

**R&D OF A HIGH-PERFORMANCE DIRC DETECTOR FOR  
USE IN AN ELECTRON-ION COLLIDER**

by

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## **ABSTRACT**

### R&D OF A HIGH-PERFORMANCE DIRC DETECTOR FOR USE IN AN ELECTRON-ION COLLIDER

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## ACKNOWLEDGEMENTS

# Contents

<b>LIST OF TABLES</b>	vii
<b>LIST OF FIGURES</b>	viii
Chapter	
<b>1. INTRODUCTION</b>	<b>2</b>
<b>2. ELECTRON-ION COLLIDER</b>	<b>3</b>
2.1 SCIENCE GOALS . . . . .	3
2.2 FACILITIES . . . . .	4
<b>3. DIRC TECHNOLOGY</b>	<b>7</b>
3.1 CHERENKOV RADIATION . . . . .	7
3.2 APPLYING THE CHERENKOV EFFECT TO PARTICLE ID . . . . .	8
3.3 RING IMAGING DETECTORS . . . . .	9
3.4 DIRC DETECTORS . . . . .	10
<b>4. HIGH-PERFORMANCE DIRC@EIC</b>	<b>12</b>
4.1 EVOLUTION OF THE DIRC@EIC DESIGN . . . . .	12
4.2 CURRENT HIGH-PERFORMANCE DIRC DESIGN . . . . .	12
4.3 SIMULATED PERFORMANCE . . . . .	12
4.4 POTENTIAL OPTIMIZATIONS . . . . .	12
<b>5. TEST BENCH EVALUATION OF DIRC@EIC COMPONENTS</b>	<b>13</b>
5.1 OPTICAL PROPERTIES OF 3-LAYER LENS . . . . .	13
5.2 RADIATION HARDNESS OF NLAK33 MATERIAL . . . . .	13
5.3 PERFORMANCE OF PHOTOSENSORS IN HIGH MAGNETIC FIELD . . . . .	13
<b>6. 3-LAYER LENS PERFORMANCE IN PARTICLE BEAM</b>	<b>14</b>
6.1 PROTOTYPE SETUP . . . . .	14
6.2 SIMULATED PERFORMANCE . . . . .	14
6.3 DATA ANALYSIS . . . . .	14
<b>7. OUTLOOK AND SUMMARY</b>	<b>15</b>

<b>BIBLIOGRAPHY</b>	<b>15</b>
APPENDICES	
A. ERROR EVALUATION . . . . .	17
<b>VITA</b>	<b>18</b>

# List of Tables

# List of Figures

2.1.	Evolution of our understanding of nucleon spin structure. <b>Left:</b> In the 1980s, a nucleons spin was naively explained by the alignment of the spins of its constituent quarks. <b>Right:</b> In the current picture, valence quarks, sea quarks and gluons, and their possible orbital motion are expected to contribute to overall nucleon spin. Figure from reference [1] . . . . .	4
2.2.	Current design of the EIC facility for JLab with the two interaction points highlighted in red. . . . .	5
2.3.	Current design of the EIC facility for BNL. . . . .	6
2.4.	Current design of the detector to be used in the JLab EIC at IP1 . . . . .	6
3.1.	Illustration of the Cherenkov cone. . . . .	8
3.2.	Particle momentum versus Cherenkov angle for different particle species in fused silica ( $n \approx 1.473$ ). . . . .	9
3.3.	Typical Ring Imaging Cherenkov (RICH) detector. . . . .	10
3.4.	Annulus ring of RICH detector. . . . .	11
3.5.	The basic components of a DIRC detector: a solid radiator, typically fused silica (green); a mirror to redirect backward-going photons (pink); optional focusing optics (purple); an expansion volume to allow photons to separate in space (cyan); and a detector surface to record the position and arrival time of Cherenkov photons (maroon). . . . .	11

# Comments

is this the correct citation to use?? . . . . .	3
maybe these are unnecessary subsections? Could probably just expand on the science goals section a little more . . . . .	4
I feel like I should say more about eRHIC here since I said so much about JLEIC . . . . .	5

## CHAPTER 1

### INTRODUCTION

## CHAPTER 2

### ELECTRON-ION COLLIDER

It has been known for nearly a century that atoms are composed of nucleons (protons and neutrons), but it took another 50 years for Murray Gell-Mann and George Zweig to independently develop a model proposing that nucleons themselves are made up of constituent components, called quarks, bound together by the exchange of gluons [2]. This lead to the development of the fundamental theory of the strong interaction, known as Quantum Chromo-Dynamics (QCD). It is now a strong goal of the nuclear physics community to understand the interactions of quarks and gluons and how those interactions make manifest both nucleons themselves, which account for nearly all the mass of the visible matter in the universe, as well as the nucleons' spin, mass, and magnetic moment.

Although it would theoretically be possible to study these properties using fixed-target electron beam experiments, it is three-fold prohibitive: it is much more costly to construct an accelerator to accelerate electrons to the necessary momentum (on the order of TeV) than to build a collider, it is more difficult and complicated to do transverse nucleon polarization studies with a fixed target due to the nature of the magnetic fields required, and it is very difficult to study the beam fragments of a fixed target reactions due to final state interactions whereas in a collider the fragments will be boosted in the same direction as the beam [3] **is this the correct citation to use??**. It was therefore deemed a priority by the 2007 Nuclear Science Advisory Committee's Long-Range Plan that an Electron-Ion Collider (EIC) be the next facility to be built in the United States.

#### 2.1 SCIENCE GOALS

One major question still pestering nuclear physicists is "What is the origin of the nucleon spin?". In the 1980s the naive answer was that the total nucleon spin was the sum of the spin of its three valance quarks, but many years of experimentation has revealed that it is much more complicated (Fig. 2.1). The EIC will be capable of much more detailed study of the contributions to the nucleon structure by enabling multi-dimensional projections of the distribution of quarks and gluons in space, longitudinal and transverse momenta, spin, and flavor.

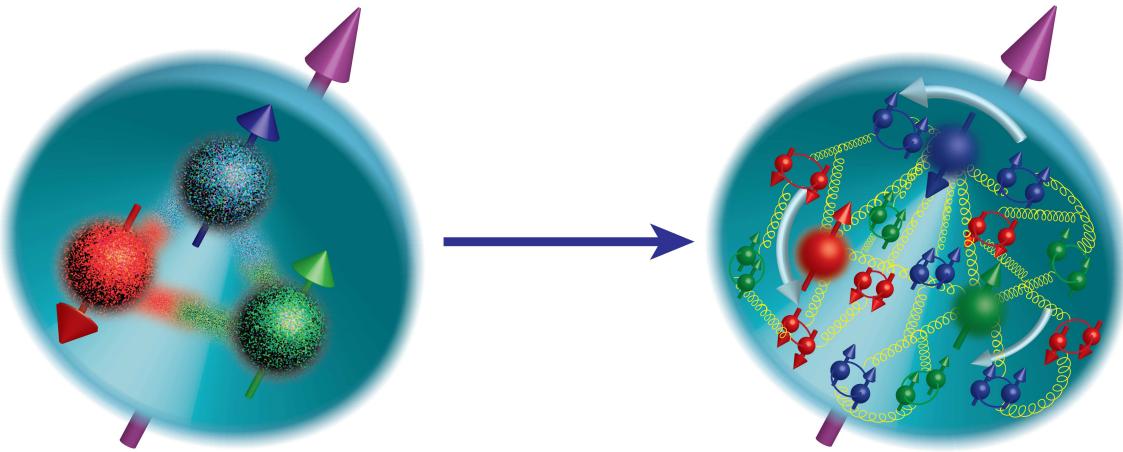


FIG. 2.1: Evolution of our understanding of nucleon spin structure. **Left:** In the 1980s, a nucleon's spin was naively explained by the alignment of the spins of its constituent quarks. **Right:** In the current picture, valence quarks, sea quarks and gluons, and their possible orbital motion are expected to contribute to overall nucleon spin. Figure from reference [1]

### 2.1.1 NUCLEON SPIN

**maybe these are unnecessary subsections? Could probably just expand on the science goals section a little more**

### 2.1.2 QUARK/GLUON DISTRIBUTION

## 2.2 FACILITIES

As of the writing of this thesis there are two competing designs for an EIC facility to be built in the United States: a figure-8 accelerator design for Thomas Jefferson National Accelerator Facility (JLab) (Figure 2.2), and a ring-ring accelerator design for Brookhaven National Lab (BNL) (Figure 2.3).

The JLab EIC (JLEIC) is planned to be approximately 1.4 km in circumference and have a footprint of roughly 500 m by 170 m. The design is a ring-ring with electrons and ions being stored in separate beam lines and collided at two interaction points (IPs) (outlined in red in Figure 2.2) on the figure-8. The JLab CEBAF SRF linac will be used as an electron injector for electrons with 3 - 11 GeV energy. The second ring will store an ion beam with energy of 20 to 100 GeV for protons or up to 40 GeV per nucleon for light to heavy ions.

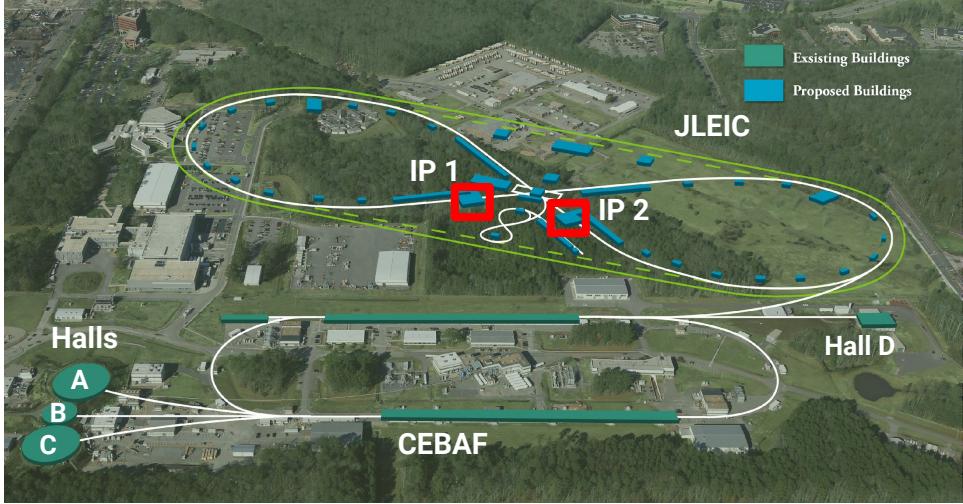


FIG. 2.2: Current design of the EIC facility for JLab with the two interaction points highlighted in red.

The ion beams are generated and accelerated in a new ion injector complex with the same figure-8 design that will be utilized to preserve ion polarization. The two rings will be stacked vertically in the same underground tunnel [4].

The BNL facility, named eRHIC, will use a new electron beam facility based on an Energy Recovery LINAC that will be built inside of the Relativistic Heavy Ion Collider (RHIC) tunnel to collide with RHIC's pre-existing polarized proton/ion beam (Fig. 2.3). **I feel like I should say more about eRHIC here since I said so much about JLEIC**

### 2.2.1 JLEIC DETECTOR DESIGN

The large center of mass energies of reactions at an EIC necessitate a very sophisticated detector system. Figure 2.4 shows the current design of the JLab EIC detector at IP1. This

### 2.2.2 PARTICLE IDENTIFICATION

The ability to accurately identify hadrons in the final state is a key requirement for the physics program at an EIC.

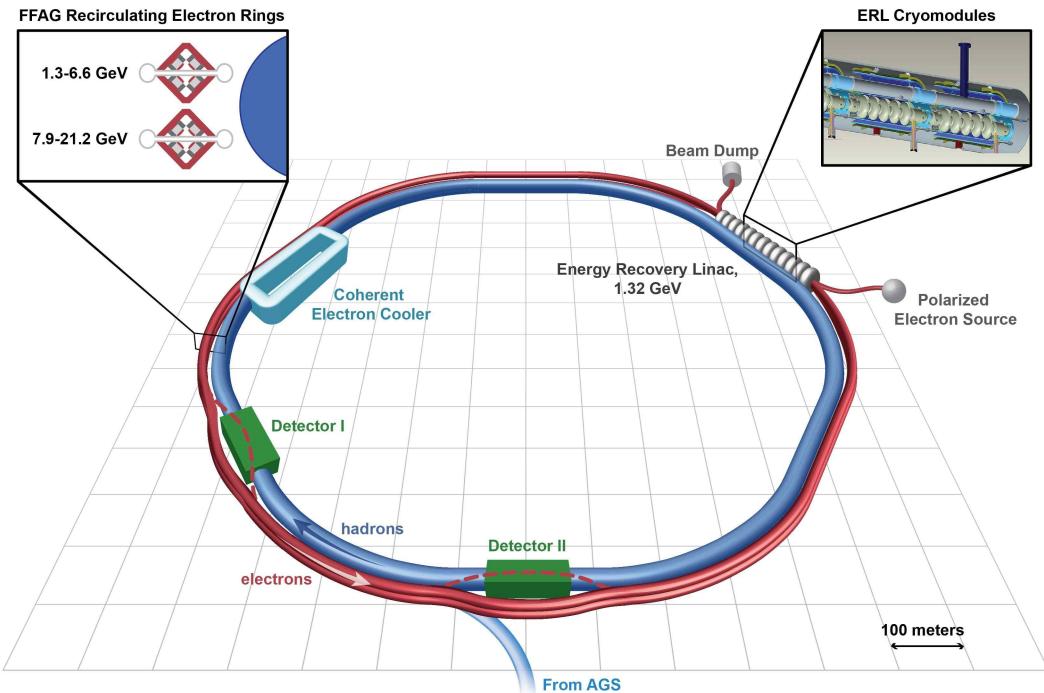


FIG. 2.3: Current design of the EIC facility for BNL.

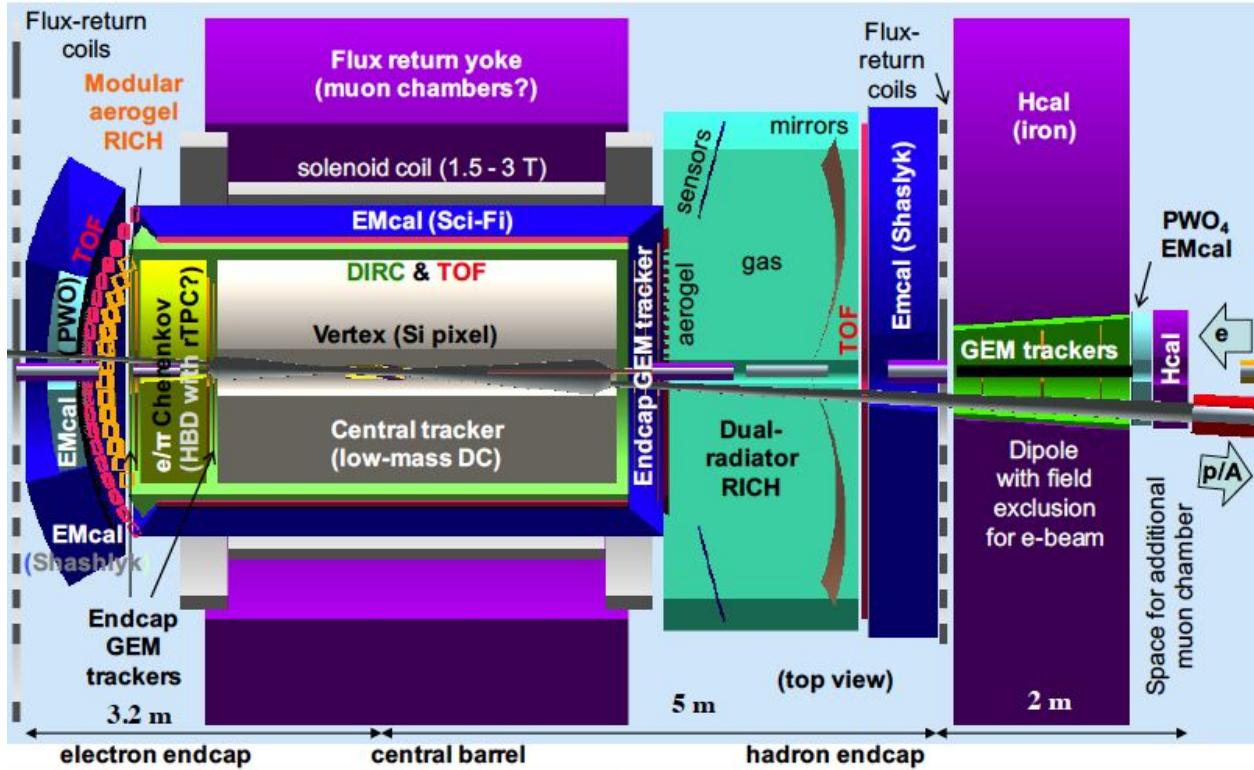


FIG. 2.4: Current design of the detector to be used in the JLab EIC at IP1

# CHAPTER 3

## DIRC TECHNOLOGY

The DIRC detector is based on the concept of the Detection of Internally Reflected Cherenkov light (DIRC) produced in a solid radiator bar to identify charged particle. It is a special type of Cherenkov counter, which uses the unique properties of the Cherenkov radiation.

### 3.1 CHERENKOV RADIATION

Einstein postulated in his Theory of Relativity that the speed of light in a vacuum,  $c$ , is the limit of the velocity of massive particles. In an optically transparent medium, however, the speed at which light propagates is modified:  $c_{med} = c/n$ , where  $n$  is the index of refraction of the medium. Pavel Cherenkov discovered in 1934 that massive particles moving through a medium faster than the speed of light in that medium emit light in the form of now-called Cherenkov radiation. Cherenkov was able to establish several interesting properties of this radiation: it is only emitted from charged particles above a certain velocity threshold  $v > c/n$ , the intensity is proportional to the particle's path length, emission is prompt, and the light is polarized with a continuous wavelength spectrum. Later in 1937 Ilya Frank and Igor Tamm theoretically formulated this radiation with fantastic agreement to Cherenkov's findings, and the three shared the 1958 Nobel Prize in Physics for their efforts [5].

Further studies confirmed that Cherenkov radiation is emitted uniformly in azimuth ( $\phi_c$ ) around the particle's direction of travel with the polar opening angle  $\theta_C$  defined as

$$\cos \theta_C = \frac{1}{\beta n(\lambda)}, \quad (3.1)$$

where  $\beta = v_p/c$ ,  $v_p$  is the particle's velocity, and the index of refraction is a function of the emitted photon wavelength. In a normal, dispersive optical medium the opening half-angle of the shock wave produced by the Cherenkov radiation,  $\eta_C$  defined in Figure 3.1, is not complementary to the Cherenkov angle. The relationship between the two is given by

$$\cot \eta_C = \left[ \frac{d}{d\omega} (\omega \tan \theta_C) \right]_{\omega_0} = \left[ \tan \theta_C + \beta^2 \omega n(\omega) \frac{dn}{d\omega} \cot \theta_C \right]_{\omega_0} \quad (3.2)$$

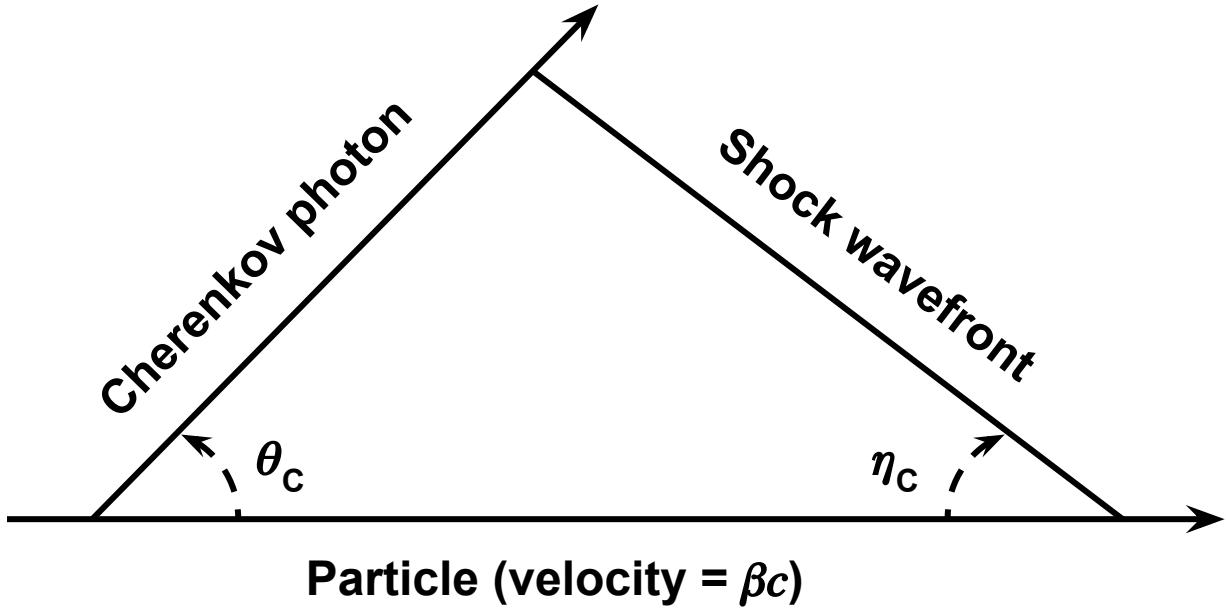


FIG. 3.1: Illustration of the Cherenkov cone.

where  $\omega_0$  is the central value of the considered frequency range. Because the second term in (3.2) is zero only for non-dispersive media the shock wave front is not perpendicular to the Cherenkov cone in real detectors.

Because particles lose very little energy when radiating Cherenkov photons the emission is very weak. The number of photons  $N_{photons}$  emitted per path length  $L$  (in cm) by a moving particle with charge  $z$  is given by the Frank-Tamm equation

$$\frac{N_{photons}}{L} = \frac{\alpha^2 z^2}{r_e m_e c^2} \int \sin^2 \theta_C(E) dE \quad (3.3)$$

where  $E$  is the photon energy in eV, the integral is taken over the region where  $n(E)$  is greater than 1, and  $\frac{\alpha^2 z^2}{r_e m_e c^2} = 370 \text{ cm}^{-1} \text{ eV}^{-1}$ .

### 3.2 APPLYING THE CHERENKOV EFFECT TO PARTICLE ID

In order to identify particle species one must know both the mass and charge of the particle in question. Because the Cherenkov angle encodes the particle's velocity it is, in principle, a simple matter to measure the particle's momentum with a tracking chamber as well as the velocity obtained from (3.1) to determine the mass and charge. Figure 3.2 shows how different particle species can be distinguished for a given momentum in fused silica.

Cherenkov counters are detectors used for particle identification (PID) by exploiting the fact that only particles above the threshold velocity  $\beta > 1/n$  will emit Cherenkov photons. The information about a particle's velocity can be combined with momentum information from a tracking system to determine the mass as [6]

$$m = \frac{p}{c} \sqrt{n^2 \cos^2 \theta_C - 1} \quad (3.4)$$

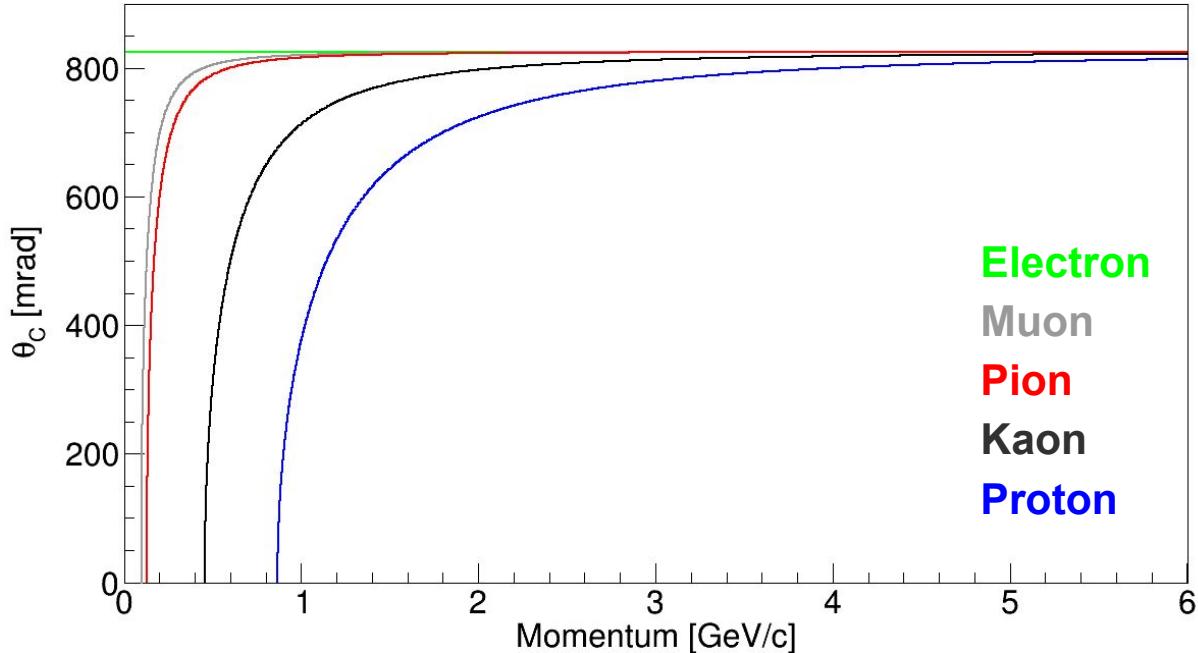


FIG. 3.2: Particle momentum versus Cherenkov angle for different particle species in fused silica ( $n \approx 1.473$ ).

### 3.3 RING IMAGING DETECTORS

Ring Imaging Cherenkov (RICH) detectors are designed to efficiently identify and separate different particle species over a wide range of momenta. A basic RICH system is shown in Figure 3.3.

A volume of radiator, either gaseous (e.g.  $C_4F_{10}$ ) or solid (e.g. aerogel), is positioned upstream of an array of photosensors. A charged particle travelling through the radiator above the threshold velocity will continuously emit Cherenkov photons in a cone. The resulting image on the photosensor array is an annulus of thickness  $L \tan \theta_C$ , where  $L$  is the distance the particle travelled inside the radiator, and  $\theta_C$  is the usual Cherenkov angle (Figure

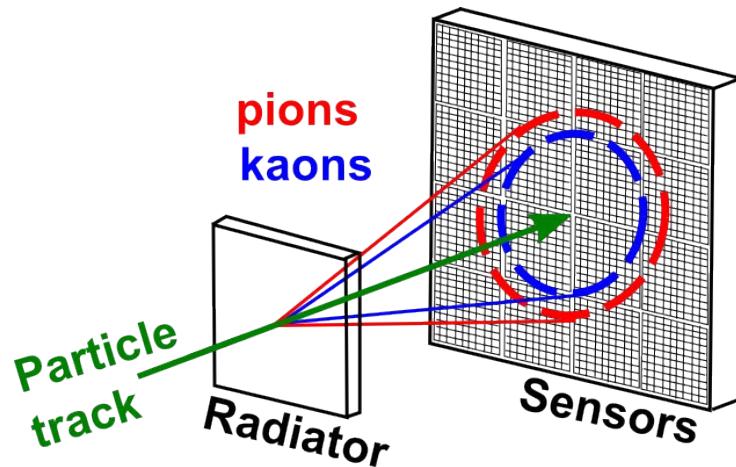


FIG. 3.3: Typical Ring Imaging Cherenkov (RICH) detector.

3.4). PID can be done by measuring the average radius of the annulus and reconstructing the Cherenkov angle geometrically.

### 3.4 DIRC DETECTORS

The basic components of a DIRC detector are shown in Figure 3.5.

#### 3.4.1 DIRCS IN FUTURE EXPERIMENTS

EIC, PANDA, Bell, ToRCH, GlueX

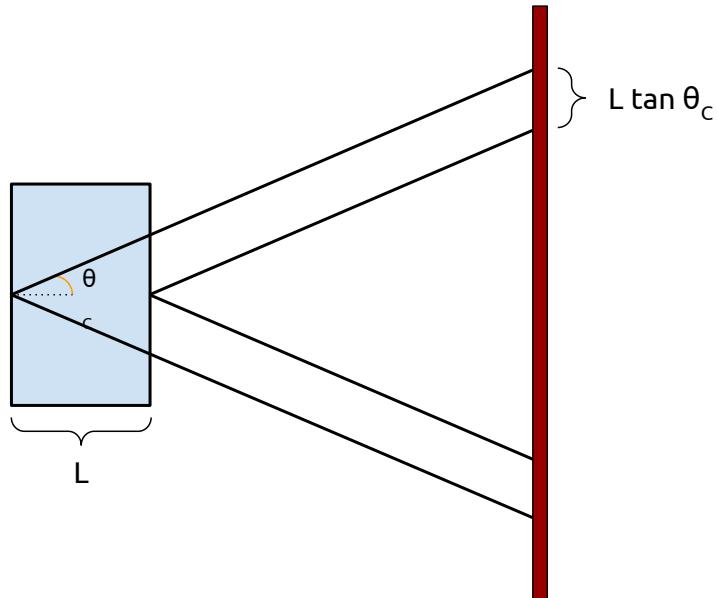


FIG. 3.4: Annulus ring of RICH detector.

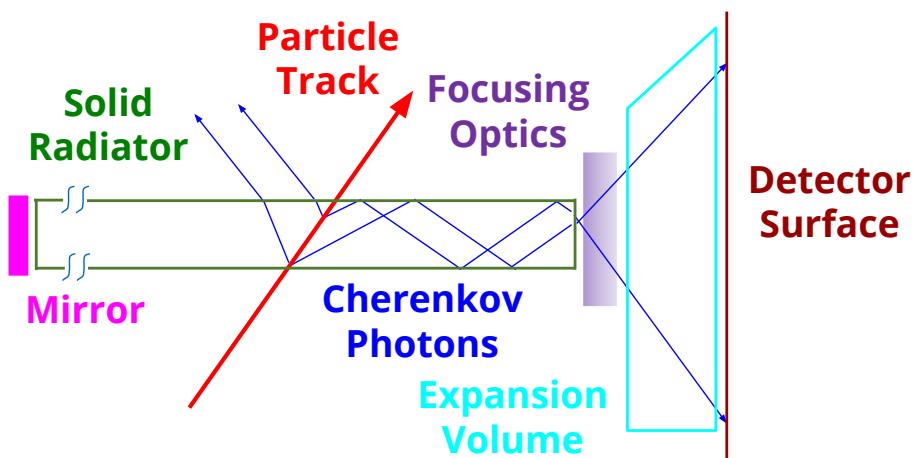


FIG. 3.5: The basic components of a DIRC detector: a solid radiator, typically fused silica (green); a mirror to redirect backward-going photons (pink); optional focusing optics (purple); an expansion volume to allow photons to separate in space (cyan); and a detector surface to record the position and arrival time of Cherenkov photons (maroon).

## CHAPTER 4

### HIGH-PERFORMANCE DIRC@EIC

4.1 EVOLUTION OF THE DIRC@EIC DESIGN

4.2 CURRENT HIGH-PERFORMANCE DIRC DESIGN

4.3 SIMULATED PERFORMANCE

4.4 POTENTIAL OPTIMIZATIONS

## CHAPTER 5

### TEST BENCH EVALUATION OF DIRC@EIC COMPONENTS

5.1 OPTICAL PROPERTIES OF 3-LAYER LENS

5.2 RADIATION HARDNESS OF NLAK33 MATERIAL

5.3 PERFORMANCE OF PHOTOSENSORS IN HIGH MAGNETIC  
FIELD

## CHAPTER 6

### 3-LAYER LENS PERFORMANCE IN PARTICLE BEAM

6.1 PROTOTYPE SETUP

6.2 SIMULATED PERFORMANCE

6.3 DATA ANALYSIS

6.3.1 ERROR EVALUATION

## CHAPTER 7

### OUTLOOK AND SUMMARY

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## APPENDIX A

### ERROR EVALUATION

**VITA**

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