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Spin Orbit coupling implementation in DFTB/GFN-xTB

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Kohn-Sham equations  $E=E[\rho(\mathbf{r})]$ 



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Expanding  $E_{\rho}$  at  $\rho_0$  to second (or third) order in fluctuation  $\delta\rho$ 

$$E_{\mathsf{tot}}[\rho_0 + \delta \rho] = E^0[\rho_0] + E^1[\rho_0, \delta \rho] + E^2[\rho_0, \delta \rho^2] + O(\delta \rho^3)$$



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- $E^0[\rho_0] = -\frac{1}{2} \int \int \frac{\rho_0(\mathbf{r})\rho_0(\mathbf{r}')}{|\mathbf{r}-\mathbf{r}'|} d\mathbf{r} d\mathbf{r}' \int V^{xc}[\rho_0]\rho_0(\mathbf{r}) d\mathbf{r} + E_{xc}[\rho_0] + \frac{1}{2} \sum_{A \neq B} \frac{Z_A Z_B}{R_{AB}}$
- $E_1[\rho_0, \delta \rho] = \sum_i^{\text{occ}} f_i \langle \psi_i | -\frac{1}{2} \nabla^2 + V_{\text{ext}} + \int \frac{\rho_0(\mathbf{r})}{|\mathbf{r} \mathbf{r}'|} + V_{xc}[\rho_0] |\psi_i \rangle$
- $E^{2}[\rho_{0}, \delta \rho^{2}] = \frac{1}{2} \int \int \left( \frac{1}{|\mathbf{r} \mathbf{r}'|} + \frac{\delta^{2} E_{xc}}{\delta \rho(\mathbf{r}) \delta \rho(\mathbf{r}')} \right) \delta \rho(\mathbf{r}) \delta \rho(\mathbf{r}') d\mathbf{r} d\mathbf{r}'$





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# Non-SCC (?) DFTB (DFTB1)

Expanding  $E_{\rho}$  at  $\rho_0$  to second (or third) order in fluctuation  $\delta \rho$ 

$$\begin{split} E_{\mathsf{tot}}[\rho_0 + \delta \rho] &= \underline{E}^{\mathsf{0}}[\rho_0] + E^{\mathsf{1}}[\rho_0, \delta \rho] + E^{\mathsf{2}}[\rho_0, \delta \rho^2] + O(\delta \rho^3) \\ &= \sum_{I < J} V_{IJ}^{\mathsf{rep}}(\mathbf{r}_{IJ}) + \sum_{i}^{\mathsf{occ}} f_i \langle \psi_i | \hat{H}[\rho_0] | \psi_i \rangle + \dots \end{split}$$



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$$= \sum_{I < J} V_{IJ}^{\text{rep}}(\mathbf{r}_{IJ}) + \sum_i^{\text{occ}} f_i \langle \psi_i | \hat{H}[\rho_0] | \psi_i \rangle + \dots$$

- ullet minimal "pseudo" atomic basis set  $\psi_i = \sum_{\mu} c_{\mu i} arphi_{\mu}$
- reference density  $\rho_0$
- tabulated as function of distance





# Non-SCC (?) DFTB (DFTB1)

$$E_{\mathsf{Non\text{-}SCC}} = \sum_{I < J} V_{IJ}^{\mathsf{rep}}(\mathbf{r}_{IJ}) + \sum_{i}^{\mathsf{occ}} f_i \langle \psi_i | \hat{H}[\rho_0] | \psi_i \rangle$$

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- reference density  $ho_0$
- tabulated as function of distance





Expanding  $E_{
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ho$ 

$$E_{\text{tot}}[\rho_0 + \delta \rho] = E^0[\rho_0] + E^1[\rho_0, \delta \rho] + E^2[\rho_0, \delta \rho^2] + O(\delta \rho^3)$$

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$$E^{2}[\rho_{0}, \delta \rho^{2}] = \frac{1}{2} \int \int \left( \frac{1}{|\mathbf{r} - \mathbf{r}'|} + \frac{\delta^{2} E_{xc}}{\delta \rho(\mathbf{r}) \delta \rho(\mathbf{r}')} \right) \delta \rho(\mathbf{r}) \delta \rho(\mathbf{r}') d\mathbf{r} d\mathbf{r}'$$



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$$E^{2}[\rho_{0}, \delta \rho^{2}] = \frac{1}{2} \int \int \left( \frac{1}{|\mathbf{r} - \mathbf{r}'|} + \frac{\delta^{2} E_{\text{XC}}}{\delta \rho(\mathbf{r}) \delta \rho(\mathbf{r}')} \right) \delta \rho(\mathbf{r}) \delta \rho(\mathbf{r}') d\mathbf{r} d\mathbf{r}'$$

$$\int \delta \rho(\mathbf{r}) = \sum_{I} \Delta q_{I} \delta \rho_{I}(\mathbf{r})$$
•  $E^{2}[\rho_{0}, \delta \rho^{2}] \approx \frac{1}{2} \sum_{IJ} \gamma_{IJ} (R_{IJ}) \Delta q_{I} \Delta q_{J}$ 





- $E_{\mathsf{DFTB2(3)}} = E_{\mathsf{Non-SCC}} + \frac{1}{2} \sum_{IJ} \gamma_{IJ} (R_{IJ}) \Delta q_I \Delta q_J + \frac{1}{3} \sum_{IJ} \Delta q_I^2 \Delta q_J \Gamma_{IJ}$
- $H_{\mu\nu} = H_{\mu\nu}^0 + H_{\mu\nu}^2 [\gamma^h, \Delta q] + H_{\mu\nu}^3 [\Gamma, \Delta q], \quad \mu \in I, \nu \in J$
- $q_I = \sum_i f_i \int_{V_I} |\psi_i(\mathbf{r})|^2 d^3r = \frac{1}{2} \sum_i^{\text{occ}} f_i \sum_{\mu \in I} \sum_{\nu} \left( c_{\mu i}^* c_{\nu i} S_{\mu \nu} + c_{\nu i}^* c_{\mu i} S_{\nu \mu} \right)$



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- $q_I = \sum_i f_i \int_{V_I} |\psi_i(\mathbf{r})|^2 d^3r = \frac{1}{2} \sum_i^{\text{OCC}} f_i \sum_{\mu \in I} \sum_{\nu} \left( c_{\mu i}^* c_{\nu i} S_{\mu \nu} + c_{\nu i}^* c_{\mu i} S_{\nu \mu} \right)$

Charges iterated until self consistency has been reached

$$q^{(0)} \xrightarrow{H^{(0)}_{\mu\nu}} c^{(0)}_{ij} \to q^{(1)} \xrightarrow{H^{(1)}_{\mu\nu}} c^{(1)}_{ij}$$

Self-consistent charge (SCC) iteration



# Spin-Orbit Coupling (SOC)

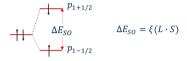






Moving charge creates magnetic field (1820)





> SOC hamiltonian in Dirac Equation

$$\hat{H}^{SOC} = -\frac{e\hbar}{4m^2c^2} \mathbf{\sigma} \cdot [E_f \times \hat{p}]$$

Spherical potential & static case

$$\widehat{H}^{SOC} = -\frac{\hbar}{4m^2c^2} \frac{1}{r} \frac{dV}{dr} \mathbf{\sigma} \cdot [\hat{r} \times \hat{p}]$$

$$\widehat{H}^{SOC} = \xi(\widehat{L} \cdot \widehat{S})$$



# Spin-Orbit Coupling (SOC)

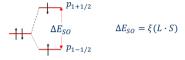


Hans C. Ørsted (1777-1851)



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See J. Chem. Theory Comput. 18, 4472 (2022)





• Single-particle on-site spin-orbit interaction

$$\hat{H}_{SO} = \frac{\zeta}{2} \mathbf{L} \cdot \mathbf{S} = \frac{\zeta}{2} (\hat{L}_x \sigma_x + \hat{L}_y \sigma_y + \hat{L}_z \sigma_z)$$

$$= \frac{\zeta}{2} \begin{pmatrix} \hat{L}_z & \hat{L}_x - i\hat{L}_y \\ \hat{L}_x + i\hat{L}_y & \hat{L}_z \end{pmatrix} = \frac{\zeta}{2} \begin{pmatrix} \hat{L}_z & \hat{L}_- \\ \hat{L}_+ & \hat{L}_z \end{pmatrix}$$

#### Pauli matrices:

$$\sigma_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

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- 
$$\hat{L}_z|Y_l^{\pm m}\rangle = \pm ml|Y_l^{\pm m}\rangle$$
,  $\hat{L}_\pm|Y_l^{\pm m}\rangle = \sqrt{l(l+1) - m(m\pm 1)}|Y_l^{\pm m}\rangle$ 

$$- \ \varphi_{\mu}(\mathbf{r} - \mathbf{R}_I) = R_{\mu}(r)\widetilde{Y}_{\mu}(\theta,\varphi)(\mu \in I), \quad \text{where } \widetilde{Y} \propto Y_{lm} \pm Y_{lm}^*$$





• 2-component spinor wavefunctions

$$\psi_i = \sum_{\mu} \binom{c_{\mu i}^{\alpha}}{c_{\mu i}^{\beta}} \varphi_{\mu}$$

Non-collinear Hamiltonian

$$\hat{H} = \left(\hat{H}_{\mu\nu}^{0} + \hat{H}_{\mu\nu}^{2} + \hat{H}_{\mu\nu}^{3}\right) \otimes \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} + \begin{bmatrix} \xi_{I}^{l} \begin{pmatrix} \hat{L}_{z} & \hat{L}_{-} \\ \hat{L}_{+} & -\hat{L}_{z} \end{pmatrix}_{l} + \xi_{J}^{l'} \begin{pmatrix} \hat{L}_{z} & \hat{L}_{-} \\ \hat{L}_{+} & -\hat{L}_{z} \end{pmatrix}_{l'} \end{bmatrix}$$



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Secular equation

$$\sum_{\mu} \begin{pmatrix} H^{\alpha\alpha}_{\mu\nu} - \epsilon_i S_{\mu\nu} & H^{\alpha\beta}_{\mu\nu} \\ H^{\beta\alpha}_{\mu\nu} & H^{\beta\beta}_{\mu\nu} - \epsilon_i S_{\mu\nu} \end{pmatrix} \begin{pmatrix} c^{\alpha}_{\mu i} \\ c^{\beta}_{\mu i} \end{pmatrix} = 0, \quad H^{\sigma\sigma'}_{\mu\nu} = \langle \varphi_{\mu} | \hat{H}^{\sigma\sigma'} | \varphi_{\nu} \rangle$$



• Density matrix

$$\rho(\mathbf{r}) = \sum_{i}^{\text{occ}} f_{i} \begin{pmatrix} \psi_{i}^{\alpha*} \\ \psi_{i}^{\beta*} \end{pmatrix} \begin{pmatrix} \psi_{i}^{\alpha} \\ \psi_{i}^{\beta} \end{pmatrix} = \begin{pmatrix} \rho^{\alpha\alpha} & \rho^{\alpha\beta} \\ \rho^{\beta\alpha} & \rho^{\beta\beta} \end{pmatrix}$$
$$= n(\mathbf{r}) \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} + m^{x}(\mathbf{r})\sigma_{x} + m^{y}(\mathbf{r})\sigma_{y} + m^{z}(\mathbf{r})\sigma_{z}$$

Pauli matrices:

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Electron and magnetization densities

$$n(\mathbf{r}) = \frac{1}{2} \operatorname{Re} \left( \rho^{\alpha \alpha} + \rho^{\beta \beta} \right), \quad m^z(\mathbf{r}) = \frac{1}{2} \operatorname{Re} \left( \rho^{\alpha \alpha} - \rho^{\beta \beta} \right)$$



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Mulliken charges

$$q_{\mu} = \frac{1}{2} \left( q_{\mu}^{\alpha \alpha} + q_{\mu}^{\beta \beta} \right), \quad q_{\mu}^{\sigma \sigma'} = \sum_{\nu} \rho^{\sigma \sigma'} S_{\mu \nu}$$





• Energy :  $E_{SO} = \text{Tr}(\rho H_{SO})$ 

$$E_{\mathsf{SO}} = \sum_{\mu\nu} \left[ H_{\mathsf{SO}}^{\alpha\alpha} \rho_{\nu\mu}^{\alpha\alpha} + H_{\mathsf{SO}}^{\beta\beta} \rho_{\nu\mu}^{\beta\beta} + 2 \operatorname{Re} \left( H_{\mathsf{SO}}^{\alpha\beta} \rho_{\nu\mu}^{\alpha\beta} \right) \right]$$

• Forces :  $\mathbf{F}_{\mathsf{SO},I} = -\sum_{\sigma\sigma'} rac{\partial E_{\mathsf{SO}}^{\sigma\sigma'}}{\partial \mathbf{R}_I}$ 

$$\mathbf{F}_{\mathsf{SO},I} = \sum_{\mu\nu,\mathbf{R}} \left[ \sum_{\sigma\sigma'} \rho_{\mu\nu}^{\sigma\sigma'}(\mathbf{R}) \frac{\partial H_{\mathsf{SO}}^{\sigma\sigma'}(\mathbf{R})}{\partial \mathbf{R}_I} - \sum_{\sigma\sigma'} \rho_{\mathsf{SO}}^{\epsilon,\sigma\sigma'}(\mathbf{R}) \frac{\partial S_{\mu\nu}(\mathbf{R})}{\partial \mathbf{R}_I} \right]$$



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No explicit contribution

| $Bi_2$  | Bond Length (Å) | Frequency (cm $^{-1}$ ) |
|---------|-----------------|-------------------------|
| w SOC   | 1.98            | 1366                    |
| w/t SOC | 2.02            | 1336                    |

Table: calculated with GFN1-xTB in AMS/DFTB





• Periodic boundary conditions

$$\phi_{\mu}^{\mathbf{k}}(\mathbf{r}) = \frac{1}{\sqrt{N}} \sum_{\mathbf{R}} e^{i\mathbf{k}\cdot\mathbf{R}} \varphi_{\mu}(\mathbf{r} - \mathbf{R})$$

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Fatbands

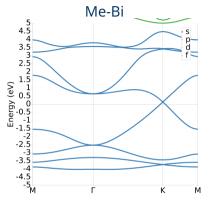
$$n_{i\mathbf{k}}^{\mu} = \sum_{\nu} \left| \langle \psi_{\mu}^{\mathbf{k}} | \phi_{\nu}^{\mathbf{k}} \rangle \right|^{2} = \frac{1}{2} \sum_{\nu} \operatorname{Re} \left( c_{\mu i \mathbf{k}}^{\alpha *} c_{\nu i \mathbf{k}}^{\alpha} S_{\mu \nu} + c_{\mu i \mathbf{k}}^{\beta *} c_{\nu i \mathbf{k}}^{\beta} S_{\mu \nu} \right)$$

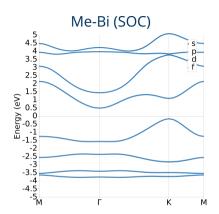
• Spin texture  $S_{i\mathbf{k}} = \langle \psi_{i\mathbf{k}} | \sigma | \psi_{i\mathbf{k}} \rangle$ 

$$m_{i\mathbf{k}}^{\mu,z} = \frac{1}{2} \sum \operatorname{Re} \left( c_{\mu i\mathbf{k}}^{\alpha*} c_{\nu i\mathbf{k}}^{\alpha} S_{\mu\nu} - c_{\mu i\mathbf{k}}^{\beta*} c_{\nu i\mathbf{k}}^{\beta} S_{\mu\nu} \right)$$



#### Benchmark calculations



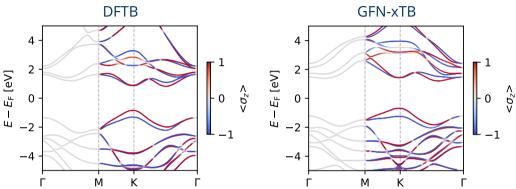


- Me functionalized Bi(111) topological insulators
- QUASINANO2013 Slater-Koster parameters
- Visualized in amsbands





#### Benchmark calculations



- 2D WS<sub>2</sub> with SOC
- Plotted with python
- Spinor Visualization in amsbands to be done





#### Benchmark calculations

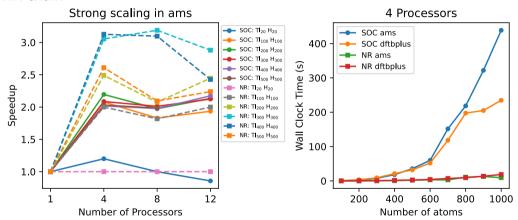
- $H_{SO}$  is **transferable** among different parameter sets and methods
- Successful benchmarked on close-shell molecules and materials such as III-V 3D semiconductors, TMDC 2D crystals, topological insulators, with comparison to DFTB+
- **Regression test** on single point calculation, geometry optimization, and frequency calculation





## **Scalability**

#### TIH chain







# THANK YOU





### Backup: GFN-xTB

$$\psi_i = \sum_{\mu} c_{\mu i} \varphi_{\mu}(\zeta, \mathsf{STO} - mG) \mathsf{b}$$

$$H_{\mu\nu} = K_{IJ} \frac{1}{2} (k_l + k_{l'}) \frac{1}{2} (h_J + h_{J'}) S_{\mu\nu} (1 + k_{EN} \Delta E_N^2) \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} + \Pi(R_{I,l'}) \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$
$$+ \frac{1}{2} S_{\mu\nu} \sum_{c} \sum_{l''} \left( Y_{IC,l''} + Y_{JC,l''} \right) P_{l''}^c \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} + \frac{1}{2} S_{\mu\nu} (q_I^2 + q_J^2) \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$
$$+ \frac{1}{2} S_{\mu\nu} \left[ \epsilon_{\mu} \begin{pmatrix} L_z & L_- \\ L_+ & -L_z \end{pmatrix} + \epsilon_{\nu} \begin{pmatrix} L_z & L_- \\ L_+ & -L_z \end{pmatrix} \right]$$

