

1 CC Density Matrix

1.1 Reduced Density Matrix

We have the one- and two-body RDM as:

$$\gamma(x, x') = \sum_{pq} \phi_q(x) \gamma_{qp} \phi_p^*(x') \quad (1)$$

$$\Gamma(x_1, x_2; x'_1, x'_2) = \sum_{pqrs} \phi_r(x_1) \phi_s(x_2) \Gamma_{rspq} \phi_p^*(x'_1) \phi_q^*(x'_2) \quad (2)$$

in which γ_{qp} and Γ_{rspq} are the elements of discrete representation of one- and two- body RDM $\gamma(x, x')$ and $\Gamma(x_1, x_2; x'_1, x'_2)$:

$$\gamma_{qp} = \frac{\langle \Psi | \hat{p}^\dagger \hat{q} | \Psi \rangle}{\langle \Psi | \Psi \rangle} = \frac{\langle 0 | e^{\hat{T}^\dagger} \hat{p}^\dagger \hat{q} e^{\hat{T}} | 0 \rangle}{\langle 0 | e^{\hat{T}^\dagger} e^{\hat{T}} | 0 \rangle} = \langle 0 | e^{\hat{T}^\dagger} \hat{p}^\dagger \hat{q} e^{\hat{T}} | 0 \rangle_C \quad (3)$$

$$\Gamma_{rspq} = \frac{\langle \Psi | \hat{p}^\dagger \hat{q}^\dagger \hat{s} \hat{r} | \Psi \rangle}{\langle \Psi | \Psi \rangle} = \frac{\langle 0 | e^{\hat{T}^\dagger} \hat{p}^\dagger \hat{q}^\dagger \hat{s} \hat{r} e^{\hat{T}} | 0 \rangle}{\langle 0 | e^{\hat{T}^\dagger} e^{\hat{T}} | 0 \rangle} = \langle 0 | e^{\hat{T}^\dagger} \hat{p}^\dagger \hat{q}^\dagger \hat{s} \hat{r} e^{\hat{T}} | 0 \rangle_C \quad (4)$$

where we used the eqn(46(?)), i.e.:

$$\begin{aligned} \langle 0 | e^{\hat{T}^\dagger} \hat{p}^\dagger \hat{q} e^{\hat{T}} | 0 \rangle &= \langle 0 | e^{\hat{T}^\dagger} \{ \hat{p}^\dagger \hat{q} \} e^{\hat{T}} | 0 \rangle + \langle 0 | e^{\hat{T}^\dagger} \{ \hat{p}^\dagger \hat{q} \} e^{\hat{T}} | 0 \rangle \\ &= \langle 0 | e^{\hat{T}^\dagger} e^{\hat{T}} | 0 \rangle \langle 0 | e^{\hat{T}^\dagger} \{ \hat{p}^\dagger \hat{q} \} e^{\hat{T}} | 0 \rangle_C + \delta_{pq} \delta_{p \in \text{occ}} \langle 0 | e^{\hat{T}^\dagger} e^{\hat{T}} | 0 \rangle \end{aligned} \quad (5)$$

Therefore:

$$\begin{aligned} \gamma_{qp} &= \langle 0 | e^{\hat{T}^\dagger} \{ \hat{p}^\dagger \hat{q} \} e^{\hat{T}} | 0 \rangle_C + \delta_{pq} \delta_{p \in \text{occ}} \\ &= (\gamma_N)_{qp} + \delta_{pq} \delta_{p \in \text{occ}} \end{aligned} \quad (6)$$

With these, we can evaluate the expected values operators. The one-body expression can be obtained as:

$$\begin{aligned} \bar{O} &= \frac{\langle \Psi | \hat{O} | \Psi \rangle}{\langle \Psi | \Psi \rangle} \\ &= \frac{\sum_{pq} o_{pq} \langle 0 | e^{\hat{T}^\dagger} \{ \hat{p}^\dagger \hat{q} \} e^{\hat{T}} | 0 \rangle}{\langle 0 | e^{\hat{T}^\dagger} e^{\hat{T}} | 0 \rangle} + \frac{\sum_{pq} o_{pq} \langle 0 | e^{\hat{T}^\dagger} \{ \hat{p}^\dagger \hat{q} \} e^{\hat{T}} | 0 \rangle}{\langle 0 | e^{\hat{T}^\dagger} e^{\hat{T}} | 0 \rangle} \\ &= \frac{\langle 0 | e^{\hat{T}^\dagger} e^{\hat{T}} | 0 \rangle \sum_{pq} o_{pq} \langle 0 | e^{\hat{T}^\dagger} \{ \hat{p}^\dagger \hat{q} \} e^{\hat{T}} | 0 \rangle_C}{\langle 0 | e^{\hat{T}^\dagger} e^{\hat{T}} | 0 \rangle} + \frac{\sum_{pq} o_{pq} \langle 0 | e^{\hat{T}^\dagger} e^{\hat{T}} | 0 \rangle \delta_{pq} \delta_{p \in \text{occ}}}{\langle 0 | e^{\hat{T}^\dagger} e^{\hat{T}} | 0 \rangle} \\ &= \sum_{pq} o_{pq} \langle 0 | e^{\hat{T}^\dagger} \{ \hat{p}^\dagger \hat{q} \} e^{\hat{T}} | 0 \rangle_C + \sum_{pq} o_{pq} \delta_{pq} \delta_{p \in \text{occ}} \\ &= \sum_{pq} o_{pq} (\gamma_N)_{qp} + \sum_i o_{ii} \\ &= \sum_{pq} o_{pq} \gamma_{qp} \end{aligned} \quad (7)$$

Similarly, for a general two-body operator $\hat{G} = \frac{1}{4} \sum_{pqrs} \langle pq | \hat{g} | rs \rangle_A \hat{p}^\dagger \hat{q}^\dagger \hat{s} \hat{r}$, where $\langle pq | \hat{g} | rs \rangle_A = \langle pq | \hat{g} | rs \rangle - \langle pq | \hat{g} | sr \rangle$, the expectation value is:

$$\bar{G} = \frac{\langle \Psi | \hat{G} | \Psi \rangle}{\langle \Psi | \Psi \rangle} = \frac{1}{4} \sum_{pqrs} \langle pq | \hat{g} | rs \rangle_A \Gamma_{rspq} \quad (8)$$

The Γ_{rspq} tensor could be partitioned as:

1.2 Response Density Matrix

1.3 Relaxed/Effective Density Matrix