

$$E_{\phi}^{(0)} = d\phi k^2 G^{(0)}(\vec{r}, \vec{r}') + (\vec{d} \cdot \vec{\nabla}) \frac{1}{\rho} \frac{\partial G^{(0)}(\vec{r}, \vec{r}')}{\partial \phi} \\ = d\phi \left( \frac{1}{\rho^2} \frac{\partial^2 G^{(0)}(\vec{r}, \vec{r}')}{\partial \phi^2} + k^2 G^{(0)}(\vec{r}, \vec{r}') \right) + d\phi \left( \frac{1}{\rho} \frac{\partial^2 G^{(0)}(\vec{r}, \vec{r}')}{\partial \rho \partial \phi} - \frac{1}{\rho^2} \frac{\partial G^{(0)}(\vec{r}, \vec{r}')}{\partial \phi} \right) + \frac{d_z}{\rho} \frac{\partial G^{(0)}(\vec{r}, \vec{r}')}{\partial z \partial \phi}$$

Using the definition that

$$d_{\pm} = \frac{1}{\sqrt{2}} (e^{-i\phi} d_{\pm} + e^{i\phi} d_{\mp}), \quad d_{\rho} = \frac{1}{\sqrt{2}} (-e^{-i\phi} d_{+} + e^{i\phi} d_{-}), \quad d_z = d_0 = d_z^{(0)}$$

$$d_{\pm} = \mp \frac{1}{\sqrt{2}} (d_{\rho}^{(0)} \pm i d_{\phi}^{(0)}), \quad d_0 = d_z^{(0)} \leftarrow \text{superscript "(0)" means at the atom position. } \vec{r}' = (\rho', \phi', z'=0).$$

$$G^{(0)}(\vec{r}, \vec{r}') = \frac{e^{-ik|\vec{r}-\vec{r}'|}}{|\vec{r}-\vec{r}'|} = \frac{1}{2} \sum_{m=-\infty}^{\infty} d_{\rho} e^{im\phi+i\beta z} J_m(\rho\rho') H_m^{(1)}(\rho\rho') \quad \text{for } (\rho < \rho').$$

we have

$$E_{\phi}^{(0)} = \frac{1}{2} d_{\rho} \left\{ \sum_{m=-\infty}^{\infty} \left( -\frac{m^2 d_{\rho}}{\rho^2} + d_{\rho} k^2 - \frac{i m d_{\rho}}{\rho^2} - \frac{\beta m d_z}{\rho} \right) e^{im\phi+i\beta z} J_m(\rho\rho') H_m^{(1)}(\rho\rho') + \sum_{m=-\infty}^{\infty} \frac{i m d_{\rho}}{\rho} e^{im\phi+i\beta z} \frac{\partial J_m(\rho\rho')}{\partial \rho} H_m^{(1)}(\rho\rho') \right\} \\ = \frac{1}{2} d_{\rho} \left\{ \sum_{m=-\infty}^{\infty} \left[ \left( -\frac{i m^2 d_{+}}{\sqrt{2} \rho^2} + \frac{i k^2 d_{+}}{\sqrt{2}} + \frac{i m d_{+}}{\sqrt{2} \rho^2} \right) e^{i(m-1)\phi} - \frac{\beta m d_z}{\rho} e^{im\phi} + \left( -\frac{i m^2 d_{-}}{\sqrt{2} \rho^2} + \frac{i k^2 d_{-}}{\sqrt{2}} - \frac{i m d_{-}}{\sqrt{2} \rho^2} \right) e^{i(m+1)\phi} \right] e^{i\beta z} J_m(\rho\rho') H_m^{(1)}(\rho\rho') + \sum_{m=-\infty}^{\infty} \frac{i m}{\sqrt{2} \rho} [-d_{+} e^{i(m-1)\phi} + d_{-} e^{i(m+1)\phi}] e^{i\beta z} \frac{\partial J_m(\rho\rho')}{\partial \rho} H_m^{(1)}(\rho\rho') \right\}.$$

We can obtain the  $E_{\phi}^{(0)}$  component associated with  $e^{im\phi+i\beta z}$  term as.

$$E_{\phi, m\beta}^{(0)} = \frac{1}{2} \left[ -\frac{i(m+1)^2 d_{+}}{\sqrt{2} \rho^2} + \frac{i k^2 d_{+}}{\sqrt{2}} + \frac{i(m+1) d_{+}}{\sqrt{2} \rho^2} \right] J_{m+1}(\rho\rho') H_{m+1}^{(1)}(\rho\rho') - \frac{i\beta m d_z^{(0)}}{2\rho} J_m(\rho\rho') H_m^{(1)}(\rho\rho') \\ + \frac{1}{2} \left[ -\frac{i(m-1)^2 d_{-}}{\sqrt{2} \rho^2} + \frac{i k^2 d_{-}}{\sqrt{2}} - \frac{i(m-1) d_{-}}{\sqrt{2} \rho^2} \right] J_{m-1}(\rho\rho') H_{m-1}^{(1)}(\rho\rho') \\ - \frac{1}{2} \left[ \frac{i(m+1) d_{+}}{\sqrt{2} \rho} \frac{\partial J_{m+1}(\rho\rho')}{\partial \rho} H_{m+1}^{(1)}(\rho\rho') - \frac{i(m-1) d_{-}}{\sqrt{2} \rho} \frac{\partial J_{m-1}(\rho\rho')}{\partial \rho} H_{m-1}^{(1)}(\rho\rho') \right] \\ = \frac{d_{+}}{2\sqrt{2} \rho} \left[ \frac{m(m+1)}{\rho^2} - k^2 \right] J_{m+1}(\rho\rho') H_{m+1}^{(1)}(\rho\rho') + \frac{d_{+} (m+1)}{2\sqrt{2} \rho} \frac{\partial J_{m+1}(\rho\rho')}{\partial \rho} H_{m+1}^{(1)}(\rho\rho') \\ - \frac{i\beta m d_z^{(0)}}{2\rho} J_m(\rho\rho') H_m^{(1)}(\rho\rho') \\ + \frac{d_{-}}{2\sqrt{2} \rho} \left[ \frac{m(m-1)}{\rho^2} - k^2 \right] J_{m-1}(\rho\rho') H_{m-1}^{(1)}(\rho\rho') - \frac{d_{-} (m-1)}{2\sqrt{2} \rho} \frac{\partial J_{m-1}(\rho\rho')}{\partial \rho} H_{m-1}^{(1)}(\rho\rho').$$

Now we can use the property that

$$\begin{cases} J_m(\rho\rho') = \frac{m+1}{\rho\rho'} J_{m+1}(\rho\rho') + \frac{\partial J_{m+1}(\rho\rho')}{\partial(\rho\rho')} \\ J_m(\rho\rho') = \frac{m-1}{\rho\rho'} J_{m-1}(\rho\rho') - \frac{\partial J_{m-1}(\rho\rho')}{\partial(\rho\rho')} \end{cases}$$

$$\Rightarrow E_{\phi, m\beta}^{(0)} = \frac{d_{+} (m+1)}{2\sqrt{2} \rho} \left[ \frac{(m+1)-1}{\rho\rho'} J_{m+1}(\rho\rho') + \frac{\partial J_{m+1}(\rho\rho')}{\partial(\rho\rho')} \right] H_{m+1}^{(1)}(\rho\rho') \\ - \frac{i\beta m d_z^{(0)}}{2\rho} J_m(\rho\rho') H_m^{(1)}(\rho\rho') \\ + \frac{d_{-} (m-1)}{2\sqrt{2} \rho} \left[ \frac{(m-1)+1}{\rho\rho'} J_{m-1}(\rho\rho') - \frac{\partial J_{m-1}(\rho\rho')}{\partial(\rho\rho')} \right] H_{m-1}^{(1)}(\rho\rho') \\ - \frac{d_{+} k^2}{2\sqrt{2} \rho} J_{m+1}(\rho\rho') H_{m+1}^{(1)}(\rho\rho') - \frac{d_{-} k^2}{2\sqrt{2} \rho} J_{m-1}(\rho\rho') H_{m-1}^{(1)}(\rho\rho') \\ = \frac{d_{+} (m+1)}{2\sqrt{2} \rho} \left[ \rho J_m(\rho\rho') - \frac{1}{\rho} J_{m+1}(\rho\rho') \right] H_{m+1}^{(1)}(\rho\rho') - \frac{i\beta m d_z^{(0)}}{2\rho} J_m(\rho\rho') H_m^{(1)}(\rho\rho') \\ + \frac{d_{-} (m-1)}{2\sqrt{2} \rho} \left[ \rho J_m(\rho\rho') + \frac{1}{\rho} J_{m-1}(\rho\rho') \right] H_{m-1}^{(1)}(\rho\rho') \\ - \frac{d_{+} k^2}{2\sqrt{2} \rho} J_{m+1}(\rho\rho') H_{m+1}^{(1)}(\rho\rho') - \frac{d_{-} k^2}{2\sqrt{2} \rho} J_{m-1}(\rho\rho') H_{m-1}^{(1)}(\rho\rho') \\ = \frac{1}{2} \left[ -\frac{\beta m d_z^{(0)}}{\rho} J_m(\rho\rho') H_m^{(1)}(\rho\rho') + \frac{i d_{+}}{\sqrt{2}} \left( k^2 + \frac{m+1}{\rho^2} \right) J_{m+1}(\rho\rho') - \frac{(m+1) \rho}{\rho} J_m(\rho\rho') \right] H_{m+1}^{(1)}(\rho\rho') \\ + \frac{i d_{-}}{\sqrt{2}} \left( k^2 - \frac{m-1}{\rho^2} \right) J_{m-1}(\rho\rho') - \frac{(m-1) \rho}{\rho} J_m(\rho\rho') \right] H_{m-1}^{(1)}(\rho\rho') ]$$

Compared to Nha's expression (Eqn. 28), our expression is equivalent if

$$\left\{ \begin{aligned} \frac{m+1}{\rho^2} J_{m+1}(\rho\rho') - \frac{\rho}{\rho} J_m(\rho\rho') &= 0 & (1) \\ \frac{m-1}{\rho^2} J_{m-1}(\rho\rho') - \frac{\rho}{\rho} J_m(\rho\rho') &= 0 & (2) \end{aligned} \right.$$

$$\Leftrightarrow \begin{cases} (m+1) J_{m+1}(\rho\rho') - \rho J_m(\rho\rho') = 0 \\ (m-1) J_{m-1}(\rho\rho') - \rho J_m(\rho\rho') = 0 \end{cases}$$

$$\Leftrightarrow \begin{cases} (m+1) J_{m+1}(x) - x J_m(x) = 0 \\ (m-1) J_{m-1}(x) - x J_m(x) = 0 \end{cases} \quad \text{but this is not true in general!}$$