



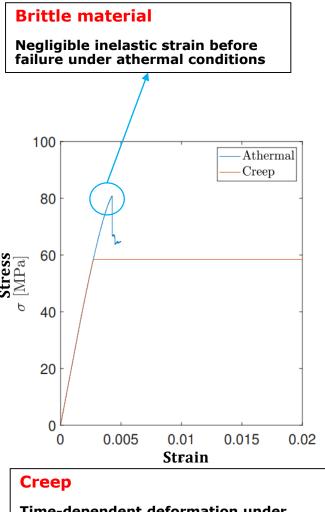


Master Project: Modelling the creep of brittle materials

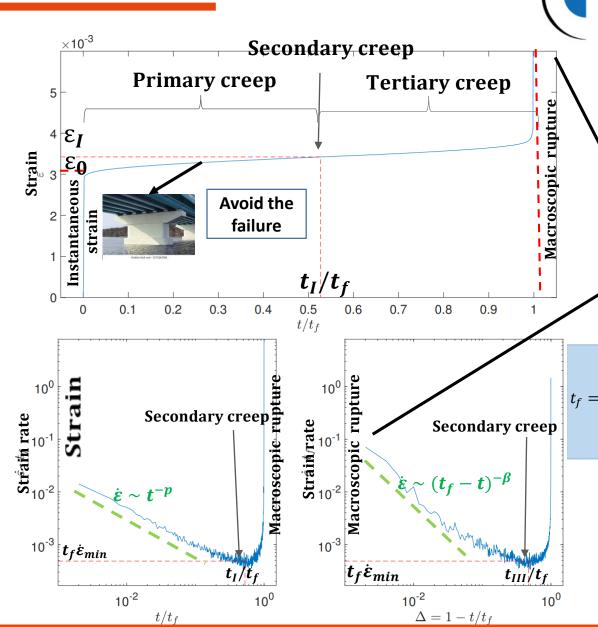
Project Advisor(s):
 Jérôme WEISS
 David AMITRANO

Juan Carlos VERANO ESPITIA

1. Presentation of the problem



Time-dependent deformation under constant applied stress





Empirical laws

$$t_f = A \dot{\varepsilon}_{min}^{-b}$$
 Monkman and Grant (1956)

$$\dot{\varepsilon} \sim \left(t_f - t\right)^{-\beta}$$
 Tertiary creep

Theoretical explanation??





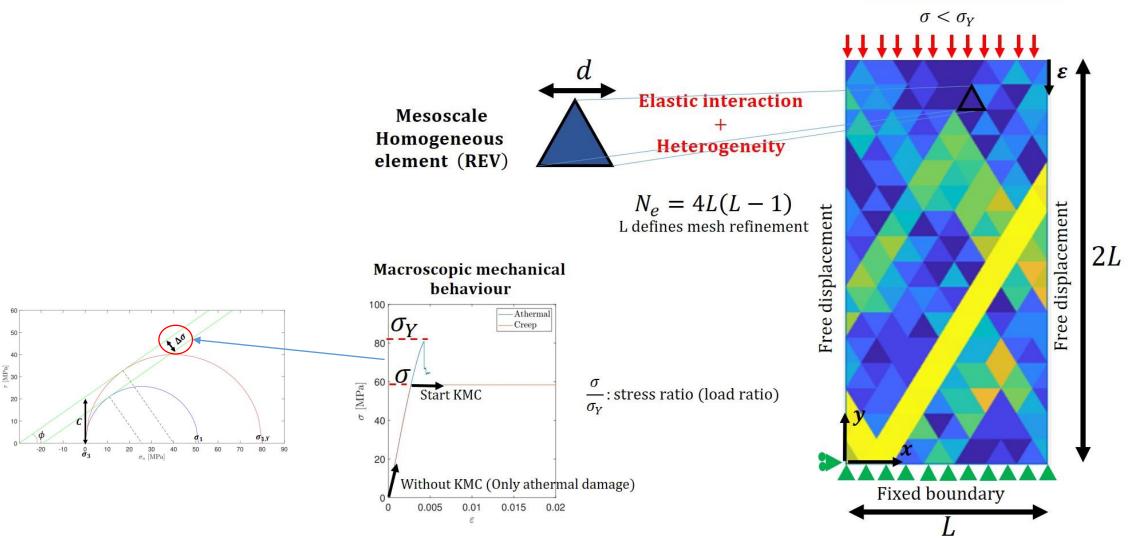


Amitrano, 1999; Amitrano and Helmstetter, 2006; Girard et al., 2010





Macroscale model 2D







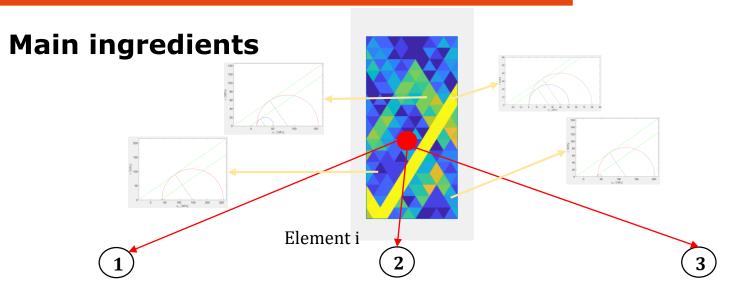


All the elements



Elastic interaction (Solved with FEM)

$$\sigma_{ej} = E_{ejkl} \varepsilon_{ej}$$



Heterogeneity

Cohesion C_i is set randomly **Uniform distribution** (only once) $C_i \in [10, 30] \text{MPa}$

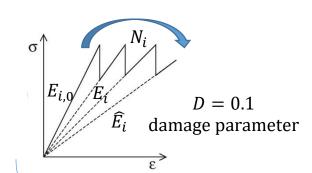
Thermal activation

$$\omega_{i} = \omega_{0} \exp\left(-\frac{E_{a,i}}{k_{B}T}\right)$$
Energy
$$E_{a,i}$$
Reaction

initial disorder (random variable) + stochastic process

Damage

$$\widehat{E_i} = (1 - D)E_i$$



Athermal damage (MC criterium)

Thermal activated damage (Probability)

ISTerre

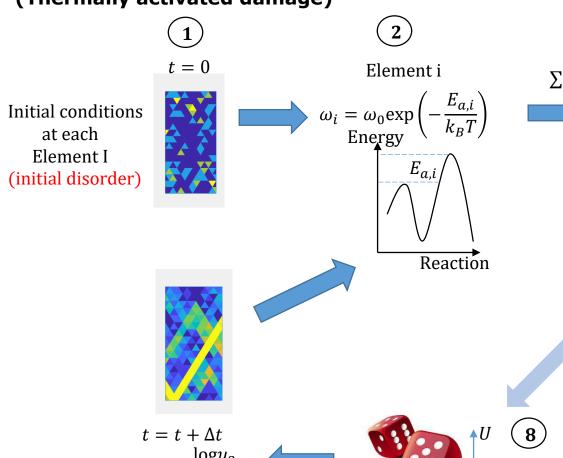


Kinetic Monte Carlo algorithm (KMC)

Bortz et al., 1975

Introduction of timescale and thermal effect

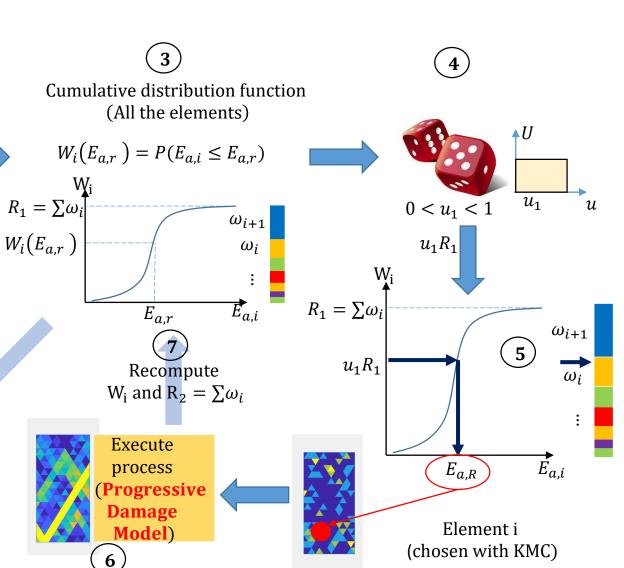
(Thermally activated damage)



 u_2

 $0 < u_2 < 1$

и



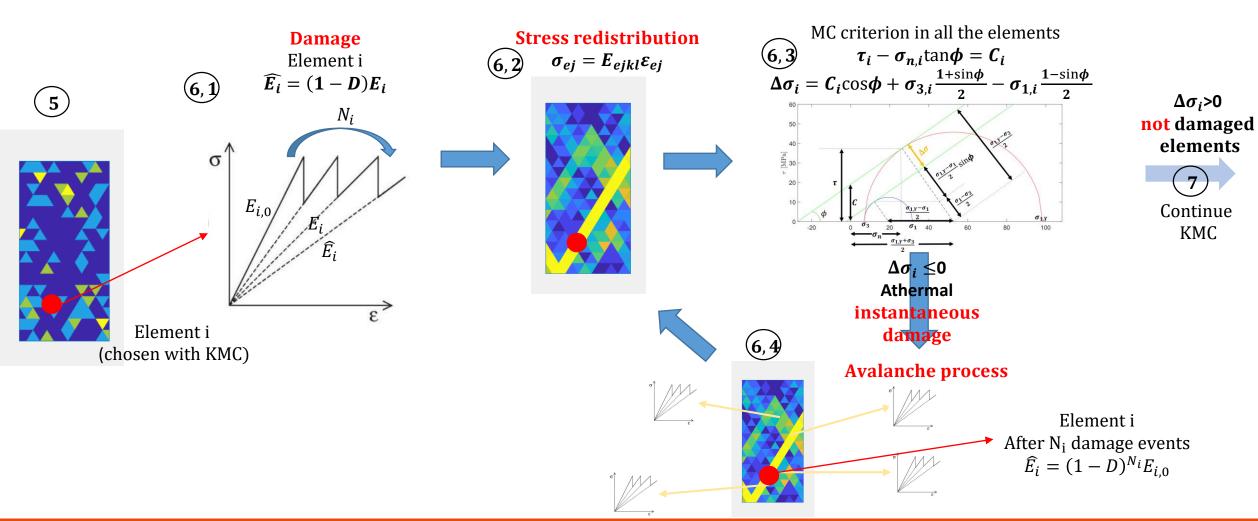
 $\Delta t = -$

(9)





Progressive isotropic damage model (PDM) (athermal damage)







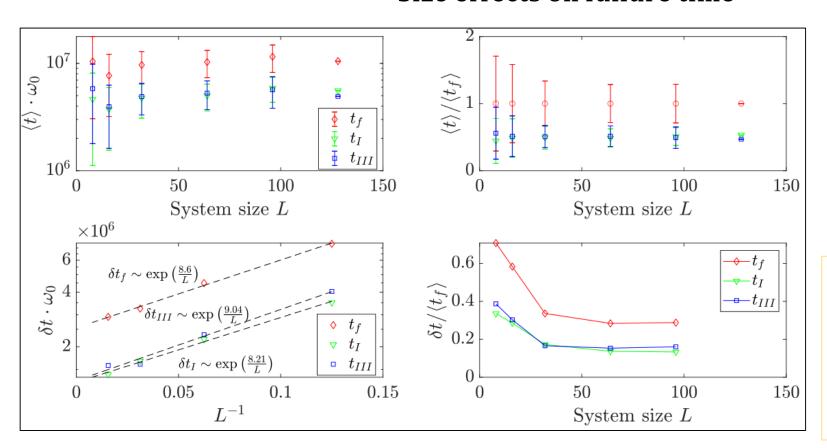
- 4.1 Size effects on failure time
- 4.2 Empirical approximations to predict rupture time
- 4.3 Response to thermally activated process at macroscale
- 4.4 Acceleration of deformation during the tertiary creep

4.1 Size effects on failure time





Size effects on failure time



$$\langle t_f \rangle (L) \sim \text{constant}$$

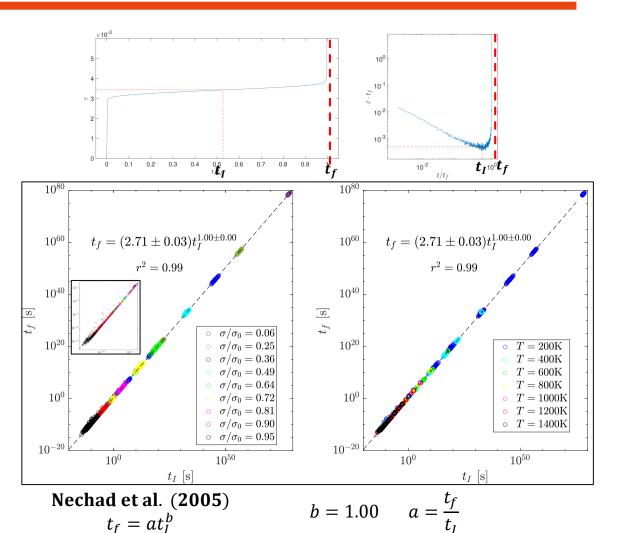
$$\delta t_f \sim \exp(L^{-1})$$

No size effects in mean

Size effects in standard deviation

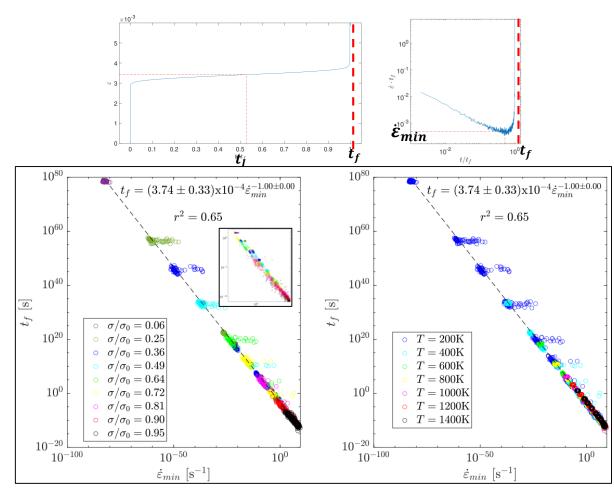
Standard deviation doesn't vanish with *larger* sizes ⇒ ???

4.2 Empirical approximations to predict rupture time









Monkman and Grant (1956) $t_f = C_1 \dot{\varepsilon}_{min}^{-m}$

m = 1.00

 $C_1 = t_f \dot{\varepsilon}_{min}$

4.4 Acceleration of deformation during the tertiary creep



(c)



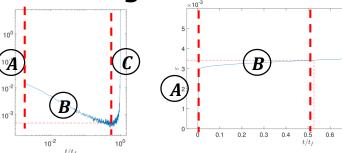
Spatial correlation of damage events

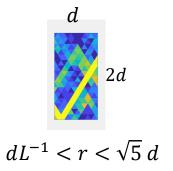
System Sizes L = 128

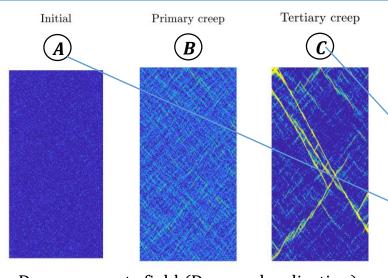
Number of simulations N = 1

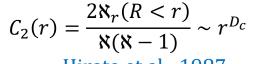
Temperature T = 800 K

Stress ratio $\sigma/\sigma_Y = 0.72 \ (\sigma = 0.72\sigma_Y)$





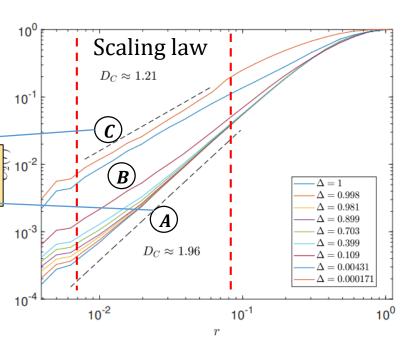




Hirata et al., 1987

Fractal behaviour of damage events??

 $D_c \approx 1$ Perfect damage localisation $D_c \approx 2$ Homogeneous damage



Damage events field (Damage localisation)

$$rac{E_{i,0}-E_i}{E_{i,0}}=1-(1-D)^{\mathcal{N}_i}$$
 0.2 0.4 0.6 0.8 1

4.4 Acceleration of deformation during the tertiary creep





Stress gap distribution

System Sizes L = 128

Number of simulations

N = 1

Temperature

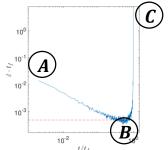
T = 800 K

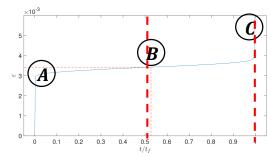
Stress ratio

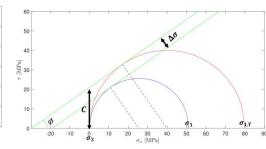
$$\sigma/\sigma_Y = 0.72 \quad (\sigma = 0.72\sigma_Y)$$

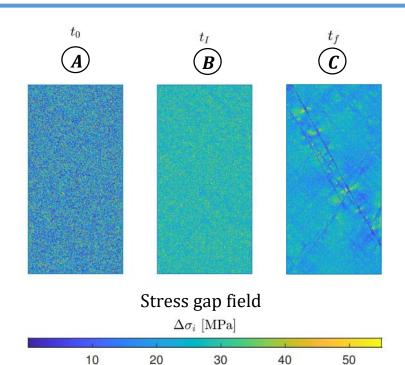
Macro stress gap

 $\Delta \sigma \approx 11 \text{ MPa}$



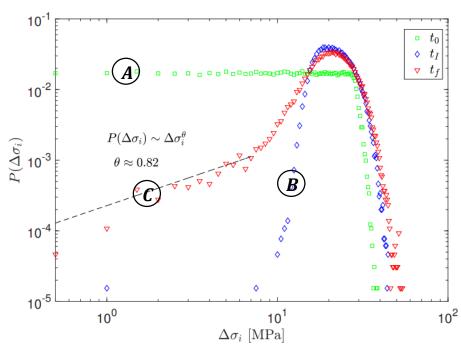






 $P(\Delta\sigma_i) \sim \Delta\sigma_i^{\theta}$ Lin et al., 2014

Power law tail towards small stress gaps



5. Conclusions





- Numerical modelling with small set of physical ingredients (disorder, thermal activation, damage, elastic interaction).
- Numerical modelling based on progressive damage coupled with KMC algorithm.
- Time dependency → through a KMC algorithm
- Numerical modelling allows to obtain information at microscopic scale during different time steps of the creep process.
- For different temperature and stress conditions the model reproduces:
 - o the different stages of creep phenomenology and its characteristics (as shown experimentally).
 - damage localisation and fractal structure of the damage field
 - o empirical laws of creep (Monkman and Grant, 1956 and Nechad et al., 2005)
- The empirical laws are a consequence of the Arrhenius law with same activation volume, i.e., importance of thermal activation in the phenomenology of creep.
- The empirical laws (MG) could be used for forecasting the failure
- Distribution of the stress gap is the key to understand what is happening during the different stages of creep.





16'

PhD: Statistical physics of creep rupture of structural materials

Thesis Directors: Jérôme WEISS, ISTerre Mikko ALAVA, Aalto University

Creep and relaxation experiments









Classic concrete

Light-weight concrete (Polystyrene: 1 < D < 5mm)

Creep and relaxation experiments





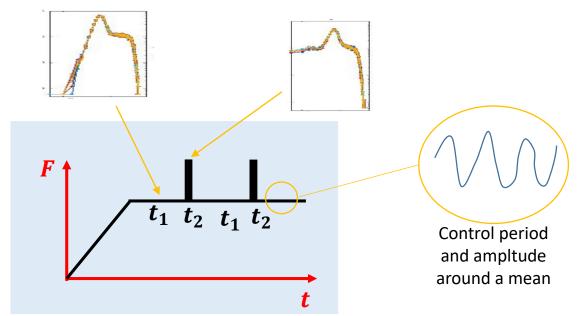
- Classic concrete
- Light-weight concrete (Polystyrene: 1 < D < 5mm)

Two kind of experiments



Acoustic emission Tomography

Creep test



Experimental setup 01

Creep and relaxation experiments





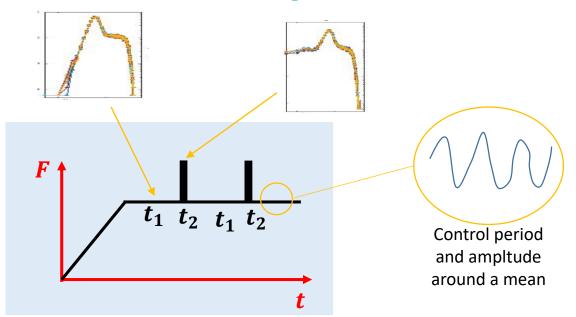
- Classic concrete
- Light-weight concrete
 (Polystyrene: 1 < D < 5mm)

Two kind of experiments

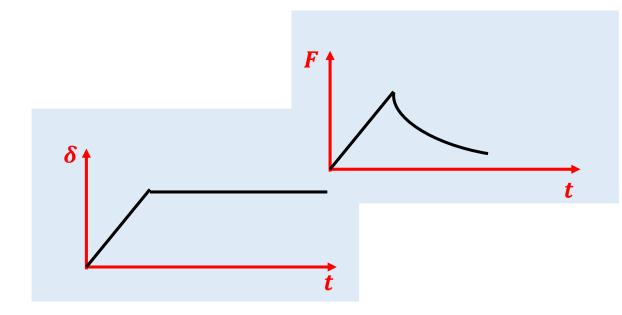


Acoustic emission Tomography

Creep test



Relaxation test



Experimental setup 01

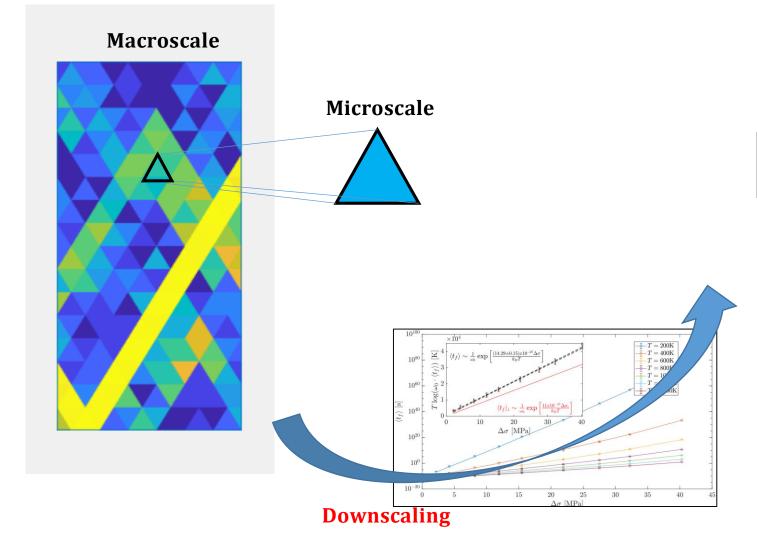


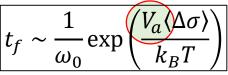
Paper in process:

The effect of quenched disorder on creep lifetimes of brittle materials

Problem:







Arrhenius law

From macro is imposed the temperature and the load, and plotting the failure time as a function of the inverse of the temperature or the stress gap is possible to obtain the activation volume at microscale

This assumption does not take into account the assumption of disorder effect

The effect of disorder on the fracture nucleation process (S. Ciliberto, A. Guarino, R. Scorreti (2001))







Physica D 158 (2001) 83-104

www.elsevier.com/locate/physd

The effect of disorder on the fracture nucleation process

S. Ciliberto*, A. Guarino, R. Scorretti¹

Ecole Normale Supérieure de Lyon, Laboratoire de Physique, CNRS UMR5672, 46, Allée d'Italie, 69364 Lyon Cedex 07, France

Received 12 January 2001; accepted 28 May 2001 Communicated by S. Fauve

Abstract

The statistical properties of failure are studied in a fiber bundle model with thermal noise. We show that the macroscopic failure is produced by a thermal activation of microcracks. Most importantly, the effective temperature of the system is amplified by the spatial disorder (heterogeneity) of the fiber bundle. The case of a time-dependent force and the validity of the Kaiser effects are also discussed. These results can give more insight to the recent experimental observations on thermally activated crack and can be useful to study the failure of electrical networks. © 2001 Elsevier Science B.V. All rights reserved.

The effect of disorder on the fracture nucleation process (S. Ciliberto, A. Guarino, R. Scorreti (2001))



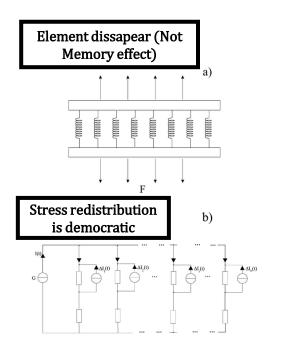




Method: DFBM

$$1 \quad F = \sum_{i=1}^{N} f_i.$$

- 2. $f_i = Ye_i$,
- 3. $f_c(i) \sim N_d(\langle f_c \rangle, KT_d)$.
- 4. $\Delta f_i(t) \sim N_{\rm T}(0, KT)$



The effect of disorder on the fracture nucleation process (S. Ciliberto, A. Guarino, R. Scorreti (2001))







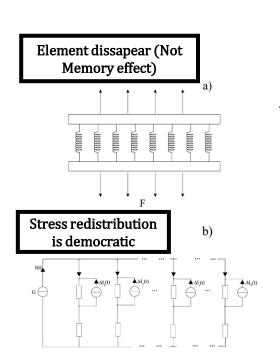
Method: DFBM

1.
$$F = \sum_{i=1}^{N} f_i$$
.

$$2 f_i = Ye_i$$

3
$$f_c(i) \sim N_d(\langle f_c \rangle, KT_d)$$
.

$$4 \Delta f_i(t) \sim N_{\rm T}(0, KT)$$



Method: DFBM

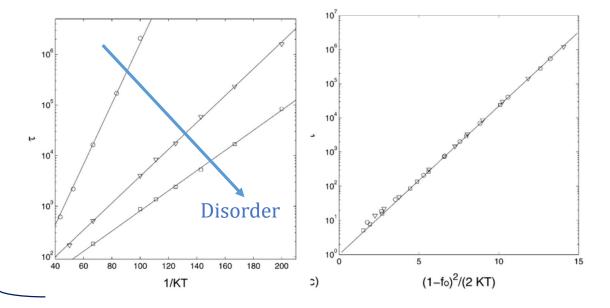
1.
$$F = \sum_{i=1}^{N} f_i$$
.

2. $f_i = Ye_i$,

3. $f_c(i) \sim N_d(\langle f_c \rangle, KT_d)$.

 $\tau \simeq \tau_0 \exp\left(\frac{(1-f_0)^2}{2KT_{eff}}\right)$
 $\tau \simeq \tau_0 \exp\left(\frac{(1-f_0)^2}{2KT_{eff}}\right)$
 $\tau \simeq \tau_0 \exp\left(\frac{(1-f_0)^2}{2KT_{eff}}\right)$
 $\tau_0 = \frac{2\sqrt{2\pi KT}}{(f_0 - \sqrt{\pi}\sigma_0)[1 + \exp(-\sqrt{\pi}\sigma_0(1-f_0)/KT)]}$.

$$au \simeq \frac{\sqrt{2\pi KT}}{f_0} \exp\left(\frac{(1-f_0)^2}{2KT}\right)$$



The effect of disorder on the fracture nucleation process (S. Ciliberto, A. Guarino, R. Scorreti (2001))







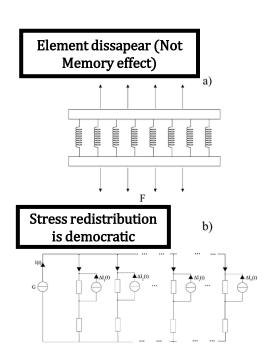
Method: DFBM

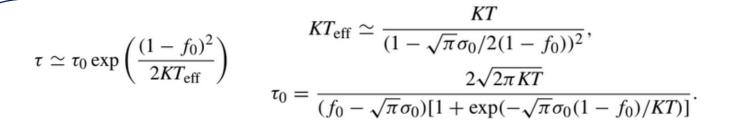
$$1 \quad F = \sum_{i=1}^{N} f_i.$$

$$2 f_i = Ye_i$$

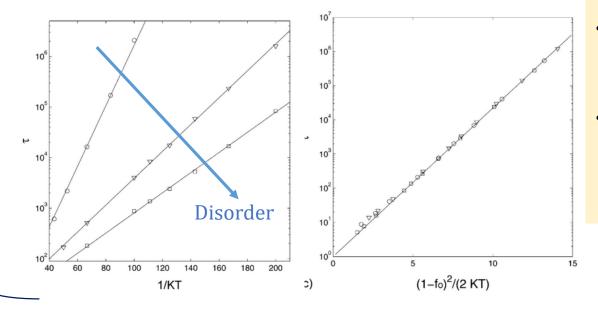
3
$$f_c(i) \sim N_d(\langle f_c \rangle, KT_d)$$
.

4.
$$\Delta f_i(t) \sim N_{\rm T}(0, KT)$$





$$\tau \simeq \frac{\sqrt{2\pi KT}}{f_0} \exp\left(\frac{(1-f_0)^2}{2KT}\right)$$



- Disorder of the material « amplifies » the thermal noise and produces an effecive temperature
- The failure time decreases as the disorder noise increases (the more the medius heterogeneous, the smaller the failure time is)
- As the disorder noise increases, the derivative of failure time with respect to KT decreases (failure time becomes less sensitive to the effective value of thermal noise.

$$kT_{eff} \ge kT$$

2. Methods: PDM + KMC

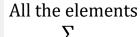








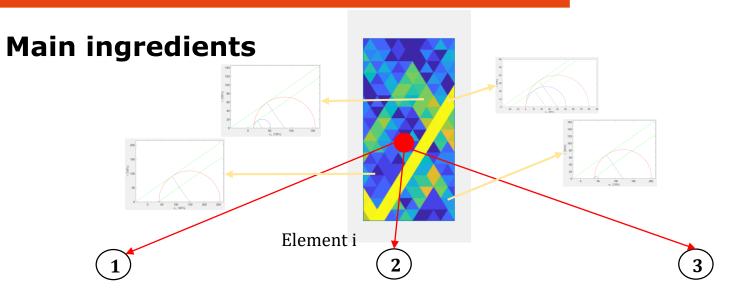








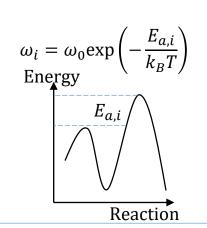
$$\sigma_{ej} = E_{ejkl} \varepsilon_{ej}$$



Heterogeneity

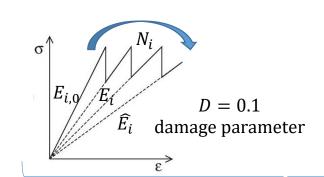
Cohesion C_i is set randomly Gaussian distribution (quenched) $C_i \in [10, 30] MPa$

Thermal activation



 $\widehat{E_i} = (1 - D)E_i$

Damage



initial disorder (random variable) + stochastic process

Athermal damage (MC criterium)

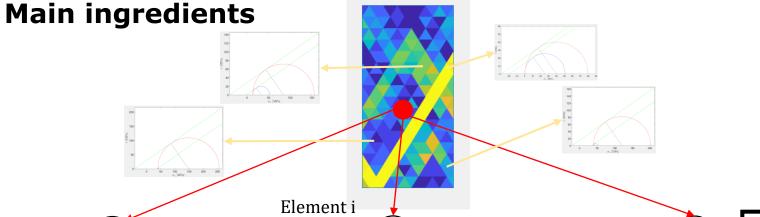
Thermal activated damage (Probability)

2. Methods: PDM + KMC





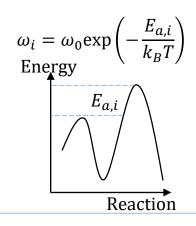




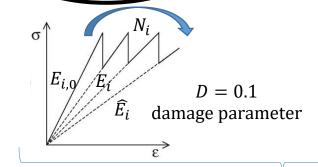


Cohesion C_i is set randomly Gaussian distribution (quenched) $C_i \in [10, 30]$ MPa

Thermal activation



Damage Element does not dissapear (Memory effect) $\widehat{E_i} = (1 - D)E_i$





All the elements

Stress redistribution is not democratic

Elastic interaction (Solved with FEM)

 $\sigma_{ej} = E_{ejkl} \varepsilon_{ej}$

Athermal damage (MC criterium)

initial disorder (random variable) + stochastic process

Thermal activated damage (Probability)



$$\omega_{i} = \omega_{0} \exp\left(-\frac{E_{a,i}}{k_{B}T}\right)$$

$$R_{2} = \sum \omega_{i}$$

$$\Delta t = -\frac{\log u_{2}}{R_{2}}$$

 $t = t + \Delta t$

$$\omega_{i} = \omega_{0} \exp\left(-\frac{E_{a,i}}{k_{B}T}\right) \qquad \qquad t_{f} = -\frac{U}{N\omega_{0}} \int_{U} \log(u) p(u) \left[\int_{\Delta\sigma(u)} \exp\left(-\frac{V_{a,0}\Delta\sigma}{k_{B}T}\right) p(\Delta\sigma) d(\Delta\sigma)\right]^{-1} du,$$

We impose the activation volume V_a



$$\omega_{i} = \omega_{0} \exp\left(-\frac{E_{a,i}}{k_{B}T}\right)$$

$$R_{2} = \sum_{i} \omega_{i}$$

$$\Delta t = -\frac{\log u_{2}}{R_{2}}$$

$$t = t + \Delta t$$

$$Disorder effect$$

$$t_{f} \sim \frac{1}{\omega_{0}} \exp\left(\frac{V_{a}\langle \Delta \sigma \rangle}{k_{B}T}\right) f(\delta)$$









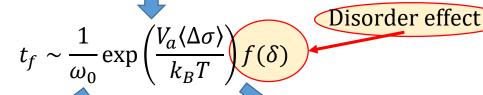
$$\omega_{i} = \omega_{0} \exp\left(-\frac{E_{a,i}}{k_{B}T}\right)$$

$$R_{2} = \sum \omega_{i}$$

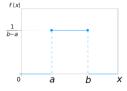
$$\Delta t = -\frac{\log u_{2}}{R_{2}}$$

$$t = t + \Delta t$$

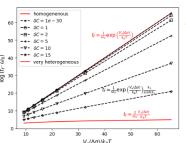
$$t_f = -\frac{U}{N\omega_0} \int_U \log(u) p(u) \left[\int_{\Delta\underline{\sigma(u)}} \exp\left(-\frac{V_{a,0}\Delta\sigma}{k_B T}\right) p(\Delta\sigma) d(\Delta\sigma) \right]^{-1} du,$$



We impose the activation volume V_a

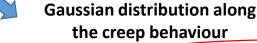


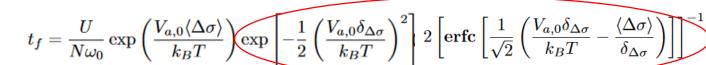
Uniform distribution along the creep behaviour



$$t_f = \frac{1}{\omega_0} \exp\left(\frac{V_{a,0}\langle \Delta \sigma \rangle}{k_B T}\right) \frac{U}{N} \frac{k_1}{\sinh(k_1)}.$$

$$k_1 = Va\delta/k_BT$$











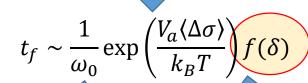
$$\omega_{i} = \omega_{0} \exp\left(-\frac{E_{a,i}}{k_{B}T}\right)$$

$$R_{2} = \sum \omega_{i}$$

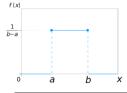
$$\Delta t = -\frac{\log u_{2}}{R_{2}}$$

$$t = t + \Delta t$$

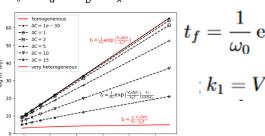
$$t_f = -\frac{U}{N\omega_0} \int_U \log(u) p(u) \left[\int_{\Delta\underline{\sigma}(u)} \exp\left(-\frac{V_{a,0}\Delta\sigma}{k_B T}\right) p(\Delta\sigma) d(\Delta\sigma) \right]^{-1} du,$$



We impose the activation volume V_a

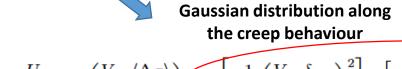


Uniform distribution along the creep behaviour

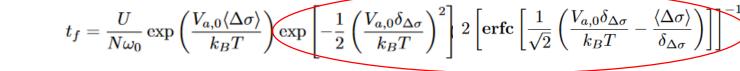


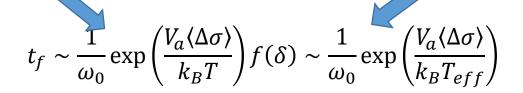
$$t_f = \frac{1}{\omega_0} \exp\left(\frac{V_{a,0}\langle \Delta \sigma \rangle}{k_B T}\right) \frac{U}{N} \frac{k_1}{\sinh(k_1)}.$$

$$k_1 = Va\delta/k_BT$$



Disorder effect

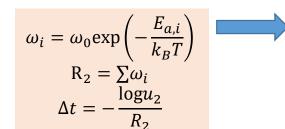












 $t = t + \Delta t$

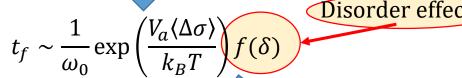
$$t_f = -\frac{U}{N\omega_0} \int_U \log(u) p(u) \left[\int_{\Delta\underline{\sigma}(u)} \exp\left(-\frac{V_{a,0}\Delta\sigma}{k_B T}\right) p(\Delta\sigma) d(\Delta\sigma) \right]^{-1} du,$$

$$t_f \sim \frac{1}{\omega_0} \exp\left(\frac{V_a \langle \Delta c \rangle}{k_B T}\right)$$

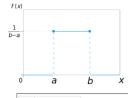
Disorder effect

Gaussian distribution along

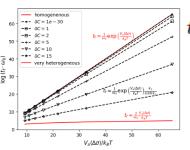
the creep behaviour





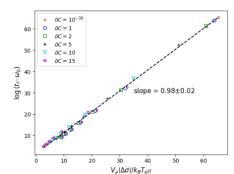


Uniform distribution along the creep behaviour

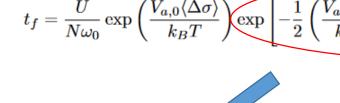


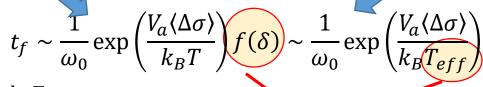
$$t_f = \frac{1}{\omega_0} \exp\left(\frac{V_{a,0}\langle \Delta \sigma \rangle}{k_B T}\right) \frac{U}{N} \frac{k_1}{\sinh(k_1)}. \qquad t_f = \frac{U}{N\omega_0} \exp\left(\frac{V_{a,0}\langle \Delta \sigma \rangle}{k_B T}\right) \exp\left[-\frac{1}{2} \left(\frac{V_{a,0}\delta_{\Delta \sigma}}{k_B T}\right)^2\right] 2 \left[\operatorname{erfc}\left[\frac{1}{\sqrt{2}} \left(\frac{V_{a,0}\delta_{\Delta \sigma}}{k_B T} - \frac{\langle \Delta \sigma \rangle}{\delta_{\Delta \sigma}}\right)\right]\right]$$

$$k_1 = Va\delta/k_BT$$
 ,

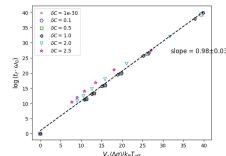


$$t_f \sim \frac{1}{\omega_0} \exp\left(\frac{k_B T}{1 - \frac{\delta}{\langle \Delta \sigma \rangle}}\right)$$







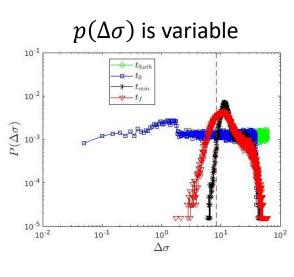


Expression for failure time? – reality: complexity to track the evolution of the stress gap along the deformation







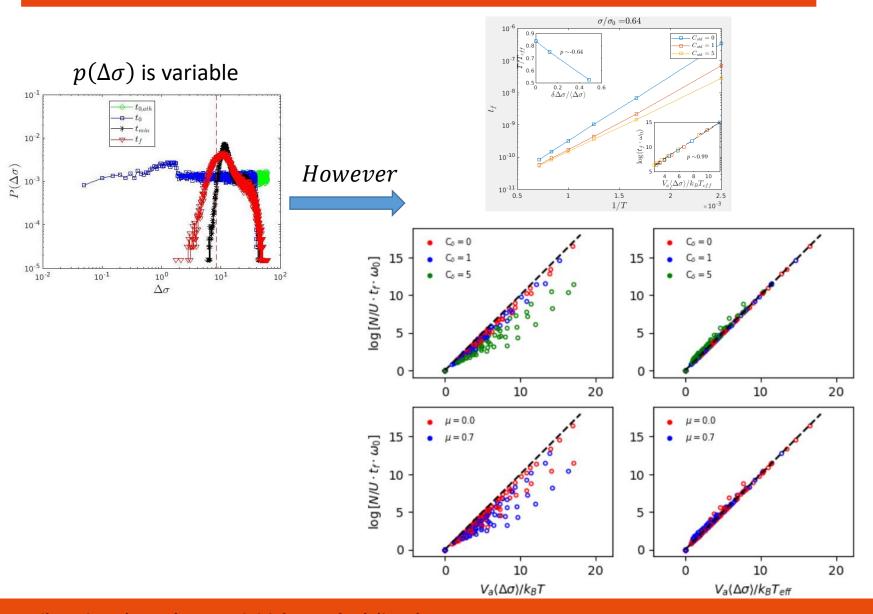


Expression for failure time? - reality: complexity to track the evolution of the stress gap along the deformation







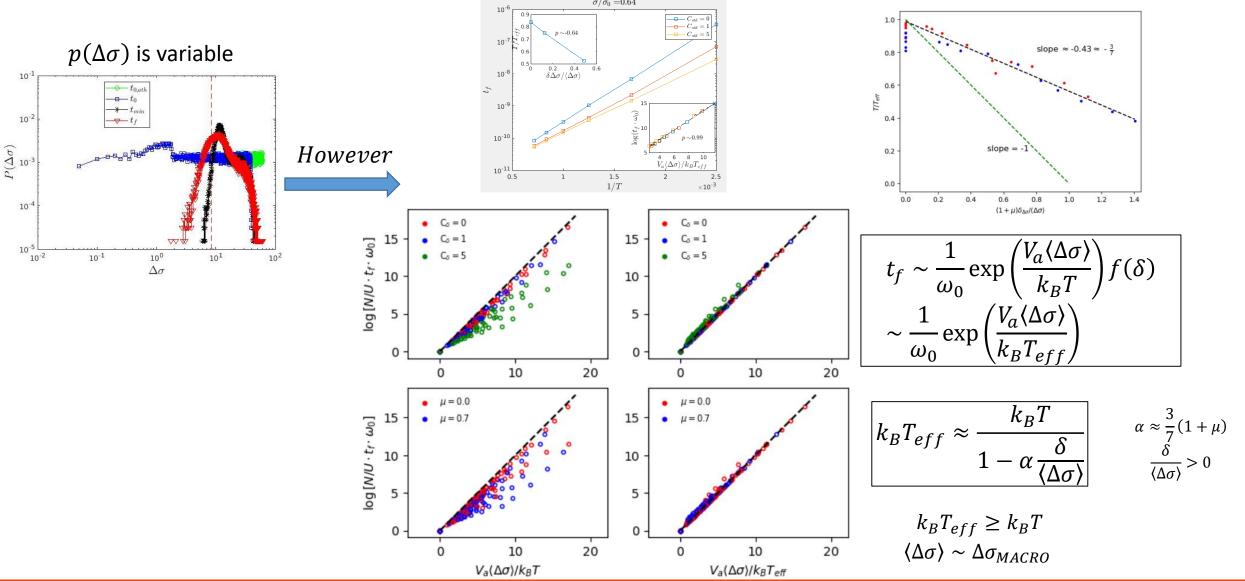


Expression for failure time? - reality: complexity to track the evolution of the stress gap along the deformation









Expression for failure time? – reality: complexity to track



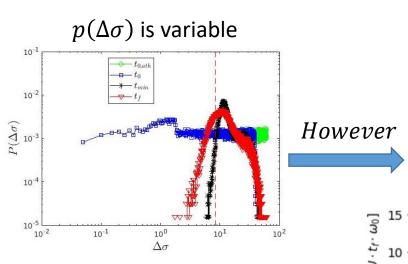




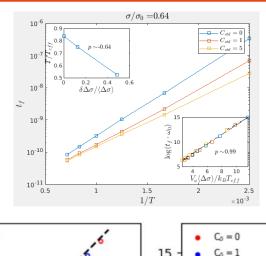


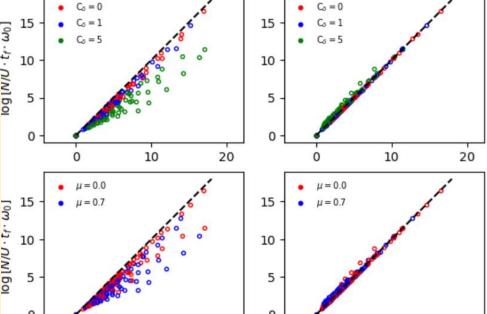
10

 $V_a(\Delta\sigma)/k_BT$



- Initial quenched disorder of the material « amplifies » the temperature and produces an effecive one
- The failure time decreases as the disorder increases (the more the medius heterogeneous, the smaller the failure time is)
- As the disorder increases, the derivative of failure time with respect to KT decreases (failure time becomes less sensitive to the effective value of thermal noise.



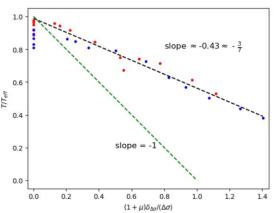


20

10

 $V_a(\Delta\sigma)/k_BT_{eff}$

20



$$t_f \sim \frac{1}{\omega_0} \exp\left(\frac{V_a \langle \Delta \sigma \rangle}{k_B T}\right) f(\delta)$$
$$\sim \frac{1}{\omega_0} \exp\left(\frac{V_a \langle \Delta \sigma \rangle}{k_B T_{eff}}\right)$$

$$k_B T_{eff} \approx \frac{k_B T}{1 - \alpha \frac{\delta}{\langle \Delta \sigma \rangle}}$$

$$\alpha \approx \frac{3}{7}(1+\mu)$$
$$\frac{\delta}{\langle \Delta \sigma \rangle} > 0$$

$$k_B T_{eff} \ge k_B T$$

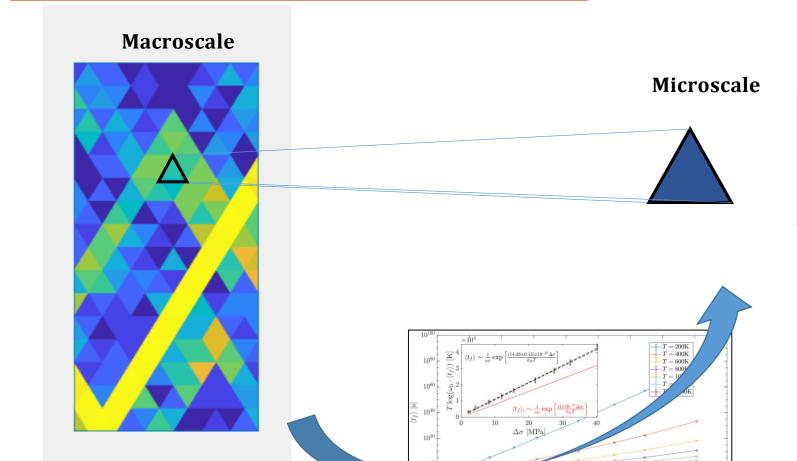
 $\langle \Delta \sigma \rangle \sim \Delta \sigma_{MACRO}$

Problem:

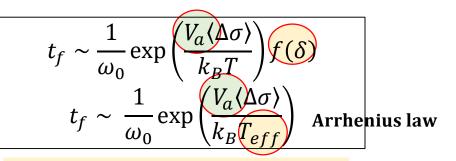








Downscaling



The obtention of the activation volume at macroscale is more complex than only use the Arrhenius expression

It is important to include a disorder effect to compute the real activation volume





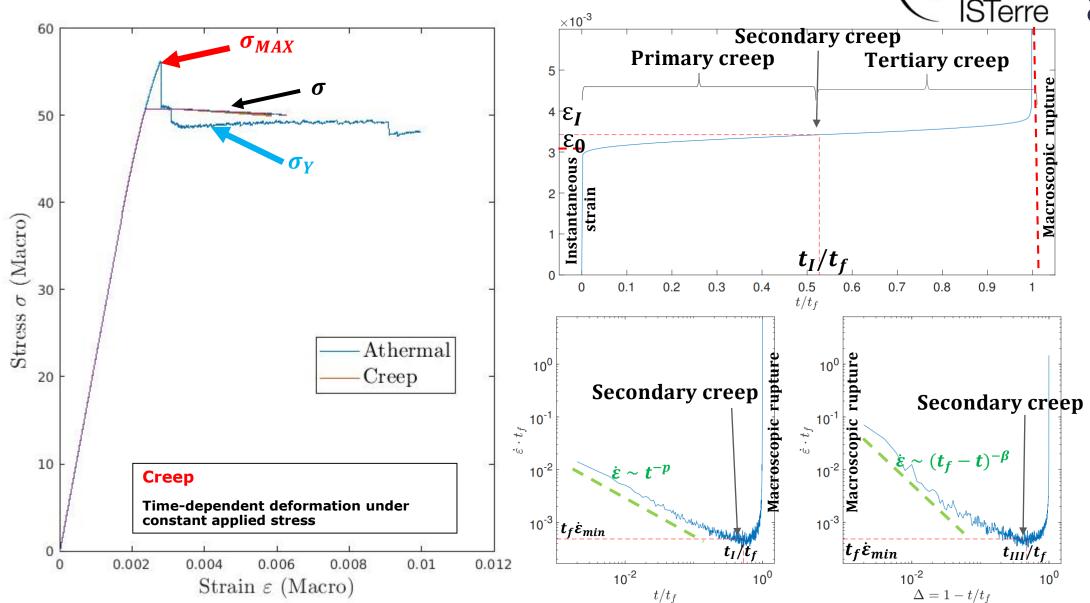


Quenched initial uniform disorder

1. Presentation of the problem







Q01

4.4 Acceleration of deformation during the tertiary creep





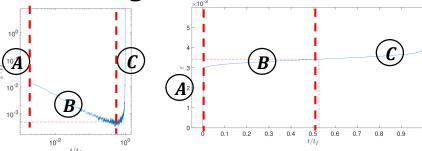
Spatial correlation of damage events

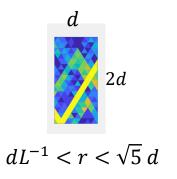
System Sizes L = 128

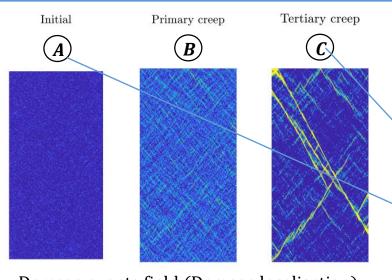
Number of simulations N=1

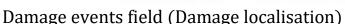
Temperature T = 800 K

 $\sigma/\sigma_V = 0.72 \quad (\sigma = 0.72\sigma_V)$ Stress ratio

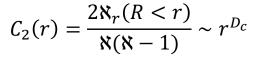








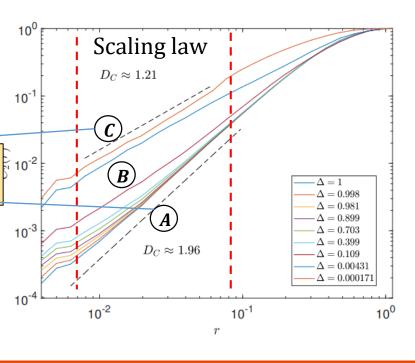
$$rac{E_{i,0}-E_i}{E_{i,0}}=1-(1-D)^{\mathcal{N}_i}$$
 0.2 0.4 0.6 0.8 1



Hirata et al., 1987

Fractal behaviour of damage events??

Perfect damage localisation Homogeneous damage



4.4 Acceleration of deformation during the tertiary creep





Stress gap distribution

System Sizes L = 128

Number of simulations N=1

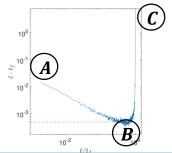
Temperature

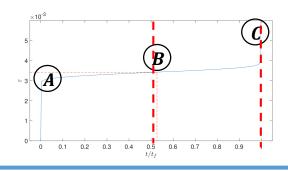
T = 800 K

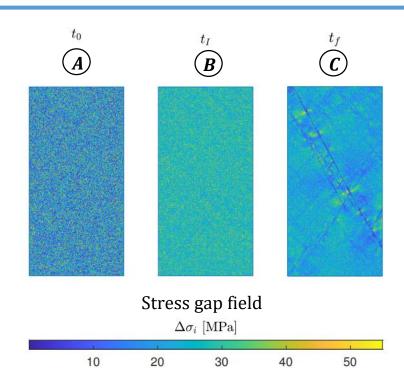
Stress ratio

 $\sigma/\sigma_Y = 0.72 \ (\sigma = 0.72\sigma_Y)$

Macro stress gap $\Delta \sigma \approx 11 \text{ MPa}$

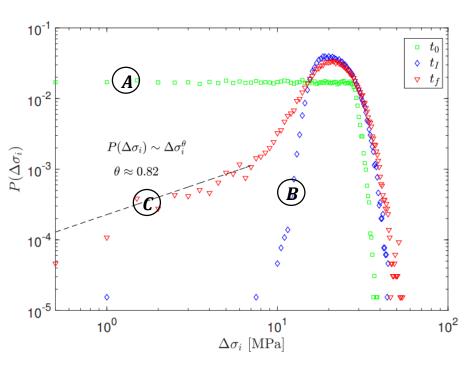






 $P(\Delta\sigma_i) \sim \Delta\sigma_i^{\theta}$ Lin et al., 2014

Power law tail towards small stress gaps



4.1 Size effect of the stress gap



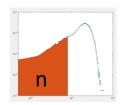


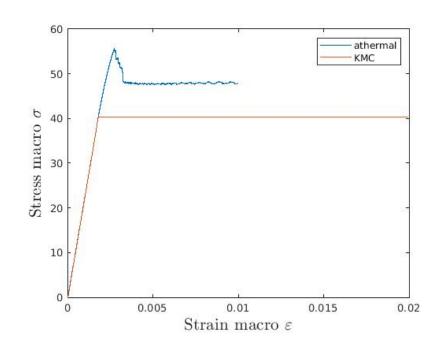
System Sizes
Number of simulations
Temperature
Stress ratio
Macro stress gap

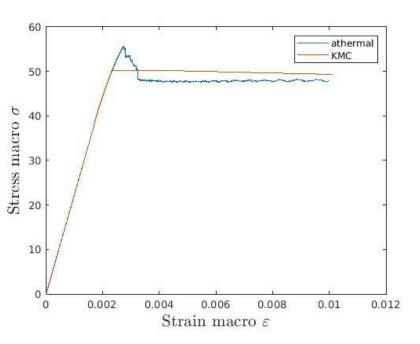
 $L = \{8,16,32,64\}$ $N = \{291,68,17,4\}$ T = 800 K

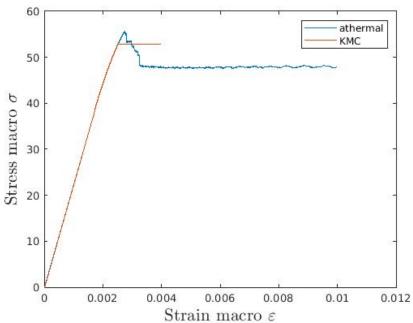
 $N = \{291, 68, 17, T = 800 \text{ K}$ $\sigma/\sigma_{MAX} \quad \sigma/\sigma_{Y}$

Size effect of the stress gap

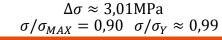




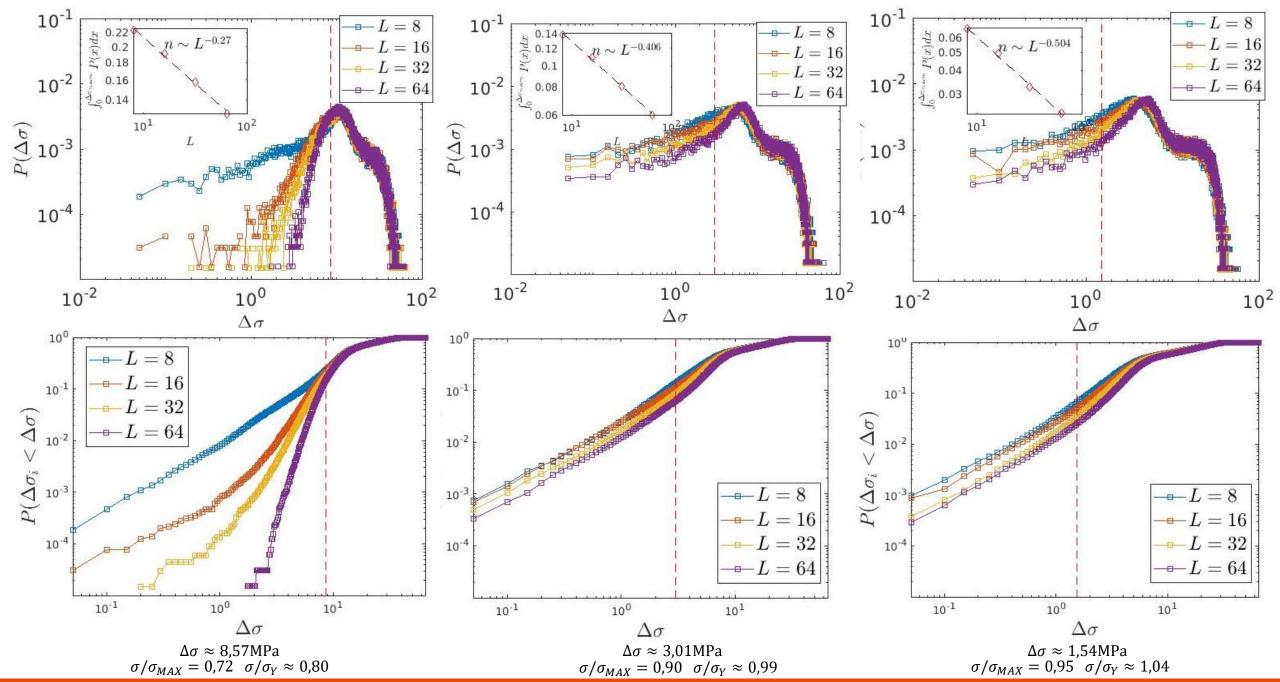




 $\Delta\sigma \approx 8.57 \mathrm{MPa}$ $\sigma/\sigma_{MAX} = 0.72 \ \sigma/\sigma_{Y} \approx 0.80$



 $\Delta\sigma \approx 1,54$ MPa $\sigma/\sigma_{MAX} = 0.95 \ \ \sigma/\sigma_{Y} \approx 1.04$



4.4 Stress ratio effect

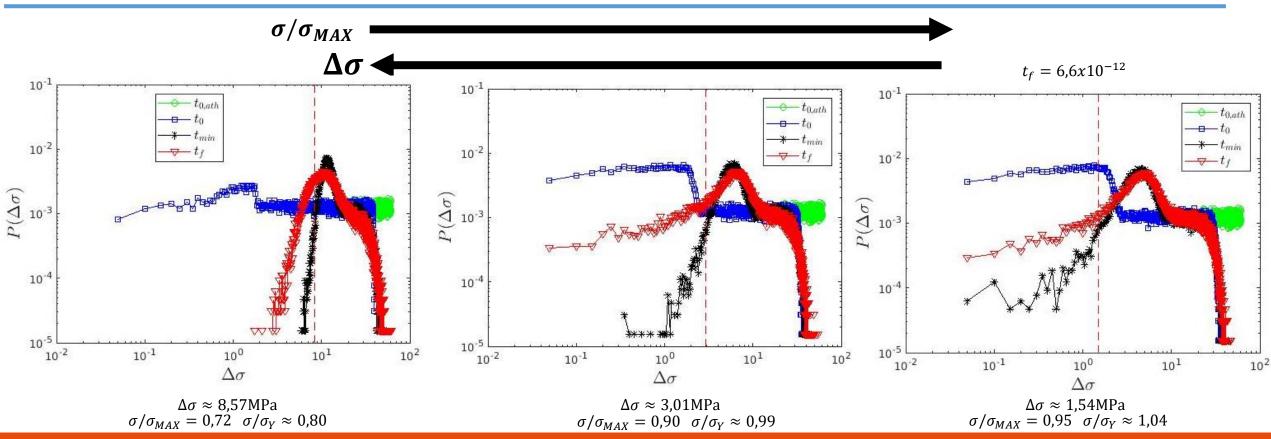


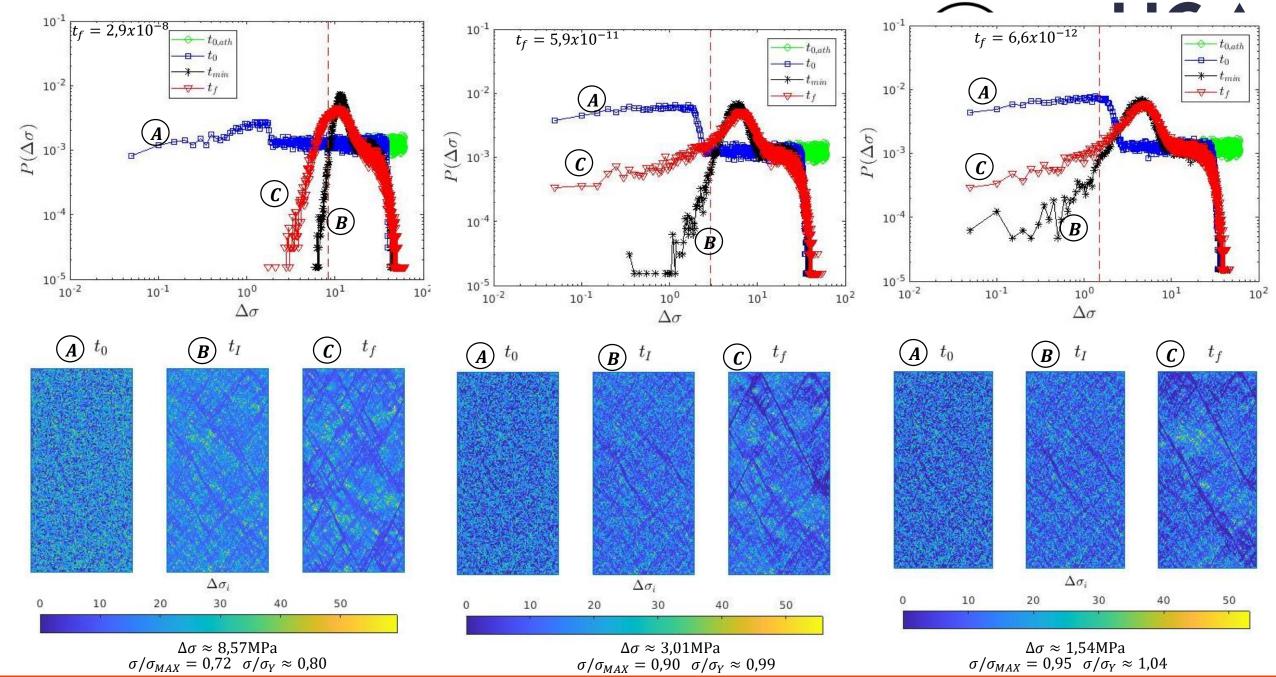


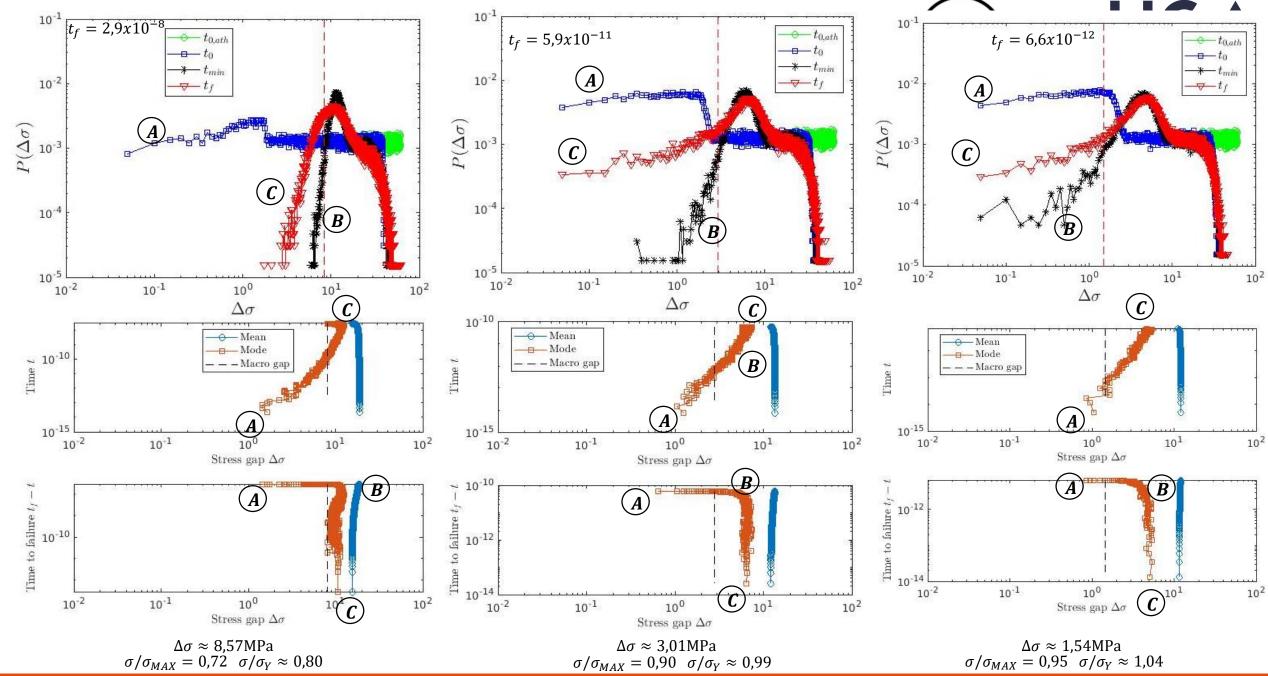
System Sizes L = 64**Number of simulations** N = 1T = 800 K**Temperature**

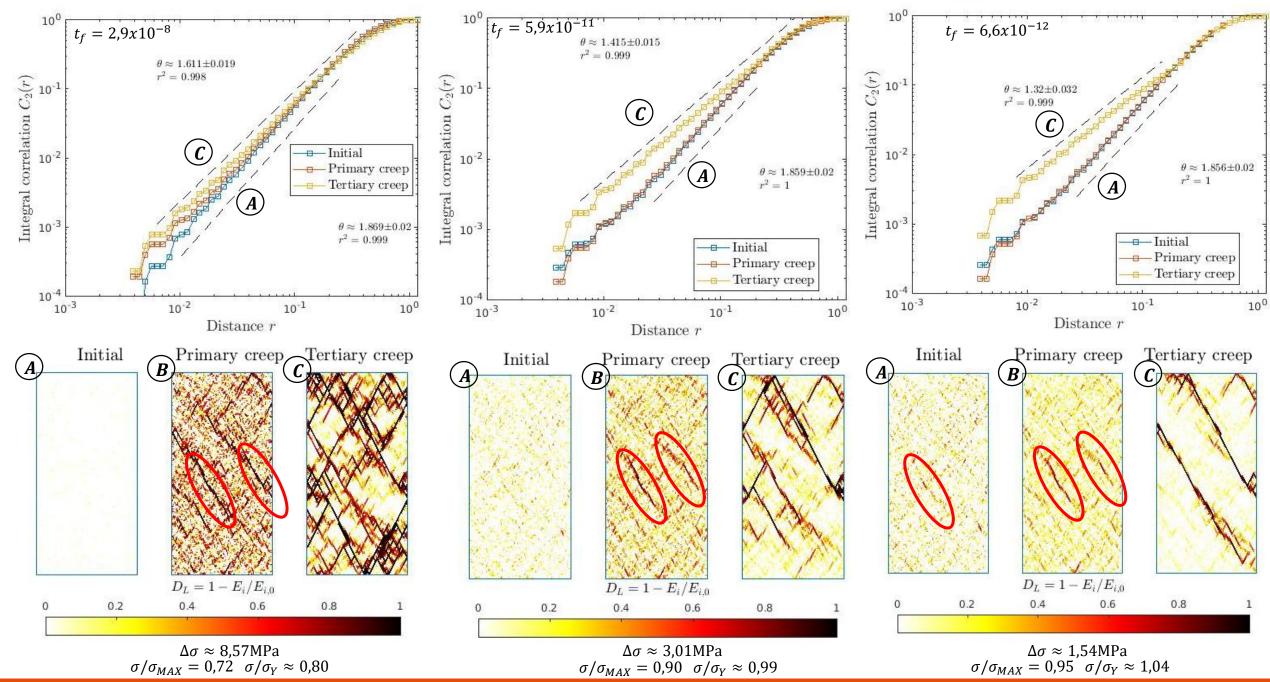
Stress ratio σ/σ_{MAX} σ/σ_{Y} Macro stress gap

Stress ratio dependence









4.4 Effect of thermal activation (same initial disorder)



KMC effect: same

initial disorder



System Sizes

L = 64

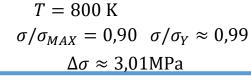
Number of simulations

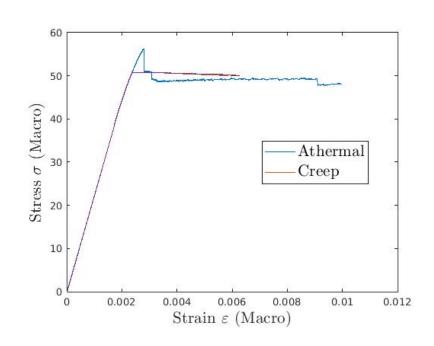
N = 10

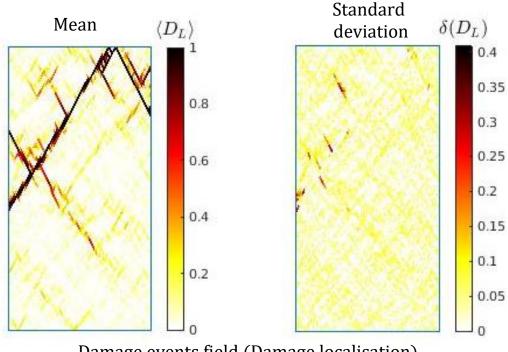
Temperature

Stress ratio

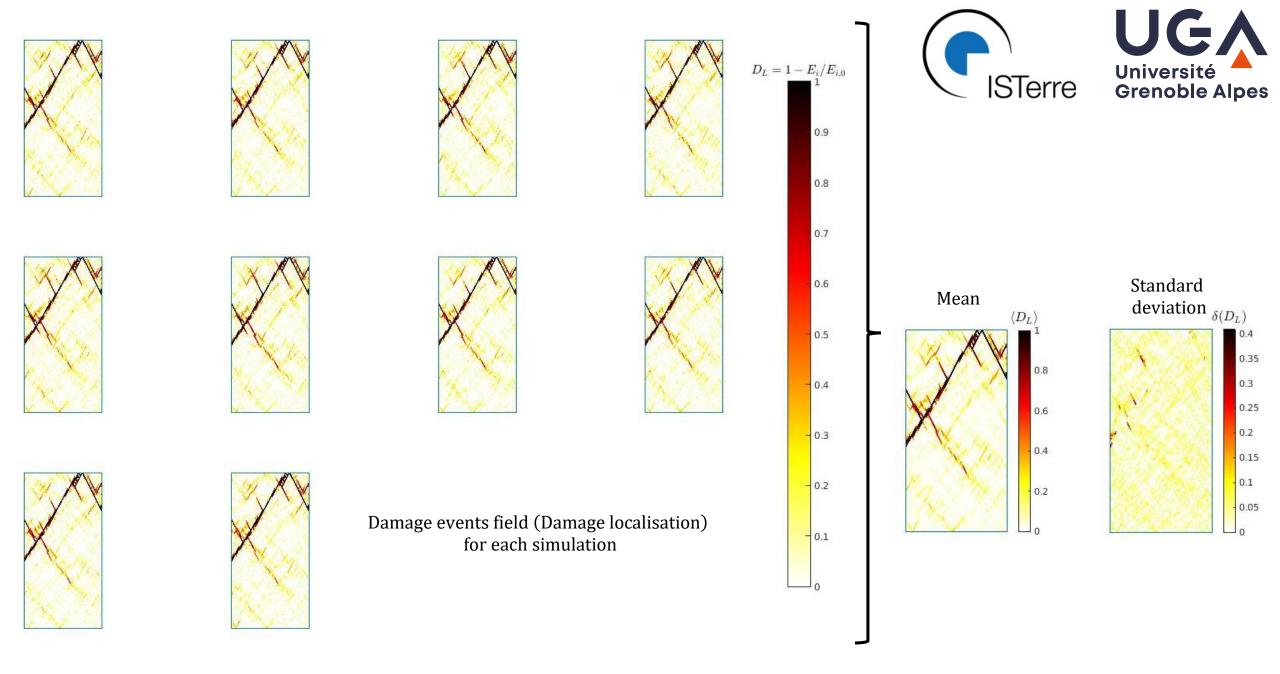
Macro stress gap

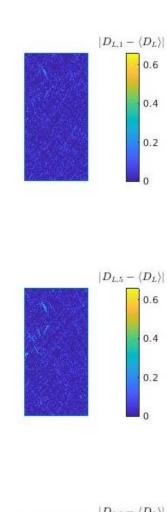


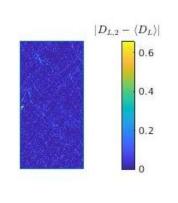


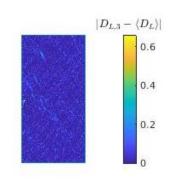


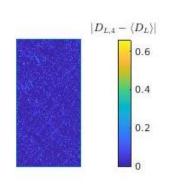
Damage events field (Damage localisation)





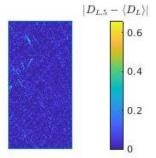


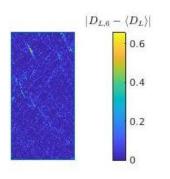


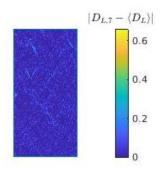


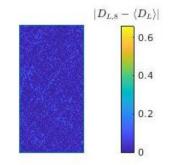


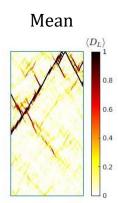


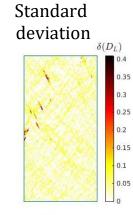


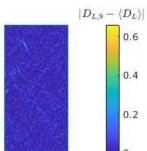


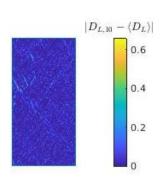




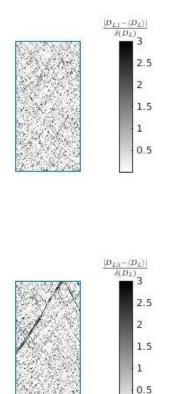


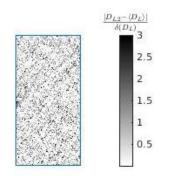


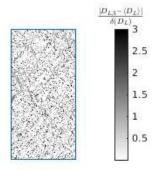


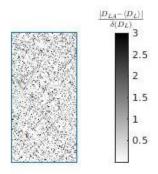


Damage events field (Damage localisation) for each simulation Damaged field — Mean of Damaged field $D_L - < D_L >$







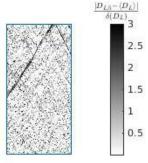


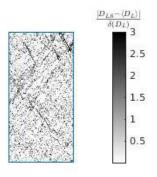


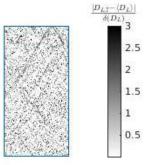


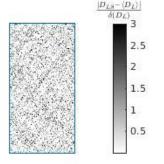
 $\delta(D_L)$

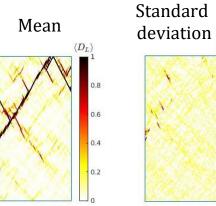
0.05

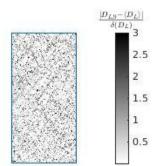


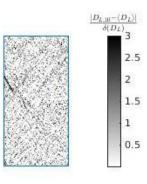






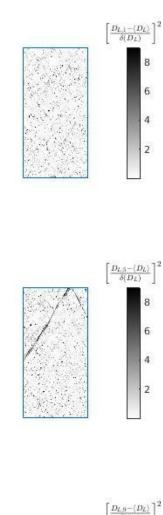


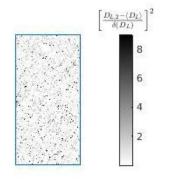


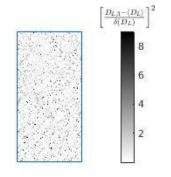


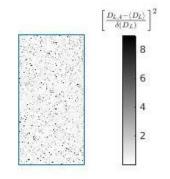
Damage events field (Damage localisation) for each simulation $\mathrm{D_L} - < \mathrm{D_L} >$

$$z = \frac{D_L - \langle D_L \rangle}{\delta D_L}$$





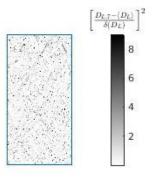


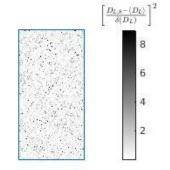


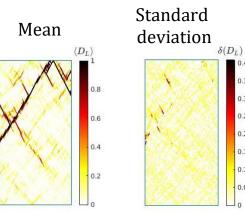


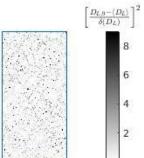


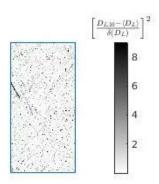












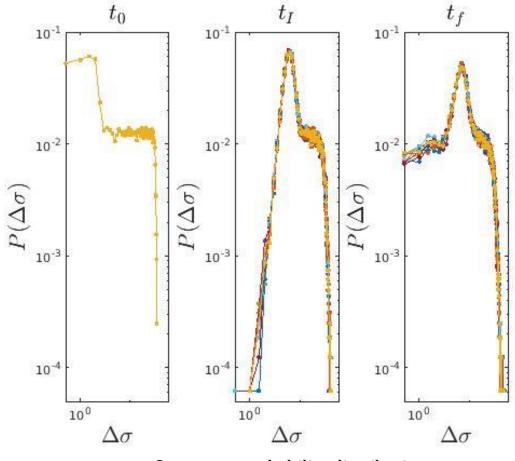
Damage events field (Damage localisation) for each simulation

$$z^2 = \left(\frac{D_L - \langle D_L \rangle}{\delta D_L}\right)^2$$

0.05





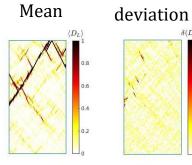


Integral correlation C_r 10-1 10⁻² 10⁻³ 10⁻² 10⁻¹ 10⁰ Distance r

Integral correlations

Stress gap probability distributions

Modelling the creep of brittle materials



10⁰

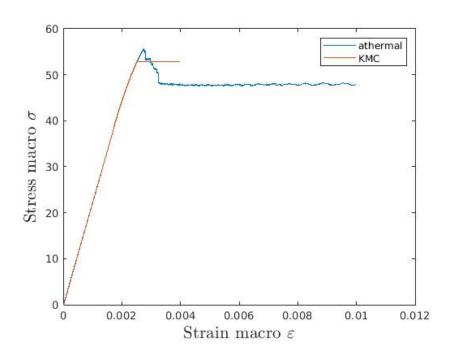
Standard

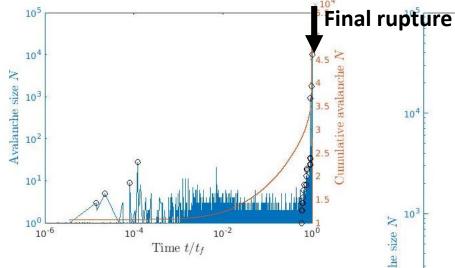
4.4 Level of damaged during the largest avalanche size

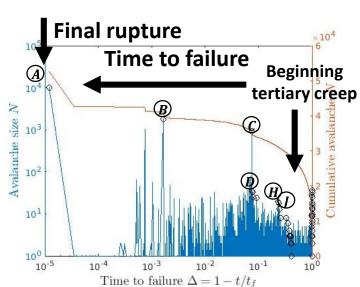


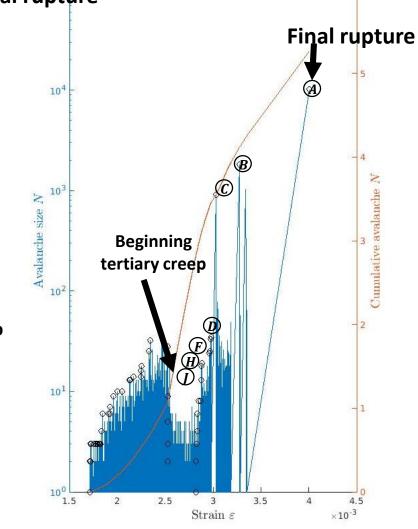


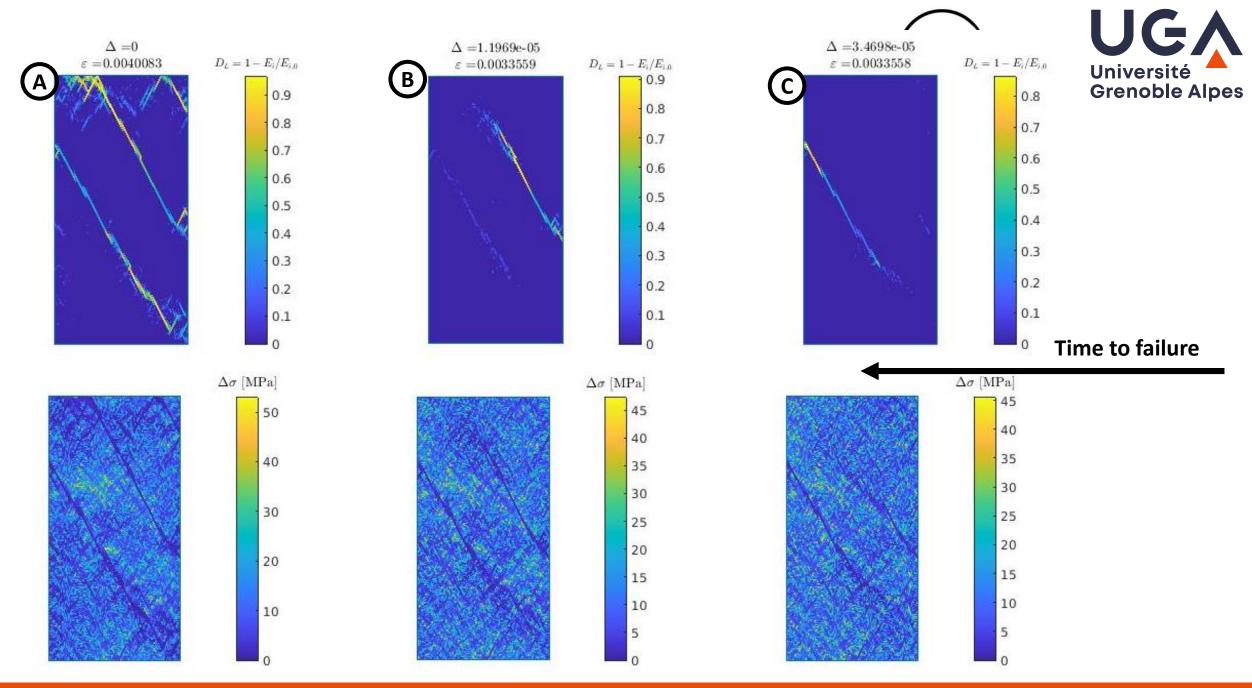
 $\begin{array}{lll} \textbf{System Sizes} & L = 64 \\ \textbf{Number of simulations} & N = 1 \\ \textbf{Temperature} & T = 800 \text{ K} \\ \textbf{Stress ratio} & \sigma/\sigma_{MAX} = 0.95 & \sigma/\sigma_{Y} \approx 1.04 \\ \textbf{Micro/Macro stress gap} & \Delta\sigma \approx 1.54 \text{MPa} \end{array}$

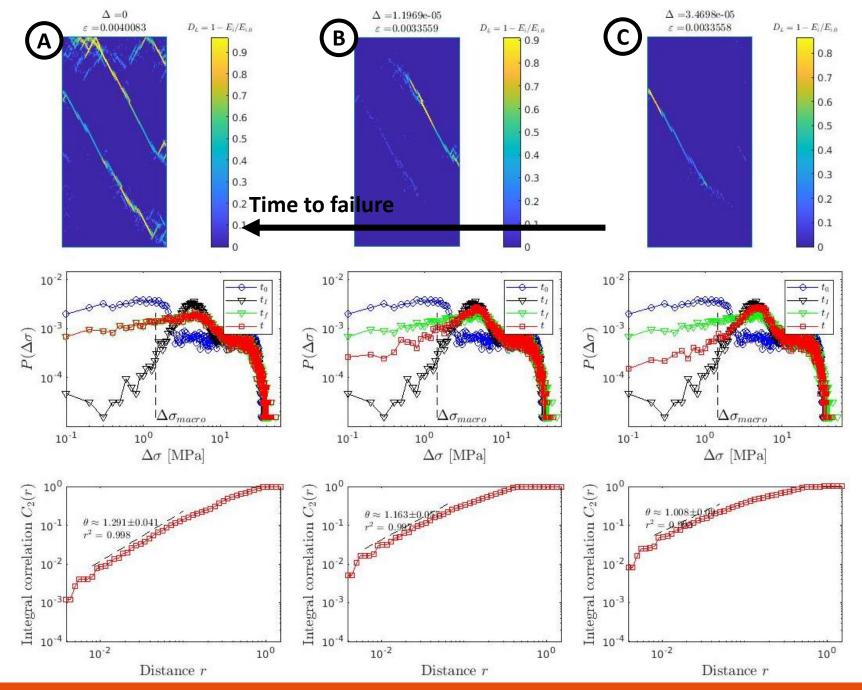






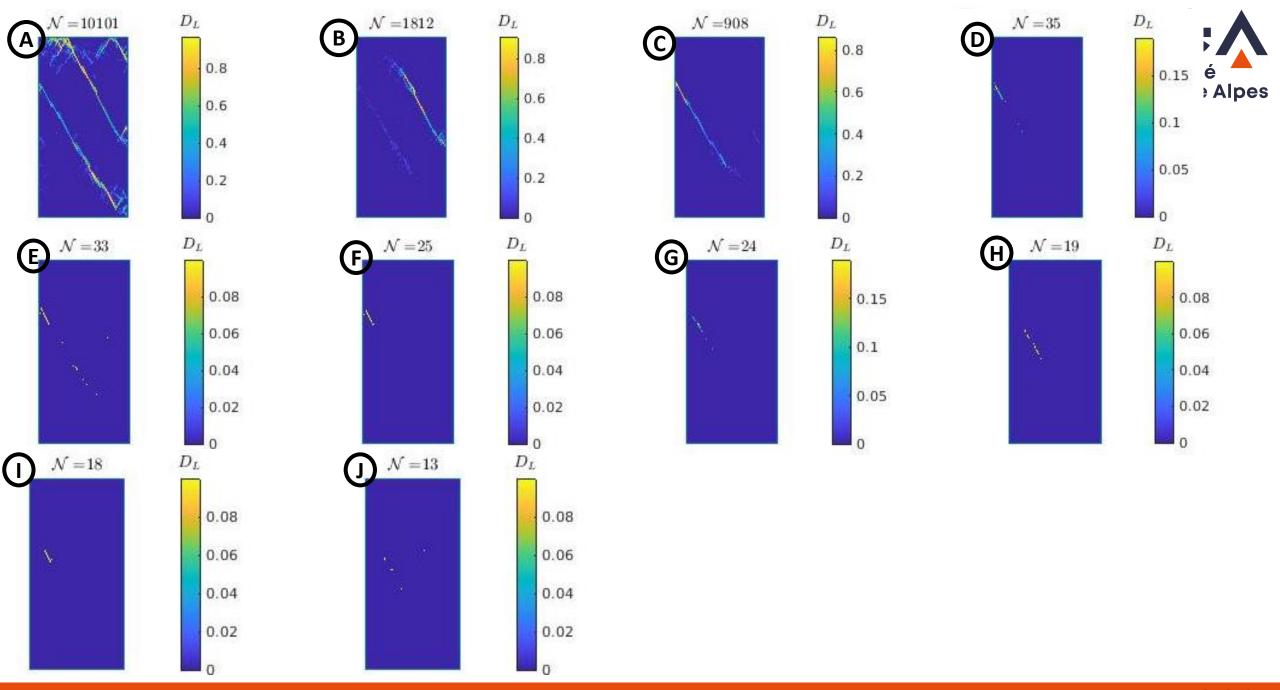












4.4 Level of damaged in different windows of time





 $\begin{array}{ccc} \textbf{System Sizes} & L = 64 \\ \textbf{Number of simulations} & N = 1 \\ \textbf{Temperature} & T = 800 \text{ K} \\ \textbf{Stress ratio} & \sigma/\sigma_{MAX} = 0.95 & \sigma/\sigma_{Y} \approx 1.04 \\ \textbf{Micro/Macro stress gap} & \Delta\sigma \approx 1.54 \text{MPa} \end{array}$

