

Chapter 7 – Magnetic Field and Magnetic Forces

UM-SJTU Joint Institute
Physics II (Fall 2020)
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Agenda

1 Magnetic Field

- Experimental Observations and Basic Features of Magnetic Field
- Sources of Magnetic Field. Examples

2 Magnetic Force

- Magnetic Force On a Moving Charge
- Examples
- Applications

3 The Hall Effect

4 Current–Carrying Conductor In Magnetic Field

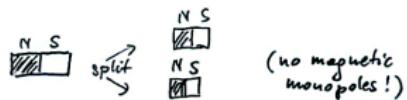
- Magnetic Force On a Current–Carrying Conductor
- Current Loop In Uniform Magnetic Field
- Magnetic Dipole Moment

5 Flux of Magnetic Field (Magnetic Flux) and Gauss's Law for Magnetic Field

Magnetic Field

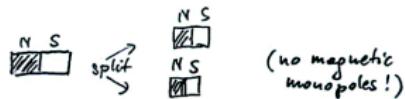
Introduction

- permanent magnets (e.g. a bar magnet, the Earth)

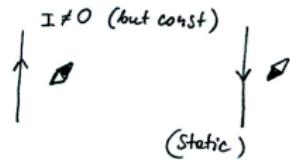


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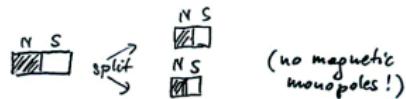


- Oersted's experiment (H.Ch. Oersted, 1820)

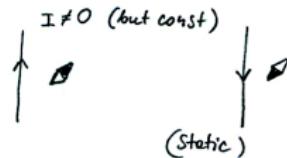


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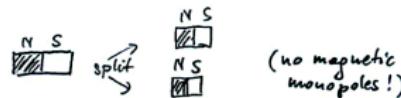


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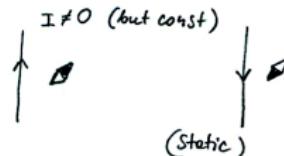
Introduction

- permanent magnets (e.g. a bar magnet, the Earth)



(no magnetic monopoles!)

- Oersted's experiment (H.Ch. Oersted, 1820)



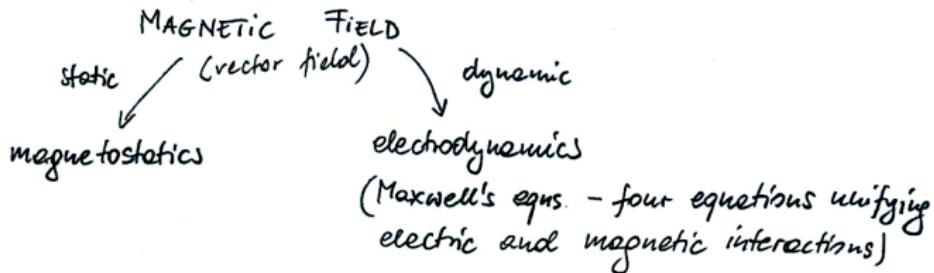
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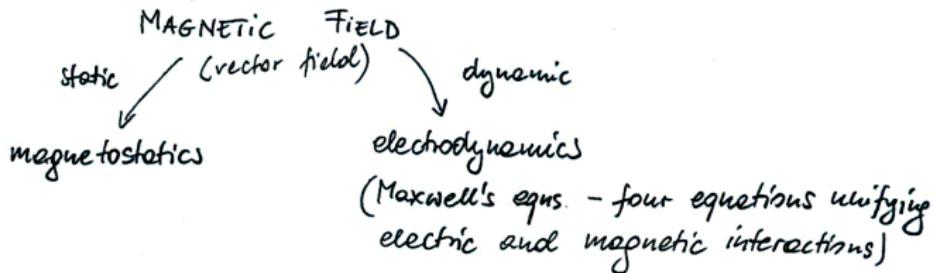
- two wires carrying electric current attract/repel each other



Features of Magnetic Field

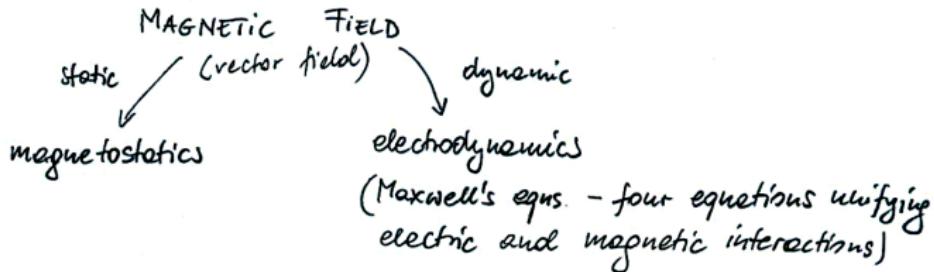


Features of Magnetic Field



Magnetic Field vs. Electric Field

Features of Magnetic Field

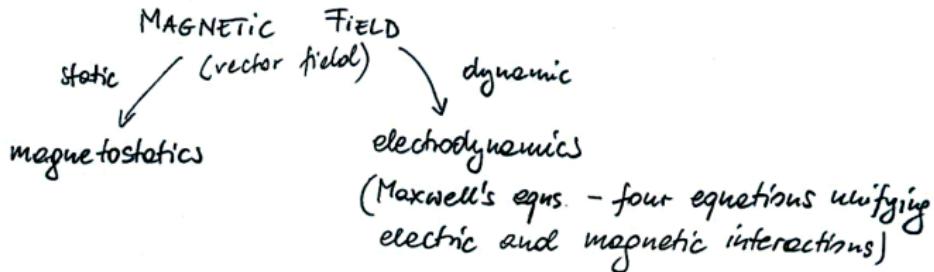


Magnetic Field vs. Electric Field

electric field (static)

- static electric charges \Rightarrow
static electric field

Features of Magnetic Field

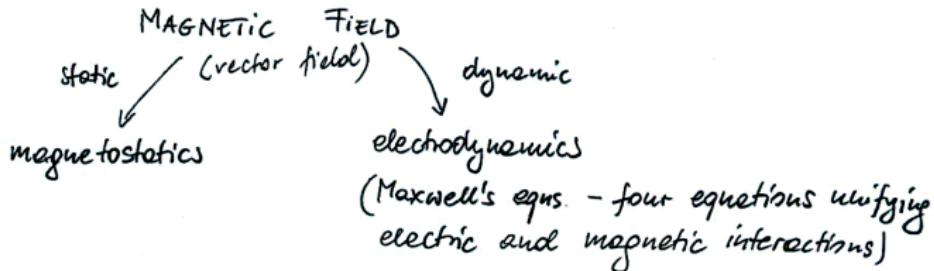


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Features of Magnetic Field



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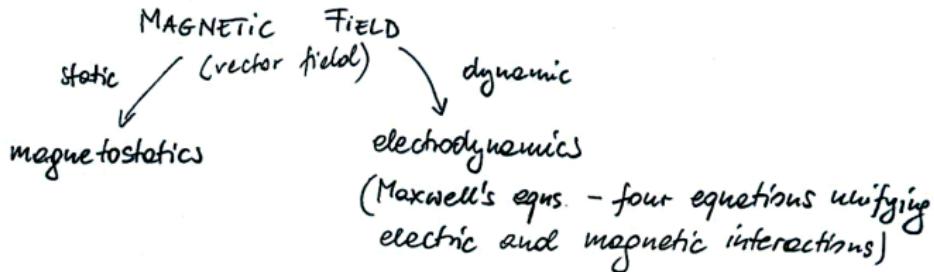
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magnetic field

- moving charges (currents)
⇒ magnetic field

Features of Magnetic Field



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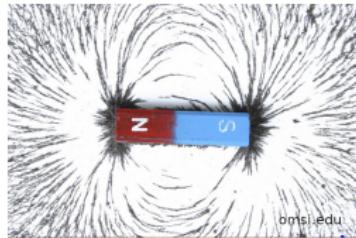
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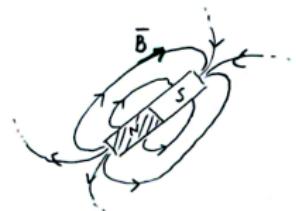
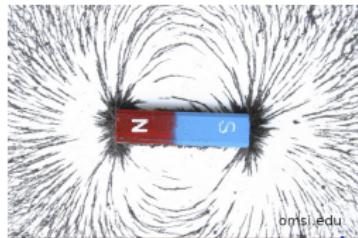
magnetic field

- moving charges (currents)
⇒ magnetic field
 - a magnetic force acts on a charge moving in the magnetic field

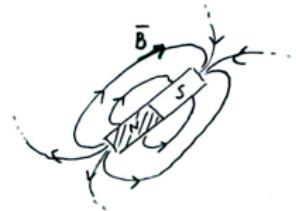
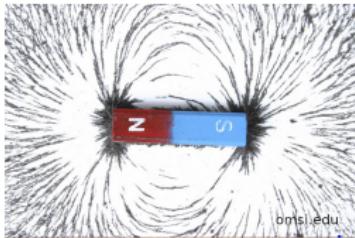
(a) permanent magnet – bar magnet



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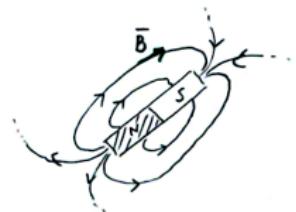
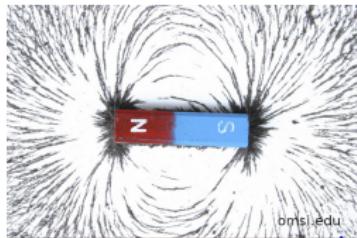


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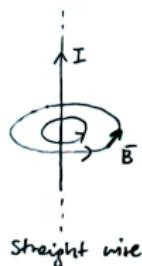
At any point, the magnetic field vector \vec{B} is tangential to the field line passing through that point.

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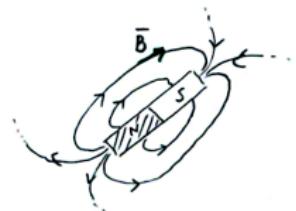
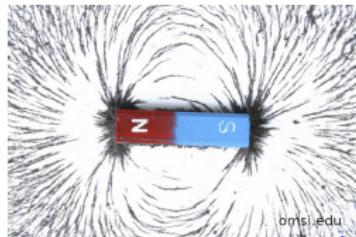


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(b) electric current

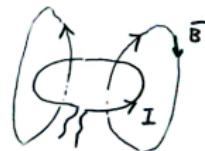
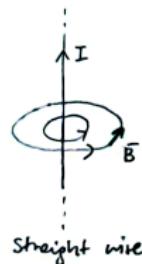


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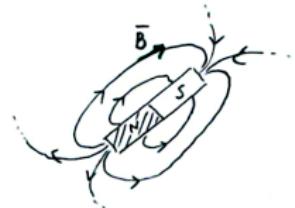
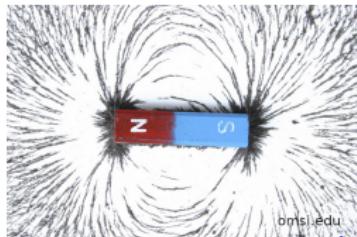
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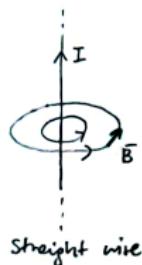
Current loop

(a) permanent magnet – bar magnet



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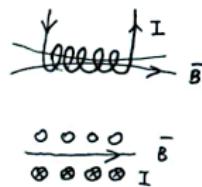
(b) electric current



straight wire

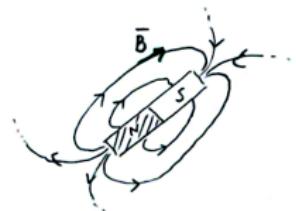
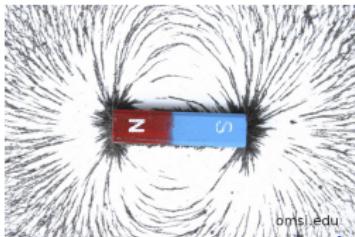


current loop



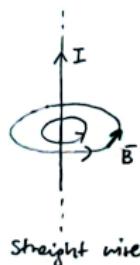
coil
(solenoid)

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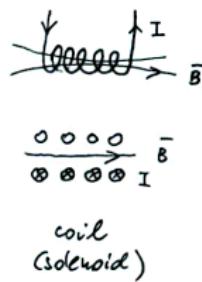


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Current loop



coil
(solenoid)

Observation. Magnetic field lines always form loops!

Magnetic Field

Magnetic Force

The Hall Effect

Current-Carrying Conductor In Magnetic Field

Flux of Magnetic Field (Magnetic Flux), and Gauss's Law for Mag

Magnetic Force On a Moving Charge

Examples

Applications

Magnetic Force

Experimental Facts

A charge q moving with velocity \vec{v} in a magnetic field \vec{B} is acted upon by a force \vec{F} with the following features

① $|\vec{F}| \propto |q|$

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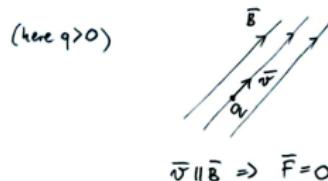
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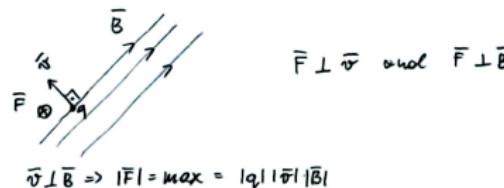
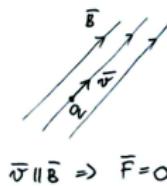
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(here $q > 0$)



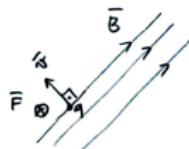
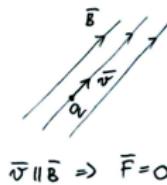
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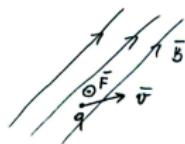
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$$\vec{v} \perp \vec{B} \Rightarrow |\vec{F}| = \text{max} = |q| |\vec{v}| |\vec{B}|$$

$$\vec{F} \perp \vec{v} \text{ and } \vec{F} \perp \vec{B}$$



$$\begin{aligned}\vec{F} &\perp \vec{v} \\ \vec{F} &\perp \vec{B}\end{aligned}$$

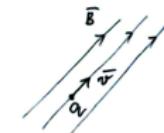
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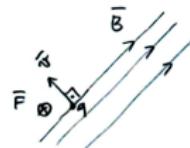
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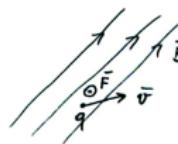


$$\vec{v} \parallel \vec{B} \Rightarrow \vec{F} = 0$$



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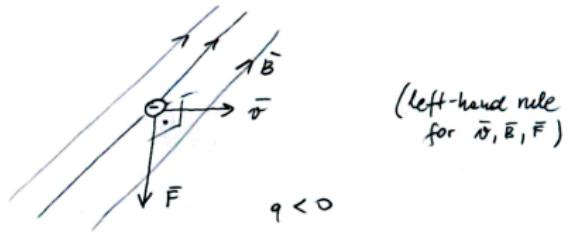
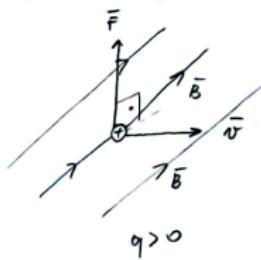
component $\perp \vec{B}$

$$|\vec{F}| = |q| |\vec{v}| |\vec{B}| \sin\theta(\vec{v}, \vec{B}) = |q| |\vec{v}_\perp| |\vec{B}| = |q| |\vec{v}| |\vec{B}_\perp|$$

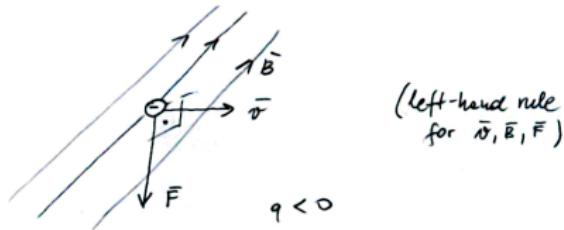
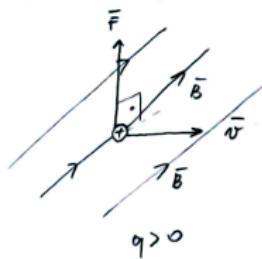
angle between \vec{v}, \vec{B}
(the smaller one)

component
 \perp to \vec{v}

The direction of \bar{F} depends on the sign of the charge.



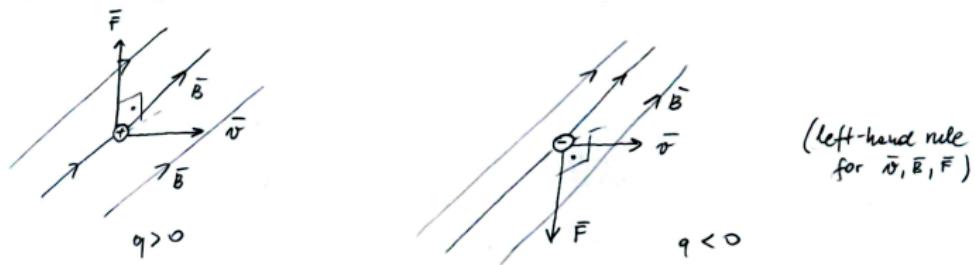
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Hence, the magnetic force on a moving charge is

$$\boxed{\bar{F} = q \bar{v} \times \bar{B}}.$$

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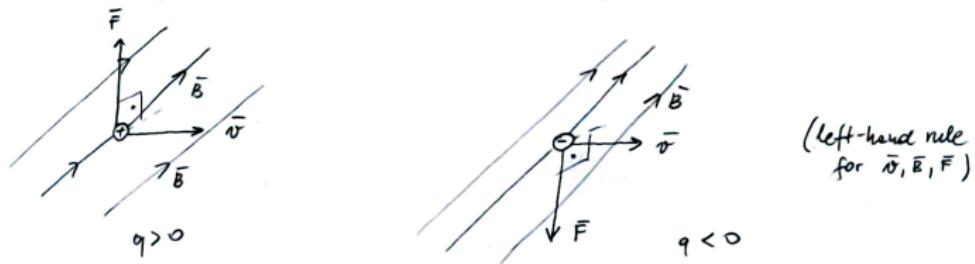
(left-hand rule
for $\bar{v}, \bar{B}, \bar{F}$)

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SI unit of the magnetic field: $\left[\frac{\text{N}}{\text{C} \cdot \text{m/s}} \right]$

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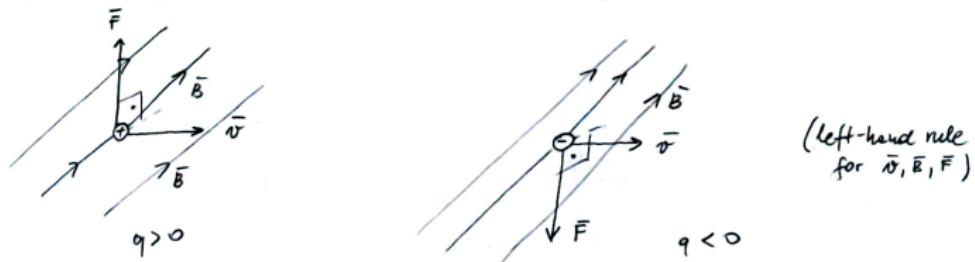
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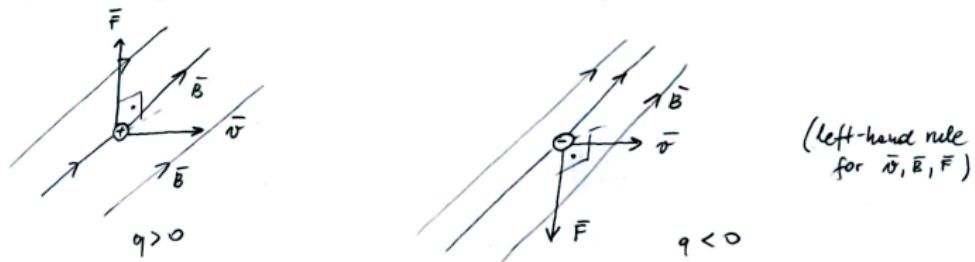


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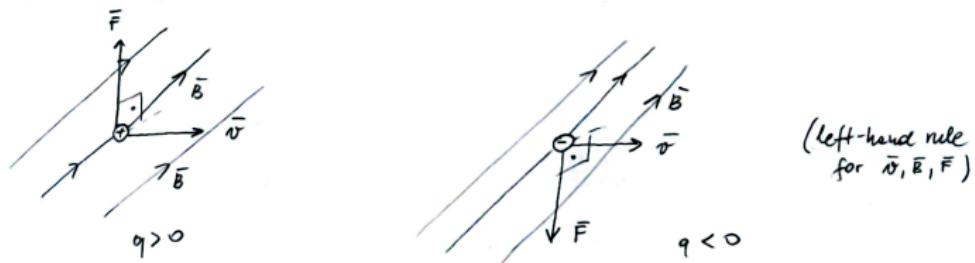


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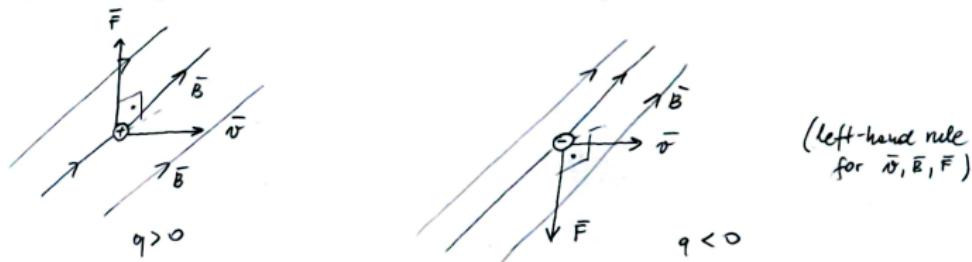
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Alternative units: Gauss [G]; $1 \text{ G} = 10^{-4} \text{ T}$.

Magnetic Force. Comments (assume $v \ll c$)

$$\bar{F} = q \bar{v} \times \bar{B} \implies \bar{F} \perp \bar{v}$$

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$$\frac{dK}{dt} = \frac{d}{dt} \left(\frac{1}{2} m \bar{v} \circ \bar{v} \right)$$

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$$\frac{dK}{dt} = \frac{d}{dt} \left(\frac{1}{2} m \bar{v} \circ \bar{v} \right) = \frac{1}{2} m (\dot{\bar{v}} \circ \bar{v} + \bar{v} \circ \dot{\bar{v}})$$

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$$\bar{F} = q\bar{v} \times \bar{B} \implies \bar{F} \perp \bar{v} \implies \bar{F} \text{ cannot change the magnitude of } \bar{v}$$

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If both, the electric \bar{E} and the magnetic \bar{B} fields are present, then

$$\bar{F} = q (\bar{E} + \bar{v} \times \bar{B})$$

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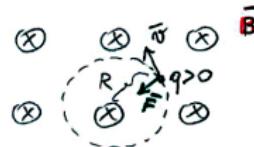
$$\boxed{\bar{F} = q (\bar{E} + \bar{v} \times \bar{B})},$$

and the force is referred to as the **Lorentz force**.

Examples

(a) cyclotron motion

$\vec{B} \neq 0$ and $\vec{v} \perp \vec{B}$



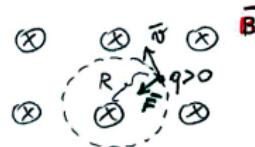
trajectory - circle

Lorentz force = centripetal
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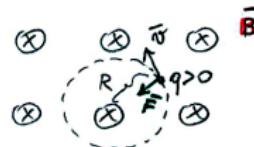
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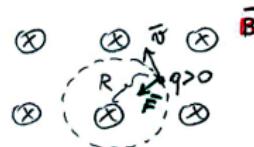
$$\frac{mv^2}{R} = |q| v B \Rightarrow v = \sqrt{\frac{|q| B}{m}} R = \omega R$$

$$\omega = \frac{|q| B}{m} \text{ cyclotron frequency}$$

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(b) motion along a helix

Examples

(a) cyclotron motion

$\vec{B} \neq 0$ and $\vec{v} \perp \vec{B}$



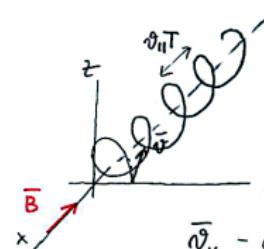
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(b) motion along a helix

$\vec{B} \neq 0$ and $\vec{v} \not\perp \vec{B}$
 $\vec{v} \parallel \vec{B}$



$\vec{v}_{||}$ - component of \vec{v} parallel to \vec{B}

Magnetic Field

Magnetic Force

The Hall Effect

Current-Carrying Conductor In Magnetic Field

Flux of Magnetic Field (Magnetic Flux), and Gauss's Law for Mag

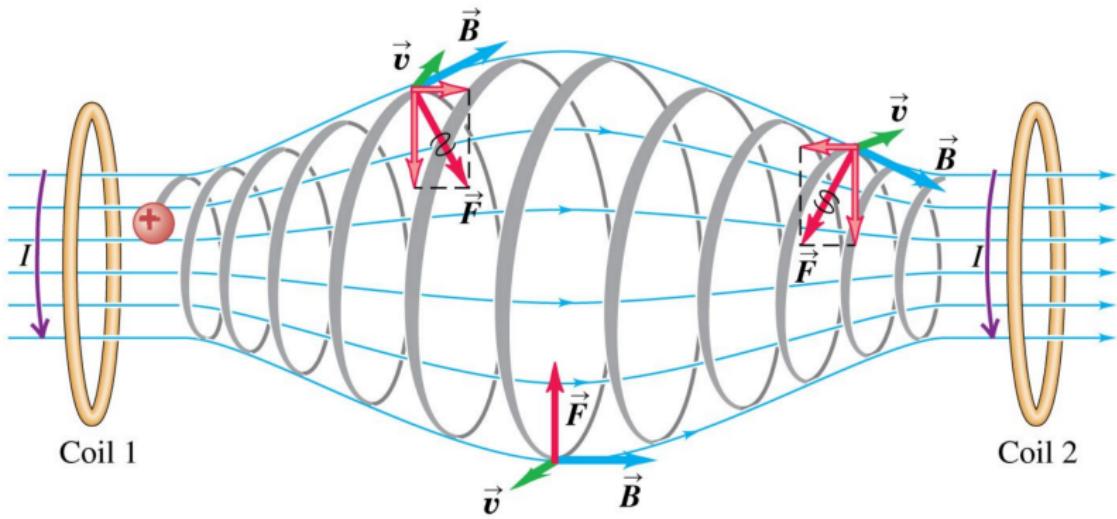
Magnetic Force On a Moving Charge

Examples

Applications

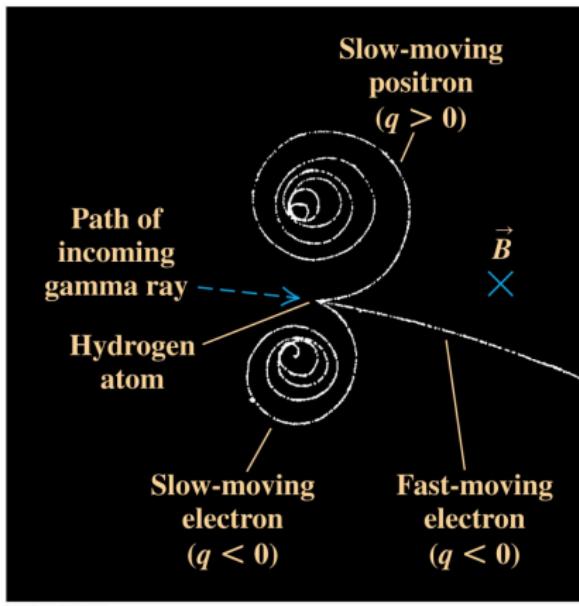
Applications

Magnetic Trap



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Bubble Chamber

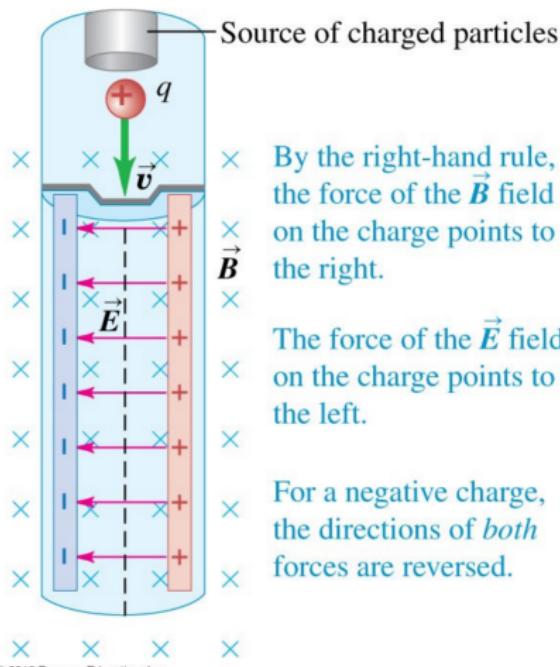


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Velocity Selector

(a) Schematic diagram of velocity selector



- × By the right-hand rule, the force of the \vec{B} field on the charge points to the right.
 - × The force of the \vec{E} field on the charge points to the left.
 - × For a negative charge, the directions of *both* forces are reversed.

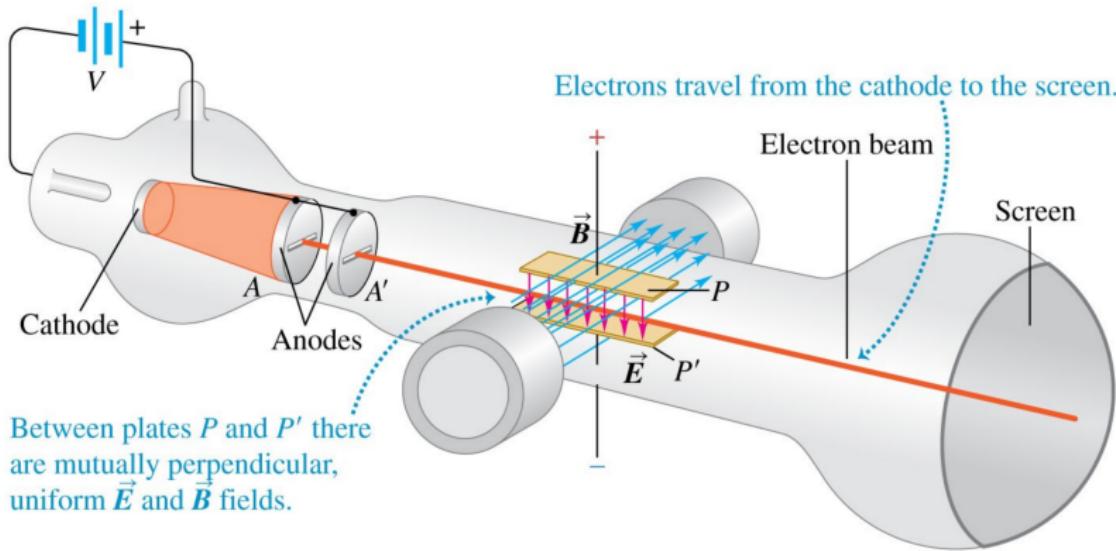
(b) Free-body diagram for a positive particle

$$F_E = qE \quad F_B =$$

Only if a charged particle has $v = E/B$ do the electric and magnetic forces cancel. All other particles are deflected.

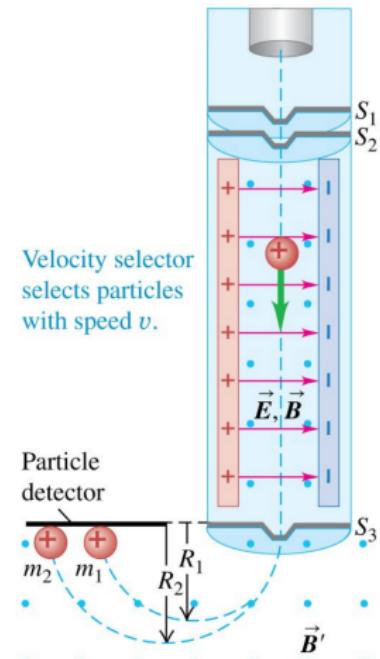
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Thomson's Experiment (e/m Ratio)



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Mass Spectrometer



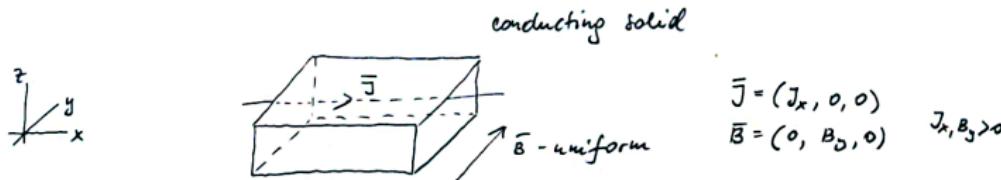
Magnetic field separates particles by mass;
the greater a particle's mass, the larger is
the radius of its path.

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The Hall Effect

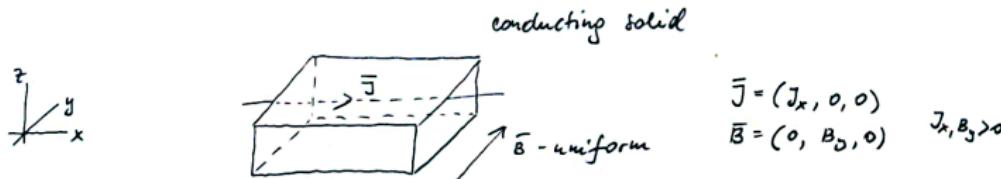
The Hall Effect

Consider electric current flowing through a conducting slab placed in a uniform magnetic field.



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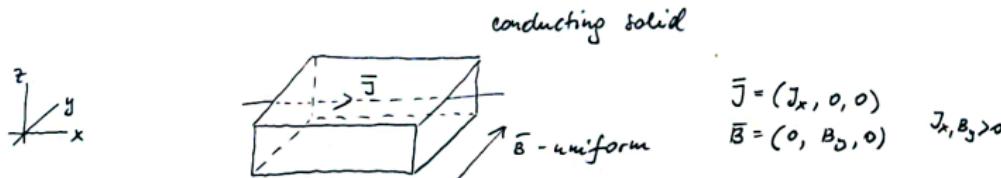


xz cross-section



The Hall Effect

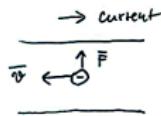
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xz cross-section



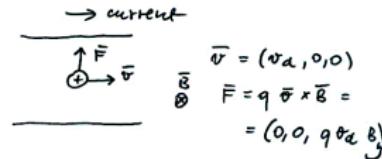
negative charge carriers



$$q_i v_d > 0$$

$$\vec{v} = (-v_d, 0, 0)$$
$$\vec{F} = -q (\vec{v} \times \vec{B}) =$$
$$= (0, 0, q v_d B)$$

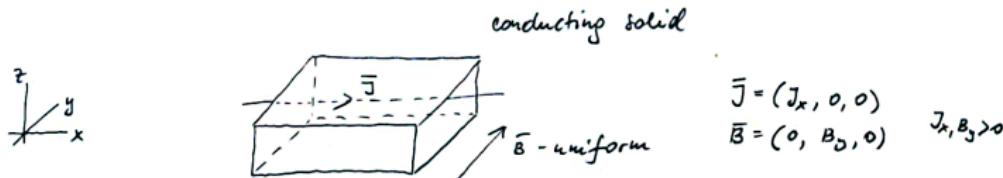
positive charge carriers



Observation. Charges are deflected towards the surface of the slab

The Hall Effect

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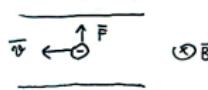


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negative charge carriers

→ current



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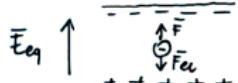


$$\vec{v} = (v_d, 0, 0)$$

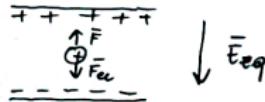
$$\begin{aligned}\vec{F} &= q \vec{v} \times \vec{B} \\ &= (0, 0, q v_d B)\end{aligned}$$

Observation. Charges are deflected towards the surface of the slab \implies electric field in the z direction builds up

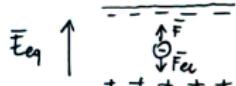
Equilibrium FBD



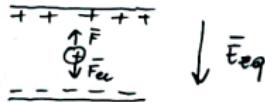
note the direction
of \bar{F}_{eq} in both cases!



Equilibrium FBD



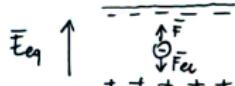
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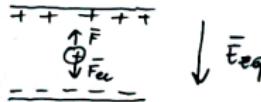
Equilibrium condition

$$\bar{F}_{eq} + \bar{F} = 0$$

Equilibrium FBD



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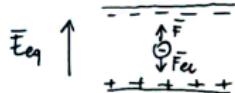


Equilibrium condition

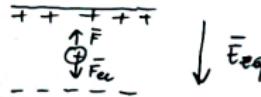
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Equilibrium FBD



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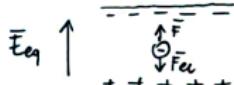
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$$\boxed{\bar{E}_{eq} = v_e B_y}$$

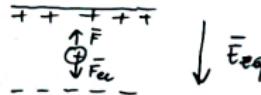
in the positive z-direction for
negative charges

in the negative z-direction for
positive charges)

Equilibrium FBD



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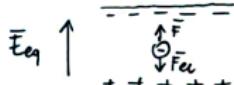
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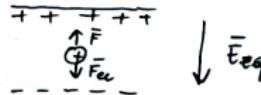
Comments

- Can be used to determine the sign of charge carriers.
- If the concentration of charge carriers is known, can be used to measure the magnetic field (Hall probe).

Equilibrium FBD



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Equilibrium condition

$$\bar{F}_{eq} + \bar{F} = 0$$

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$$\boxed{\bar{E}_{eq} = v_a B_y}$$

in the positive z-direction for
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$$\bar{J}_x = nq \bar{v}_a$$

$$nq = \frac{\bar{J}_x}{\bar{v}_a} = \frac{\bar{J}_x}{\bar{E}_{eq}} B_y \quad \Rightarrow$$

$$\boxed{B_y = \frac{nq \bar{E}_{eq}}{\bar{J}_x}}$$

Current–Carrying Conductor In Magnetic Field

Magnetic Force On a Current–Carrying Conductor

Starting point

$$\overline{F} = q \overline{v} \times \overline{B}$$

Magnetic Force On a Current–Carrying Conductor

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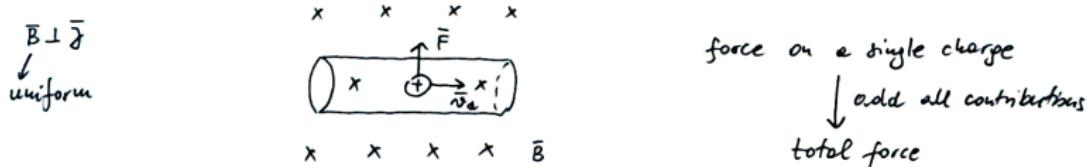
Assume positive charges, all move with \overline{v}_d along a straight-line conductor; uniform magnetic field.

Magnetic Force On a Current-Carrying Conductor

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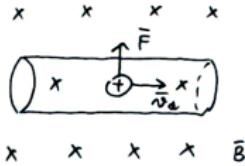
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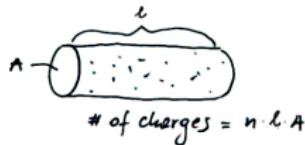
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$\overline{B} \perp \overline{v}$
uniform



force on a single charge
↓ add all contributions
total force



$$\text{total force} = \# \text{ of charges} \cdot \text{force on a single charge}$$

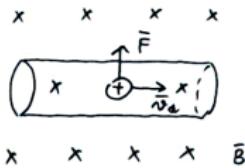
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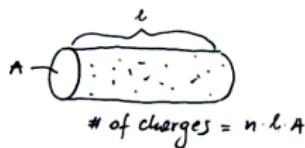
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force on a single charge
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total force = # of charges · force on a single charge

$$F = (n l A) q v_d B$$

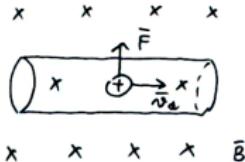
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Starting point

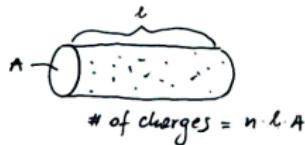
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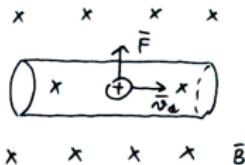
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Starting point

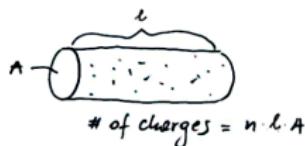
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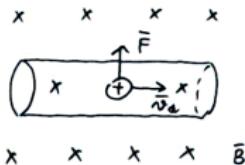
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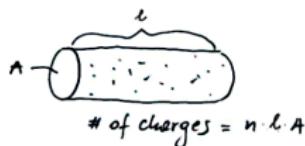
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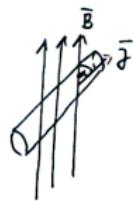
force on a single charge
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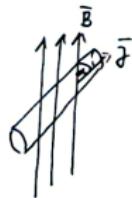
$$F = (n l A) q v_d B = \underbrace{(n q v_d) A l B}_I = I l B$$

$\bar{B}, \bar{x}, \bar{y}$

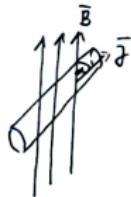


$$F = (n\ell A)qv_d B \sin \angle(\vec{J}, \vec{B})$$

$\vec{B} \times \vec{J}$

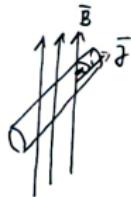


$\bar{B} \times \bar{J}$



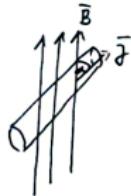
$$F = (n\ell A)qv_d B \sin \angle(\bar{J}, \bar{B}) = \underbrace{(nqv_d)_l}_I \ell B \sin \angle(\bar{J}, \bar{B})$$

$\bar{B} \times \bar{J}$



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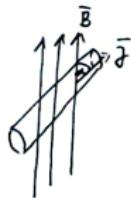


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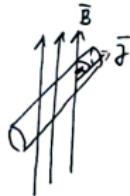
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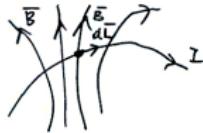


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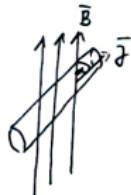
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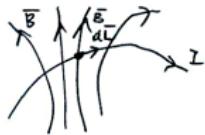


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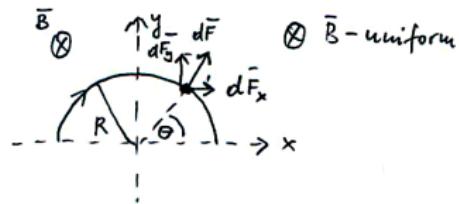
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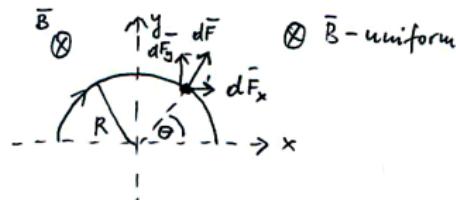


$$\bar{F} = \int_{\text{conductor}} I d\bar{\ell} \times \bar{B}.$$

Example



Example

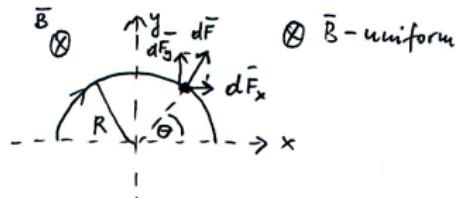


$$dF_x = \cos\theta \ dF$$

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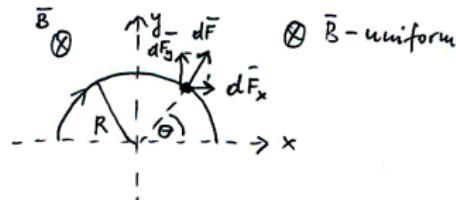
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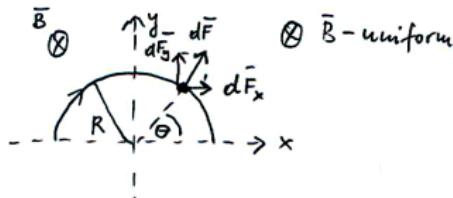
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$$d\theta = \frac{d\ell}{R}$$

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$$d\theta = \frac{d\bar{l}}{R}$$

$$d\bar{l} = R d\theta$$

$$= \int_0^\pi I R B \sin\theta d\theta = I B R \cdot 2$$

$$F_x = 0 \quad (\text{same method, or symmetry})$$

Current Loop In Uniform Magnetic Field

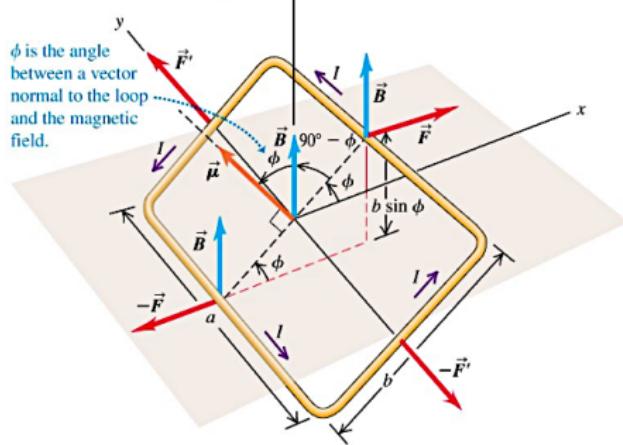
Current Loop in Uniform Magnetic Field

(a)

The two pairs of forces acting on the loop cancel, so no net force acts on the loop.

However, the forces on the a sides of the loop (\vec{F} and $-\vec{F}$) produce a torque

$$\tau = (IBa)(b \sin\phi)$$
 on the loop.



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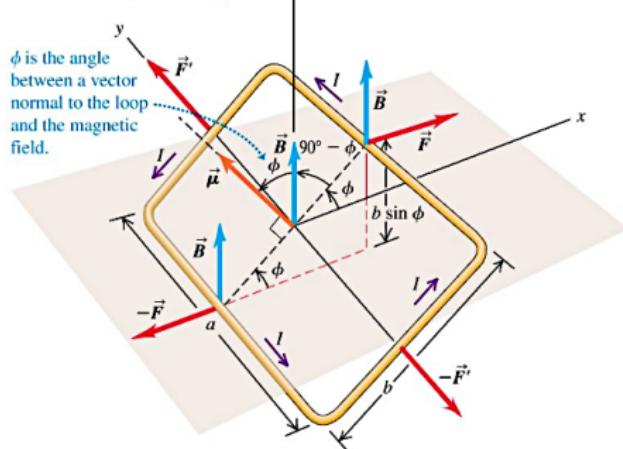
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Observation. The net force on the loop is always zero.

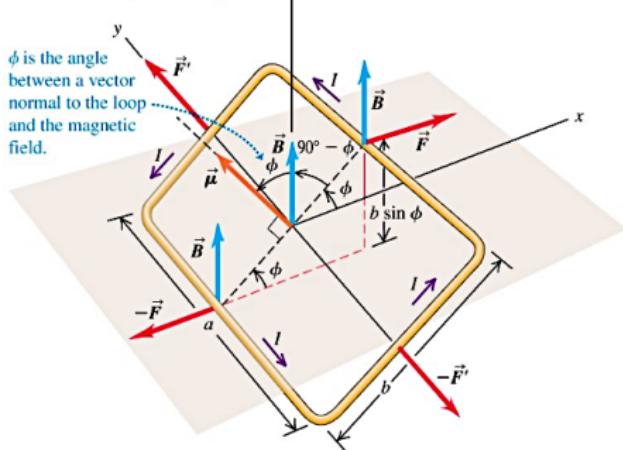
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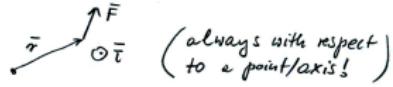
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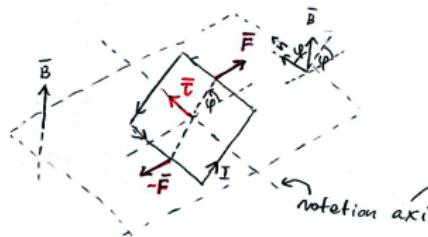
Observation. The net force on the loop is always zero. However, the net torque is not.

Torque $\bar{\tau} = \bar{r} \times \bar{F}$

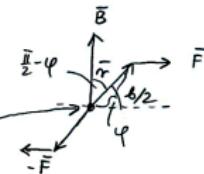


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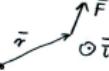
$\bar{\tau} \circlearrowleft \bar{F}$ (always with respect to a point/axis!)

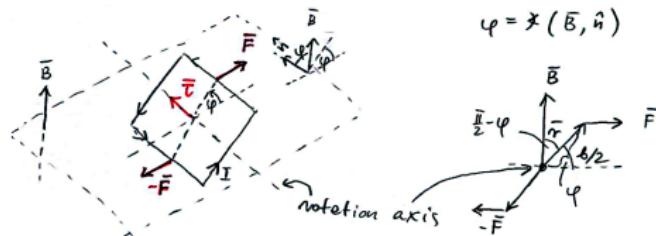


$$\varphi = \chi(\bar{B}, \hat{n})$$



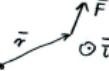
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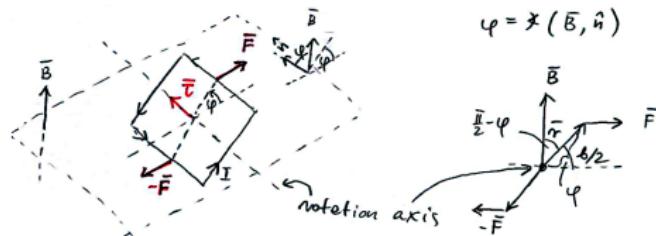
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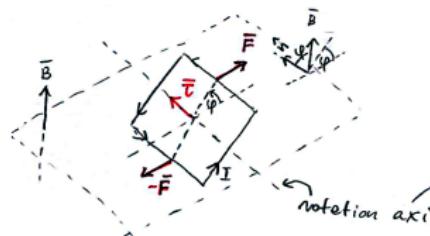
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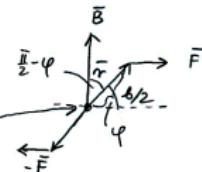
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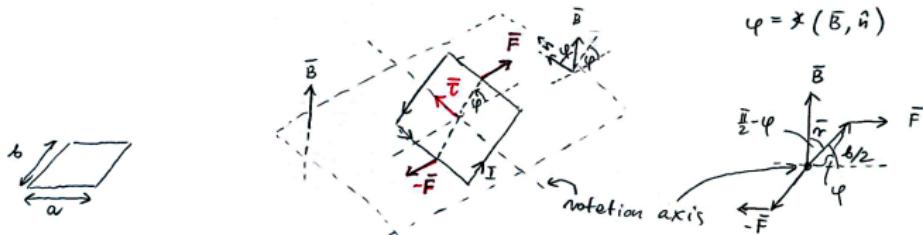
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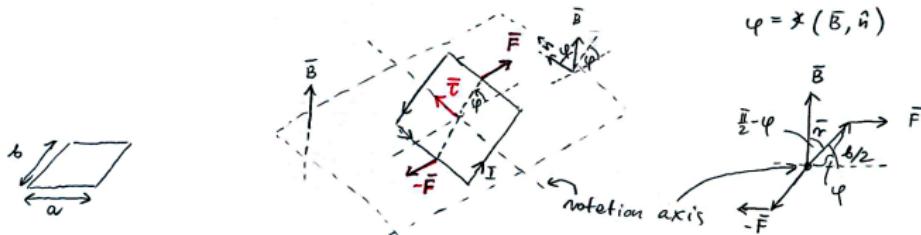
Stable equilibrium

unstable equilibrium

τ_1 positive direction of \vec{v}_p

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\vec{F}
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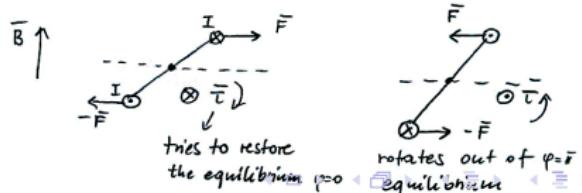


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Magnetic Dipole Moment

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Torque on a current loop placed in uniform magnetic field

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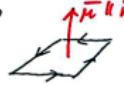
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Compare: Electric dipole moment.

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Potential energy (see derivation for the electric dipole)

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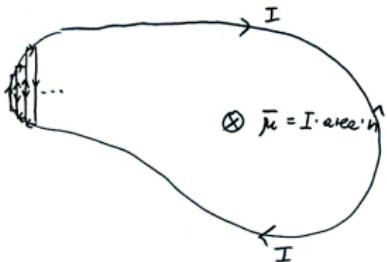
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The definition of the magnetic dipole moment and all discussions for uniform magnetic field can be generalized to an **arbitrarily shaped loop**.

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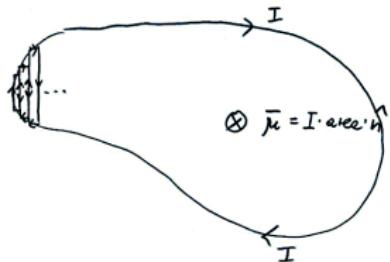
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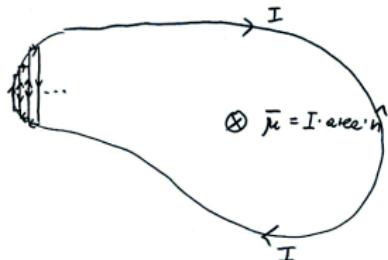


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Example. Solenoid (coil)

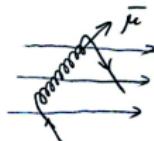
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Example. Solenoid (coil)



collection of N loops of area A

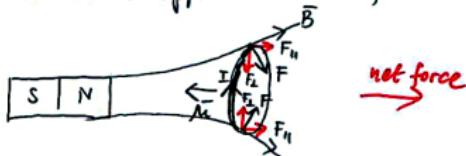
$$\boxed{\vec{\mu} = N \cdot I A \hat{n}}$$

right-hand rule

Generalization (II). Non-Uniform Magnetic Field

What happens if a magnetic dipole moment (current loop) is placed in a non-uniform magnetic field?

- (1) $\vec{\mu}$ in the direction opposite to the field



net components

$$F_{\perp} = 0$$

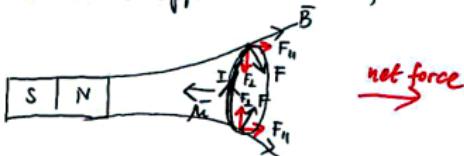
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\Downarrow
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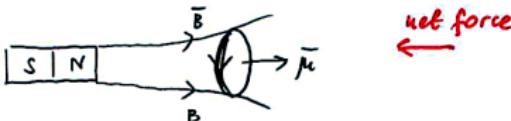
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\Downarrow
repulsion

- (2) $\bar{\mu}$ in the direction of the field



$$F_{\perp} = 0$$

$$F_{\parallel} \neq 0 \text{ (to the left)} \Rightarrow \text{attraction}$$

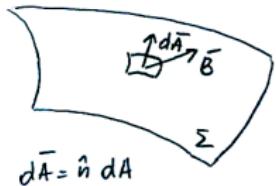
Flux of Magnetic Field (Magnetic Flux) and Gauss's Law for Magnetic Field

Magnetic Flux

Definition of the flux of the magnetic field (magnetic flux) is fully analogous to that of the electric flux.

Magnetic Flux

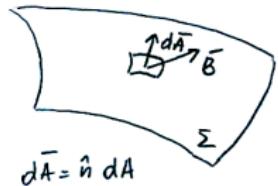
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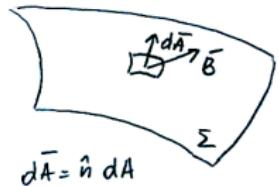
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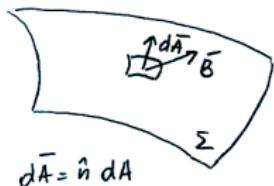
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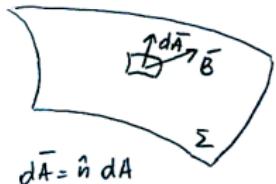
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Total flux through the surface Σ

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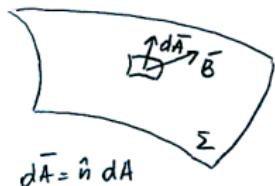
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SI units of the magnetic flux: Wb (Weber)

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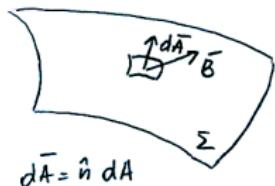
Total flux through the surface Σ

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Gauss's Law for Magnetic Field

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Gauss's Law for Magnetism

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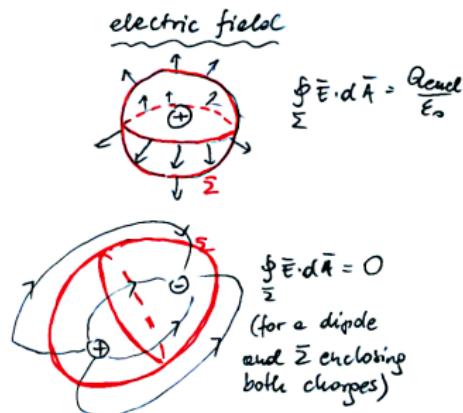
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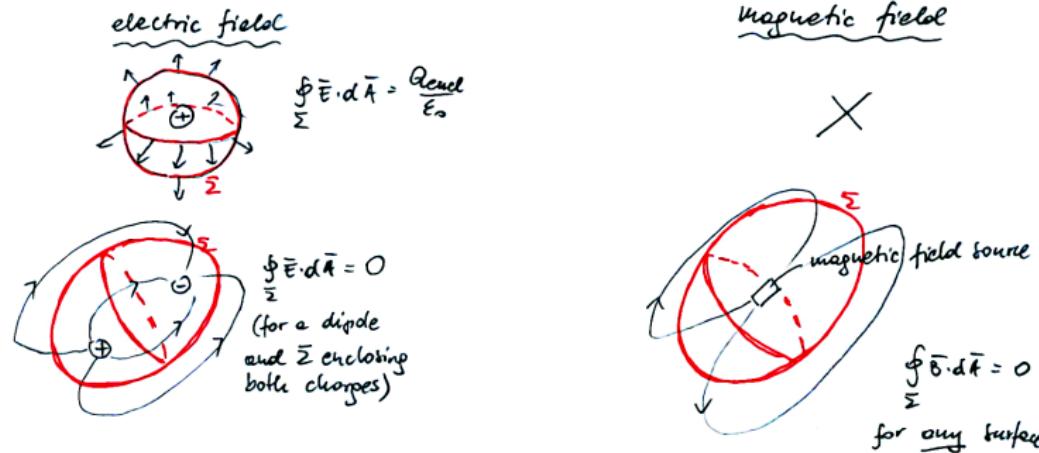


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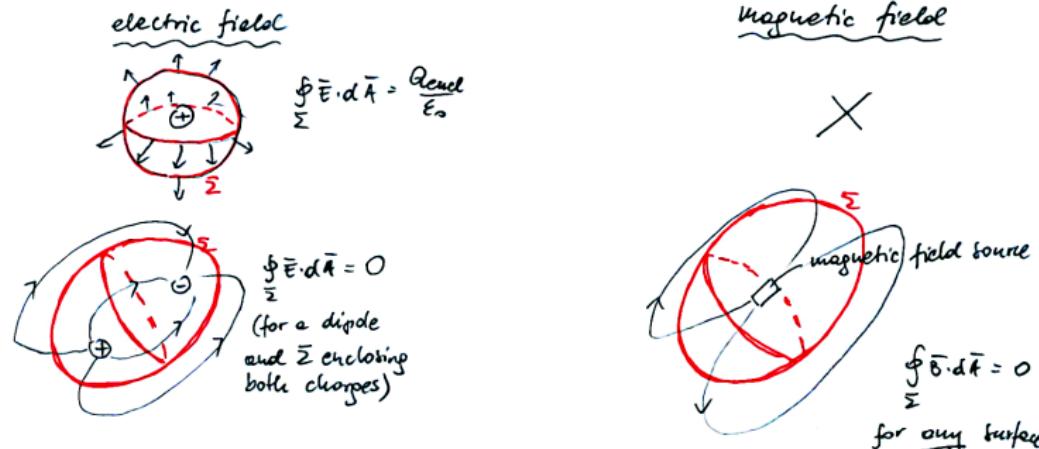


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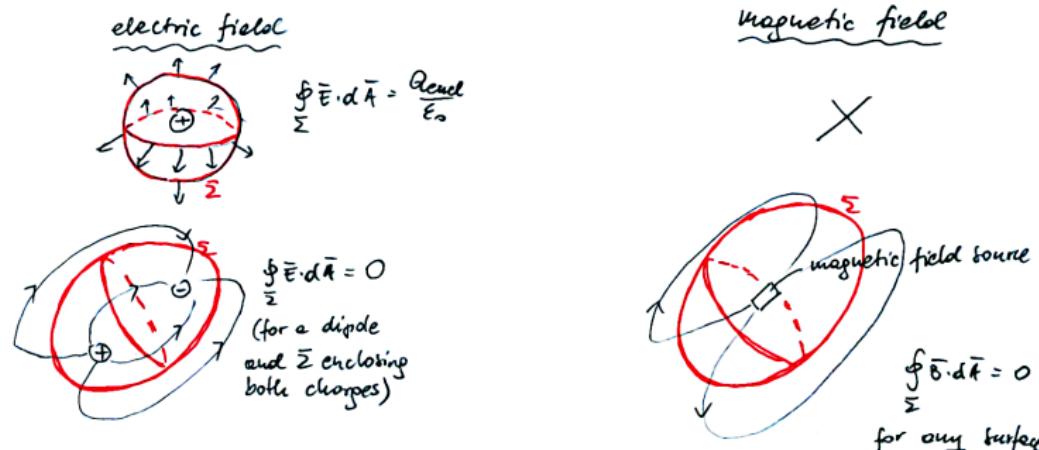
Electric field lines begin and end on electric charges

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Electric field lines begin and end on electric charges; magnetic field lines always form loops.

Gauss's Law For Magnetic Field. Differential Form

The differential form of Gauss's Law for the magnetic field can be derived in the same manner as for the electric field

$$\oint_{\Sigma} \overline{B} \circ d\overline{A} \stackrel{\text{divergence thm}}{=} \int_{\Omega} \operatorname{div} \overline{B} d\tau$$

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Recall the interpretation of divergence

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Recall the interpretation of divergence; there are no magnetic monopoles (no magnetic "point charges").