

# **Family Floer mirror space for local SYZ singularities**

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Consider a Lagrangian fibration with singularities

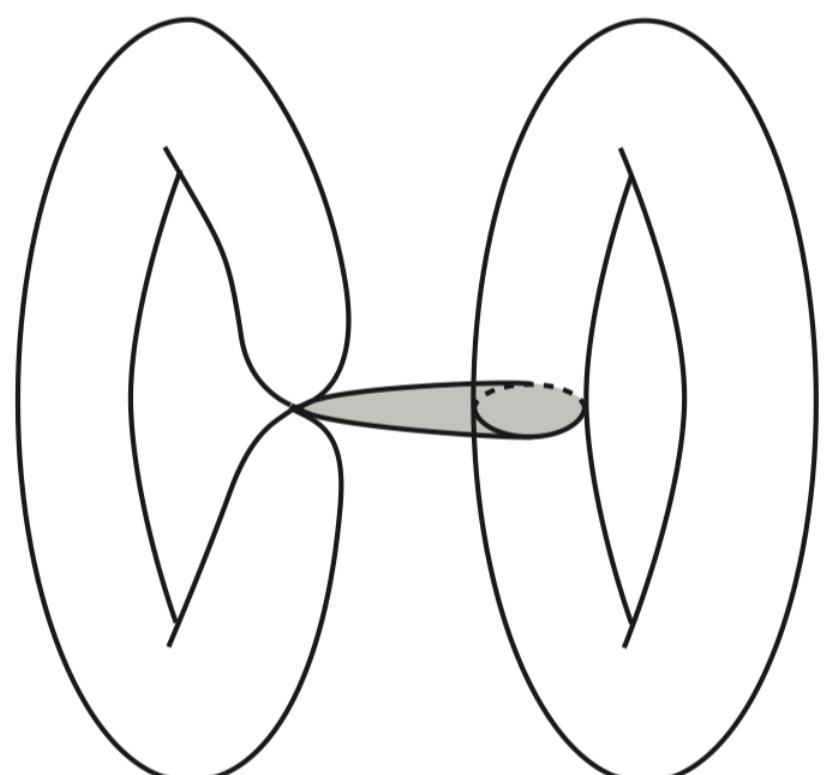
$$\pi : X \rightarrow B$$

on a Kähler manifold  $(X, \omega)$ . Denote its smooth part by

$$\pi_0 : X_0 \rightarrow B_0$$

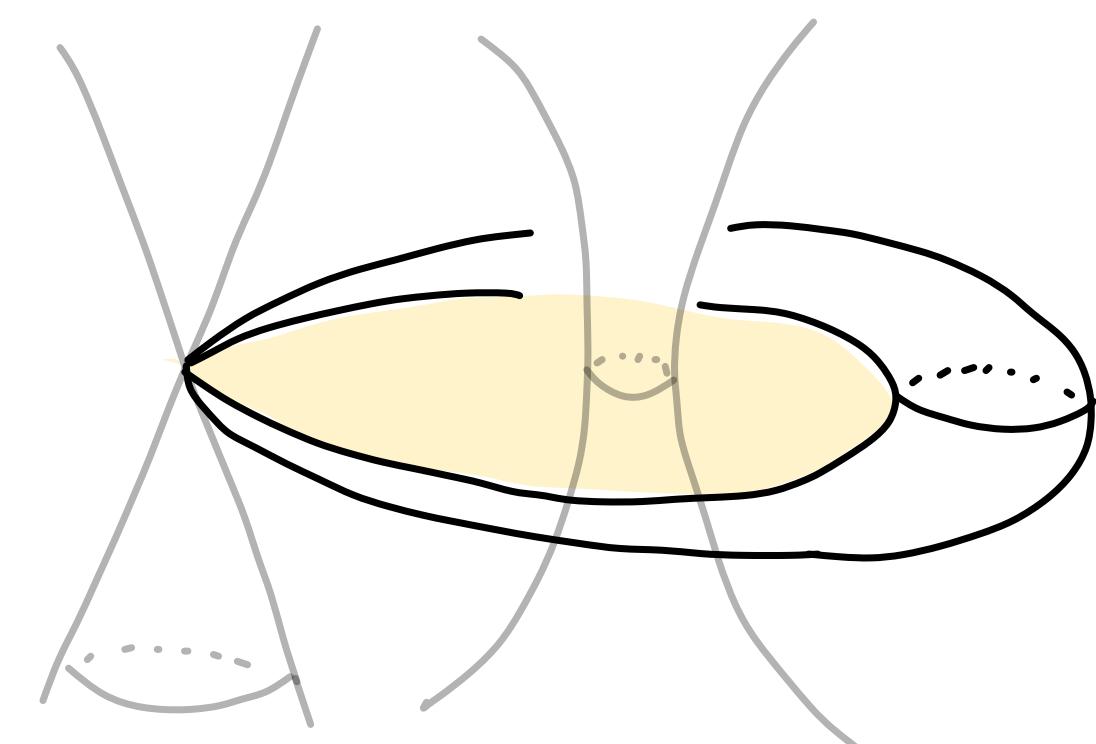
But, we must place  $\pi_0$  in  $X$  not just in  $X_0$ , because the holomorphic disks we consider are all sweeping in  $X$

**For the singular Lagrangian fibers, we study two different types of holomorphic disks**



Type (I)

- (I) Disks emanating from various singular fibers but eventually bounded by a smooth fiber
- (II) Disks bounded by a singular fiber  
Alternatively, the disks meet the singular fibers at interior vs boundary
- We must deal with them in two different ways, respectively emphasizing
  - the Floer aspect for (I) (Left figure: Done in my thesis)
  - the NA analytic / topological aspect for (II)  
(Right figure: discuss today)



Type (II)

Let's first briefly review the mirror construction in my thesis

## Theorem (Y.)

We can associate to  $(X, \pi_0)$  a triple  $(X_0^\vee, \pi_0^\vee, W_0^\vee)$  consisting of

- (a) a  $\Lambda$ -analytic space  $X_0^\vee$
- (b) an affinoid torus fibration  $\pi_0^\vee : X_0^\vee \rightarrow B_0$
- (c) a global function  $W_0^\vee$

unique up to isomorphism of analytic spaces

$\Lambda = \mathbb{C}((T^\mathbb{R}))$  - **Novikov field** - NA valuation  $v$  or norm  $|z| = e^{-v(z)}$

We also set  $\Lambda_0 = \{ |z| \leq 1 \}$ ,  $\Lambda_+ = \{ |z| < 1 \}$

$U_\Lambda = \{ |z| = 1 \}$ , Novikov unitary group, similar to  $U(1) \cong S^1$

- Set-theoretically, the mirror space is not very interesting:

$$X_0^\vee = \bigcup_{q \in B_0} H^1(L_q; U_\Lambda)$$

- Then, a main point in my thesis is that on this set, we can further give an analytic space structure by considering the Maslov-0 disks.
- This analytic topology on  $X_0^\vee$  already contains (partial) information of singularities. These disks usually meet the singular fibers.

### Affinoid torus fibration:

It is simply a continuous map with respect to analytic topology and the manifold topology on  $B_0$ , and it is locally modeled on the *tropicalization map*

$$\text{trop} : (\Lambda^*)^n \rightarrow \mathbb{R}^n \quad y_i \mapsto v(y_i)$$

I think it is first introduced by Kontsevich-Soibelman. It is further studied and is given this name by Nicaise-Xu-Yu. Here ‘continuous’ is really a strong condition, since we use the analytic topology.

- A brief picture of the mirror is as follows:
- In  $B_0$ , let  $\chi : (U, q_0) \rightarrow (V, c) \subset \mathbb{R}^n$  be a (**pointed**) integral affine chart. We allow  $q_0 \notin U$ . Then, we have an **affinoid tropical chart**

$$\bigcup_{\mathfrak{U} \in U} H^1(L_{\mathfrak{U}}; U_\Lambda) \xrightarrow{\text{set}} (\pi_0^\vee)^{-1}(U) \xrightarrow{\cong} \text{trop}^{-1}(V - c)$$

$$\downarrow \quad \quad \quad \downarrow$$

$$U \xrightarrow{\chi} V$$

- Given two such affinoid tropical charts for  $i = 1, 2$

$$\tau_i : (\pi_0^\vee)^{-1}(U) \rightarrow \text{trop}^{-1}(V_i - c_i)$$

There is a transition map (or say a gluing map)

$$\Phi : \text{trop}^{-1}(V_1 - c_1) \rightarrow \text{trop}^{-1}(V_2 - c_2)$$

between the two *analytic open domains* in  $(\Lambda^*)^n$ . It is decided by **some**  $A_\infty$  homomorphism associated to an isotopy from  $L_{q_1}$  to  $L_{q_2}$  (roughly). But,  $\Phi$  is the same for **any**  $A_\infty$  homomorphism obtained in this way. This is carefully proved in my thesis.

# Theorem (Y.)

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unique up to isomorphism of analytic spaces

- We aim to develop an analytic extension  $\pi^\vee$  over  $B$  (rather than just  $B_0$ )
- All possible (analytic) continuous extension may be too much.
- We add extra conditions to control the extension.
- Based on our computations for the Gross's special Lagrangian fibration, we propose to use:

## tropically continuous maps:

in the sense of Chambert-Loir and Ducros. See Section (3.1.6) of their famous paper in which they develop  $(p,q)$ -forms on analytic space:

### **Formes différentielles réelles et courants sur les espaces de Berkovich**

- ▶ Under this condition, the topological extension from  $B_0$  to  $B$  can somehow control the analytic extension from  $\pi_0^\vee$  to some potential extension.
- ▶ I'm inspired by Gross's Topological Mirror Symmetry to think like this. Also, I'm inspired by Kontsevich-Soibelman's singular model in  
(‘Affine structures and non-archimedean analytic spaces’, Section 8)

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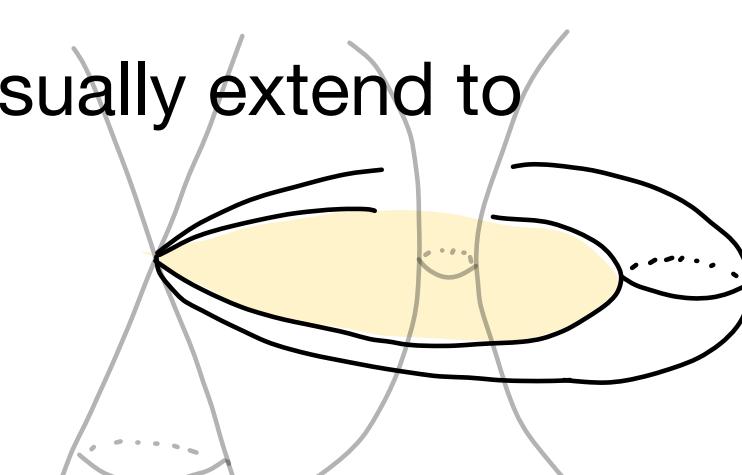
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- Roughly speaking, a tropically continuous map  $F$  is *locally* in the following form:

$$F|_{\mathcal{U}} = \varphi(v(f_1), \dots, v(f_n))$$

- $\mathcal{U}$  is an analytic open subset
- $f_1, \dots, f_n : \mathcal{U} \rightarrow \Lambda^*$  are invertible analytic functions.  
(e.g. local coordinates for analytification of an algebraic variety)
- $\varphi : \mathbb{R}^n \rightarrow \mathbb{R}^m$  is just a continuous map for Euclidean topology
- Let's give a naive idea, we will be more specific soon.
- Very intuitively, imagine  $v(f_1), \dots, v(f_n)$  are ‘action coordinates’, so more or less correspond to symplectic areas.
- The symplectic areas, as functions on  $B_0$ , can usually extend to the singular locus. See the right side figure:
- Anyway, we will just focus on an example later.



## Why the conventional Maurer-Cartan idea is not enough

It's been long expected that the mirror space should be the union of the Maurer-Cartan sets (with singular fibers)

$$\bigcup_{q \in B} \mathcal{MC}(L_q)$$

Over  $B_0$ , this roughly gives the correct picture. Indeed, the  $\mathcal{MC}(L_q)$  is very close to  $H^1(L_q; U_\Lambda)$  (slightly different)

In short, the **set-theoretic cocycle condition** is essentially straightforward by the homotopy invariance of Maurer-Cartan sets.

(a property well-known for the classic homotopy theory of  $A_\infty$  structures).

Over  $B \setminus B_0$ , the Maurer-Cartan picture fails unfortunately !

### Maurer-Cartan set of a singular Lagrangian

#### $\subsetneq$ ‘Dual singular fiber’

Roughly, the ‘dual singular fibers’ here will extend  $\pi_0^\vee$  *tropically continuously*. We will see this more clearly soon.

*Admittedly, the Maurer-Cartan picture may offer some good ideas, or inspirations, etc.*

*But unfortunately, it is eventually not the correct picture.*

Nevertheless, for the **analytic cocycle condition**, we must introduce more powerful ideas and tools as in my thesis: (at least 4 points below)

1. Study a uniform version of *Groman-Solomon’s reverse isoperimetric inequalities* for the non-archimedean convergence.
2. Establish a minimal model version of Fukaya’s trick. With nontrivial Maslov-0 disks, this creates further difficulties. Why minimal model? Very roughly, want  $H^1(L_q; U_\Lambda)$  rather than  $\Omega^1(L_q; U_\Lambda)$ .
3. Prove the transition maps are well-defined. Otherwise, what the cocycle conditions mean is very ambiguous. This requires the following ud-homotopy.
4. Upgrade the classic homotopy to the **ud-homotopy** for the  $A_\infty$  structures. The point is, the analytic gluing needs stronger homotopy.
  - For example, individual  $A_\infty$  maps satisfy the divisor axiom are not enough. We want the homotopies between them also satisfy the divisor axiom in a very specific sense.
  - This requires lots of difficult homological algebras, and finally we need to introduce the so-called category  $\mathcal{UD}$  in my thesis.

In a word, the ud-homotopy theory enables us to **upgrade the ‘classic Maurer-Cartan idea’ to a higher and more precise level**, matching NA adic-convergent formal power series rather than just set bijections.

This is a totally different story, and is crucial for the analytic topology.

- ❖ By the way, the inclusion of Landau-Ginzburg superpotential will be also indispensable for our results later.

## Theorem: $Y$ is SYZ mirror to $X$

Let's go to a fundamental example which has been long predicted by Gross-Siebert program. Define

$$X = \mathbb{C}^n \setminus \{z_1 \cdots z_n = 1\}$$

(equipped with the standard symplectic form), and define

$$Y = \{(x, y) \in \Lambda^2 \times (\Lambda^*)^{n-1} \mid x_0 x_1 = 1 + y_1 + \cdots + y_{n-1}\}$$

**Definition:** We say an algebraic variety  $Y$  over  $\Lambda$  is **SYZ mirror** to a complex manifold  $X$  over  $\mathbb{C}$  if

- there exists a proper tropically continuous analytic fibration  $f: \mathcal{Y} \rightarrow B$  on a Zariski-dense analytic open domain  $\mathcal{Y}$  in  $Y$
- there exists a Lagrangian fibration  $\pi: X \rightarrow B$  onto the same base manifold  $B$  for some Kähler form  $\omega$  on  $X$

such that the following conditions hold

- 1) The  $\pi$  (resp.  $f$ ) restricts to a Lagrangian torus fibration  $\pi_0$  (resp. an affinoid torus fibration  $f_0$ ) over a common open subset  $B_0 \subset B$  such that the **two** induced integral affine structures agree with each other and  $\Delta = B \setminus B_0$  is codimesion-2.

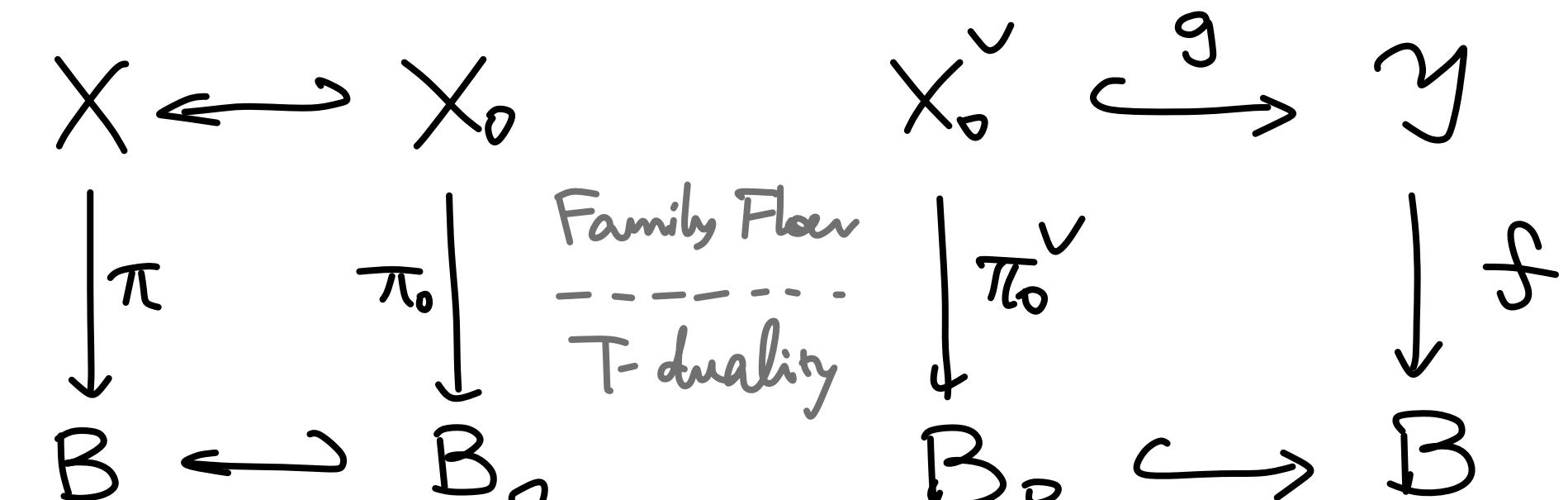
(same smooth/singular locus, integral affine str. It's already very nontrivial)

- 2) There is an isomorphism of affinoid torus fibration  $\pi_0^\vee \cong f_0$ . Here  $\pi_0^\vee$  is the family Floer mirror fibration for  $\pi_0$  (a sort of T-duality)
- 3) The set  $\mathcal{Y}_0 := f_0^{-1}(B_0)$  is Zariski dense in  $Y$   
(possibly redundant, but useful for a folklore conjecture later)

- Kontsevich-Soibelman proved that any affinoid torus fibration also induces an integral affine structure on the base.
- The existence of affinoid torus fibration, parallel to that of Lagrangian fibration, should be also a nontrivial problem in NA geometry.
- In general,  $\mathcal{Y}$  depends on  $\omega$ , and vice versa. This gives a picture of “Kähler moduli v.s. complex moduli”:  $(Y, \mathcal{Y})$  is SYZ mirror to  $(X, \omega)$ .
- For example, we can simply run the family Floer T-duality for the toric moment map. The analytic domain is  $\mathcal{Y} = \mathcal{Y}_0 \cong \text{trop}^{-1}(P)$  in  $(\Lambda^*)^n$  for the moment polytope  $P = P_\omega$  relying on  $\omega$ . But, the Zariski-closure is the same algebraic variety  $Y = (\Lambda^*)^n$  not relying on  $\omega$ .
- We focus only on SYZ now. Hopefully, we could achieve Abouzaid's family Floer functor to prove HMS in the future.

**Remark** if you allow me to remove the condition 2), still nontrivial to get 1), then

- Nothing about Floer/Fukaya theories
- Nothing about the moduli spaces of pseudo-holomorphic disks.
- It may be very unmotivated without some Floer-theoretic considerations.
- But, as we will see, the construction of  $f$  on  $Y$  itself is very elementary



**Theorem:**  $Y = \{x_0x_1 = 1 + y_1 + \cdots + y_{n-1}\}$  is SYZ mirror to  $X = \mathbb{C}^n \setminus \{z_1 \cdots z_n = 1\}$

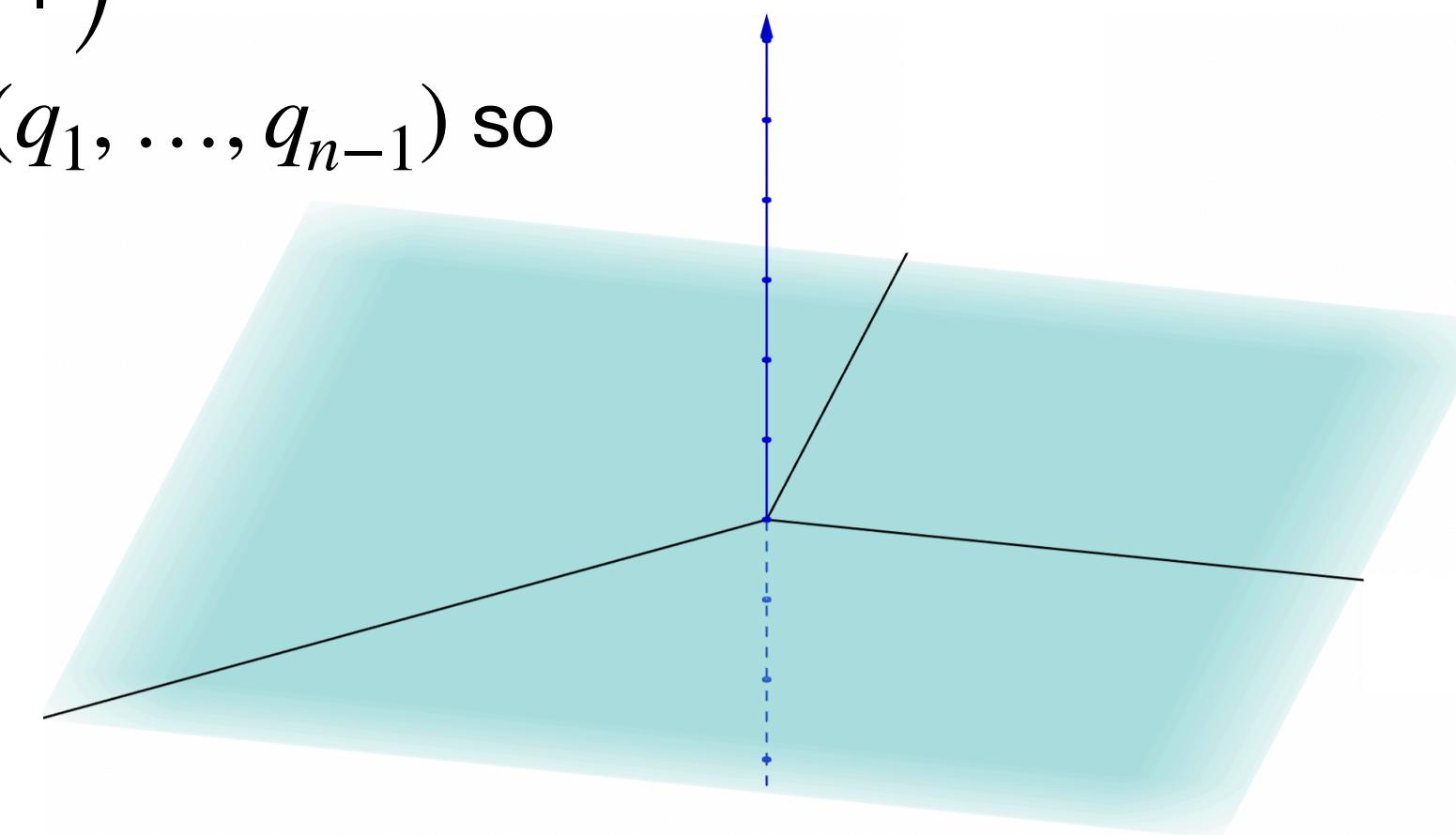
- On the A-side, we consider a Gross's special Lagrangian fibration.

$$\pi : X \rightarrow B \cong \mathbb{R}^n \quad z \mapsto \left( |z_1|^2 - |z_n|^2, \dots, |z_{n-1}|^2 - |z_n|^2, \log |z_1 \cdots z_n - 1| \right)$$

- The singular locus  $\Delta = B \setminus B_0$  is a tropical hypersurface in  $\mathbb{R}^{n-1} \times \{0\}$  given by those  $\bar{q} = (q_1, \dots, q_{n-1})$  so that  $\min\{0, q_1, \dots, q_{n-1}\}$  attains at least twice. (See the right figure when  $n = 3$ )
- Let  $B_0 := B \setminus \Delta$ . By the family Floer theory, there is an ‘abstract’ affinoid torus fibration

$$\pi_0^\vee : X_0^\vee \rightarrow B_0$$

on the ‘abstract’ set  $X_0^\vee \equiv \bigcup_{q \in B_0} H^1(L_q; U_\Lambda)$

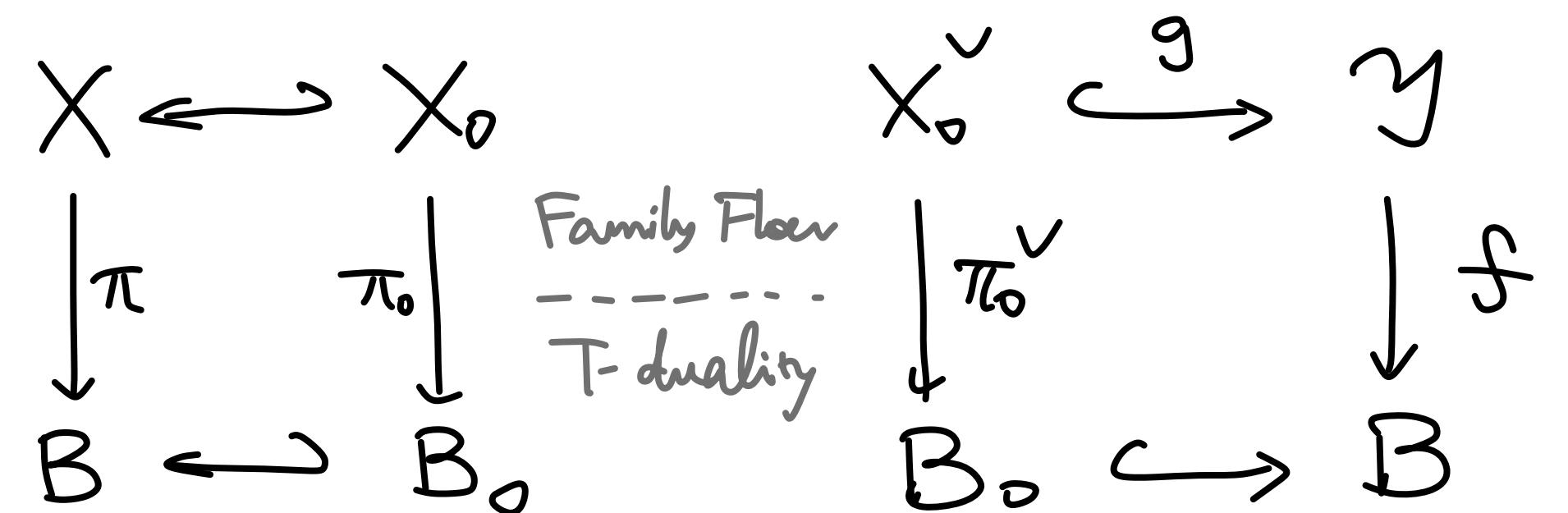


- Our Zariski-dense analytic open domain  $\mathcal{Y}$  in  $Y = \{x_0x_1 = 1 + y_1 + \cdots + y_{n-1}\}$  is defined by setting the NA norm  $|x_1| < 1$ . This is very explicit. We will explain why we choose it later.

(It is also fine to consider  $|x_1| < R$  for other  $R \neq 1$  or replace  $x_1$  by  $x_0$ . Just a convention)

Now, our major task is to find both the constructions of (see the diagram)

- (i) the **analytic embedding**  $g : X_0^\vee \rightarrow \mathcal{Y}_0$  from the abstract to the concrete
- (ii) the **dual singular fibration**  $f : \mathcal{Y} \rightarrow B$  so that  $f_0 \cong \pi_0^\vee$  via the above  $g$



**(i) Construction of  $g$ :** First, study the wall-crossing of  $\pi$  (use the family Floer theory), and we can finally show a simple **identification**:

(\*)

$$X_0^\vee \cong T_+ \cup T_- / \sim$$

$T_\pm$  are analytic open subdomains  $\subsetneq (\Lambda^*)^n$  (wrt the action coordinates for  $\omega$ )

$T_\pm$  correspond to the Clifford and Chekanov tori respectively.

the gluing relation  $\sim$  can be written down explicitly.

Under the identification (\*), the analytic embedding  $g$  is obtained by gluing

$$g_+ : T_+ \rightarrow Y \quad (y_1, \dots, y_n) \mapsto \left( \frac{1}{y_n}, y_n h, y_1, \dots, y_{n-1} \right)$$

$$g_- : T_- \rightarrow Y \quad (y_1, \dots, y_n) \mapsto \left( \frac{h}{y_n}, y_n, y_1, \dots, y_{n-1} \right)$$

where  $h = 1 + y_1 + \dots + y_{n-1}$

❖ The formula of  $g$  is due to GHK and GS, but we further add NA picture (KS).

- Gross-Hacking-Keel ( see Lemma 3.1 in ‘Birational Geometry of Cluster Algebras’ )
- Kontsevich-Soibelman ( see Page 44 in ‘Affine structures and non-archimedean analytic spaces’ )

Because GHK is over  $\mathbb{C}$ , we discuss some ways of reduction from  $\Lambda$  to  $\mathbb{C}$  :

- If an analytic space over  $\Lambda$  is the generic fiber of a formal scheme over  $\Lambda_0$ , then the special fiber is the so-called analytic reduction which is a variety over  $\mathbb{C}$ . (not unique)
- $(\Lambda^*)^n$  ‘contains’ infinity copies of  $(\mathbb{C}^*)^n$  c.f. *exploded tropicalization map* (Sam Payne).
- Study the *Maslov’s dequantization*. (Mikhalkin, Abouzaid-Ganatra-Iritani-Sheridan)

Anyway, the Novikov field is good for the **T-duality idea** :

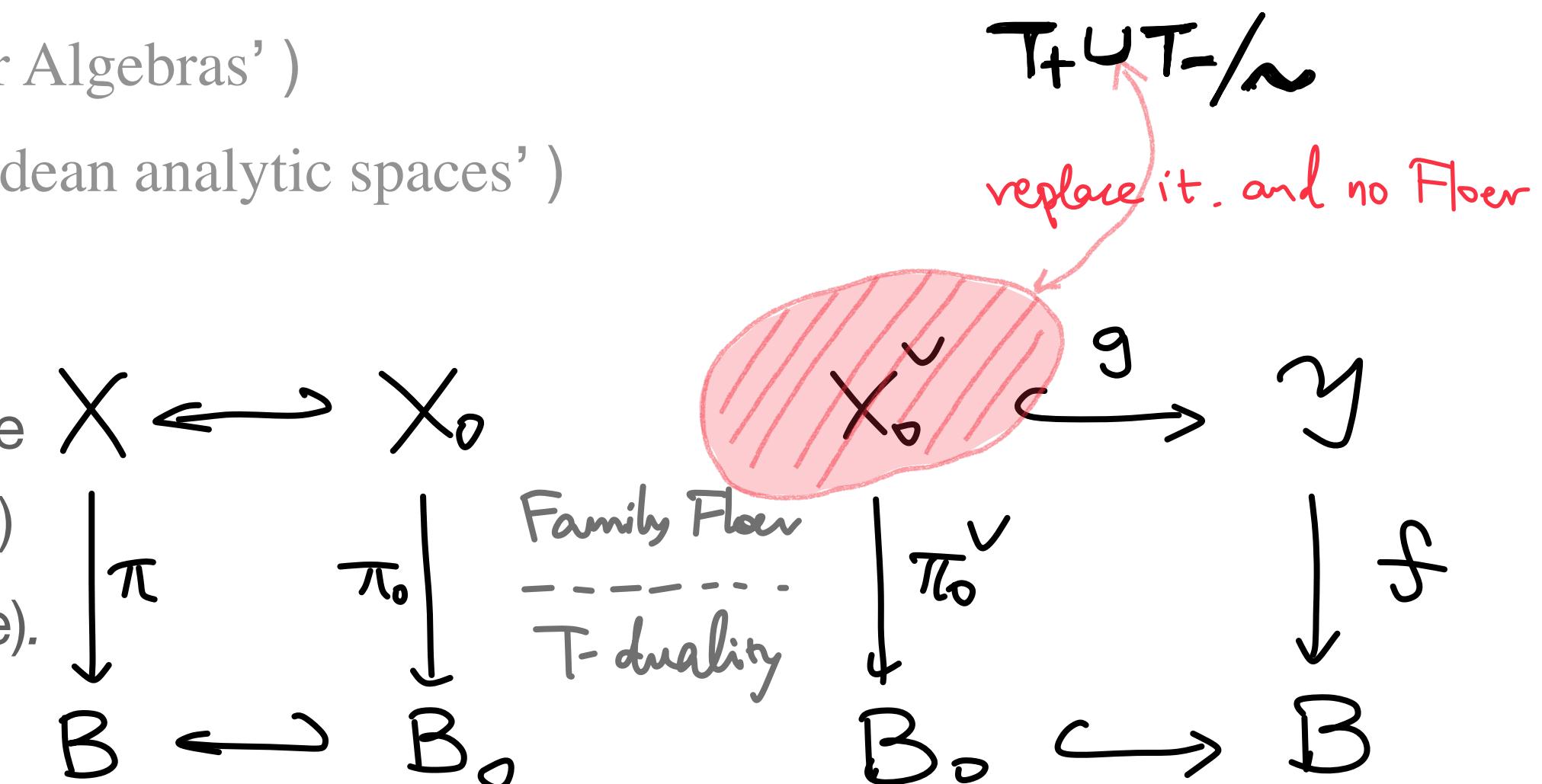
More intrinsically,  $T_\pm \subset \bigcup_{q \in B_0} H^1(L_q; U_\Lambda)$  as sets, and

$$(\Lambda^*)^n \ni (y_1, \dots, y_n) \leftrightarrow (T^{\chi_1} \nabla(\sigma_1), \dots, T^{\chi_n} \nabla(\sigma_n)) \leftrightarrow \nabla$$

where  $\nabla \in H^1(L_q; U_\Lambda)$  and  $\chi = (\chi_1, \dots, \chi_n)$  is the action coordinates for a basis  $\{\sigma_1, \dots, \sigma_n\} \subset \pi_1(L_q)$ .

**Remark:** The only place we use Floer theory is the identification (\*)

- In fact, all of  $T_\pm$ , the gluing relation  $\sim$ , the embedding  $g$  have nothing to do with the Floer theory or the moduli space business.
- If we’re content with the main theorem without the T-duality condition  $\pi_0^\vee \cong f_0$ , then we can entirely exclude the Floer theory.



**(ii) Construction of  $f$ :** This is very difficult ! We obtain  $f = j^{-1} \circ F$  by decomposing it into a topological embedding  $j : B \equiv \mathbb{R}^n \rightarrow \mathbb{R}^{n+1}$  and an analytic tropically continuous map  $F : Y \rightarrow \mathbb{R}^{n+1}$ . (I will give a very beautiful picture for the image of  $j$  )

This decomposition imitates Kontsevich-Soibelman's model, and I would like to thank Tony Yue Yu for his suggestion to KS's paper.

Anyway, after a lot of trials, we find the following pair  $(j, F)$  by hand, and it finally works !

$$j(q) = (\theta_0(q), \theta_1(q), \bar{q})$$

- ❖  $q = (q_1, \dots, q_{n-1}, q_n) = (\bar{q}, q_n)$  are in  $B \equiv \mathbb{R}^n$
- ❖  $\theta_0(q) = \min\{-\psi(q), -\psi(\bar{q}, 0)\} + \min\{0, \bar{q}\}$
- ❖  $\theta_1(q) = \min\{\psi(q), \psi(\bar{q}, 0)\}$
- ❖  $\psi : B \rightarrow \mathbb{R}_+$  is the  $\omega$ -areas of holomorphic disks

$$F = (F_0, F_1, G_1, \dots, G_{n-1})$$

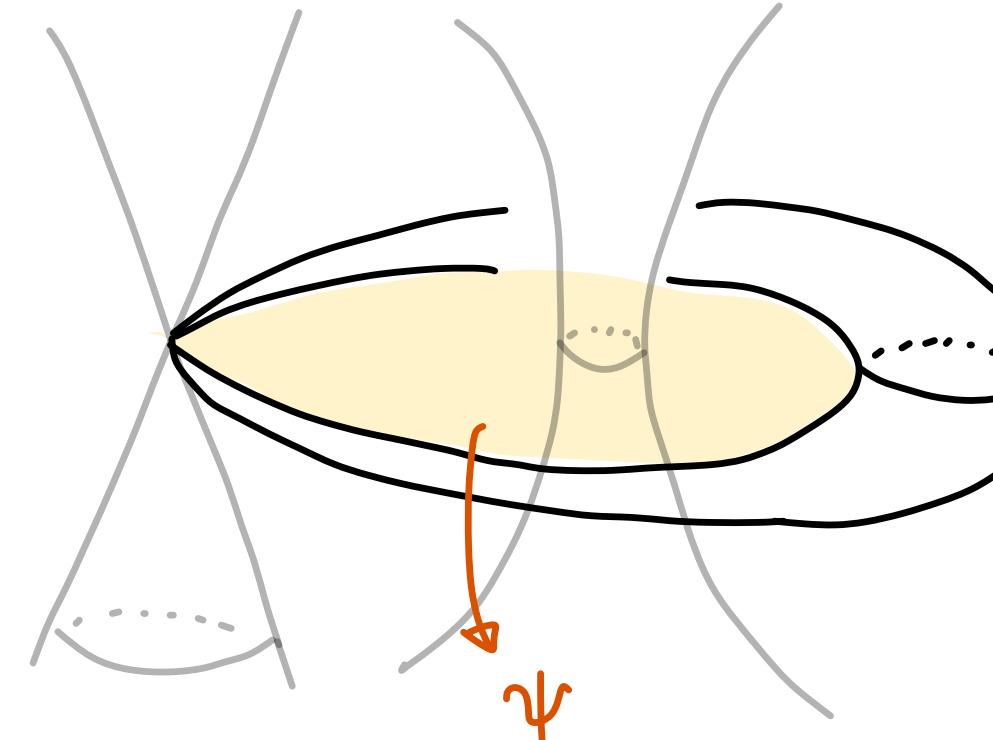
Given  $z = (x_0, x_1, y_1, \dots, y_{n-1})$  in  $Y$ , we define

- ❖  $F_0(z) = \min \left\{ v(x_0), -\psi(v(y_1), \dots, v(y_{n-1}), 0) + \min\{0, v(y_1), \dots, v(y_{n-1})\} \right\}$
- ❖  $F_1(z) = \min \left\{ v(x_1), \psi(v(y_1), \dots, v(y_{n-1}), 0) \right\}$
- ❖  $G_k(z) = v(y_k) \quad 1 \leq k \leq n-1$

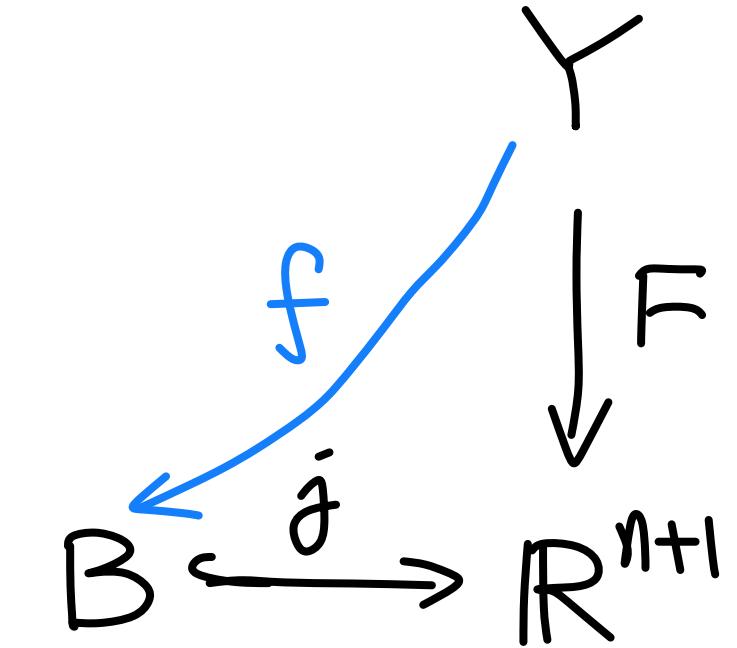
where  $v$  is the non-archimedean valuation and the  $\psi$  is the same as above.

This depends on the Kähler form in general.

This includes ‘singular’ analytic fibers in the sense of Kontsevich-Soibelman.



$\psi(q)$  for  $q \in \Delta \equiv B \setminus B_0$

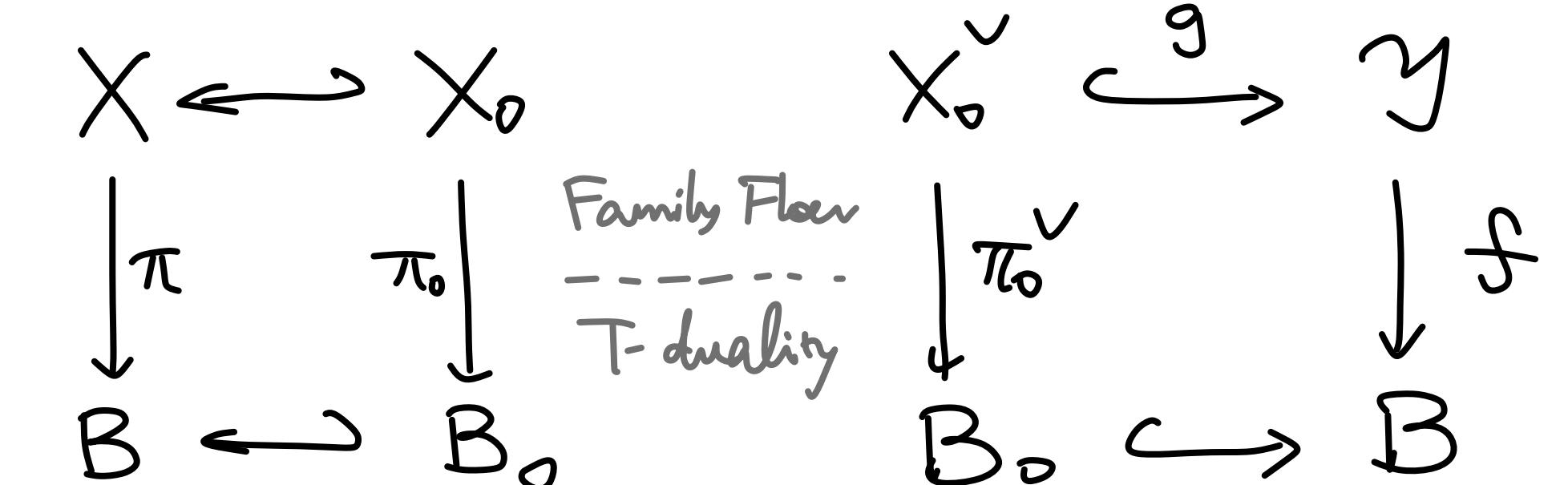


The Zariski-dense analytic subdomain  $\mathcal{Y}$  exactly satisfies that  $j(B) = F(\mathcal{Y})$ . Define

$$f = j^{-1} \circ F|_{\mathcal{Y}}$$

**It is explicit, elementary, and has singular fibers**

Although not obvious now, it will meet all the conditions for our definition of ‘SYZ mirror’ (e.g. match singular locus, integral affine str)  $\square$



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- ▶ A prototype is found by KS many years ago in a very different context.
- ▶ KS use 3 charts, but we only use 2 charts  $T_\pm$  which geometrically correspond to Clifford/Chekanov tori. We simplify it to 2 charts, exactly inspired by GHK's work.
- ▶ **As we see, the formula of  $f$  is very elementary itself (maybe still complicated)**
- ▶ In fact, if we are content with the result without the T-duality condition  $\pi_0^\vee \cong f_0$  there is even no need to know anything about Floer theory

$$F = (F_0, F_1, G_1, \dots, G_{n-1})$$

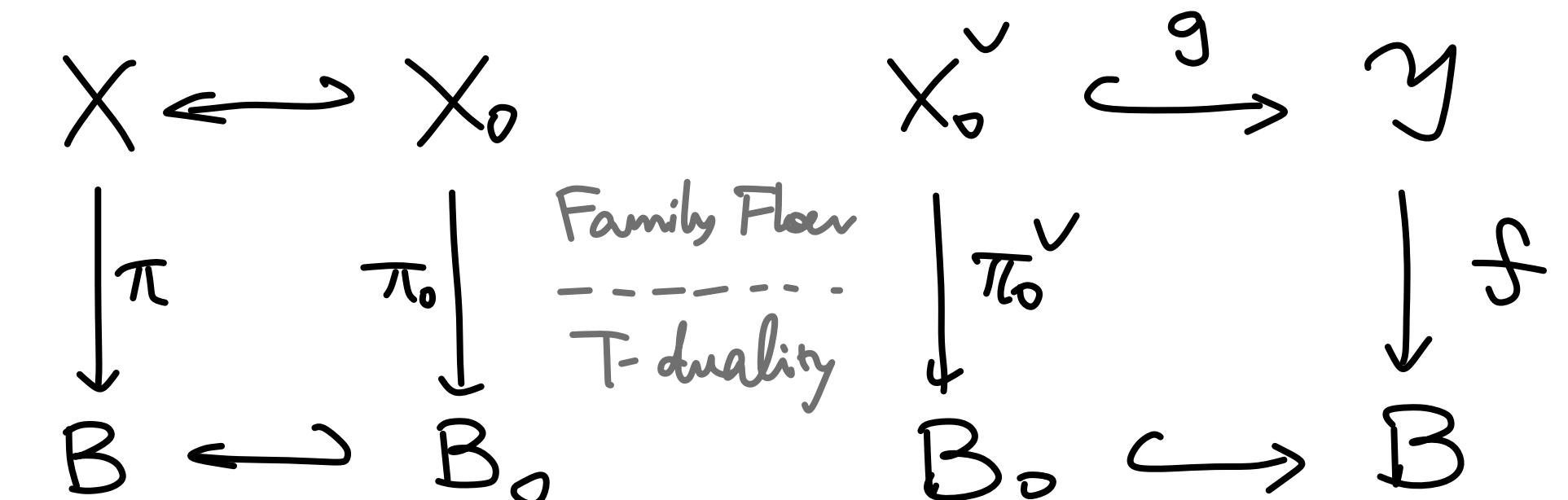
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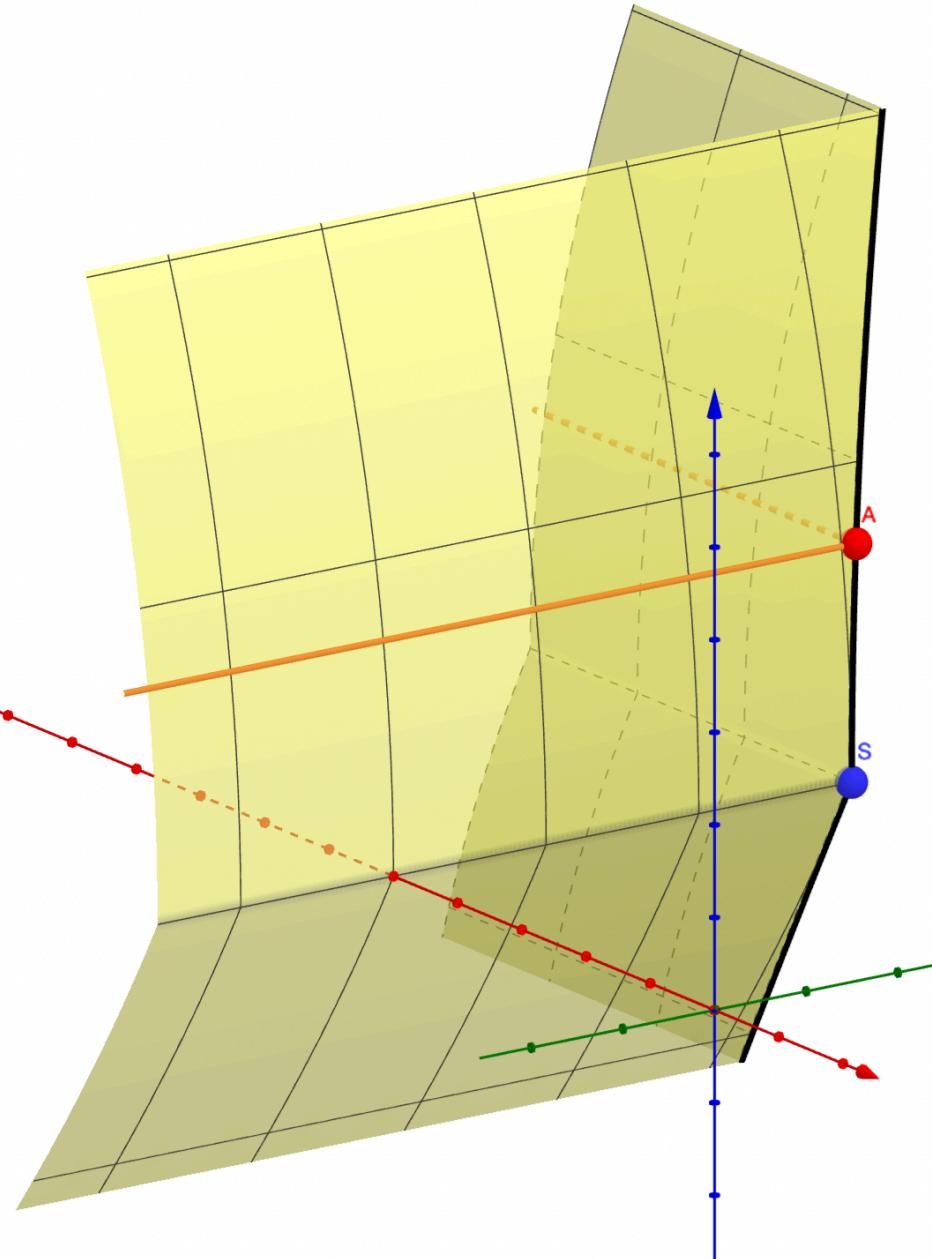
This includes 'singular' analytic fibers in the sense of Kontsevich-Soibelman.



# Visualization for $j : \mathbb{R}^n \rightarrow \mathbb{R}^{n+1}$

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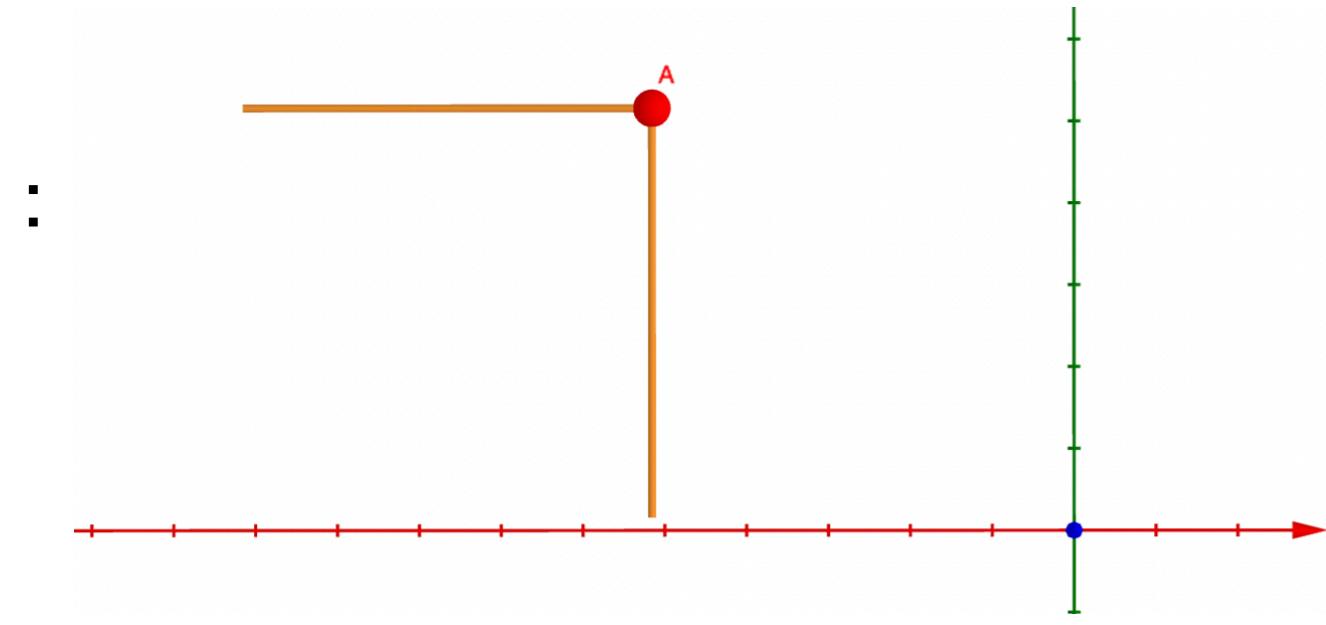
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- ❖  $\theta_1(q) = \min\{\psi(q), \psi(\bar{q}, 0)\}$
- ❖  $\psi : B \rightarrow \mathbb{R}_+$  is the  $\omega$ -area of a holomorphic disk



In  $\mathbb{R}^2$ , consider a broken line  $R_{\bar{q}}$  (in orange) with a corner point :

$$A = (a_0(\bar{q}), a_1(\bar{q})) := (\min\{0, \bar{q}\} - \psi(\bar{q}, 0), \psi(\bar{q}, 0)).$$

Get a family of broken lines  $R_{\bar{q}}$  parametrized by  $A(\bar{q})$  or  $\bar{q}$ .



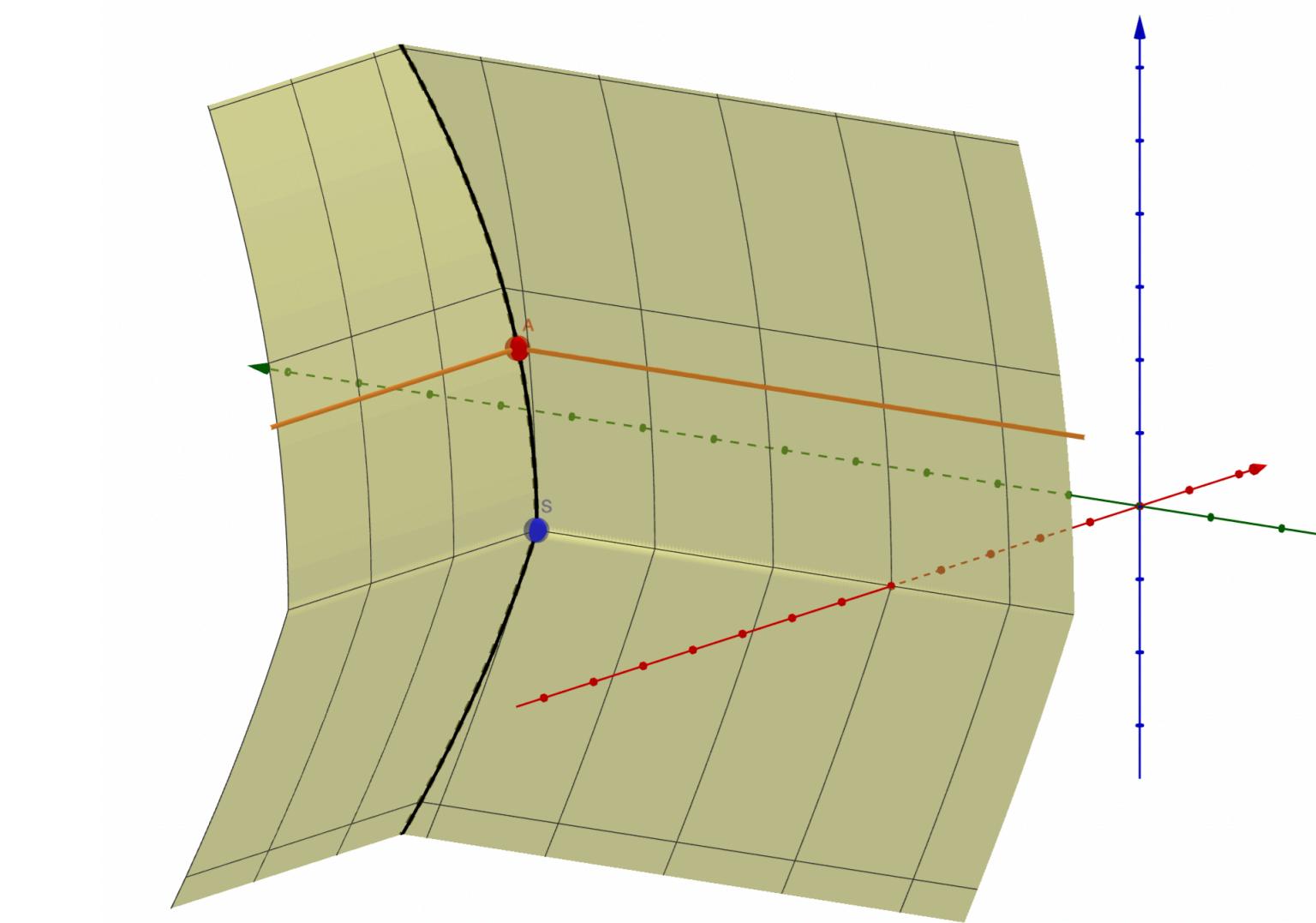
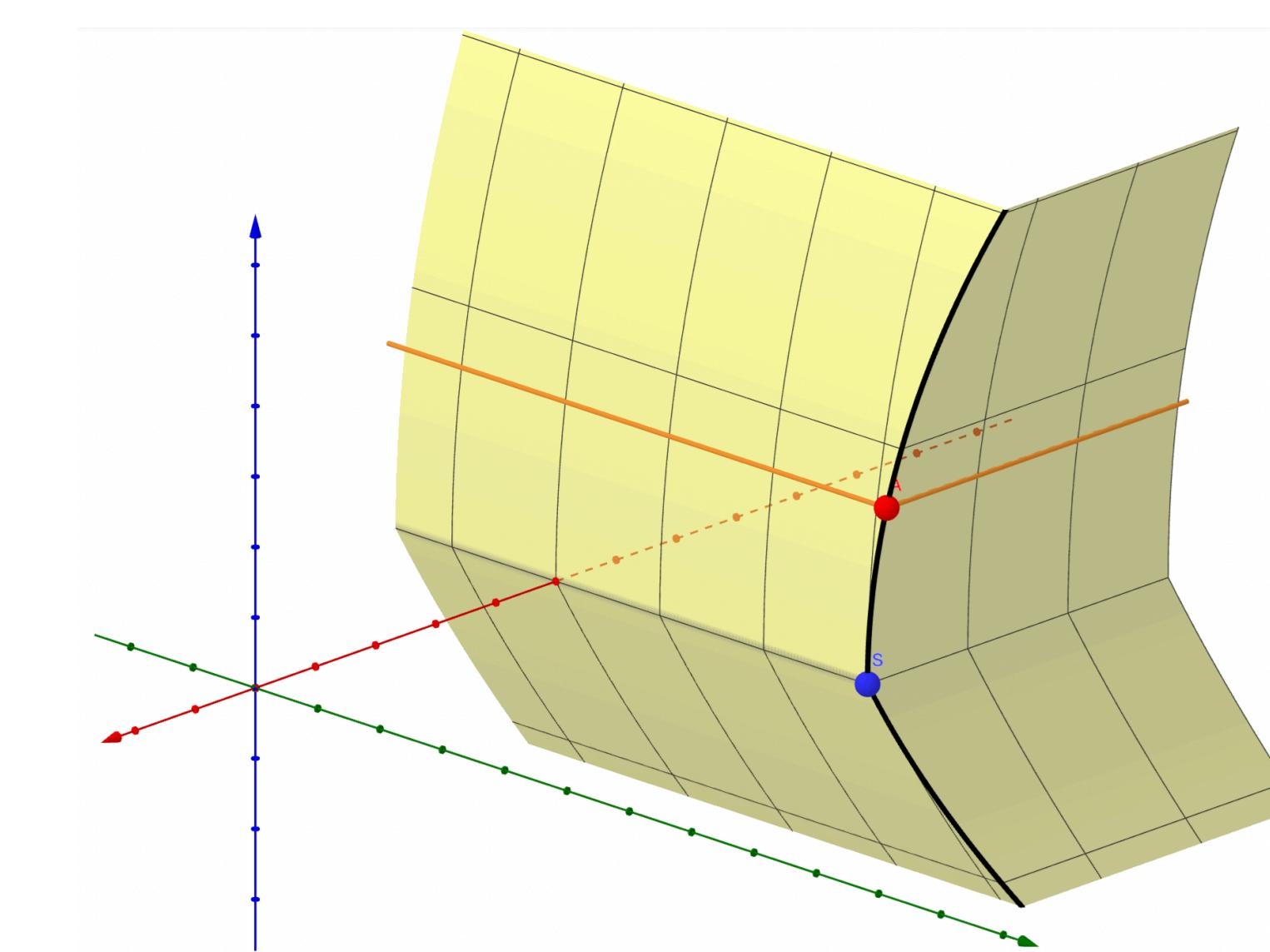
Now, the image  $j(B)$  is simply the union of these broken lines:

$$j(B) = \bigcup_{\bar{q} \in \mathbb{R}^{n-1}} R_{\bar{q}} \times \{\bar{q}\}$$

The black curve is the trace of the point  $A(\bar{q})$  and relies on  $\omega$  in general. Also,

$$j(\Delta) = \bigcup_{\bar{q} \in \Pi} \{(A(\bar{q}), \bar{q})\}$$

where  $\Pi = V(\min\{0, \bar{q}\})$ . If  $n = 2$ , only one singular point, the blue point below.



# Visualization for $j : \mathbb{R}^n \rightarrow \mathbb{R}^{n+1}$

$$j(q) = (\theta_0(q), \theta_1(q), \bar{q})$$

- ❖  $q = (q_1, \dots, q_{n-1}, q_n) = (\bar{q}, q_n)$  are in  $B \equiv \mathbb{R}^n$
- ❖  $\theta_0(q) = \min\{-\psi(q), -\psi(\bar{q}, 0)\} + \min\{0, \bar{q}\}$
- ❖  $\theta_1(q) = \min\{\psi(q), \psi(\bar{q}, 0)\}$
- ❖  $\psi : B \rightarrow \mathbb{R}_+$  is the  $\omega$ -area of a holomorphic disk

Maybe very surprisingly, the  $F = (F_0, F_1, G_1, \dots, G_{n-1})$  has  
**exactly the same image in  $\mathbb{R}^{n+1}$ ,** i.e.  $j(B) = F(\mathcal{Y})$

Given  $z = (x_0, x_1, y_1, \dots, y_{n-1})$  in  $Y$ , we define

- ❖  $F_0(z) = \min \left\{ v(x_0), -\psi(v(y_1), \dots, v(y_{n-1}), 0) + \min\{0, v(y_1), \dots, v(y_{n-1})\} \right\}$
- ❖  $F_1(z) = \min \left\{ v(x_1), \psi(v(y_1), \dots, v(y_{n-1}), 0) \right\}$
- ❖  $G_k(z) = v(y_k) \quad 1 \leq k \leq n-1$

where  $v$  is the non-archimedean valuation and the  $\psi$  is the same as above.

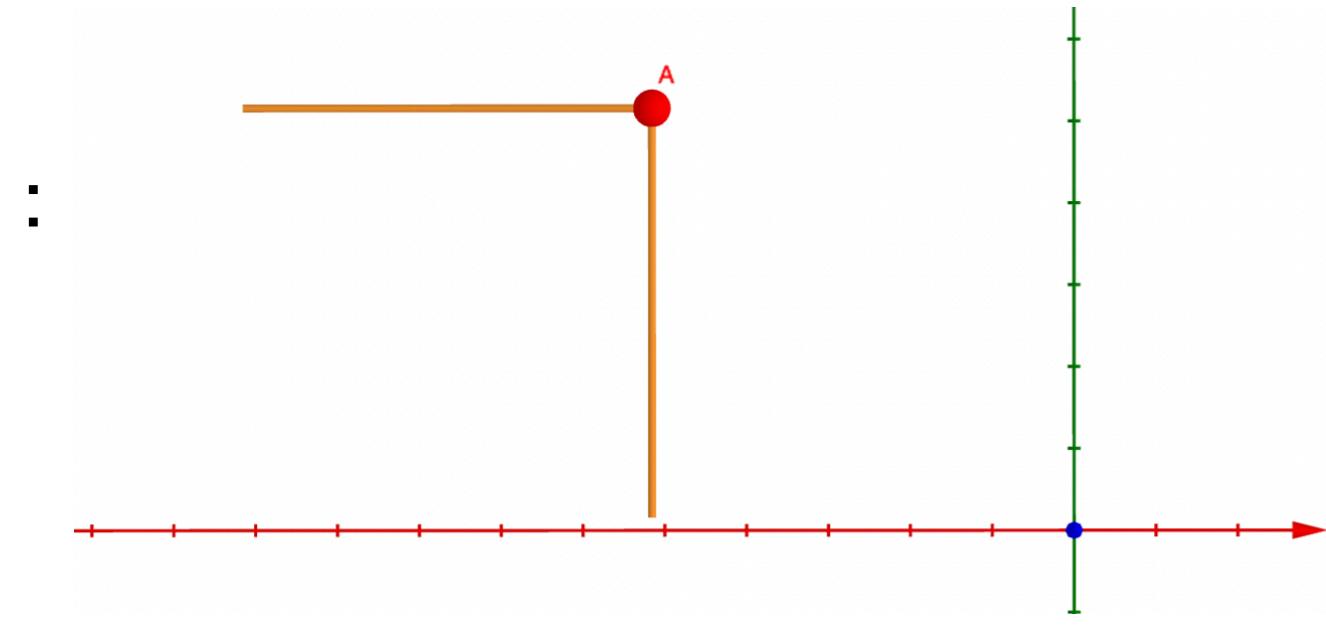
This depends on the Kähler form in general.

This includes ‘singular’ analytic fibers in the sense of Kontsevich-Soibelman.

In  $\mathbb{R}^2$ , consider a broken line  $R_{\bar{q}}$  (in orange) with a corner point :

$$A = (a_0(\bar{q}), a_1(\bar{q})) := (\min\{0, \bar{q}\} - \psi(\bar{q}, 0), \psi(\bar{q}, 0)).$$

Get a family of broken lines  $R_{\bar{q}}$  parametrized by  $A(\bar{q})$  or  $\bar{q}$ .



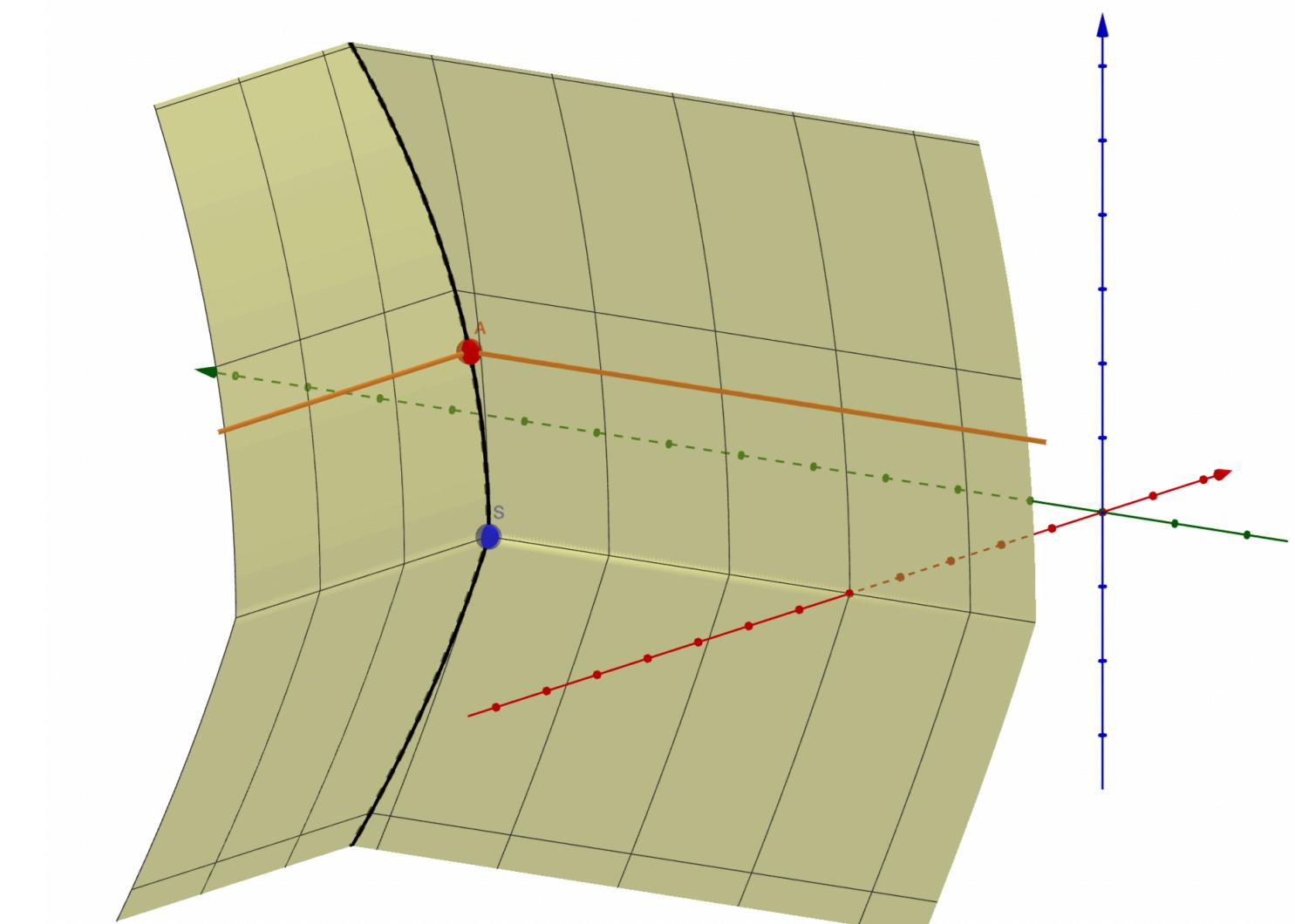
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- ❖ The motivation to find  $(j, F)$  is difficult for me to explain. It is really found by hand.
- ❖ But, let me try to explain the idea: the key observation is

$$v(1 + y_1 + \dots + y_{n-1}) = \min\{0, v(y_1), \dots, v(y_{n-1})\}$$

whenever the minimum is attained only once.

But if the minimum is attained twice, it may happen that

$$v(1 + y_1 + \dots + y_{n-1}) > \min\{0, v(y_1), \dots, v(y_{n-1})\}$$

which gives some ambiguity.

- ❖ The pair  $(j, F)$  is very carefully designed to ‘eliminate’ this ambiguity.
- ❖ This ambiguity also plays the crucial role to ‘create’ the singularity of the dual fibration  $f$ .
- ❖ Regardless of how we find the  $(j, F)$  from the family Floer picture, it is really elementary to define them and check a key lemma as follows :

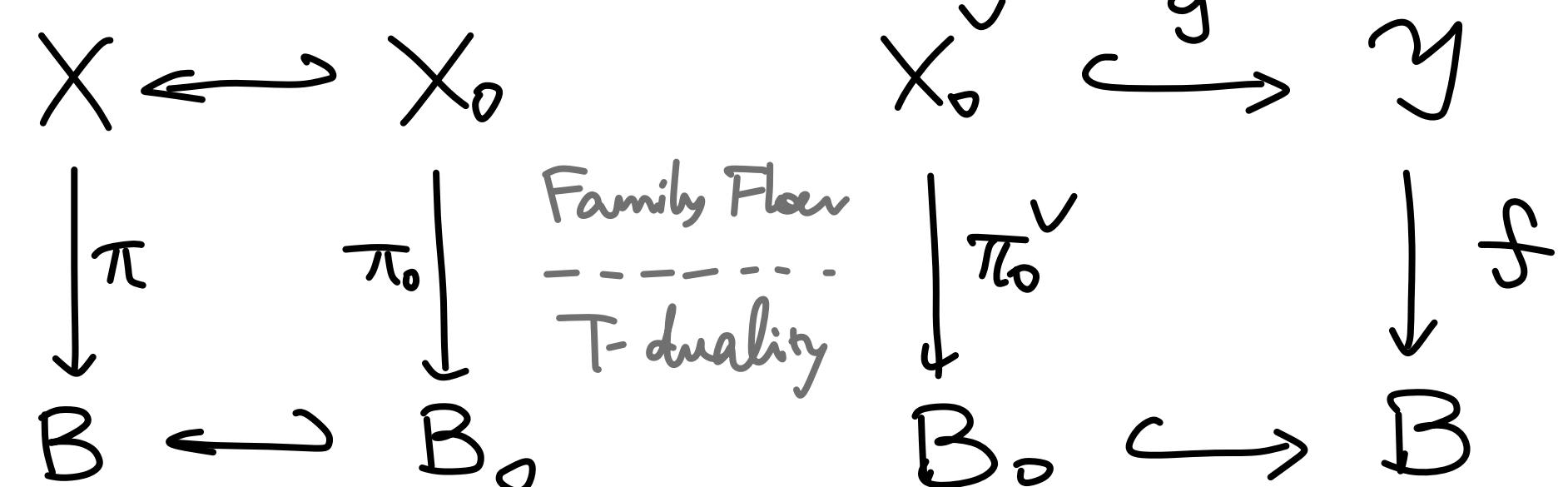
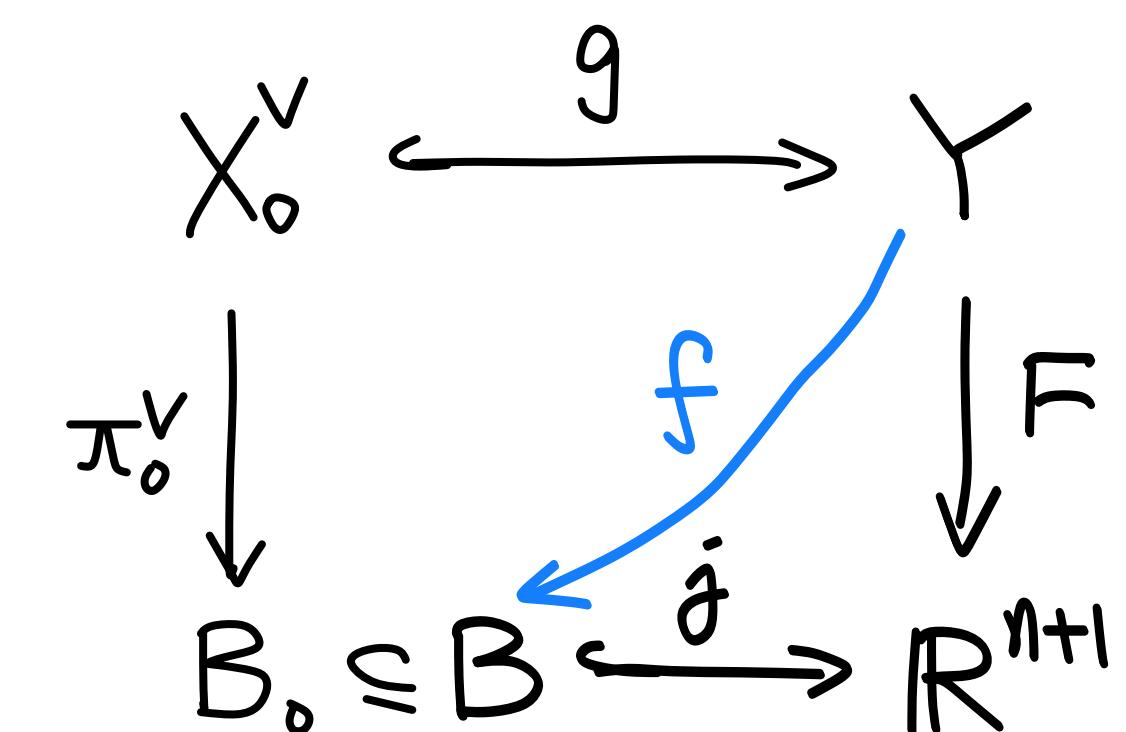
$$j \circ \pi_0^\vee = F \circ g$$

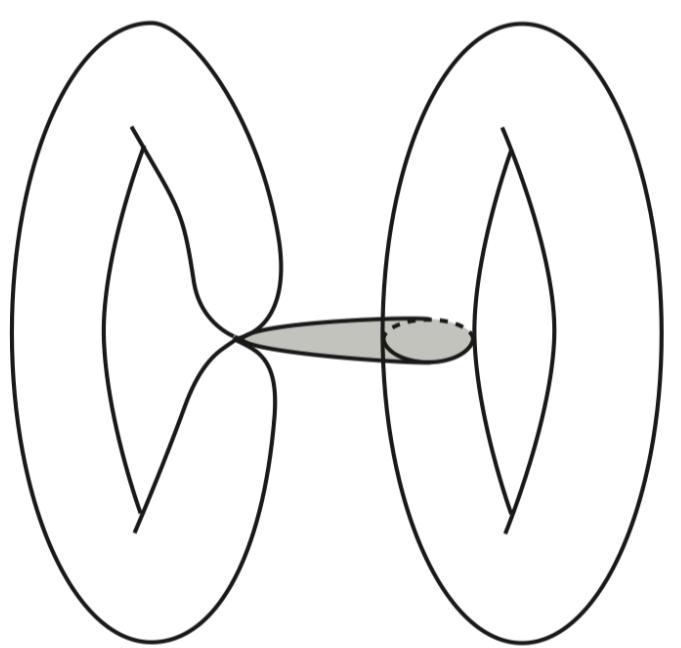
This implies the T-duality condition  $\pi_0^\vee \cong f_0$

- The only place we need family Floer is the identification (not easy to get)

$$(*) \quad X_0^\vee \cong T_+ \cup T_- / \sim$$

- By (\*), we can also directly define the affinoid torus fibration  $\pi_0^\vee$  on the domains  $T_\pm$  regardless of any Floer-theoretic considerations.
- Except (\*), all of  $T_\pm, \sim, j, F, Y, f, g$  can be defined **directly** in the pure NA world.





follows Gross-Hacking-Keel-Siebert's principle based on the identification

agrees with many previous results like Auroux, Abouzaid-Auroux-Katzarkov, Abouzaid-Sylvan, Gammage, Gross-Siebert, etc.

**Type (I)**

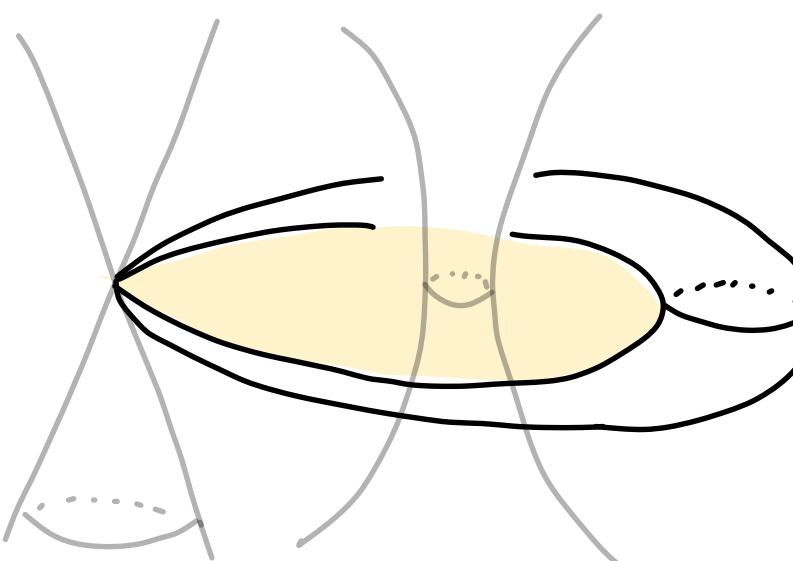
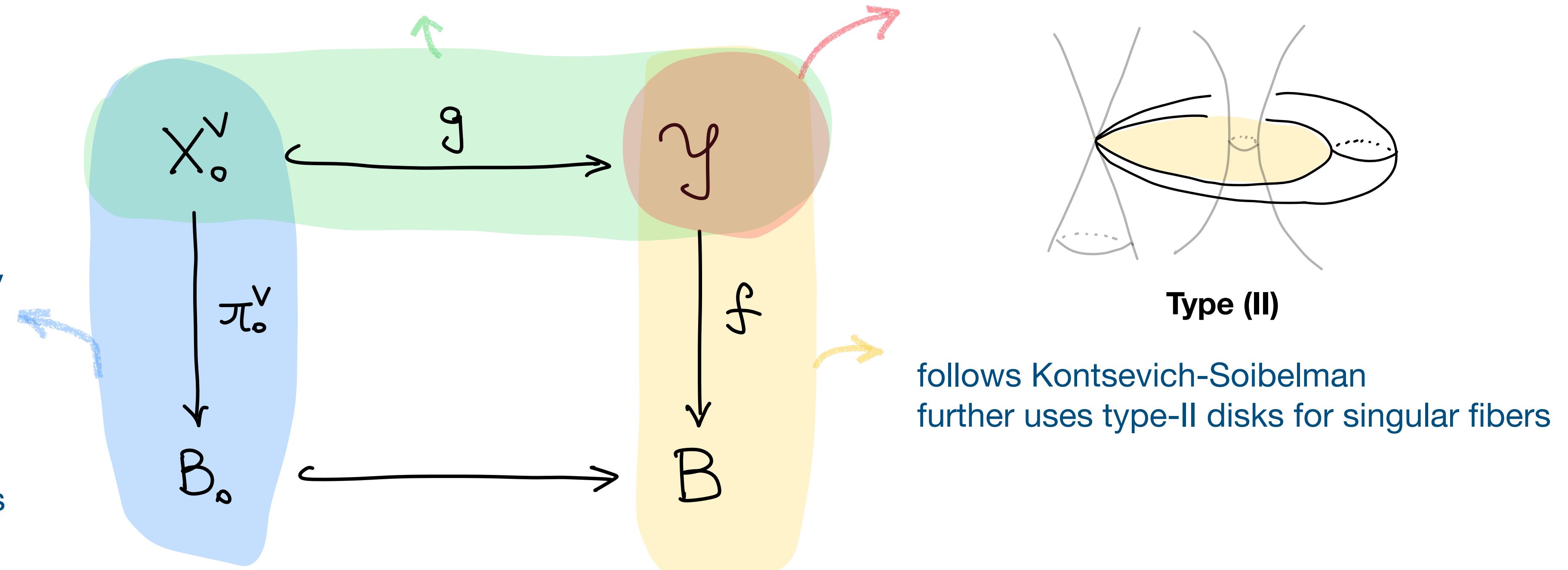
follows family Floer theory and T-duality

We put an analytic topology on

$$X_0^\vee \equiv \bigcup_{q \in B_0} H^1(L_q; U_\Lambda)$$

and the wall-crossing of Maslov-0 disks (type (I)) can imply the identification

$$X_0^\vee \equiv T_+ \sqcup T_- / \sim$$



**Type (II)**

follows Kontsevich-Soibelman  
further uses type-II disks for singular fibers

- Given the above, our very explicit T-duality picture is compatible with so many previous mirror symmetry results. Thus, it is very reasonable to believe the analytic fibers of  $f$  over  $\Delta = B \setminus B_0$  should be the correct **dual singular fibers**
- Meanwhile, by the original family Floer picture, we should expect that  $f^{-1}(q)$  is the Maurer-Cartan set  $\mathcal{MC}(L_q)$  for  $q \in \Delta$
- But unfortunately, the two approaches only have some partial agreements.
- The Maurer-Cartan set is only a strict subset of the corresponding dual singular  $f$ -fiber.

## Dual singular fiber is not a Maurer-Cartan set !

For simplicity, we assume  $n = 2$ . Then  $B = \mathbb{R}^2$ ,  $\Delta = \{0\}$ . Recall  $X = \mathbb{C}^2 \setminus \{z_1 z_2 = 1\}$ ,  $Y = \{x_0 x_1 = 1 + y \text{ in } \Lambda^2 \times \Lambda^*\}$

We have a pinched sphere Lagrangian  $L_0$  as the fiber over the singular point 0

Then, the dual singular fiber is (recall that  $f$  has an explicit formula)

$$\begin{aligned} S := f^{-1}(0) &= \{(x_0, x_1, y) \in Y \mid v(x_0) \geq -\psi(0), v(x_1) \geq \psi(0), v(y) = 0\} \\ &\cong \{(x_0, x_1, y) \in Y \mid v(x_0) \geq 0, v(x_1) \geq 0, v(y) = 0\} \quad x_0 \mapsto T^{-\psi(0)} x_0, x_1 \mapsto T^{\psi(0)} x_1 \end{aligned}$$

- 1) If  $1 + y \in \Lambda_+$  then  $v(x_0) + v(x_1) = v(1 + y) > 0$ , so  $(x_0, x_1) \in \Lambda_0 \times \Lambda_+ \cup \Lambda_+ \times \Lambda_0$  vice versa.
- 2) If  $1 + y \notin \Lambda_+$  then  $v(x_0) = v(x_1) = 0$ , so  $(x_0, x_1)$  is a pair in  $U_\Lambda^2$  such that  $\bar{x}_0 \bar{x}_1 - 1 \neq 0$

In other words, we conclude that

$$S = S_1 \sqcup S_2$$

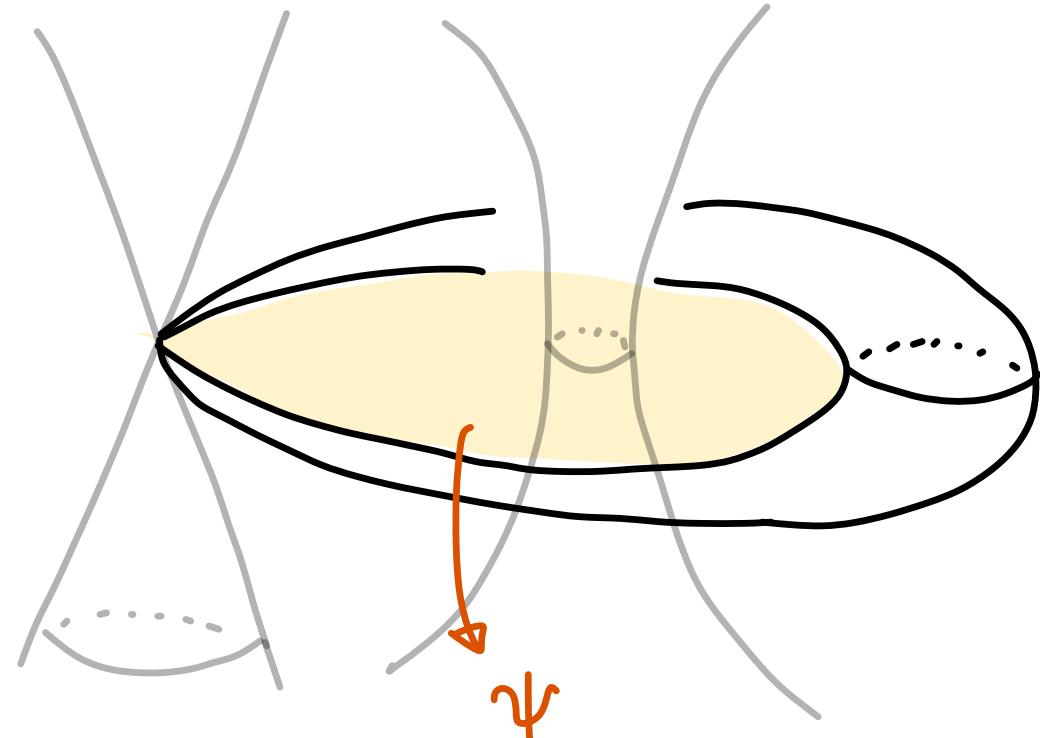
where

$$S_1 = \Lambda_0 \times \Lambda_+ \cup \Lambda_+ \times \Lambda_0$$

$$S_2 = \{(x_0, x_1) \in U_\Lambda^2 \mid \bar{x}_0 \bar{x}_1 \neq 1\} \cong U_\Lambda \times (\mathbb{C}^* \setminus \{-1\} \oplus \Lambda_+)$$

On the other hand, Hong, Kim, and Lau have proved that the Maurer-Cartan set for the singular Lagrangian  $L_0$  is **exactly** given by

$$\mathcal{MC}(L_0) \cong S_1 \subsetneq S$$



- ▶ Therefore,  $f^{-1}(0) \supsetneq \mathcal{MC}(L_0)$
- ▶ There are extra points in  $S_2$  beyond the scope of the conventional MC picture.
- ▶ One possibility is we need additional ‘deformation data’ of MC sets of singular Lagrangians.
- ▶ The NA analytic topology more or less enforces us to handle singular Lagrangian fibers differently.

## Further evidence: a folklore conjecture

### Conjecture: (Kontsevich, Seidel, Auroux, ...)

The critical values of the mirror Landau-Ginzburg superpotential on  $X^\vee$  are the eigenvalues of the quantum multiplication by the first Chern class on  $X$ .

- Recall that  $X = \mathbb{C}^n \setminus \{z_1 \cdots z_n = 1\}$  and  $Y = \{x_0 x_1 = 1 + y_1 + \cdots + y_{n-1}\}$
- The Gross's Lagrangian fibration  $\pi$  can be placed in not only  $X$  but also possibly **a larger ambient manifold  $\bar{X}$**
- Often, we can check **the Maslov-0 disks keep the same**. Then, we will have the same **analytic topology** for  $(X_0^\vee, \pi_0^\vee)$   
Besides, we can also use the same  $g : X_0^\vee \rightarrow \mathcal{Y}_0$  and  $f : \mathcal{Y} \rightarrow B$  as before.
- On the other hand, there are no Maslov-2 disks in  $X$ , but **there will be new Maslov-2 disks** in  $\bar{X}$ .  
It gives a global superpotential  $W_0$  on the analytic open domain  $X_0^\vee \cong \mathcal{Y}_0$  (using the embedding  $g$ ).  
Moreover,  $W_0$  is polynomial for our example.
- By our definition of 'SYZ mirror', the analytic domain  $\mathcal{Y}_0$  is **Zariski dense** in the algebraic  $\Lambda$ -variety  $Y$ .  
Hence, it can be extended on the whole algebraic variety  $Y$ , denoted by  $W$ .  
(In general, it depends on the Kähler form  $\omega$ )
- Choosing various ambient space  $\bar{X}$  will produce various different Landau-Ginzburg superpotential  $W$  on  $Y$

## Further evidence: a folklore conjecture

**ambient space**

$\bar{X} = \mathbb{CP}^n$  Let  $H \in \pi_2(\bar{X})$  be the class of a complex line.

**LG superpotential**

$W = x_1 + \frac{T^{E(H)} x_0^n}{y_1 \cdots y_{n-1}}$  defined on  $Y = \{x_0 x_1 = 1 + y_1 + \cdots + y_{n-1}\}$ , where  $E(H) = \frac{1}{2\pi} \omega \cap H$

**Critical points**

$$\begin{cases} x_0 = T^{-\frac{E(H)}{n+1}} e^{-\frac{2\pi i s}{n+1}} \\ x_1 = nT^{\frac{E(H)}{n+1}} e^{\frac{2\pi i s}{n+1}} \\ y_1 = \cdots = y_{n-1} = 1 \end{cases} \quad s \in \{0, 1, \dots, n\}$$

**Remark:** There may be many other examples by thinking

- (1) other toric CY variety than  $\mathbb{C}^n$
- (2) other compactification  $\bar{X}$

**Critical values**

$(n+1)T^{\frac{\omega(H)}{n+1}} e^{\frac{2\pi i}{n+1}s}$  for  $s \in \{0, 1, \dots, n\}$  One can check the folklore conjecture holds

## Further evidence: a folklore conjecture

ambient space

$$\bar{X} = \mathbb{CP}^m \times \mathbb{CP}^{n-m} \text{ for } 0 < m < n$$

This is also a compactification of  $\mathbb{C}^n$

Let  $H_1, H_2 \in \pi_2(\bar{X})$  be the classes of a complex line in  $\mathbb{CP}^m \times pt$  and in  $pt \times \mathbb{CP}^{n-m}$

LG superpotential

$$1) \quad W = x_1 + \frac{T^{E(H_1)} x_0^m}{y_1 \cdots y_m} + \frac{T^{E(H_2)} x_0^{n-m}}{y_{m+1} \cdots y_{n-1}} \text{ defined on the same } Y = \{x_0 x_1 = 1 + y_1 + \cdots + y_{n-1}\}$$

$$2) \quad W = x_1 + \frac{T^{E(H_1)} x_0}{y_1} + T^{E(H_2)} x_0 \text{ defined on } Y = \{x_0 x_1 = 1 + y_1\}, \text{ (for } n = 2, m = 1\text{)}$$

We have four critical points  $\begin{cases} x_0 = \pm T^{\frac{-E(H_2)}{2}} \\ x_1 = \pm T^{\frac{E(H_1)}{2}} \pm T^{\frac{E(H_2)}{2}} \\ y_1 = \pm T^{\frac{E(H_1)-E(H_2)}{2}} \end{cases}$  in the fiber of  $f$  over  $\hat{q} = \left( \frac{E(H_1) - E(H_2)}{2}, a_\omega \right) \in B = \mathbb{R}^2$

Critical points

**Remark:** It may happen that for some Kähler form  $\omega$ , the number  $a_\omega = 0$ ;

If  $E(H_1) = E(H_2)$ , the  $\hat{q}$  is a singular point; we don't know how to prove the folklore conjecture.

If  $E(H_1) \neq E(H_2)$ , the  $\hat{q}$  lies on the wall; the conventional proof fails for the Maslov-0 disks, but we can still prove it in my other paper. In general, we don't know if the critical points always avoid the walls or singular locus. It seems the **walls** might be dispersed in an *open subset in  $B$*  (e.g. blowup of  $\mathbb{C}^n$  along a hyper)

Critical values

$$(m+1)T^{\frac{E(H_1)}{m+1}} e^{\frac{2\pi i}{m+1}r} + (n-m+1)T^{\frac{E(H_2)}{n-m+1}} e^{\frac{2\pi i}{n-m+1}s}$$

for  $r \in \{0, 1, \dots, m\}$  and  $s \in \{0, 1, \dots, n-m\}$  One can also check the folklore conjecture.

## Generalizations

We can repeat the proof almost verbatim for more general examples.

- Let  $\mathcal{X}_P$  (e.g.  $\mathbb{C}^n$ ) be a toric Calabi-Yau manifold equipped with a toric Kähler form  $\omega$  corresponding to the (unbounded) moment polytope

$$P : \quad \langle m, v_i \rangle + \lambda_i \geq 0 \quad \text{where } m \in M_{\mathbb{R}} \cong \mathbb{R}^n \text{ and } v_i \text{'s are the rays in the fan}$$

We may assume  $v_1, \dots, v_n$  form a basis of  $N = M^*$ . Let  $v_{n+a} = k_{a1}v_1 + \dots + k_{an}v_n$  be the remaining rays. The Calabi-Yau condition means there is  $m_0 \in M$  such that  $\langle m_0, v_i \rangle = 1$ ; in particular,  $k_{a1} + \dots + k_{an} = 1$

Note that  $v_s - v_n$  ( $1 \leq s < n$ ) form a basis in the sub-lattice  $\langle m_0, \cdot \rangle = 0$ . We define

$$X_P = \mathcal{X}_P \setminus (z^{m_0} = 1)$$

It admits a Gross's special Lagrangian fibration  $\pi$  as before.

- On the other side, we define a Laurent polynomial

$$h(y_1, \dots, y_{n-1}) = \sum_{s=1}^{n-1} T^{\lambda_s}(1 + \delta_s)y_s + T^{\lambda_n}y_n(1 + \delta_n) + \sum_a T^{\lambda_{n+a}}(1 + \delta_{n+a}) \prod_{s=1}^{n-1} y_s^{k_{as}}$$

The singular locus  $\Delta$  of  $\pi$  is precisely decided by the tropicalization  $h_{\text{trop}}$  of  $h$ . Here each  $\delta_i \in \Lambda_+$  is given by the counts of stable disks with sphere bubbles. Sometimes they are not zero; the valuation  $v(\delta_i)$  is the smallest area of sphere bubble. Finally, we define

$$Y_h = \{(x, y) \in \Lambda^2 \times (\Lambda^*)^{n-1} \mid x_0x_1 = h(y)\}$$

**Theorem:**  $Y_h$  is SYZ mirror to  $X_P$

**Example:** Take  $h(y) = 1 + y_1 + \dots + y_{n-1}$ , and the corresponding tropical polynomial is

$$h_{\text{trop}} = \min\{0, q_1, \dots, q_{n-1}\}$$

which somehow plays the leading role. For instance

- describes the walls on the A side
- appears in the formula of  $f$  i.e. the pair  $(j, F)$  on the B side
- gives the singular locus  $\Delta = B \setminus B_0$

**“SYZ converse”:** Given  $h_{\text{trop}}$  and  $h$ , we conversely have  $P' := \{(\bar{q}, q_n) \mid q_n + h_{\text{trop}}(\bar{q}) \geq 0\} \cong P$

This picture will be lost if we only work over  $\mathbb{C}$ .

**Example:**  $h(y) = y_1 + T^{-1}y_2 + T^{3.14} + T^2y_1^2 + y_1y_2 + T^2y_2^2$

Then, we can recover  $P$ :  $v_1 = (1, 0, 0)$   $\lambda_1 = 0$

Have infinite such examples.  $v_2 = (0, 1, 0)$   $\lambda_2 = -1$

$$v_3 = (0, 0, 1) \quad \lambda_3 = 3.14$$

$$v_4 = (2, 0, -1) \quad \lambda_4 = 2$$

$$v_5 = (1, 1, -1) \quad \lambda_5 = 0$$

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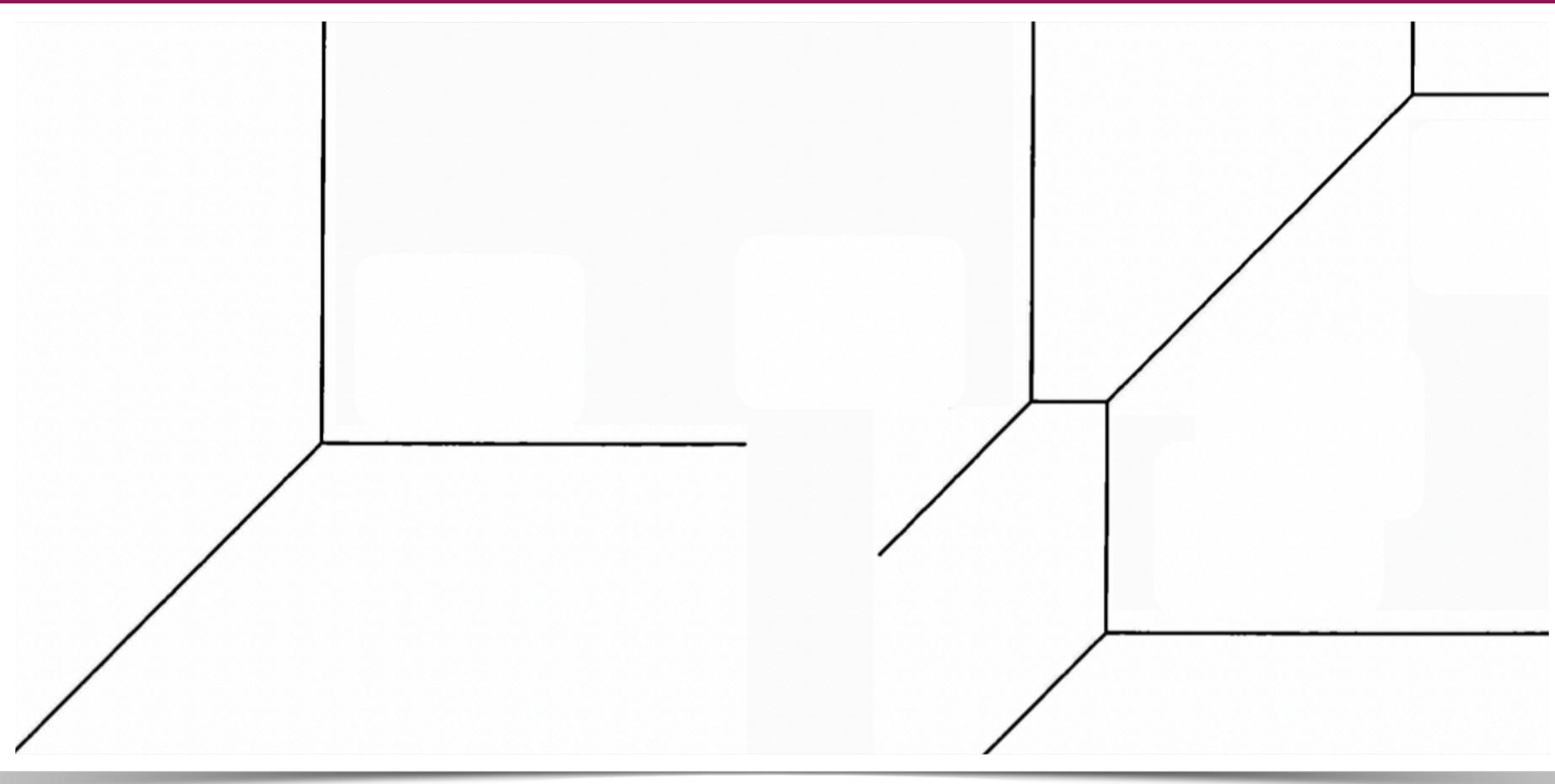
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$v_6 = (0, 2, -1) \quad \lambda_6 = 2$

**Thanks for your attention !**