Introduction and Motivation QFI based on expectation values Case study Conclusion and outlook

# Optimal bound on the quantum Fisher Information

Based on few initial expectation values of the prove state.

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Recent Advances in Quantum Metrology; Warsaw - 2016

#### Outline

- Introduction and Motivation
- 2 QFI based on expectation values: Are they optimal?
  - Optimisation problem
- Case study
  - Spin squeezed states
  - Unpolarised Dicke states
- Conclusion and outlook

Many inequalities have been proposed to lower bound the quantum Fisher Information.

#### Bounds for qFI

$$F_{Q}[\varrho, J_{z}] \geq \frac{\langle J_{x} \rangle^{2}}{(\Delta J_{y})^{2}}, \qquad F_{Q}[\varrho, J_{y}] \geq \beta^{-2} \frac{\langle J_{x}^{2} + J_{z}^{2} \rangle}{(\Delta J_{z})^{2} + \frac{1}{4}},$$

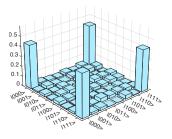
$$F_{Q}[\varrho, J_{z}] \geq \frac{4(\langle J_{x}^{2} + J_{y}^{2} \rangle)^{2}}{2\sqrt{(\Delta J_{x}^{2})^{2} (\Delta J_{y}^{2})^{2} + \langle J_{x}^{2} \rangle - 2\langle J_{y}^{2} \rangle(1 + \langle J_{x}^{2} \rangle) + 6\langle J_{y}J_{x}^{2}J_{y} \rangle}}$$

[ I.A., B. Lücke, J. Peise, C. Klempt & G. Toth, NJP 17, 083027 (2015) ]

[ L. Pezzé & A. Smerzi, PRL 102, 100401 (2009) ]

[ Z. Zhang & L.-M. Duan, NJP 16, 103037 (2014) ]

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- ② For large systems, we only have a couple of expectation values to characterise the state.



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- The archetypical criteria that demonstrates metrologicaly useful entanglement on the state.

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- Many inequalities have been proposed to lower bound the quantum Fisher Information.
- For large systems, we only have a couple of expectation values to characterise the state.
- The archetypical criteria that demonstrates metrologicaly useful entanglement on the state.
- It is essential either to verify them or find new ones for different set of expectation values.

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## The non-trivial exercise of computing the qFI

• Different forms of the qFI

$$F_{Q}[\varrho, J_{z}] = 2 \sum_{\lambda, \gamma} \frac{(\rho_{\lambda} - \rho_{\gamma})^{2}}{\rho_{\lambda} + \rho_{\gamma}} |\langle \lambda | J_{z} | \gamma \rangle|^{2}$$
$$F_{Q}[\varrho, J_{z}] = \min_{\{\rho_{k}, |\Psi_{k}\rangle\}} 4 \sum_{k} p_{k} (\Delta J_{z})^{2}_{|\Psi_{k}\rangle}$$

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[ M.G.A. Paris, Int. J. Quant. Inf. 7, 125 (2009) ]
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• In the general case, usually *lower bounded* by its "classical" counterparts.

## Optimisation: Legendre Transform

 For a convex function of the state, we construct a thight lower bound as follows,

$$g(\varrho) \geq \mathcal{B}\big(\{w_k := \langle W_k \rangle\}\big) = \sup_{\{r_k\}} \big(r \cdot w - \sup_{\varrho} [r \cdot \langle W \rangle_{\varrho} - g(\varrho)]\big).$$

The scalar product reads as follows 
$$r \cdot w \equiv \sum_k r_k w_k$$
 and  $r \cdot \langle W \rangle_{\varrho} \equiv \sum_k r_k \langle W_k \rangle_{\varrho}$ .

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• When  $g(\varrho)$  is defined as *infimum over the convex roof*, the  $2^{\rm nd}$  optimisation simplified to pure states only,

$$\mathcal{B}(\lbrace w_k \rbrace) = \sup_{\lbrace r_k \rbrace} \big( r \cdot w - \sup_{|\psi\rangle} [r \cdot \langle W \rangle - g(|\psi\rangle)] \big).$$

- O. Gühne, M. Reimpell & R.F. Werner, PRL 98, 110502 (2007)
- [ J. Eisert, F.G.S.L. Brandão & K.M.R. Audenaert, NJP **9**, 46 (2007) ]

#### Optimisation for the qFI

The *simplicity* of qFI for pure states allows to rewrite the problem:

$$\mathcal{F}(\{w_k\}) = \sup_{\{r_k\}} \left(r \cdot w - \sup_{\mu} [\lambda_{\max}(r \cdot W - 4(J_z - \mu)^2)]\right).$$

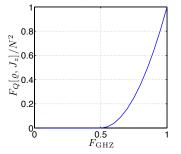
• Therefore, we have parametrised the optimisation, which leads to a *more efficient finding* of the solution.

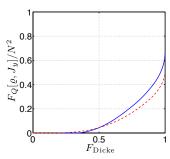
[ I.A., M. Kleinmann, O. Güne & G. Tóth, arXiv:1511.05203 ]

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• One of the cases using projector operators, *i.e.*, using the *fidelity*, leads to *analytic soulution*!

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- Hence, the optimisation can be accomplished *only using 2* operators  $\{J_z, J_x^2\}$  while the resulting bound is mapped directly to  $\langle J_z \rangle, (\Delta J_x)^2$ .

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- **1** Hence, the optimisation can be accomplished *only using 2* operators  $\{J_z, J_x^2\}$  while the resulting bound is mapped directly to  $\langle J_z \rangle$ ,  $(\Delta J_x)^2$ .
- **3** Since  $F_Q \ge \langle J_z \rangle^2 / (\Delta J_x)^2$  is a valid lower bound, we compare it with our numerical result.

• We numerically optimise the lower bound of qFI for a 4 particle system for all possible values of  $\langle J_z \rangle$  and  $(\Delta J_x)^2$ .

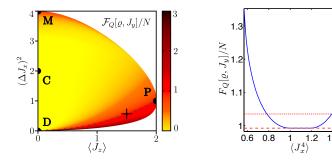


Figure: Below the dashed line,  $\mathcal{F}$  surpasses the shot noise limit. Cross point enhanced with extra parameter. The result shows an extreme similarity with respect to  $F_Q \geq \langle J_z \rangle^2/\left(\Delta J_x\right)^2$ .

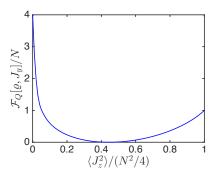
[ I.A., M. Kleinmann, O. Güne & G. Tóth, arXiv:1511.05203 ]

## Metrology with unpolarised Dicke states

• For unpolarised Dicke state we use these 3 operators,  $\{J_x^2, J_y^2, J_z^2\}$ , with the following constraint,  $\langle J_x^2 \rangle = \langle J_y^2 \rangle$ .

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- ② For  $\sum_{I}\langle J_{I}^{2}\rangle=\frac{N}{2}(\frac{N}{2}+1)$ , *i.e.* maximal, and for 6 particles:



#### Realistic characterisation of Dicke state

Experiment  $\rightarrow$  [ B. Lücke et al., PRL 112, 155304 (2014) ]

 We now consider an interesting experimental case with BECs where the following initial values are measured.

#### Experimental details of unpolarized Dicke state with BEC

$$N = 7900 \qquad \langle J_z^2 \rangle = 112 \pm 31$$
$$\langle J_x^2 \rangle = \langle J_y^2 \rangle = 6 \times 10^6 \pm 0.6 \times 10^6$$

 For that large system we developed a successful extrapolation tool.

## Extrapolation:

• First, we equivalently increase the expectation values and the unknown bound *to the symmetric subspace*.

$$\langle J_z^2 \rangle_{\mathrm{sym},N} = 145.69 \qquad \langle J_x^2 \rangle_{\mathrm{sym},N} = 7.8 \times 10^6$$

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② We freeze  $\langle J_z^2 \rangle$  while we reduce the particle number of an auxiliary system. Here, we perform the optimisation.

$$2\langle J_x^2\rangle_{\mathrm{sym},N'}=\frac{N'}{2}(\frac{N'}{2}+1)-\langle J_z^2\rangle_{\mathrm{sym},N'}$$

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The extrapolation is directly obtained by

$$\frac{\mathcal{F}}{\sum_I \langle J_I^2 \rangle} \gtrsim \frac{\mathcal{F}_{\mathrm{sym},N}}{N^2/4} \approx \frac{\mathcal{F}_{\mathrm{sym},N'}}{N'^2/4}$$

• Obtaining a numerical lower bound this *large system*.

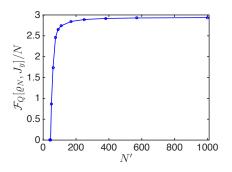


Figure: The extrapolated lower bound *approaches rapidly* to the desired value. The points over the line indicate the systems used for optimisation.

[ I.A., M. Kleinmann, O. Güne & G. Tóth, arXiv:1511.05203 ]

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- We have shown that the lower bounds can be improved with few extra considerations.
- For large systems the optimisation method can be complemented with scaling considerations, such as the extrapolation for Dicke states case.
- This method is constructed to be universal and it could be used in many other situations.

## Thank you for your attention!

Preprint  $\rightarrow$  arXiv:1511.05203

#### Groups' home pages

- → https://sites.google.com/site/gedentqopt
- → http://www.physik.uni-siegen.de/tqo/



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