

# Optimal bound on the quantum Fisher Information

*Based on few initial expectation values of the prove state.*

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# Outline

- 1 Introduction and Motivation
- 2 QFI based on expectation values: Are they optimal?
  - Optimization problem
- 3 Case study
  - Spin squeezed states
  - Unpolarized Dicke states
- 4 Conclusion and outlook

- ① **Many inequalities** have been proposed to lower bound the quantum Fisher Information.

### Bounds for qFI

$$\mathcal{F}[\varrho, J_z] \geq \frac{\langle J_x \rangle^2}{(\Delta J_y)^2}, \quad \mathcal{F}[\varrho, J_y] \geq \beta^{-2} \frac{\langle J_x^2 + J_z^2 \rangle}{(\Delta J_z)^2 + \frac{1}{4}},$$

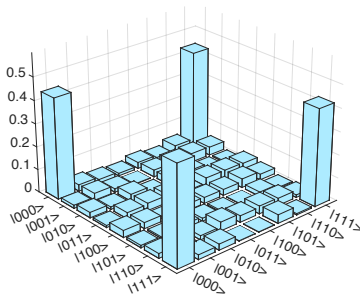
$$\mathcal{F}[\varrho, J_z] \geq \frac{4(\langle J_x^2 + J_y^2 \rangle)^2}{2\sqrt{(\Delta J_x^2)^2 (\Delta J_y^2)^2 + \langle J_x^2 \rangle - 2\langle J_y^2 \rangle (1 + \langle J_x^2 \rangle)} + 6}$$

[ I.A., B. Lücke, J. Peise, C. Klempt & G. Toth, New J. Phys. **17**, 083027 (2015) ]

[ L. Pezzé & A. Smerzi, Phys. Rev. Lett. **102**, 100401 (2009) ]

[ Z. Zhang & L.-M. Duan, 2014 New J. Phys. **16** 103037 (2014) ]

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- ① *Many inequalities* have been proposed to lower bound the quantum Fisher Information.
- ② Typically, we only have *a couple of expectation values* to characterize the state.
- ③ The archetypical criteria that demonstrates *useful entanglement* on the state.
- ④ It is essential either to *verify them or find new ones* for different set of expectation values.

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# The non-trivial exercise of computing the qFI

## 1 Different forms of the qFI

$$\mathcal{F}[\varrho, J_z] = 2 \sum_{\lambda, \gamma} \frac{(p_\lambda - p_\gamma)^2}{p_\lambda + p_\gamma} |\langle \lambda | J_z | \gamma \rangle|^2$$

$$\mathcal{F}[\varrho, J_z] = \min_{\{p_k, |\psi_k\rangle\}} 4 \sum_k p_k (\Delta J_z)_{|\psi_k\rangle}^2$$

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## 3 In the general case, usually *lower bounded* by its "classical" counterparts.

# Optimization: Legendre Transform

- For any kind of function of the state, we construct a *tight lower bound* as follows,

$$g(\varrho) \geq \mathcal{B}(\{w_k := \langle W_k \rangle\}) = \sup_{\{r_k\}} (r \cdot w - \sup_{\varrho} [r \cdot \langle W \rangle - g(\varrho)]).$$

- When  $g(\varrho)$  is the infimum over the convex roof, the 2<sup>nd</sup> optimization simplified to pure states only,

$$\mathcal{B}(\{w_k\}) = \sup_{\{r_k\}} (r \cdot w - \sup_{|\psi\rangle} [r \cdot \langle W \rangle - g(|\psi\rangle)]).$$

TODO: Cite Otfried.

## Optimization for the qFI

It is slightly different than for other functions because of the *simplicity of the qFI for pure states*.

$$\mathcal{F}(\{w_k\}) = \sup_{\{r_k\}} (r \cdot w - \sup_{\mu} [\lambda_{\max}(r \cdot W - 4(J_z - \mu)^2)]).$$

Therefore, we have parametrised the optimization, which leads to a *more efficient finding* of the solution.

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# Measuring $\langle J_z \rangle$ and $(\Delta J_x)^2$ for Spin Squeezed States

- 1 We use the following 3 operators  $\{J_z, J_x, J_x^2\}$  to characterize the input state with their respective expectation values.
- 2 In the direction of  $\langle J_x \rangle$  the worst case is it take the value zero
- 3 Therefore the optimisation can be done only for 2 operators  $\{J_z, J_x^2\}$  and it can be mapep directly to  $\langle J_z \rangle, (\Delta J_x)^2$ .





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