathcal{L}(\mathcal{H}), \text{Tr}(\Phi(x)) = \text{Tr } x$   
 $\Leftrightarrow \Phi$  completely positive     $\Leftrightarrow \forall n \in \mathbb{N} \quad \Phi \otimes \text{Id}_n \geq 0.$

Pour  $n \in \mathbb{N}$ ,

$$\mathcal{M}_n(\mathcal{L}(\mathcal{H})) := \{(x_{ij})_{i,j \in \llbracket 1, n \rrbracket} \mid \forall i,j, x_{ij} \in \mathcal{L}(\mathcal{H})\}.$$

$$\begin{aligned} \text{Id}_n \otimes \Phi: \mathcal{M}_n(\mathcal{L}(\mathcal{H})) &\longrightarrow \mathcal{M}(\mathcal{L}(\mathcal{H}')) \\ X = (x_{ij})_{ij} &\longmapsto (\Phi(x_{ij}))_{ij} \end{aligned}$$

Pour  $X \in \mathcal{M}_n(\mathcal{L}(\mathcal{H}))$ ,

$$X \geq 0 \Leftrightarrow \text{Spec}(X) \subseteq \mathbb{R}^+ \Leftrightarrow \exists Y \in \mathcal{M}_n(\mathcal{L}(\mathcal{H})), X = Y^*Y.$$

Pour  $\rho \in \mathcal{D}(\mathcal{H})$  et  $\sigma \in \mathcal{D}(\mathcal{H}')$  où  $\dim \mathcal{H}' = m$

$$\text{alors } \sigma \otimes \rho \in \mathcal{M}_m(\mathcal{L}(\mathcal{H}))$$

$$\text{et } \text{Id}_m \otimes \Phi(\sigma \otimes \rho) = \sigma \otimes \Phi(\rho) \geq 0.$$

2) Avec  $\Phi: x \mapsto \text{Tr}(x) \rho_0$  où  $\rho_0 \in \mathcal{D}(\mathbb{C}^n)$ ,

on a donc

$$\mathcal{D}(\Phi(\rho) \| \Phi(\sigma)) = \mathcal{D}(\rho_0 \| \rho_0) = \text{Tr}(\rho_0 (\log \rho_0 - \log \rho_0)) = 0$$

Si  $U$  est unitaire, alors  $\Phi: x \mapsto UxU^*$ .

$(\mathcal{L}(\mathcal{H}), \langle \cdot, \cdot \rangle_{HS})$  où  $\langle x, y \rangle = \text{tr}(x y^\dagger)$ .

↳ Hilbert-Schmidt

Si  $\Phi$  est un canal, alors  $\Phi^*$  est unital & completely positive.

### Traces partielles

Pour deux systèmes  $A, B$  joints,

$$\mathcal{L}(\mathcal{H}_A \otimes \mathcal{H}_B) = \underbrace{\text{span}}_{\Delta} \{ x_A \otimes x_B \mid x_A \in \mathcal{L}(\mathcal{H}_A), x_B \in \mathcal{L}(\mathcal{H}_B) \}.$$

$$\text{Trace partielle } \text{tr}_B : \mathcal{L}(\mathcal{H}_A \otimes \mathcal{H}_B) \longrightarrow \mathcal{L}(\mathcal{H}_A)$$

$$x_A \otimes x_B \longmapsto \text{tr}(x_B) x_A$$

$$\text{Ainsi, } \text{tr}(\text{tr}_B(x_A \otimes x_B)) = \text{tr}(x_A) \text{tr}(x_B) = \text{tr}(x_A \otimes x_B).$$

$\text{tr}_B$  est un canal quantique qui donne la marginale sur A de l'état biparti  $\rho_{AB}$ .

$$H(A|B) = -D(\rho_{AB} \parallel \mathbb{1}_A \otimes \text{tr}_B(\rho_{AB}))$$

$$3) \rho_A = \text{tr}_B(\rho_{AB}) = \text{tr}_B(\mathbb{1}_{A/d} \otimes \mathbb{1}_{B/d}) = \underbrace{\text{tr}\left(\frac{\mathbb{1}_B}{d}\right)}_1 \mathbb{1}_A/d.$$

$$\text{De même, } \rho_B = \mathbb{1}_B/d.$$

$$I(A:B) = D(\rho_{AB} \parallel \rho_A \otimes \rho_B) = D\left(\frac{\mathbb{1}_{AB}}{d^2} \parallel \frac{\overbrace{\mathbb{1}_A \otimes \mathbb{1}_B}^{1_{AB}}}{d^2}\right) = 0$$

$$\rho_{AB} \text{ produit} \Leftrightarrow I(A:B) = 0$$

$$H(A|B) = -D(\rho_{AB} \parallel \mathbb{1}_A \otimes \rho_B) = -D\left(\frac{\mathbb{1}_{AB}}{d^2} \parallel \frac{\mathbb{1}_{AB}}{d}\right)$$

$$= -\text{tr}\left(\frac{\mathbb{1}_{AB}}{d^2} \left( \log\left(\frac{\mathbb{1}_{AB}}{d^2}\right) - \log\left(\frac{\mathbb{1}_{AB}}{d}\right) \right)\right)$$

$$= -\frac{1}{d^2} \sum_{i=1}^{d^2} \left( \log \frac{1}{d^2} - \log \frac{1}{d} \right) = -\log(1/d)$$

$$= \log d$$

$$\rho_{AB} = |\Psi_{AB}\rangle \langle \Psi_{AB}| = \frac{1}{d} \sum_{i,j} |ii\rangle \langle jj|$$

$$= \frac{1}{d} \sum_{ij} |i\rangle \langle j|_A \otimes |i\rangle \langle j|_B$$

$$\rho_A = \text{tr}_B \left( \frac{1}{d} \sum_{ij} |i\rangle \langle j|_A \otimes |i\rangle \langle j|_B \right)$$

$$= \frac{1}{d} \sum_{ij} \text{tr}(|i\rangle \langle j|_B) |i\rangle \langle j|_A$$

$$= \frac{\mathbb{1}_A}{d}$$

$$\text{et } \rho_B = \mathbb{1}_B / d$$

$$I(A:B)_{|\Psi_{AB}\rangle \langle \Psi_{AB}|} = D(|\Psi\rangle \langle \Psi| \parallel \mathbb{1}_{AB} / d^2)$$

$$= \text{tr} \left[ |\Psi\rangle \langle \Psi| \left( \log (|\Psi\rangle \langle \Psi|) - \log (\mathbb{1}_{AB} / d^2) \right) \right]$$

$$= -\text{tr} \left[ |\Psi\rangle \langle \Psi| \log (\mathbb{1}_{AB} / d^2) \right]$$

$$= -\text{tr} \left[ \frac{1}{d} \left( \sum |ii\rangle \langle jj| \right) \sum \log (1/d^2) |kk\rangle \langle kk| \right]$$

$$= -\frac{1}{d} \text{tr} \left[ \sum \log \frac{1}{d^2} |ii\rangle \langle ii| \right] = 2 \log d$$

$$H(A|B) = \dots = -\log d <_o \text{Quantique! Intrication!}$$

$$\left( H(A|B) \in [-\log d, \log d] \quad I(A:B) \in [0, 2\log d] \right)$$

$$4) H(A|Bc) = -D(p_{ABC} \parallel \mathbb{1}_A \otimes p_{Bc})$$

$$\leq -D(p_{AB} \parallel \mathbb{1}_A \otimes p_B) = H(A|B).$$

DPI

# TD M

## I Continuity of the entropy

$$1) \quad \rho - \tau = (\rho - \tau)_+ - (\rho - \tau)_- \quad \text{donc } \rho + (\rho - \tau)_+ = (\rho - \tau)_- + \tau = V$$

$$\Delta(\rho, \tau) = \text{tr}((\rho - \tau)_+) + \text{tr}((\rho - \tau)_-)$$

$$= 2 \text{tr} V - \text{tr} \rho - \text{tr} \tau$$

Let  $(t_i)$  be the eigenvalues of  $V$  ( $\lambda_i$ ) for  $\rho$ , ( $\delta_i$ ) for  $\tau$ .

$$t_i \geq \max(\lambda_i, \delta_i)$$

$$2t_i - \lambda_i - \delta_i \geq |\lambda_i - \delta_i|$$

$$\Delta(\rho, \tau) = \sum (2t_i - \lambda_i - \delta_i) \geq \sum |\lambda_i - \delta_i|$$

2.  $S(\rho) = -\text{tr}(\rho \log \rho)$ . Define  $\eta : x \mapsto x \log x$ .

$$\text{If } |\lambda - \delta| \leq \frac{1}{2}, \quad |\eta(\lambda) - \eta(\delta)| \leq \eta(|\lambda - \delta|).$$

$$\begin{aligned} |S(\rho) - S(\tau)| &= \left| \sum (\eta(\lambda_i) - \eta(\delta_i)) \right| \\ &\leq \sum |\eta(\lambda_i) - \eta(\delta_i)| \\ &\leq \sum \eta(|\lambda_i - \delta_i|) \end{aligned}$$

$$\text{Define } p_i = |\lambda_i - \delta_i| / s$$

$$\begin{aligned} \eta(|\lambda_i - \delta_i|) &= s \eta\left(\frac{|\lambda_i - \delta_i|}{s}\right) - |\lambda_i - \delta_i| \log s \\ \text{so, } \sum \eta(|\lambda_i - \delta_i|) &= s \sum \left( \eta\left(\frac{|\lambda_i - \delta_i|}{s}\right) / s - |\lambda_i - \delta_i| \log s \right) \\ &= s \sum \eta(p_i) - s \log s \\ &\leq s \log d - s \log s \end{aligned}$$

3. If  $\sigma$  and  $\rho$  have the same eigenvalues,

$$|S(\sigma) - S(\rho)| = 0 \leq 0 = \Delta(\sigma, \rho)$$

$$\text{Otherwise, } \delta > 0 \text{ and so } |S(\sigma) - S(\rho)| \leq \delta \log d - \delta \log \delta \\ \leq \Delta(\rho, \sigma) \log_2 d \\ - \Delta(\rho, \sigma) \log \Delta(\rho, \sigma)$$

$$\text{as } \delta \leq \Delta(\rho, \sigma) \leq \frac{1}{2}.$$

## II Density matrices and von Neumann entropy

a.  $\rho_{XA} = \begin{pmatrix} p_1 \rho_1 & \dots \\ & \ddots \\ & & p_m \rho_m \end{pmatrix}$

$$\begin{aligned} b. \text{tr}_X(\rho_{XA}) &= \text{tr}_X \left( \sum p_x |x\rangle\langle x| \otimes \rho_x \right) \\ &= \sum p_x \text{tr}_X(|x\rangle\langle x| \otimes \rho_x) \\ &= \sum p_x \text{tr}(|x\rangle\langle x|)_{\rho_x} = \sum p_x \rho_x \end{aligned}$$

$$\begin{aligned} \text{tr}_A(\rho_{XA}) &= \text{tr}_A \left( \sum p_x |x\rangle\langle x| \otimes \rho_x \right) \\ &= \sum p_x \text{tr}(\rho_x) |x\rangle\langle x| = \sum p_x |x\rangle\langle x|. \end{aligned}$$

c) Let  $Q_i^{(bc)}$  be the eigenvalues of  $\rho_x$ :  $\rho_x = U_x^\dagger D_x U_x$

$$U_{tot} = \sum |x\rangle\langle x| \otimes U_x$$

$$\begin{aligned} \text{so } U_{tot} U_{tot}^\dagger &= \left( \sum |x\rangle\langle x| \otimes U_x \right) \left( \sum |x\rangle\langle x| \otimes U_x^\dagger \right) \\ &= \sum |x\rangle\langle x| \otimes U_x U_x^\dagger \\ &= \sum |x\rangle\langle x| \otimes \mathbb{1} = \mathbb{1} \end{aligned}$$

$$\begin{aligned} U_{\text{tot}} \rho_{XA} U_{\text{tot}}^+ &= \sum p_x |x\rangle\langle x| \otimes U_x \rho_x U_x^+ \\ &= \sum p_x |x\rangle\langle x| \otimes \text{diag}(\lambda_1^{(x)}, \dots, \lambda_m^{(x)}) \end{aligned}$$

d)  $H(\{p_1, \dots, p_m\}) = S\left(\sum p_x |x\rangle\langle x|\right).$

$$\begin{aligned} S(\rho_{XA}) &= - \sum_{x,i} p_x \lambda_i^{(x)} \log(p_x \lambda_i^{(x)}) \\ &= - \sum p_x \lambda_i^{(x)} (\log p_x + \log \lambda_i^{(x)}) \\ &= - \sum p_x (\sum \lambda_i^{(x)}) \log p_x - \sum p_x \sum \lambda_i^{(x)} \log \lambda_i^{(x)} \\ &= H(\{p_1, \dots, p_m\}) + \sum p_x S(\rho_x). \end{aligned}$$

2. a)

$$\begin{aligned} \rho &= \frac{1}{2} \left( \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} + \begin{pmatrix} \cos^2 \theta & \cos \theta \sin \theta \\ \cos \theta \sin \theta & \sin^2 \theta \end{pmatrix} \right) \\ &= \frac{1}{2} \begin{pmatrix} 1 + \cos^2 \theta & \cos \theta \sin \theta \\ \cos \theta \sin \theta & \sin^2 \theta \end{pmatrix} \\ &= \frac{\sin \theta \cos \theta}{2} \underbrace{\begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}}_{X} + \frac{1}{2} \begin{pmatrix} 1 + \cos^2 \theta & 0 \\ 0 & 1 - \cos^2 \theta \end{pmatrix} \\ &= \frac{\sin \theta \cos \theta}{2} X + \frac{1}{2} I + \frac{\cos^2 \theta}{2} Z \quad Z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \end{aligned}$$

b.  $\text{tr}(\rho) = 1$

$$\det(\rho) = \frac{1}{4} \left( (1 + \cos^2 \theta) \sin^2 \theta - \cos^2 \sin^2 \theta \right) = \frac{\sin^2 \theta}{4}$$

$$\lambda_1, \lambda_2 = \frac{1}{4} \sin^2 \theta \pm \frac{1}{2} \sqrt{1 - \sin^2 \theta} = \frac{1}{4} \sin^2 \theta \pm \frac{1}{2} \cos \theta$$

  $\text{Sp}(\rho) = \left\{ \frac{1 \pm \cos \theta}{2} \right\}.$

### III Entanglement & nonlocality

a)  $\rho_{WW}(p) = \frac{1}{2} (|00\rangle\langle 00| + |00\rangle\langle 11| + |11\rangle\langle 00| + |11\rangle\langle 11|)$   
 $+ \frac{1-p}{u} (|01\rangle\langle 00| + |01\rangle\langle 10| + |10\rangle\langle 01| + |10\rangle\langle 11|)$

$$\rho_{WW}(p) = \begin{pmatrix} \frac{1+p}{u} & 0 & 0 & p/2 \\ 0 & \frac{1-p}{u} & 0 & 0 \\ 0 & 0 & \frac{1-p}{u} & 0 \\ p/2 & 0 & 0 & 1+p/u \end{pmatrix}$$

b)  $\left\{ \begin{array}{l} x \in \mathcal{L}(\mathcal{H}_A \otimes \mathcal{H}_B) \\ x = \sum x_i^{(A)} \otimes x_i^{(B)} \end{array} \right.$   
 $x^{T_B} = \sum x_i^{(A)} \otimes (x_i^{(B)})^T$

$$\rho_{WW}(p)^{T_B} = \begin{pmatrix} \frac{1+p}{u} & 0 & 0 & 0 \\ 0 & \frac{1-p}{u} & p/2 & 0 \\ 0 & p/2 & \frac{1-p}{u} & 0 \\ 0 & 0 & 0 & 1+p/u \end{pmatrix}$$

c) Let  $p$  s.t.  $\rho_{WW}(p)^{T_B} \geq 0$

$\begin{pmatrix} 1 \\ 0 \\ 0 \\ 0 \end{pmatrix}, \begin{pmatrix} 0 \\ 0 \\ 0 \\ 1 \end{pmatrix}$  are eigen vectors of  $\frac{1+p}{u}$

$\begin{pmatrix} \frac{1-p}{u} & p/2 \\ p/2 & \frac{1-p}{u} \end{pmatrix}$  has eigen vectors  $\begin{pmatrix} 1 \\ 1 \end{pmatrix}$  w/ eigenvalue  $\frac{1+p}{u}$

$\begin{pmatrix} 1 \\ -1 \end{pmatrix}$  w/ eigenvalue  $\frac{1-3p}{u}$

d)  $(\rho_A \otimes \rho_B)^{T_B} = \underbrace{\rho_A}_{\geq 0} \otimes \underbrace{\rho_B^T}_{\geq 0} \geq 0$

c) if  $p \leq 1/3$ ,  $\rho_{AB}(p)^T_B \geq 0$  so  $\rho_{AB}(p)^T_B$  product  
and so  $\rho_{AB}(p)^T_B$  entangled when  $p > 1/3$

## IV Quantum Channels

If  $\phi$  is a quantum channel then there exists an environment E and  
 $V: A \rightarrow BE$  st

$$\begin{aligned}\phi(p) &= \text{tr}_E(V_p V^\dagger) \\ &= \text{tr}_E(U_p \otimes |0\rangle\langle 0|_E \otimes U^\dagger)\end{aligned}$$