

# 1 Direct Differentiation

MP2 energy in spin-orbital formalism:

$$\begin{aligned} E^{(2)} &= \frac{1}{4} \sum_{ijab} \frac{\langle ab||ij\rangle \langle ij||ab\rangle}{f_{ii} + f_{jj} - f_{aa} - f_{bb}} \\ &= \frac{1}{4} \sum_{ijab} T_{ij}^{ab} \langle ij||ab\rangle \end{aligned} \quad (1)$$

in which we denote the MP2 amplitude:

$$T_{ij}^{ab} = \frac{\langle ab||ij\rangle}{\Delta_{ab}^{ij}} = \frac{\langle ab||ij\rangle}{f_{ii} + f_{jj} - f_{aa} - f_{bb}} \quad (2)$$

Directly differentiate the second-order energy w.r.t. perturbation parameter  $\lambda$ :

$$\begin{aligned} \frac{\partial E^{(2)}}{\partial \lambda} &= \frac{1}{4} \sum_{ijab} \frac{\partial}{\partial \lambda} (T_{ij}^{ab} \langle ij||ab\rangle) \\ &= \frac{1}{4} \sum_{ijab} \left( \frac{\partial T_{ij}^{ab}}{\partial \lambda} \right) \langle ij||ab\rangle + \frac{1}{4} \sum_{ijab} T_{ij}^{ab} \left( \frac{\partial \langle ij||ab\rangle}{\partial \lambda} \right) \end{aligned} \quad (3)$$

in which:

$$\frac{\partial \langle ij||ab\rangle}{\partial \lambda} = \langle i^\lambda j||ab\rangle + \langle ij^\lambda||ab\rangle + \langle ij||a^\lambda b\rangle + \langle ij||ab^\lambda\rangle \quad (4)$$

exploiting the permutational symmetry and equivalence of the dummy indices:

$$\sum_{ijab} \langle i^\lambda j||ab\rangle = \sum_{ijab} \langle j^\lambda i||ba\rangle = \sum_{jiba} \langle i^\lambda j||ab\rangle \quad (5)$$

Therefore:

$$\sum_{ijab} \langle ij||ab\rangle^\lambda = 2 \sum_{ijab} (\langle i^\lambda j||ab\rangle + \langle ij||a^\lambda b\rangle) \quad (6)$$

Then:

$$\begin{aligned} \frac{\partial T_{ij}^{ab}}{\partial \lambda} &= \frac{\partial}{\partial \lambda} \left( \frac{\langle ab||ij\rangle}{\Delta_{ab}^{ij}} \right) \\ &= \frac{\langle ab||ij\rangle^\lambda}{\Delta_{ab}^{ij}} - \frac{\langle ab||ij\rangle}{(\Delta_{ab}^{ij})^2} \left( \frac{\partial \Delta_{ab}^{ij}}{\partial \lambda} \right) \\ &= \frac{\langle ab||ij\rangle^\lambda}{\Delta_{ab}^{ij}} - \frac{\langle ab||ij\rangle}{(\Delta_{ab}^{ij})^2} (\varepsilon_i^\lambda + \varepsilon_j^\lambda - \varepsilon_a^\lambda - \varepsilon_b^\lambda) \end{aligned} \quad (7)$$

Putting eqns (6) and (7) together:

$$\begin{aligned} \frac{\partial E^{(2)}}{\partial \lambda} &= \frac{1}{4} \sum_{ijab} \langle ij||ab\rangle (T_{ij}^{ab})^\lambda + \frac{1}{4} \sum_{ijab} T_{ij}^{ab} \langle ij||ab\rangle^\lambda \\ &= \frac{1}{4} \sum_{ijab} \langle ij||ab\rangle \left( \frac{\langle ab||ij\rangle^\lambda}{\Delta_{ab}^{ij}} - \frac{\langle ab||ij\rangle}{(\Delta_{ab}^{ij})^2} (\varepsilon_i^\lambda + \varepsilon_j^\lambda - \varepsilon_a^\lambda - \varepsilon_b^\lambda) \right) + \frac{1}{4} \sum_{ijab} \langle ij||ab\rangle^\lambda T_{ij}^{ab} \\ &= \frac{1}{4} \sum_{ijab} \frac{\langle ij||ab\rangle \langle ab||ij\rangle^\lambda}{\Delta_{ab}^{ij}} + \frac{1}{4} \sum_{ijab} \frac{\langle ij||ab\rangle^\lambda \langle ab||ij\rangle}{\Delta_{ab}^{ij}} - \frac{1}{4} \sum_{ijab} \frac{\langle ij||ab\rangle \langle ab||ij\rangle}{(\Delta_{ab}^{ij})^2} (2\varepsilon_i^\lambda - 2\varepsilon_a^\lambda) \end{aligned}$$

$$\begin{aligned}
&= \frac{1}{2} \sum_{ijab} (T_{ij}^{ab})^* \langle a^\lambda b || ij \rangle + \frac{1}{2} \sum_{ijab} (T_{ij}^{ab})^* \langle ab || i^\lambda j \rangle + \frac{1}{2} \sum_{ijab} T_{ij}^{ab} \langle i^\lambda j || ab \rangle + \frac{1}{2} \sum_{ijab} T_{ij}^{ab} \langle ij || a^\lambda b \rangle \\
&\quad - \frac{1}{2} \sum_{ijab} T_{ij}^{ab} (T_{ij}^{ab})^* (\varepsilon_i^\lambda - \varepsilon_a^\lambda)
\end{aligned} \tag{8}$$

Using the expression for  $|i^\lambda\rangle$  and  $|a^\lambda\rangle$ :

$$\begin{aligned}
\frac{\partial E^{(2)}}{\partial \lambda} &= \frac{1}{2} \sum_{ijab} (T_{ij}^{ab})^* \left( \sum_k (U_{ka}^\lambda)^* \langle kb || ij \rangle + \sum_{f \neq a} (U_{fa}^\lambda)^* \langle fb || ij \rangle + \sum_\mu C_{\mu a}^* \langle \mu^\lambda b || ij \rangle \right) \\
&\quad + \frac{1}{2} \sum_{ijab} (T_{ij}^{ab})^* \left( \sum_{k \neq i} U_{ki}^\lambda \langle ab || kj \rangle + \sum_f U_{fi}^\lambda \langle ab || fj \rangle + \sum_\mu C_{\mu i} \langle ab || \mu^\lambda j \rangle \right) \\
&\quad + \frac{1}{2} \sum_{ijab} T_{ij}^{ab} \left( \sum_{k \neq i} (U_{ki}^\lambda)^* \langle kj || ab \rangle + \sum_f (U_{fi}^\lambda)^* \langle fj || ab \rangle + \sum_\mu C_{\mu i}^* \langle \mu^\lambda j || ab \rangle \right) \\
&\quad + \frac{1}{2} \sum_{ijab} T_{ij}^{ab} \left( \sum_k U_{ka}^\lambda \langle ij || kb \rangle + \sum_{f \neq a} U_{fa}^\lambda \langle ij || fb \rangle + \sum_\mu C_{\mu a} \langle ij || \mu^\lambda b \rangle \right) \\
&\quad - \frac{1}{2} \sum_{ijab} T_{ij}^{ab} (T_{ij}^{ab})^* (\varepsilon_i^\lambda - \varepsilon_a^\lambda)
\end{aligned} \tag{9}$$

Using the orthonormality condition on  $(U_{ka}^\lambda)^*$ ,  $(U_{fa}^\lambda)^*$ ,  $(U_{ki}^\lambda)^*$  and  $U_{ka}^\lambda$ , and omitting the AO terms (terms involving  $|\phi_\mu\rangle$ ) for now:

$$\begin{aligned}
\frac{\partial E^{(2)}}{\partial \lambda} &= \frac{1}{2} \sum_{ijab} \sum_k (T_{ij}^{ab})^* \langle kb || ij \rangle (-S_{ak}^\lambda - U_{ak}^\lambda) + \frac{1}{2} \sum_{ijab} \sum_{f \neq a} (T_{ij}^{ab})^* \langle fb || ij \rangle (-S_{af}^\lambda - U_{af}^\lambda) \\
&\quad + \frac{1}{2} \sum_{ijab} \sum_{k \neq i} (T_{ij}^{ab})^* \langle ab || kj \rangle U_{ki}^\lambda + \frac{1}{2} \sum_{ijab} \sum_f (T_{ij}^{ab})^* \langle ab || fi \rangle U_{fi}^\lambda \\
&\quad + \frac{1}{2} \sum_{ijab} \sum_{k \neq i} T_{ij}^{ab} \langle kj || ab \rangle (-S_{ik}^\lambda - U_{ik}^\lambda) + \frac{1}{2} \sum_{ijab} \sum_f T_{ij}^{ab} \langle fj || ab \rangle (U_{fi}^\lambda)^* \\
&\quad + \frac{1}{2} \sum_{ijab} \sum_k T_{ij}^{ab} \langle ij || kb \rangle (-S_{ak}^\lambda - (U_{ak}^\lambda)^*) + \frac{1}{2} \sum_{ijab} \sum_{f \neq a} T_{ij}^{ab} \langle ij || fb \rangle U_{fa}^\lambda \\
&\quad - \frac{1}{2} \sum_{ijab} T_{ij}^{ab} (T_{ij}^{ab})^* (\varepsilon_i^\lambda - \varepsilon_a^\lambda) + \text{AO terms}
\end{aligned} \tag{10}$$

collecting the  $U_{af}^\lambda$  and  $U_{fa}^\lambda$  terms and swapping some dummy indices:

$$\begin{aligned}
A &= -\frac{1}{2} \sum_{ijab} \sum_{f \neq a} (T_{ij}^{ab})^* \langle fb || ij \rangle U_{af}^\lambda + \frac{1}{2} \sum_{ijab} \sum_{f \neq a} T_{ij}^{ab} \langle ij || fb \rangle U_{fa}^\lambda \\
&= -\frac{1}{2} \sum_{ijfb} \sum_{a \neq f} (T_{ij}^{fb})^* \langle ab || ij \rangle U_{fa}^\lambda + \frac{1}{2} \sum_{ijab} \sum_{f \neq a} T_{ij}^{ab} \langle ij || fb \rangle U_{fa}^\lambda \\
&= \frac{1}{2} \sum_{ijab} \sum_{f \neq a} \langle ij || fb \rangle \langle ab || ij \rangle U_{fa}^\lambda \left( \frac{1}{\Delta_{ab}^{ij}} - \frac{1}{\Delta_{fb}^{ij}} \right) \\
&= \frac{1}{2} \sum_{ijab} \sum_{f \neq a} \langle ij || fb \rangle \langle ab || ij \rangle U_{fa}^\lambda \frac{\epsilon_a - \epsilon_f}{\Delta_{ab}^{ij} \Delta_{fb}^{ij}}
\end{aligned} \tag{11}$$

using the expression for  $U_{fa}^\lambda$ :

$$A = \frac{1}{2} \sum_{ijab} \sum_{f \neq a} \frac{\langle ij || fb \rangle \langle ab || ij \rangle (Q_{fa}^\lambda + \sum_{gm} [U_{gm}^\lambda \langle fm || ag \rangle + (U_{gm}^\lambda)^* \langle fg || am \rangle])}{\Delta_{ab}^{ij} \Delta_{fb}^{ij}}$$

$$= \frac{1}{2} \sum_{ijab} \sum_{f \neq a} T_{ij}^{ab} (T_{ij}^{fb})^* \left( Q_{fa}^\lambda + \sum_{gm} [U_{gm}^\lambda \langle fm || ag \rangle + (U_{gm}^\lambda)^* \langle fg || am \rangle] \right) \quad (12)$$

evaluating the  $\varepsilon_a^\lambda$  term in  $\frac{\partial E^{(2)}}{\partial \lambda}$  expression:

$$\begin{aligned} \frac{1}{2} \sum_{ijab} T_{ij}^{ab} (T_{ij}^{ab})^* \varepsilon_a^\lambda &= \frac{1}{2} \sum_{ijab} T_{ij}^{ab} (T_{ij}^{ab})^* \left( Q_{aa}^\lambda + \sum_{gm} [U_{gm}^\lambda \langle am || ag \rangle + (U_{gm}^\lambda)^* \langle ag || am \rangle] \right) \\ &= \frac{1}{2} \sum_{ijab} \sum_{f=a} T_{ij}^{ab} (T_{ij}^{fb})^* \left( Q_{fa}^\lambda + \sum_{gm} [U_{gm}^\lambda \langle fm || ag \rangle + (U_{gm}^\lambda)^* \langle fg || am \rangle] \right) \end{aligned} \quad (13)$$

add this term into  $A$ :

$$A + \frac{1}{2} \sum_{ijab} T_{ij}^{ab} (T_{ij}^{ab})^* \varepsilon_a^\lambda = \frac{1}{2} \sum_{ijabf} T_{ij}^{ab} (T_{ij}^{fb})^* \left( Q_{fa}^\lambda + \sum_{gm} [U_{gm}^\lambda \langle fm || ag \rangle + (U_{gm}^\lambda)^* \langle fg || am \rangle] \right) \quad (14)$$

Similarly, collecting the  $U_{ki}^\lambda$  and  $U_{ik}^\lambda$  terms, and add the  $\varepsilon_i^\lambda$  term into them we get:

$$- \frac{1}{2} \sum_{ijabk} (T_{ij}^{ab})^* T_{kj}^{ab} \left( Q_{ki}^\lambda + \sum_{gm} [U_{gm}^\lambda \langle km || ig \rangle + (U_{gm}^\lambda)^* \langle kg || im \rangle] \right) \quad (15)$$

Now looking at the  $S_{ak}^\lambda$  terms:

$$\begin{aligned} & - \frac{1}{2} \sum_{ijab} \sum_k (T_{ij}^{ab})^* \langle kb || ij \rangle S_{ak}^\lambda - \frac{1}{2} \sum_{ijab} \sum_k T_{ij}^{ab} \langle ij || kb \rangle S_{ak}^\lambda \\ &= - \frac{1}{2} \sum_{ijkab} S_{ak}^\lambda \left( \frac{\langle ij || ab \rangle \langle kb || ij \rangle + \langle ab || ij \rangle \langle ij || kb \rangle}{\Delta_{ab}^{ij}} \right) \end{aligned} \quad (16)$$

no further simplification for general (complex) orbitals, but could be further simplified if assumed real orbitals

Now putting these all back into eqn (10):

$$\begin{aligned} \frac{\partial E^{(2)}}{\partial \lambda} &= - \frac{1}{2} \sum_{ijkab} S_{ak}^\lambda \left( T_{ij}^{ab} \langle ij || kb \rangle + (T_{ij}^{ab})^* \langle kb || ij \rangle \right) \\ & - \frac{1}{2} \sum_{ijab} \sum_{f \neq a} (T_{ij}^{ab})^* \langle fb || ij \rangle S_{af}^\lambda - \frac{1}{2} \sum_{ijab} \sum_{k \neq i} T_{ij}^{ab} \langle kj || ab \rangle S_{ik}^\lambda \\ & + \frac{1}{2} \sum_{ijab} \sum_f (T_{ij}^{ab})^* \langle ab || fi \rangle U_{fi}^\lambda + \frac{1}{2} \sum_{ijab} \sum_f T_{ij}^{ab} \langle fj || ab \rangle (U_{fi}^\lambda)^* \\ & - \frac{1}{2} \sum_{ijab} \sum_k (T_{ij}^{ab})^* \langle kb || ij \rangle U_{ak}^\lambda - \frac{1}{2} \sum_{ijab} \sum_k T_{ij}^{ab} \langle ij || kb \rangle (U_{ak}^\lambda)^* \\ & + \frac{1}{2} \sum_{ijabf} T_{ij}^{ab} (T_{ij}^{fb})^* \left( Q_{fa}^\lambda + \sum_{gm} [U_{gm}^\lambda \langle fm || ag \rangle + (U_{gm}^\lambda)^* \langle fg || am \rangle] \right) \\ & - \frac{1}{2} \sum_{ijabk} (T_{ij}^{ab})^* T_{kj}^{ab} \left( Q_{ki}^\lambda + \sum_{gm} [U_{gm}^\lambda \langle km || ig \rangle + (U_{gm}^\lambda)^* \langle kg || im \rangle] \right) \end{aligned} \quad (17)$$

By defining:

$$D_{ki} = - \frac{1}{2} \sum_{jab} (T_{ij}^{ab})^* T_{kj}^{ab} \quad (18)$$

$$D_{fa} = \frac{1}{2} \sum_{ijb} (T_{ij}^{fb})^* T_{ij}^{ab} \quad (19)$$

$$I_{ik} = -\frac{1}{2} \sum_{jab} T_{ij}^{ab} \langle kj || ab \rangle \quad (20)$$

$$I_{af} = -\frac{1}{2} \sum_{ijb} (T_{ij}^{ab})^* \langle fb || ij \rangle \quad (21)$$

$$I_{ak} = -\frac{1}{2} \sum_{ijb} (T_{ij}^{ab} \langle ij || kb \rangle + (T_{ij}^{ab})^* \langle kb || ij \rangle) \quad (22)$$

the derivative becomes:

$$\begin{aligned} \frac{\partial E^{(2)}}{\partial \lambda} &= \sum_{ak} S_{ak}^\lambda I_{ak} + \sum_{a \neq f} S_{af}^\lambda I_{af} + \sum_{i \neq k} S_{ik}^\lambda I_{ik} \\ &+ \sum_{af} D_{fa} Q_{fa}^\lambda + \sum_{ik} D_{ki} Q_{ki}^\lambda \\ &+ \frac{1}{2} \sum_{ijabf} (T_{ij}^{fb})^* \langle fb || ai \rangle U_{ai}^\lambda + \frac{1}{2} \sum_{ijabf} T_{ij}^{fb} \langle aj || fb \rangle (U_{ai}^\lambda)^* \\ &- \frac{1}{2} \sum_{ijkab} (T_{kj}^{ab})^* \langle ib || kj \rangle U_{ai}^\lambda - \frac{1}{2} \sum_{ijkab} T_{kj}^{ab} \langle kj || ib \rangle (U_{ai}^\lambda)^* \\ &+ \frac{1}{2} \sum_{ijmabfg} T_{mj}^{gb} (T_{mj}^{fb})^* U_{ai}^\lambda \langle fi || ga \rangle + \frac{1}{2} \sum_{ijmabfg} T_{mj}^{gb} (T_{mj}^{fb})^* (U_{ai}^\lambda)^* \langle fa || gi \rangle \\ &- \frac{1}{2} \sum_{ijkmabg} (T_{mj}^{gb})^* T_{kj}^{gb} U_{ai}^\lambda \langle ki || ma \rangle - \frac{1}{2} \sum_{ijkmabg} (T_{mj}^{gb})^* T_{kj}^{gb} (U_{ai}^\lambda)^* \langle ka || mi \rangle \\ &= \sum_{ak} S_{ak}^\lambda I_{ak} + \sum_{a \neq f} S_{af}^\lambda I_{af} + \sum_{i \neq k} S_{ik}^\lambda I_{ik} \\ &+ \sum_{af} D_{fa} Q_{fa}^\lambda + \sum_{ik} D_{ki} Q_{ki}^\lambda \\ &+ \frac{1}{2} \sum_{ijabf} (T_{ij}^{fb})^* \langle fb || ai \rangle U_{ai}^\lambda + \frac{1}{2} \sum_{ijabf} T_{ij}^{fb} \langle aj || fb \rangle (U_{ai}^\lambda)^* \\ &- \frac{1}{2} \sum_{ijkab} (T_{kj}^{ab})^* \langle ib || kj \rangle U_{ai}^\lambda - \frac{1}{2} \sum_{ijkab} T_{kj}^{ab} \langle kj || ib \rangle (U_{ai}^\lambda)^* \\ &+ \sum_{fg} \sum_{ai} D_{fg} \langle fi || ga \rangle U_{ai}^\lambda + \sum_{fg} \sum_{ai} D_{fg} \langle fa || gi \rangle (U_{ai}^\lambda)^* \\ &+ \sum_{km} \sum_{ai} D_{km} \langle ki || ma \rangle U_{ai}^\lambda + \sum_{km} \sum_{ai} D_{km} \langle ka || mi \rangle (U_{ai}^\lambda)^* \\ &= \sum_{ai} S_{ai}^\lambda I_{ai} + \sum_{a \neq b} S_a^\lambda I_{ab} + \sum_{i \neq j} S_{ij}^\lambda I_{ij} + \sum_{ab} D_{ab} Q_{ab}^\lambda + \sum_{ij} D_{ij} Q_{ij}^\lambda \\ &+ \frac{1}{2} \sum_{ai} \sum_{jbc} (T_{ij}^{bc})^* \langle bc || ai \rangle U_{ai}^\lambda + \frac{1}{2} \sum_{ai} \sum_{jbc} T_{ij}^{bc} \langle aj || bc \rangle (U_{ai}^\lambda)^* \\ &- \frac{1}{2} \sum_{ai} \sum_{jkb} (T_{kj}^{ab})^* \langle ib || kj \rangle U_{ai}^\lambda - \frac{1}{2} \sum_{ai} \sum_{jkb} T_{kj}^{ab} \langle kj || ib \rangle (U_{ai}^\lambda)^* \\ &+ \sum_{ai} \sum_{bc} D_{bc} \langle bi || ca \rangle U_{ai}^\lambda + \sum_{ai} \sum_{bc} D_{bc} \langle ba || ci \rangle (U_{ai}^\lambda)^* \end{aligned}$$

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$$\begin{aligned}
& + \sum_{ai} \sum_{jk} D_{jk} \langle ji || ka \rangle U_{ai}^\lambda + \sum_{ai} \sum_{jk} D_{jk} \langle ja || ki \rangle (U_{ai}^\lambda)^* \\
& = \sum_{ai} S_{ai}^\lambda I_{ai} + \sum_{ab} S_{ab}^\lambda I_{ab} + \sum_{ij} S_{ij}^\lambda I_{ij} + \sum_{ab} D_{ab} Q_{ab}^\lambda + \sum_{ij} D_{ij} Q_{ij}^\lambda \\
& + \text{something like } \sum_{ai} X_{ai} U_{ai}^\lambda
\end{aligned} \tag{23}$$

$\sum_{a \neq b}$  and  $\sum_{i \neq j}$  terms, what about the  $a = b$  and  $i = j$  terms?

Should I drop the complex conjugate? but still can't see how to merge  $\langle bi || ca \rangle$  and  $\langle ba || ci \rangle$  terms hence could not get the same  $X_{ai}$  intermediate as in the article.

## 2 Some Identities

From CPHF orthonormality condition:

$$U_{pq}^\lambda + (U_{qp}^\lambda)^* + S_{pq}^\lambda = 0 \quad (24)$$

Consider the spin-orbital (and using  $\mathbf{C}(\lambda) = \mathbf{C}(0)\mathbf{U}(\lambda)$ ):

$$\begin{aligned} |a\rangle &= |\psi_a\rangle = \sum_{\mu} C_{\mu a}(\lambda) |\phi_{\mu}\rangle \\ &= \sum_{\mu} \left( \sum_q C_{\mu q}(0) U_{qa}(\lambda) \right) |\phi_{\mu}\rangle \\ &= \sum_{\mu} \left( \sum_k C_{\mu k}(0) U_{ka}(\lambda) \right) |\phi_{\mu}\rangle + \sum_{\mu} \left( \sum_{f \neq a} C_{\mu f}(0) U_{fa}(\lambda) \right) |\phi_{\mu}\rangle \end{aligned} \quad (25)$$

**Question: Why not include  $U_{aa}$  into this sum?**

Taking derivative w.r.t.  $\lambda$ :

$$\begin{aligned} |a^\lambda\rangle &= \sum_{\mu} \left( \sum_k C_{\mu k}(0) U_{ka}^\lambda \right) |\phi_{\mu}\rangle + \sum_{\mu} \left( \sum_{f \neq a} C_{\mu f}(0) U_{fa}^\lambda \right) |\phi_{\mu}\rangle \\ &\quad + \sum_{\mu} \left( \sum_k C_{\mu k}(0) U_{ka}^\lambda \right) |\phi_{\mu}^\lambda\rangle + \sum_{\mu} \left( \sum_{f \neq a} C_{\mu f}(0) U_{fa}^\lambda \right) |\phi_{\mu}^\lambda\rangle \end{aligned} \quad (26)$$

Noticing that  $\mathbf{U}(0) = \mathbf{I}$  hence  $\sum_{\mu} C_{\mu k}(0) |\phi_{\mu}\rangle = |\psi_k\rangle$ :

$$\begin{aligned} |a^\lambda\rangle &= \sum_k U_{ka}^\lambda |\psi_k\rangle + \sum_{f \neq a} U_{fa}^\lambda |\psi_f\rangle + \sum_{\mu q} C_{\mu q}(0) U_{qa}^\lambda |\phi_{\mu}^\lambda\rangle \\ &= \sum_k U_{ka}^\lambda |k\rangle + \sum_{f \neq a} U_{fa}^\lambda |f\rangle + \sum_{\mu} C_{\mu a} |\mu^\lambda\rangle \end{aligned} \quad (27)$$

Similarly:

$$|i^\lambda\rangle = \sum_{k \neq i} U_{ki}^\lambda |k\rangle + \sum_f U_{fi}^\lambda |f\rangle + \sum_{\mu} C_{\mu i} |\mu^\lambda\rangle \quad (28)$$

Expression for CPHF coefficients (c.f. Pople et al. 1979):

$$U_{fa}^\lambda = \frac{1}{\varepsilon_a - \varepsilon_f} (Q_{fa}^\lambda + \sum_{gm} [U_{gm}^\lambda \langle fm || ag \rangle + (U_{gm}^\lambda)^* \langle fg || am \rangle]) \quad (29)$$

$$U_{ki}^\lambda = \frac{1}{\varepsilon_i - \varepsilon_k} (Q_{ki}^\lambda + \sum_{gm} [U_{gm}^\lambda \langle fm || ag \rangle + (U_{gm}^\lambda)^* \langle kg || im \rangle]) \quad (30)$$

$$\varepsilon_a^\lambda = Q_{aa}^\lambda + \sum_{gm} [U_{gm}^\lambda \langle am || ag \rangle + (U_{gm}^\lambda)^* \langle ag || am \rangle] \quad (31)$$

$$\varepsilon_i^\lambda = Q_{ii}^\lambda + \sum_{gm} [U_{gm}^\lambda \langle im || ig \rangle + (U_{gm}^\lambda)^* \langle ig || im \rangle] \quad (32)$$

### 3 Lagrangian Method

The MP2 Lagrangian could be written as:

$$\begin{aligned}\mathcal{L}_{\text{MP2}} &= E_{\text{MP2}} + C_{\text{Bri}} + C_{\text{ON}} \\ &= E_{\text{HF}} + E_{\text{H}} + C_{\text{Bri}} + C_{\text{ON}}\end{aligned}\quad (33)$$

in which  $E_{\text{H}}$  is the Hylleraas functional,  $C_{\text{Bri}}$  is the Brillouin condition, and  $C_{\text{ON}}$  is the orthonormality condition.

#### 3.1 Hartree Fock Energy

The Hartree-Fock energy has contribution from zeroth- and first-order energies in MP2:

$$\begin{aligned}E_{\text{HF}} &= E^{(0)} + E^{(1)} \\ &= \sum_i h_{ii} + \sum_{ij} \langle ij || ij \rangle - \frac{1}{2} \sum_{ij} \langle ij || ij \rangle \\ &= \sum_i h_{ii} + \frac{1}{2} \sum_{ij} \langle ij || ij \rangle\end{aligned}\quad (34)$$

#### 3.2 Hylleraas Functional

The Hylleraas functional is defined as:

$$\begin{aligned}E_{\text{H}} &= \langle \Psi^{(1)} | \hat{V} - E^{(1)} | \Phi_0 \rangle + \langle \Phi_0 | \hat{V} - E^{(1)} | \Psi^{(1)} \rangle + \langle \Psi^{(1)} | \hat{H}^{(0)} - E^{(0)} | \Psi^{(1)} \rangle \\ &= 2 \text{Re} \langle \Psi^{(1)} | \hat{V} - E^{(1)} | \Phi_0 \rangle + \langle \Psi^{(1)} | \hat{H}^{(0)} - E^{(0)} | \Psi^{(1)} \rangle\end{aligned}\quad (35)$$

in which the relevant operators and functions are:

$$\hat{V} - E^{(1)} = \frac{1}{4} \sum_{pqrs} \langle pq || rs \rangle \{ \hat{p}^\dagger \hat{q}^\dagger \hat{s} \hat{r} \} \quad (36)$$

$$\hat{H}^{(0)} - E^{(0)} = \sum_{pq} f_{pq} \{ \hat{p}^\dagger \hat{q} \} = \sum_{pq} h_{pq} \{ \hat{p}^\dagger \hat{q} \} + \sum_{pqi} \langle pi || qi \rangle \{ \hat{p}^\dagger \hat{q} \} \quad (37)$$

$$| \Psi^{(1)} \rangle = \frac{1}{4} \sum_{ijab} T_{ij}^{ab} | \Phi_{ij}^{ab} \rangle \quad (38)$$

Therefore the Hylleraas functional could be written as:

$$\begin{aligned}E_{\text{H}} &= \frac{1}{8} \text{Re} \left\{ \sum_{ijab} (T_{ij}^{ab})^\dagger \sum_{pqrs} \langle pq || rs \rangle \langle \Phi_{ij}^{ab} | \{ \hat{p}^\dagger \hat{q}^\dagger \hat{s} \hat{r} \} | \Phi_0 \rangle \right\} \\ &\quad + \frac{1}{16} \sum_{ijab} (T_{ij}^{ab})^\dagger \sum_{klcd} T_{kl}^{cd} \sum_{pq} \langle \Phi_{ij}^{ab} | \{ \hat{p}^\dagger \hat{q} \} | \Phi_{kl}^{cd} \rangle\end{aligned}\quad (39)$$

First, we need to work out the following expectations:

$$\langle \Phi_{ij}^{ab} | \{ \hat{p}^\dagger \hat{q}^\dagger \hat{s} \hat{r} \} | \Phi_0 \rangle = \langle \Phi_0 | \{ \hat{i}^\dagger \hat{j}^\dagger \hat{b} \hat{a} \} \{ \hat{p}^\dagger \hat{q}^\dagger \hat{s} \hat{r} \} | \Phi_0 \rangle \quad (40)$$

$$\langle \Phi_{ij}^{ab} | \{ \hat{p}^\dagger \hat{q} \} | \Phi_{kl}^{cd} \rangle = \langle \Phi_0 | \{ \hat{i}^\dagger \hat{j}^\dagger \hat{b} \hat{a} \} \{ \hat{p}^\dagger \hat{q} \} \{ \hat{c}^\dagger \hat{d}^\dagger \hat{l} \hat{k} \} | \Phi_0 \rangle \quad (41)$$

Using GWT, the non-zero contributions come from the fully contracted terms:

$$\{ \hat{i}^\dagger \hat{j}^\dagger \hat{b} \hat{a} \} \{ \hat{p}^\dagger \hat{q}^\dagger \hat{s} \hat{r} \} = \{ \hat{i}^\dagger \hat{j}^\dagger \hat{b} \hat{a} \hat{p}^\dagger \hat{q}^\dagger \hat{s} \hat{r} \} + \{ \hat{i}^\dagger \hat{j}^\dagger \hat{b} \hat{a} \hat{p}^\dagger \hat{q}^\dagger \hat{s} \hat{r} \} + \{ \hat{i}^\dagger \hat{j}^\dagger \hat{b} \hat{a} \hat{p}^\dagger \hat{q}^\dagger \hat{s} \hat{r} \} + \{ \hat{i}^\dagger \hat{j}^\dagger \hat{b} \hat{a} \hat{p}^\dagger \hat{q}^\dagger \hat{s} \hat{r} \} + \dots$$

$$= \delta_{ir} \delta_{js} \delta_{bq} \delta_{ap} + \delta_{is} \delta_{jr} \delta_{bp} \delta_{aq} - \delta_{ir} \delta_{js} \delta_{bp} \delta_{aq} - \delta_{is} \delta_{jr} \delta_{bq} \delta_{ap} + \dots \quad (42)$$

$$\begin{aligned} & \{\hat{i}^\dagger \hat{j}^\dagger \hat{b} \hat{a}\} \{\hat{p}^\dagger \hat{q}\} \{\hat{c}^\dagger \hat{d} \hat{l} \hat{k}\} \\ &= \{\hat{i}^\dagger \hat{j}^\dagger \hat{b} \hat{a} \hat{p}^\dagger \hat{q} \hat{c}^\dagger \hat{d} \hat{l} \hat{k}\} + \{\hat{i}^\dagger \hat{j}^\dagger \hat{b} \hat{a} \hat{p}^\dagger \hat{q} \hat{c}^\dagger \hat{d} \hat{l} \hat{k}\} + \{\hat{i}^\dagger \hat{j}^\dagger \hat{b} \hat{a} \hat{p}^\dagger \hat{q} \hat{c}^\dagger \hat{d} \hat{l} \hat{k}\} + \{\hat{i}^\dagger \hat{j}^\dagger \hat{b} \hat{a} \hat{p}^\dagger \hat{q} \hat{c}^\dagger \hat{d} \hat{l} \hat{k}\} \\ &+ \{\hat{i}^\dagger \hat{j}^\dagger \hat{b} \hat{a} \hat{p}^\dagger \hat{q} \hat{c}^\dagger \hat{d} \hat{l} \hat{k}\} + \{\hat{i}^\dagger \hat{j}^\dagger \hat{b} \hat{a} \hat{p}^\dagger \hat{q} \hat{c}^\dagger \hat{d} \hat{l} \hat{k}\} + \{\hat{i}^\dagger \hat{j}^\dagger \hat{b} \hat{a} \hat{p}^\dagger \hat{q} \hat{c}^\dagger \hat{d} \hat{l} \hat{k}\} + \{\hat{i}^\dagger \hat{j}^\dagger \hat{b} \hat{a} \hat{p}^\dagger \hat{q} \hat{c}^\dagger \hat{d} \hat{l} \hat{k}\} \\ &+ \{\hat{i}^\dagger \hat{j}^\dagger \hat{b} \hat{a} \hat{p}^\dagger \hat{q} \hat{c}^\dagger \hat{d} \hat{l} \hat{k}\} + \{\hat{i}^\dagger \hat{j}^\dagger \hat{b} \hat{a} \hat{p}^\dagger \hat{q} \hat{c}^\dagger \hat{d} \hat{l} \hat{k}\} + \{\hat{i}^\dagger \hat{j}^\dagger \hat{b} \hat{a} \hat{p}^\dagger \hat{q} \hat{c}^\dagger \hat{d} \hat{l} \hat{k}\} + \{\hat{i}^\dagger \hat{j}^\dagger \hat{b} \hat{a} \hat{p}^\dagger \hat{q} \hat{c}^\dagger \hat{d} \hat{l} \hat{k}\} + \dots \\ &= \delta_{iq} \delta_{jk} \delta_{bd} \delta_{ac} \delta_{pl} - \delta_{iq} \delta_{jk} \delta_{bc} \delta_{ad} \delta_{pl} - \delta_{iq} \delta_{jl} \delta_{bd} \delta_{ac} \delta_{pk} + \delta_{iq} \delta_{jl} \delta_{bc} \delta_{ad} \delta_{pk} \\ &- \delta_{ik} \delta_{jq} \delta_{bd} \delta_{ac} \delta_{pl} + \delta_{il} \delta_{jq} \delta_{bd} \delta_{ac} \delta_{pk} + \delta_{ik} \delta_{jq} \delta_{bc} \delta_{ad} \delta_{pl} - \delta_{il} \delta_{jq} \delta_{bc} \delta_{ad} \delta_{pk} \\ &+ \delta_{ik} \delta_{jl} \delta_{bp} \delta_{ac} \delta_{qd} - \delta_{il} \delta_{jk} \delta_{bp} \delta_{ac} \delta_{qd} - \delta_{ik} \delta_{jl} \delta_{bp} \delta_{ad} \delta_{qc} + \delta_{il} \delta_{jk} \delta_{bp} \delta_{ad} \delta_{qc} \\ &+ \delta_{ik} \delta_{jl} \delta_{bd} \delta_{ap} \delta_{qc} - \delta_{il} \delta_{jk} \delta_{bd} \delta_{ap} \delta_{qc} - \delta_{ik} \delta_{jl} \delta_{bc} \delta_{ap} \delta_{qd} + \delta_{il} \delta_{jk} \delta_{bc} \delta_{ap} \delta_{qd} + \dots \end{aligned} \quad (43)$$

the terms not fully contrated are ommitted.

Therefore, the two parts of Hylleraas functional could be simplified as:

$$\begin{aligned} & \frac{1}{8} \text{Re} \left\{ \sum_{ijab} (T_{ij}^{ab})^\dagger \sum_{pqrs} \langle pq || rs \rangle \langle \Phi_{ij}^{ab} | \{\hat{p}^\dagger \hat{q}^\dagger \hat{s} \hat{r}\} | \Phi_0 \rangle \right\} \\ &= \frac{1}{8} \text{Re} \left\{ \sum_{ijab} (T_{ij}^{ab})^\dagger (\langle ab || ij \rangle + \langle ba || ji \rangle - \langle ba || ij \rangle - \langle ab || ji \rangle) \right\} \\ &= \frac{1}{2} \text{Re} \left\{ \sum_{ijab} (T_{ij}^{ab})^\dagger \langle ab || ij \rangle \right\} \end{aligned} \quad (44)$$

$$\begin{aligned} & \frac{1}{16} \sum_{ijab} (T_{ij}^{ab})^\dagger \sum_{klcd} T_{kl}^{cd} \sum_{pq} \langle \Phi_{ij}^{ab} | \{\hat{p}^\dagger \hat{q}\} | \Phi_{kl}^{cd} \rangle \\ &= \frac{1}{16} \left( \sum_{ijlab} f_{li} (T_{ij}^{ab})^\dagger T_{jl}^{ab} - \sum_{ijlab} f_{li} (T_{ij}^{ab})^\dagger T_{jl}^{ba} - \sum_{ijkab} f_{ki} (T_{ij}^{ab})^\dagger T_{kj}^{ab} + \sum_{ijkab} f_{ki} (T_{ij}^{ab})^\dagger T_{kj}^{ba} \right. \\ &- \sum_{ijlab} f_{lj} (T_{ij}^{ab})^\dagger T_{il}^{ab} + \sum_{ijlab} f_{lj} (T_{ij}^{ab})^\dagger T_{il}^{ba} + \sum_{ijkab} f_{kj} (T_{ij}^{ab})^\dagger T_{ki}^{ab} - \sum_{ijkab} f_{kj} (T_{ij}^{ab})^\dagger T_{ki}^{ba} \\ &+ \sum_{ijabd} f_{bd} (T_{ij}^{ab})^\dagger T_{ij}^{ad} - \sum_{ijabd} f_{bd} (T_{ij}^{ab})^\dagger T_{ji}^{ad} - \sum_{ijabc} f_{bc} (T_{ij}^{ab})^\dagger T_{ij}^{ca} + \sum_{ijabc} f_{bc} (T_{ij}^{ab})^\dagger T_{ji}^{ca} \\ &+ \sum_{ijabc} f_{ac} (T_{ij}^{ab})^\dagger T_{ij}^{cb} - \sum_{ijabc} f_{ac} (T_{ij}^{ab})^\dagger T_{ji}^{cb} - \sum_{ijabd} f_{ad} (T_{ij}^{ab})^\dagger T_{ij}^{bd} + \sum_{ijabd} f_{ad} (T_{ij}^{ab})^\dagger T_{ji}^{bd} \Big) \\ &= \frac{1}{4} \left( \sum_{ijkab} f_{ki} (T_{ij}^{ab})^\dagger T_{jk}^{ab} - \sum_{ijkab} f_{ki} (T_{ij}^{ab})^\dagger T_{jk}^{ba} + \sum_{ijabc} f_{ac} (T_{ij}^{ab})^\dagger T_{ij}^{bc} - \sum_{ijabc} f_{ac} (T_{ij}^{ab})^\dagger T_{ij}^{cb} \right) \\ &= \frac{1}{2} \left( \sum_{ijkab} f_{ki} (T_{ij}^{ab})^\dagger T_{jk}^{ab} - \sum_{ijabc} f_{ac} (T_{ij}^{ab})^\dagger T_{ij}^{bc} \right) \end{aligned} \quad (45)$$



Therefore, the Hylleraas functional, written in spin-orbital form, is:

$$E_H = \frac{1}{2} \text{Re} \left\{ \sum_{ijab} (T_{ij}^{ab})^\dagger \langle ab || ij \rangle \right\} + \frac{1}{2} \left( \sum_{ijkab} f_{ki} (T_{ij}^{ab})^\dagger T_{jk}^{ab} - \sum_{ijabc} f_{ac} (T_{ij}^{ab})^\dagger T_{ij}^{bc} \right) \quad (46)$$

To formulate the Hylleraas functional into density matrix representation, we write out dependencies on one- and two-electron integrals, i.e.,  $h_{pq}$  and  $\langle pq || rs \rangle$  explicitly.

$$\begin{aligned} E_H &= \frac{1}{2} \text{Re} \left\{ \sum_{ijab} (T_{ij}^{ab})^\dagger \langle ab || ij \rangle \right\} + \frac{1}{2} \left( \sum_{ijkab} h_{ki} (T_{ij}^{ab})^\dagger T_{jk}^{ab} - \sum_{ijabc} h_{ac} (T_{ij}^{ab})^\dagger T_{ij}^{bc} \right) \\ &\quad + \frac{1}{2} \left( \sum_{ijklab} \langle kl || il \rangle (T_{ij}^{ab})^\dagger T_{jk}^{ab} - \sum_{ijkabc} \langle ak || ck \rangle (T_{ij}^{ab})^\dagger T_{ij}^{bc} \right) \\ &= \sum_{ki} h_{ij} \gamma_{ij} + \sum_{ab} h_{ab} \gamma_{ab} + \sum_{ijab} \text{Re} \left\{ \langle ab || ij \rangle \Gamma_{ij}^{ab} \right\} + \sum_{ijkl} \langle ij || kl \rangle \Gamma_{kl}^{ij} + \sum_{ijab} \langle ai || bj \rangle \Gamma_{bj}^{ai} \end{aligned} \quad (47)$$

in which

$$\gamma_{ij} = \frac{1}{2} \sum_{kab} (T_{jk}^{ab})^\dagger T_{ki}^{ab} \quad (48)$$

$$\gamma_{ab} = -\frac{1}{2} \sum_{ijc} (T_{ij}^{ac})^\dagger T_{ij}^{cb} \quad (49)$$

$$\Gamma_{ij}^{ab} = \frac{1}{2} (T_{ij}^{ab})^\dagger \quad (50)$$

$$\Gamma_{kl}^{ij} = \frac{1}{2} \sum_{mab} (T_{km}^{ab})^\dagger T_{mi}^{ab} \delta_{jl} \quad (51)$$

$$\Gamma_{bj}^{ai} = -\frac{1}{2} \sum_{klc} (T_{kl}^{ac})^\dagger T_{kl}^{cb} \delta_{ij} \quad (52)$$

### 3.3 Perturbed Orthonormality Condition

We have the general orthonormality condition, subject to perturbation, as:

$$S_{pq} = \langle p | q \rangle = \delta_{pq} \quad (53)$$

$$\sum_{\mu\nu} C_{\mu p}^* S_{\mu\nu} C_{\nu q} = \delta_{pq} \quad (54)$$

Parameterization of the MO coefficients:

$$\mathbf{C}(\lambda) = \mathbf{C}(0) \mathbf{U}(\lambda) \quad (55)$$

$$C_{\mu p}(\lambda) = \sum_r C_{\mu r}(0) U_{rp}(\lambda) \quad (56)$$

in which  $\mathbf{U}(\lambda)$  is the solution to the CPHF equations.

Now the orthonormality condition using this parameterization:

$$\sum_{\mu\nu} \left( \sum_r U_{rp}^*(\lambda) C_{\mu r}^*(0) \right) S_{\mu\nu}(\lambda) \left( \sum_s C_{\nu s}(0) U_{sq}(\lambda) \right) = \delta_{pq} \quad (57)$$

Introducing the transformed overlap matrix:

$$\mathcal{S}_{pq}(\lambda) = \sum_{\mu\nu} C_{\mu p}^*(0) S_{\mu\nu}(\lambda) C_{\nu q}(0) \quad (58)$$

we have:

$$\sum_{rs} U_{rp}^*(\lambda) \mathcal{S}_{rs}(\lambda) U_{sq}(\lambda) = \delta_{pq} \quad (59)$$

differentiating both sides of the equation gives:

$$\sum_{rs} \frac{dU_{rp}^*(\lambda)}{d\lambda} \mathcal{S}_{rs}(\lambda) U_{sq}(\lambda) + \sum_{rs} U_{rp}^*(\lambda) \frac{d\mathcal{S}_{rs}(\lambda)}{d\lambda} U_{sq}(\lambda) + \sum_{rs} U_{rp}^*(\lambda) \mathcal{S}_{rs}(\lambda) \frac{dU_{sq}(\lambda)}{d\lambda} = 0 \quad (60)$$

Noting that  $\mathcal{S}(0) = \mathbf{I}$  because the unperturbed spin-orbitals are orthonormal, and it is trivial that  $\mathbf{U}(0) = \mathbf{I}$ .

Therefore evaluating the derivative at  $\lambda = 0$ , and denoting  $A^\lambda = \left(\frac{dA}{d\lambda}\right)\big|_{\lambda=0}$  results in:

$$\sum_{rs} (U_{rp}^\lambda)^* \delta_{rs} \delta_{sq} + \sum_{rs} \delta_{rp} \mathcal{S}_{rs}^\lambda \delta_{sq} + \sum_{rs} \delta_{rp} \delta_{rs} U_{sq}^\lambda = 0 \quad (61)$$

contracting the Kronecker delta tensors we get the perturbed orthonormality condition:

$$(U_{qp}^\lambda)^* + \mathcal{S}_{pq}^\lambda + U_{pq}^\lambda = 0 \quad (62)$$

### 3.4 Perturbed Brillouin Condition

The SCF density matrix is defined as:

$$D_{\mu\nu}^{\text{SCF}} = \sum_i^N C_{\mu i}^* C_{\nu i} \quad (63)$$

in which the MO coefficients are parameterized as:

$$C_{\mu p}(\lambda) = \sum_q C_{\mu q}(0) U(\lambda)_{qp} \quad (64)$$

$$\mathbf{C}(\lambda) = \mathbf{C}(0) \mathbf{U}(\lambda) \quad (65)$$

Define the one- and two-electron parts of the fock matrix, in AO and MO basis, as:

$$\begin{aligned} h_{pq} &= \langle p | \hat{h} | q \rangle = \sum_{\mu\nu} C_{\mu p}^* h_{\mu\nu}^{\text{AO}} C_{\nu q} & g_{pq} &= \sum_i \langle pi || qi \rangle = \sum_i \sum_{\mu\nu} C_{\mu p}^* \langle \mu i || \nu i \rangle C_{\nu q} \\ h_{\mu\nu}^{\text{AO}} &= \langle \mu | \hat{h} | \nu \rangle & g_{\mu\nu}^{\text{AO}} &= \sum_i \langle \mu i || \nu i \rangle = \sum_{\rho\sigma} D_{\rho\sigma} \langle \mu \rho || \nu \sigma \rangle \\ \mathbf{h} &= \mathbf{C}^\dagger \mathbf{h}^{\text{AO}} \mathbf{C} & \mathbf{g} &= \mathbf{C}^\dagger \mathbf{g}^{\text{AO}} \mathbf{C} \end{aligned} \quad (66)$$

Therefore, the Fock matrix could be expressed as (with the dependency on SCF density explicitly addressed):

$$\begin{aligned} F_{pq} &= h_{pq} + \sum_i \langle pi || qi \rangle \\ &= \sum_{\mu\nu} C_{\mu p}^* h_{\mu\nu}^{\text{AO}} C_{\nu q} + \sum_{\mu\nu} C_{\mu p}^* g_{\mu\nu}^{\text{AO}} C_{\nu q} \\ &= \sum_{\mu\nu} C_{\mu p}^* \langle \mu | \hat{h} | \nu \rangle C_{\nu q} + \sum_{\mu\nu} C_{\mu p}^* \left( \sum_{\rho\sigma} D_{\rho\sigma} \langle \mu \rho || \nu \sigma \rangle \right) C_{\nu q} \end{aligned} \quad (67)$$

$$F_{\mu\nu}^{\text{AO}} = h_{\mu\nu}^{\text{AO}} + g_{\mu\nu}^{\text{AO}} = \langle \mu | \hat{h} | \nu \rangle + \sum_{\rho\sigma} D_{\rho\sigma} \langle \mu \rho | | \nu \sigma \rangle \quad (68)$$

$$\begin{aligned} \mathbf{F} &= \mathbf{h} + \mathbf{g}[\mathbf{D}^{\text{SCF}}] \\ &= \mathbf{C}^\dagger \mathbf{h}^{\text{AO}} \mathbf{C} + \mathbf{C}^\dagger \mathbf{g}^{\text{AO}} [\mathbf{D}^{\text{SCF}}] \mathbf{C} \\ &= \mathbf{C}^\dagger \mathbf{F}^{\text{AO}} [\mathbf{D}^{\text{SCF}}] \mathbf{C} \end{aligned} \quad (69)$$

$$\mathbf{F}^{\text{AO}} [\mathbf{D}^{\text{SCF}}] = \mathbf{h}^{\text{AO}} + \mathbf{g}^{\text{AO}} [\mathbf{D}^{\text{SCF}}] \quad (70)$$

Evaluating the derivative at  $\lambda = 0$ , noting that  $\mathbf{U}(0) = \mathbf{I}$ :

$$\begin{aligned} \mathbf{F}^\lambda &= \left. \frac{d\mathbf{F}(\lambda)}{d\lambda} \right|_{\lambda=0} = \left( \mathbf{C}^\dagger(\lambda) \mathbf{F}^{\text{AO}} [\mathbf{D}^{\text{SCF}}(\lambda)] (\lambda) \mathbf{C}(\lambda) \right)^\lambda \\ &= \mathbf{C}^{\lambda\dagger}(\lambda) \mathbf{F}^{\text{AO}} [\mathbf{D}^{\text{SCF}}(\lambda)] (\lambda) \mathbf{C}(\lambda) + \mathbf{C}^\dagger(\lambda) \mathbf{F}^{\text{AO}} [\mathbf{D}^{\text{SCF}}(\lambda)] (\lambda) \mathbf{C}^\lambda(\lambda) \\ &\quad + \mathbf{C}^\dagger(\lambda) \left( \mathbf{h}^{\text{AO}}(\lambda) + \mathbf{g}^{\text{AO}} [\mathbf{D}^{\text{SCF}}(\lambda)] (\lambda) \right)^\lambda \mathbf{C}(\lambda) \\ &= \mathbf{U}^{\lambda\dagger}(\lambda) \underbrace{\mathbf{U}^\dagger(0) \mathbf{C}^\dagger(0)}_{\mathbf{C}^\dagger(\lambda=0)} \mathbf{F}^{\text{AO}} [\mathbf{D}^{\text{SCF}}(\lambda)] (\lambda) \mathbf{C}(\lambda) + \mathbf{C}^\dagger(\lambda) \mathbf{F}^{\text{AO}} [\mathbf{D}^{\text{SCF}}(\lambda)] (\lambda) \underbrace{\mathbf{C}(0) \mathbf{U}(0)}_{\mathbf{C}(\lambda=0)} \mathbf{U}^\lambda(\lambda) \\ &\quad + \mathbf{C}^\dagger(\lambda) \left( \mathbf{h}^{\text{AO},\lambda}(\lambda) + \mathbf{g}^{\text{AO},\lambda} [\mathbf{D}^{\text{SCF}}(\lambda)] (\lambda) + \mathbf{g}^{\text{AO}} [\mathbf{D}^{\text{SCF},\lambda}(\lambda)] (\lambda) \right) \mathbf{C}(\lambda) \\ &= \mathbf{U}^{\lambda\dagger} \mathbf{F} + \mathbf{F} \mathbf{U}^\lambda + \mathbf{C}^\dagger \mathbf{h}^{\text{AO},\lambda} \mathbf{C} + \mathbf{C}^\dagger \mathbf{g}^{\text{AO},\lambda} [\mathbf{D}^{\text{SCF}}] \mathbf{C} + \mathbf{C}^\dagger \mathbf{g}^{\text{AO}} [\mathbf{D}^{\text{SCF},\lambda}] \mathbf{C} \end{aligned} \quad (71)$$

The perturbed Brillouin condition is:

$$F_{ai}^\lambda = 0 \quad (72)$$

To evaluate the perturbed Fock matrix, we write the perturbed quantities in suffix notation as (assuming canonical orbitals, i.e.  $F_{pq} = \delta_{pq} \varepsilon_p$ ):

$$(\mathbf{U}^{\lambda\dagger} \mathbf{F})_{ai} = U_{ia}^{\lambda*} \varepsilon_i = \left. \frac{dU_{ia}^*}{d\lambda} \right|_{\lambda=0} \varepsilon_i \quad (73)$$

$$(\mathbf{F} \mathbf{U}^\lambda)_{ai} = \varepsilon_a U_{ai}^\lambda = \varepsilon_a \left. \frac{dU_{ai}}{d\lambda} \right|_{\lambda=0} \quad (74)$$

$$(\mathbf{C}^\dagger \mathbf{h}^{\text{AO},\lambda} \mathbf{C})_{ai} = \sum_{\mu\nu} C_{\mu a}^* h_{\mu\nu}^\lambda C_{\nu i} = \sum_{\mu\nu} C_{\nu a}^* \left. \frac{dh_{\mu\nu}}{d\lambda} \right|_{\lambda=0} C_{\nu i} \quad (75)$$

$$\begin{aligned} (\mathbf{C}^\dagger \mathbf{g}^{\text{AO},\lambda} [\mathbf{D}^{\text{SCF}}] \mathbf{C})_{ai} &= \sum_{\mu\nu} C_{\mu a}^* \left( \sum_{\rho\sigma} D_{\rho\sigma} \langle \mu \rho | | \nu \sigma \rangle^\lambda \right) C_{\nu i} \\ &= \sum_{\mu\nu} C_{\mu a}^* \left( \sum_{\rho\sigma} D_{\rho\sigma} \left. \frac{d\langle \mu \rho | | \nu \sigma \rangle}{d\lambda} \right|_{\lambda=0} \right) C_{\nu i} \end{aligned} \quad (76)$$

$$\begin{aligned} (\mathbf{C}^\dagger \mathbf{g}^{\text{AO}} [\mathbf{D}^{\text{SCF},\lambda}] \mathbf{C})_{ai} &= \sum_{\mu\nu} C_{\mu a}^* \left( \sum_{\rho\sigma} D_{\rho\sigma}^\lambda \langle \mu \rho | | \nu \sigma \rangle \right) C_{\nu i} \\ &= \sum_{\mu\nu} C_{\mu a}^* \left( \sum_{\rho\sigma} \left. \frac{dD_{\rho\sigma}}{d\lambda} \right|_{\lambda=0} \langle \mu \rho | | \nu \sigma \rangle \right) C_{\nu i} \end{aligned} \quad (77)$$