
1 NMR Shielding

Shielding Tensor:

$$\sigma_{\beta\alpha}^K = \left. \frac{d^2 E}{dB_\alpha dm_{K\beta}} \right|_{\mathbf{B}, \mathbf{m}_K=0} \quad (1)$$

How do I parameterize energy E with \mathbf{B} and \mathbf{m}_K ?

The one-electronic Hamiltonian in magnetic field:

$$h(\mathbf{r}, \mathbf{B}, \mathbf{m}) = \frac{1}{2} \boldsymbol{\pi}^2 - \phi(\mathbf{r}) \quad (2)$$

in which:

$$\boldsymbol{\pi} = -i\nabla + \mathbf{A} \quad (3)$$

is the kinetic momentum operator.

Vector potential:

$$\mathbf{A}_i = \mathbf{A}_0(\mathbf{r}_i) + \sum_K \mathbf{A}_K(\mathbf{r}_i) \quad (4)$$

with:

$$\mathbf{A}_0(\mathbf{r}_i) = \frac{1}{2} \mathbf{B} \times \mathbf{r}_0 \quad \mathbf{B} = \nabla \times \mathbf{A}(\mathbf{r}) \quad (5)$$

$$\mathbf{A}_K(\mathbf{r}_i) = \alpha^2 \frac{\mathbf{M}_K \times \mathbf{r}_K}{r_K^3} \quad \mathbf{B}_K(\mathbf{r}) = \nabla \times \mathbf{A}_K(\mathbf{r}) \quad (6)$$

The first part is contribution from the external magnetic field, the second part from the nuclear magnetic moments.

Now parameterize with MO coefficients / densities?

2 SCF Level

$$\begin{aligned}
 E^{\text{SCF}} &= \sum_i^N h_{ii} + \frac{1}{2} \sum_{ij}^N \langle ij || ij \rangle \\
 &= \sum_i \sum_{\mu\nu} C_{\mu i}^* h_{\mu\nu} C_{\nu i} + \frac{1}{2} \sum_{ij} \langle ij || ij \rangle
 \end{aligned} \tag{7}$$

$$D_{\mu\nu}^{\text{SCF}} = \sum_i C_{\mu i}^* C_{\nu i} \tag{8}$$

$$C_{\mu p}(\lambda) = \sum_q C_{\mu q}(0) U_{qp}(\lambda) \tag{9}$$

At SCF level, the NMR shielding tensor is given as:

$$\sigma_{\beta\alpha}^{\text{SCF},K} = \left. \frac{d^2 E^{\text{SCF}}}{dB_\alpha dm_{K_\beta}} \right|_{\mathbf{B}, \mathbf{m}_K=0} \tag{10}$$

Taking the first derivative against the nuclear magnetic moment gives:

$$\begin{aligned}
 \frac{dE^{\text{SCF}}}{dm_{K_\beta}} &= \frac{d}{dm_{K_\beta}} \left(\sum_{i\mu\nu} C_{\mu i}^* h_{\mu\nu} C_{\nu i} \right) \\
 &= \sum_{i\mu\nu} C_{\mu i}^* C_{\nu i} \frac{dh_{\mu\nu}}{dm_{K_\beta}} \\
 &= \sum_{\mu\nu} D_{\mu\nu}^{\text{SCF}} \frac{dh_{\mu\nu}}{dm_{K_\beta}}
 \end{aligned} \tag{11}$$

Note that the MO coefficients are variationally determined so $\frac{d\mathbf{C}}{dm_{K_\beta}} = \mathbf{0}$, and the basis function does not depend on the nuclear magnetic moment, i.e. $\frac{d\phi_\mu}{dm_{K_\beta}} = 0$.

Now taking the second derivative w.r.t. the external magnetic field:

$$\begin{aligned}
 \sigma_{\beta\alpha}^{\text{SCF},K} &= \frac{d^2 E^{\text{SCF}}}{dB_\alpha dm_{K_\beta}} \\
 &= \frac{d}{dB_\alpha} \left(\sum_{\mu\nu} D_{\mu\nu}^{\text{SCF}} \frac{dh_{\mu\nu}}{dm_{K_\beta}} \right) \\
 &= \sum_{\mu\nu} D_{\mu\nu}^{\text{SCF}} \frac{d^2 h_{\mu\nu}}{dB_\alpha dm_{K_\beta}} + \sum_{\mu\nu} \frac{dD_{\mu\nu}^{\text{SCF}}}{dB_\alpha} \frac{dh_{\mu\nu}}{dm_{K_\beta}}
 \end{aligned} \tag{12}$$

The response of SCF density to the magnetic field perturbation is:

$$\begin{aligned}
 \frac{dD_{\mu\nu}^{\text{SCF}}}{dB_\alpha} &= \frac{d}{dB_\alpha} \left(\sum_i C_{\mu i}^*(\mathbf{B}) C_{\nu i}(\mathbf{B}) \right) \\
 &= \frac{d}{dB_\alpha} \left(\sum_{ipq} C_{\mu p}^*(0) U_{pi}^*(\mathbf{B}) C_{\nu q}(0) U_{qi}(\mathbf{B}) \right) \\
 &= \sum_{ip} C_{\mu p}^*(0) \frac{dU_{pi}^*(\mathbf{B})}{dB_\alpha} C_{\nu i}(\mathbf{B}) + \sum_{iq} C_{\mu i}^*(\mathbf{B}) \frac{dU_{qi}(\mathbf{B})}{dB_\alpha} C_{\nu q}(0)
 \end{aligned} \tag{13}$$

$$= \sum_{ip} C_{\mu p}^* (U_{pi}^{B_\alpha})^* C_{\nu i} + \sum_{ip} C_{\mu i}^* U_{pi}^{B_\alpha} C_{\nu p} \quad (14)$$

The virtual-occupied block of \mathbf{U}^B is obtained from the CPSCF equations, and the occupied-occupied block is chosen according to the orthonormality condition.

3 MP2 Level

$$\mathcal{L}_{\text{MP2}} = E_{\text{HF}} + E_{\text{H}} + \frac{1}{2} \sum_{ai} (z_{ai} f_{ai} + z_{ai}^* f_{ai}^*) \quad (15)$$

$$E_{\text{HF}} = \sum_{\mu\nu} h_{\mu\nu} D_{\mu\nu}^{\text{HF}} + \frac{1}{2} \sum_{ij} \langle ij || ij \rangle \quad (16)$$

$$\begin{aligned} E_{\text{H}} &= \sum_{ij} h_{ij} \gamma_{ij}^{\text{H}} + \sum_{ab} h_{ab} \gamma_{ab}^{\text{H}} + \sum_{pqrs} \langle pq || rs \rangle \Gamma_{rs}^{pq} \\ &= \sum_{\mu\nu} D_{\mu\nu}^{\text{H}} h_{\mu\nu}^{\text{AO}} + \sum_{pqrs} \langle pq || rs \rangle \Gamma_{rs}^{pq} \end{aligned} \quad (17)$$

$$D_{\mu\nu}^{\text{H}} = \sum_{ij} C_{\mu i}^* C_{\nu j} \gamma_{ij}^{\text{H}} + \sum_{ab} C_{\mu a}^* C_{\nu b} \gamma_{ab}^{\text{H}} \quad (18)$$

$$\begin{aligned} \sum_{ai} z_{ai} f_{ai} &= \sum_{ai} z_{ai} h_{ai} + \sum_{aij} z_{ai} \langle aj || ij \rangle \\ &= \sum_{\mu\nu ai} z_{ai} C_{\mu a}^* C_{\nu i} h_{\mu\nu}^{\text{AO}} + \sum_{aij} z_{ai} \langle aj || ij \rangle \end{aligned} \quad (19)$$

$$\begin{aligned} \sum_{ai} z_{ai}^* f_{ai}^* &= \sum_{ai} z_{ai}^* h_{ai}^* + \sum_{aij} z_{ai}^* \langle aj || ij \rangle^* \\ &= \sum_{ai} z_{ia} h_{ia} + \sum_{aij} z_{ia} \langle ij || aj \rangle \\ &= \sum_{\mu\nu ai} z_{ia} C_{\mu i}^* C_{\nu a} h_{\mu\nu}^{\text{AO}} + \sum_{aij} z_{ia} \langle ij || aj \rangle \end{aligned} \quad (20)$$

3.1 First Derivative

We take the first derivative w.r.t. the nuclear magnetic moment m_{K_β} as only one-electron Hamiltonian $h_{\mu\nu}$ depends on m_{K_β} :

$$\begin{aligned} \left. \frac{d\mathcal{L}_{\text{MP2}}}{dm_{K_\beta}} \right|_{\mathbf{m}_K=0} &= \left. \frac{\partial \mathcal{L}_{\text{MP2}}}{\partial m_{K_\beta}} \right|_{\mathbf{m}_K=0} + \sum_{bj} \left(\frac{\partial \mathcal{L}_{\text{MP2}}}{\partial \kappa_{bj}} \right) \left(\frac{\partial \kappa_{bj}}{\partial m_{K_\beta}} \right) \Big|_{\mathbf{m}_K=0} \\ &\quad + \sum_{ai} \left(\frac{\partial \mathcal{L}_{\text{MP2}}}{\partial z_{ai}} \right) \left(\frac{\partial z_{ai}}{\partial m_{K_\beta}} \right) \Big|_{\mathbf{m}_K=0} \\ &= \left. \frac{\partial \mathcal{L}_{\text{MP2}}}{\partial m_{K_\beta}} \right|_{\mathbf{m}_K=0} \end{aligned} \quad (21)$$

$$\begin{aligned} \left. \frac{\partial \mathcal{L}_{\text{MP2}}}{\partial m_{K_\beta}} \right|_{\mathbf{m}_K=0} &= \sum_{\mu\nu} D_{\mu\nu}^{\text{HF}} \frac{\partial h_{\mu\nu}^{\text{AO}}}{\partial m_{K_\beta}} \Big|_{\mathbf{m}_K=0} + \sum_{\mu\nu} D_{\mu\nu}^{\text{H}} \frac{\partial h_{\mu\nu}^{\text{AO}}}{\partial m_{K_\beta}} \Big|_{\mathbf{m}_K=0} \\ &\quad + \frac{1}{2} \sum_{\mu\nu ai} z_{ai} C_{\mu a}^* C_{\nu i} \frac{\partial h_{\mu\nu}^{\text{AO}}}{\partial m_{K_\beta}} \Big|_{\mathbf{m}_K=0} + \frac{1}{2} \sum_{\mu\nu ai} z_{ia} C_{\mu i}^* C_{\nu a} \frac{\partial h_{\mu\nu}^{\text{AO}}}{\partial m_{K_\beta}} \Big|_{\mathbf{m}_K=0} \\ &= \sum_{\mu\nu} \frac{\partial h_{\mu\nu}^{\text{AO}}}{\partial m_{K_\beta}} \Big|_{\mathbf{m}_K=0} \left(D_{\mu\nu}^{\text{HF}} + D_{\mu\nu}^{\text{H}} + \frac{1}{2} \sum_{ai} z_{ai} C_{\mu a}^* C_{\nu i} + \frac{1}{2} \sum_{ai} z_{ia} C_{\mu i}^* C_{\nu a} \right) \end{aligned} \quad (22)$$

note that need to turn z_{ai}^* into $z_{ia}...$ what about z_{ia} block?

Trying to resolve z_{ia}

$$\begin{aligned} & \sum_{\mu\nu ai} z_{ia} C_{\mu i}^* C_{\nu a} h_{\mu\nu}^{\text{AO}} \\ & \stackrel{\mu \leftrightarrow \nu}{=} \sum_{\mu\nu ai} (z_{ai} C_{\mu a}^* C_{\nu i})^* h_{\nu\mu}^{\text{AO}} \end{aligned} \quad (23)$$