

TraceWin

(V2)

D. Uriot

CEA Saclay
DSM/IRFU/SACM/LEDA
CEN Saclay 91191 Gif sur Yvette

Updated, 06th September 2010

Abstract

TRACE_WIN is a program calculating envelopes evolution of a bunched or DC beam through different transport elements. It's a linear matrix code including effects of linear space-charge forces. It has been first developed for proton linac, but now its vocation is much more general. It's a full linac-developing tool, including envelope calculation, multiparticle codes link, error studies with or without diagnostic and correction.

It has been completely written in C++ and Qt4.4 for Windows and Linux operating system. This choice makes it very easy to use. The user can change any parameter and observe the effect very easily due to the very powerful graphics display which allows to visualize most of the useful parameters of the simulation (envelopes, beam ellipses, emittances, phase advances...). All these outputs can be easily stored on disk, save in several image formats and inserted into reports (using copy and paste tools). Several matching options are available. Moreover, it generates different input files for multiparticle codes like PARTRAN, TOUTATIS from Saclay and directly run them and process their results.

You can find here, way of using, the way the beam dynamics has been modeled and gives detailed information required for a friendly use of the code and its outputs

[Installation & support](#)

[News in version 2](#)

[News and bugs](#)

[Ways of using](#)

[TraceWin files](#)

[Elements](#)

[Commands](#)

[Definitions used](#)

[Errors study](#)

Software installation & support

Technical support and code updates

This version of TRACE_WIN was written and is supported by **Didier URIOT** and **Nicolas PICHOFF**. We would appreciate hearing from you if you discover a bug or if you cannot find the information you need in this documentation to run the code. We will be pleased to try to solve your problem with the latest version of the codes. Your questions or remarks can be sent at the following Emails address:

didier.uriot@cea.fr

For copyright reasons the program is 6 months limited version. A new version can be asked at the same Email address.

Memory

The code requires at least 7.0 MB of free memory. Envelope charts needs a lot of memory, depending on the calculation step, then, it is not recommended to open too many charts at the same time especially with a very small calculation step if you don't have enough memory. I suggest 40 MB at least for an easy using.

Installation: No installation is necessary; all the extra files used by Trace_Win are directly extracted from main code and installed if a process requires it.

TRACEWIN definitions

[General description](#)

[Twiss parameters](#)

[Twiss parameters and acceleration](#)

[Conversions](#)

[Normalisation](#)

[Space charge](#)

[Transit time factor](#)

[Phase advance](#)

[Halo parameter](#)

TRACEWIN Elements

[Alpha magnet](#)

[Beam current](#)

[Beam rotation](#)

[Bending magnet](#)

[Bunched cavity or thin gap](#)

[Cavity multi-gap](#)
[Circular or rectangular aperture](#)
[Diagnostic elements](#)
[Drift](#)
[DTL cell](#)
[Edge angle on bending magnet](#)
[Electrostatic Acceleration](#)
[Electrostatic quadrupole](#)
[Field Map](#)
[Funneling gap](#)
[Multipole Field Map](#)
[RFQ cell](#)
[Thin lens](#)
[Thin matrix](#)
[Thin steering magnet](#)
[Sinus cavity or CCL](#)
[Solenoid](#)
[Space charge compensation](#)
[Quadrupole](#)

TRACEWIN Commands

Change element parameters:

[Chopper](#)
[Change structure frequency](#)
[Duplicate element](#)

[Magnetic steerer](#)
[Set RFQ coupling gap](#)
[Set RFQ electrode type](#)
[Set RFQ vane geometry](#)
[Shift](#)
[Structure file end](#)
[Superpose field map](#)

Lattice commands:

[Begin of lattice](#)
[End of lattice](#)
[Set phase advance](#)

Matching commands:

[Minimize beam envelope variation](#)
[Matching element commands](#)
[Minimize emittance growth](#)
[Minimize field variation](#)
[Minimize phase variation](#)
[Set achromatic line](#)
[Set centroid position](#)
[Set beam energy and phase](#)
[Set beam phase advance](#)
[Set beam phase error](#)
[Set beam separation](#)
[Set beam size](#)
[Set maximum beam size](#)
[Set synchronous phase](#)
[Set Twiss parameters](#)

Errors:

[Input beam errors](#)
[Bend errors](#)
[Cavity errors](#)
[RFQ errors](#)
[Quadrupole errors](#)

[Adjust commands](#)
[Gas pressure](#)
[Plot distribution](#)
[Read a multiparticle output file](#)
[Read a particle file](#)
[Set PARTRAN steps](#)

TRACEWIN files

[Data file](#) (*.dat)
[Init project file](#) (*.ini)
[Results file](#) (*.cal)
[Sigma0 file](#) (*.sig0)
[Input file for multiparticle program](#) (*.par, *.dat)
[Particle distribution](#) (*.dat)
[Particle losses](#) (*.loss)
[Error file results](#) (*.txt)
[Input & Output particle distribution](#) (*.dst, *.plt)
[Partran or Toutatis output](#) (*.out)
[Electric or magnetic field map](#)
[Current or space charge compensation map](#) (*.scc)
[Aperture map](#) (*.ouv)

Drift

| Mnemonic | Parameter | Definition |
|----------|-----------|---------------|
| DRIFT | L | Length (mm) |
| | R | Aperture (mm) |
| | R_y | Aperture (mm) |

If R_y equal 0 aperture is circular with R radius.

If R_y not equal 0 aperture is rectangular with R for x plane and R_y for Y plane.

[Drift matrix](#)

Quadrupole

| Mnemonic | Parameter | Definition |
|----------|-----------|---|
| QUAD | L | Length (mm) |
| | G | Magnetic field gradient (T/m) |
| | R | Aperture (mm) |
| | Θ | Skew Angle (°) |
| | G_3 | Sextupole gradient (T/m ²) |
| | G_4 | Octupole gradient (T/m ³) |
| | G_5 | Decapole gradient (T/m ⁴) |
| | G_6 | Dodecapole gradient (T/m ⁵) |

If L equal 0, quadrupole element are simulated as thin lens and all gradient components have to be replace by the gradient integral values.
 If L not equal to 0, multipole kicks are applied at the half of the quadrupole.

Remark: Short MULTIPOLE elements (based on field map) have to be replaced by QUAD element with multipole component. Quad element is much more efficient in terms of speed.

[Quadrupole matrix](#)

Beam Rotation

| Mnemonic | Parameter | Definition |
|----------|---------------|------------|
| BEAM_ROT | θ_{xy} | Angle (°) |
| | θ_{xz} | Angle (°) |
| | θ_{yz} | Angle (°) |
| | dx | X shift |
| | dy | Y shift |
| | dxp | Xp shift |
| | dyp | Yp shift |

Rotation are performed on after the other (in the order : XY, XZ, YZ) and finally the shift.

[Beam rotation matrix](#)

Thin Lens

| Mnemonic | Parameter | Definition |
|-----------|-----------|-------------------|
| THIN_LENS | f_x | Focal Length (mm) |
| | f_y | Focal Length (mm) |
| | R | Aperture (mm) |

[Thin lens matrix](#)

Thin Matrix

| Mnemonic | Parameter | Definition |
|-------------|----------------------|---------------------------|
| THIN_MATRIX | lg | Length (mm) |
| | a_{00} to a_{55} | Matrix terms, row per row |

The 36 terms of matrix have to be set row by row from a00 to a55. Length, Lg, is just used in graphic view.

[See matrix format R](#)

The units are respectively: x (m), x' (rad), y (m), y' (rad), z (m), dp/p (rad)

Quadrupole Example: **THIN_MATRIX 10.0 1 0 0 0 0 0 -2.5 1 0 0 0 0 0 0 1. 0 0 0 0 0 0 2.5 1 0 0 0 0 0 0 1 0 0 0 0 0 1**

Electrostatic Quadrupole

| Mnemonic | Parameter | Definition |
|----------|-----------|--------------------------------|
| QUAD_ELE | L | Length (mm) |
| | V_o | Voltage between electrodes (V) |
| | R | Aperture (mm) |

| | | |
|--|----------|--|
| | θ | Skew Angle (°) |
| | V_3 | Sextupole voltage component (V/m) |
| | V_4 | Octupole voltage component (V/m ²) |
| | V_5 | Decapole voltage component (V/m ³) |
| | V_6 | Dodecapole voltage component (V/m ⁴) |

V_0 between electrodes means $+V_0/2$ on one electrode and $-V_0/2$ on the near electrode for a total voltage of V_0 .

In order to keep coherence with old TraceWin version, all electrostatic components are in voltage: (V) and not in field: E (V/m)

If L equal 0, electrostatic quadrupole element are simulated as thin lens and all voltage components have to be replace by the voltage integral values.

If L not equal to 0, multipole kicks are applied at the half of the electrostatic quadrupole.

[Quadrupole matrix](#)

Bunched cavity or thin gap

| Mnemonic | Parameter | Definition |
|----------|-------------|--|
| GAP | $E_0 TL$ | Effective gap voltage (V) |
| | θ_s | RF phase (deg) (absolute or relative) |
| | R | Aperture (mm) |
| | P | 1: θ_s is absolute phase, 0: θ_s is relative |
| | β_s | Particle reduced velocity |
| | T_s | Transit time factor |
| | kT'_s | (*) |
| | $k^2 T''_s$ | (*) |
| | kS' | (*) |
| | $k^2 S''$ | (*) |

Red parameters are optional

(*) See [Transit time factor definition](#)

[Bunched cavity or thin gap matrix](#)

Sinus cavity or CCL

| Mnemonic | Parameter | Definition |
|----------|------------|--|
| CAVSIN | L | Length (mm) |
| | N | Cell number |
| | $E_0 T$ | Average accelerating field (V/m) |
| | θ_s | Phase of the synchronous particle at the entrance (deg) |
| | r | Aperture (mm) |
| | p | 1: θ_s is absolute phase, 0: θ_s is relative |
| | | |

[Sinus cavity or CCL matrix](#)

Bending magnet

| Mnemonic | Parameter | Definition |
|----------|-----------|---|
| BEND | α | Bend angle in the rotation plane (deg) |
| | $ \rho $ | Curvature radius of central trajectory (mm) |
| | n | Field gradient index |
| | r | Aperture (mm) |

| | | |
|--|-----------|------------------------------|
| | <i>HV</i> | 0 : horizontal, 1 : vertical |
|--|-----------|------------------------------|

Warning: A bend always have to be preceded and followed by edge element, even if edge is set to zero.

A positive bend (denoted $\alpha > 0$) bends the particles to the right in the horizontal plane, regardless of the sign of the charge.

The bend with field gradient index not equal to zero are only treated in the first order and for horizontal bending

[Bending magnet matrix](#)

Edge angle on bending magnet

| Mnemonic | Parameter | Definition |
|-------------|-----------|--------------------------------------|
| EDGE | β | Pole face rotation angle (deg) |
| | $ \rho $ | Curvature radius of bend (mm) |
| | g | Total gap of magnet (mm) |
| | $K1$ | Fringe-field factor (default = 0.45) |
| | $K2$ | Fringe-field factor (default = 2.80) |
| | r | Aperture (mm) |
| | <i>HV</i> | 0 : horizontal, 1 : vertical |

By definition: An edge focalise the transverse x plane if the pole $\beta < 0$, whatever the curvature radius sign or the bend angle sign and the particle charge.

Set $K1$ and $K2$ to a very small value to disable fringe field estimation.

The gap of magnet is used for particle loss estimation in bend.

[Edge angle on bending magnet matrix](#)

Circular or rectangular aperture

| Mnemonic | Parameter | Definition |
|-----------------|-----------|--|
| APERTURE | dx | X half width (mm) or radius hole |
| | dy | Y half width (mm) or distance between hole |
| | N | 0 : Rectangular aperture 1 : Circular aperture 2 : Perperpot mode 3 : Rectangular aperture with dx & dy corresponding to a beam fraction. (aperture is adjusted with 0.1 mm step) 4 : Horizontal finger with dx =total finger width and dy =finger position. 5 : vertical finger with dx =total finger width and dy =finger position. |
| | | |

Space charge compensation

| Mnemonic | Parameter | Definition |
|--------------------------|-----------|---|
| SPACE_CHARGE_COMP | k | Beam current is compensated by a factor k |

Beam current

| Mnemonic | Parameter | Definition |
|----------------|-----------|-------------------|
| CURRENT | Ib | Beam current (mA) |

Ib (mA) is the new current beam.

Solenoid

| Mnemonic | Parameter | Definition |
|----------|-----------|--------------------|
| SOLENOID | L | Length (mm) |
| | B | Magnetic field (T) |
| | R | Aperture (mm) |
| | | |

[Solenoid matrix](#)

Thin steering magnet

| Mnemonic | Parameter | Definition |
|---------------|-----------|-------------------|
| THIN_STEERING | BL_x | x-component (T.m) |
| | BL_y | y-component (T.m) |
| | r | Aperture (mm) |
| | | |

[Thin steering magnet matrix](#)

DTL cell

| Mnemonic | Parameter | Definition |
|----------|------------|--|
| DTL_CEL | L | Cell length (mm) |
| | $Lq1$ | First ½ quadrupole length (mm) |
| | $Lq2$ | Second ½ quadrupole length (mm) |
| | g | Cell center (mm) |
| | $B1'$ | First magnetic field gradient (T/m) |
| | $B2'$ | Second magnetic field gradient (T/m) |
| | E_{oTL} | Effective gap voltage (V) |
| | θ_s | RF phase (deg) |
| | R | Aperture (mm) |
| | p | 1: θ_s is absolute phase, 0: θ_s is relative |
| | β_s | Particle reduced velocity |
| | T_s | Transit time factor |
| | kT'_s | (*) |
| | $k^2T''_s$ | (*) |
| | | |
| | | |
| | | |

Red parameters are optional

If β_s is not equal to 0 T_s become needed

(*) See [Transit time factor definition](#)

[DTL cell matrix](#)

Copy following examples:

[examples\DTL_1.ini](#)

[examples\DTL_1.dat](#)

[examples\DTL_1.cal](#)

[examples\DTL_2.ini](#)

[examples\DTL_2.dat](#)

[examples\DTL_2.cal](#)

[examples\DTL_3.ini](#)

[examples\DTL_3.dat](#)

[examples\DTL_3.cal](#)

[examples\DTL_4.ini](#)
[examples\DTL_4.dat](#)
[examples\DTL_4.cal](#)

Cavity multi-gap

| Mnemonic | Parameter | Definition |
|----------|-------------------------------------|---|
| NCELLS | <i>Mode</i> | (0)2 π , (1) π , (2) π & 2 π |
| | <i>N_c</i> | Number of cell |
| | <i>β_s</i> | Geometric β |
| | <i>E_oT</i> | Effective gap voltage (V/m) |
| | <i>θ</i> | RF phase at the first gap position (deg) |
| | <i>R</i> | Aperture (mm) |
| | <i>P</i> | 1: θ is absolute phase, 0: θ is relative |
| | <i>kE_oT_i</i> | Input field correction, E _o = E _o *(1+k) |
| | <i>kE_oT_o</i> | Output field correction, E _o = E _o *(1+k) |
| | <i>dz_i</i> | First gap displacement (mm) |
| | <i>dz_o</i> | Last gap displacement (mm) |
| | <i>β_s</i> | Particle reduced velocity |
| | <i>T_s</i> | Transit time factor of the middle gaps |
| | <i>kT'_s</i> | (*) |
| | <i>k²T''_s</i> | (*) |
| | <i>T_i</i> | Transit time factor of the input gaps |
| | <i>kT'_i</i> | (*) |
| | <i>k²T''_i</i> | (*) |
| | <i>T_o</i> | Transit time factor of the output gaps |
| | <i>kT'_o</i> | (*) |
| | <i>k²T''_o</i> | (*) |

Red parameters are optional

If β_s is not equal to 0 *T_s* become needed

(*) See [Transit time factor definition](#)

Cavity multi-gap matrix

Copy following examples:

[examples\Supra_1.ini](#)
[examples\Supra_1.dat](#)
[examples\Supra_1.cal](#)

[examples\Supra_2.ini](#)
[examples\Supra_2.dat](#)
[examples\Supra_2.cal](#)

[examples\Supra_3.ini](#)
[examples\Supra_3.dat](#)
[examples\Supra_3.cal](#)

[examples\Supra_4.ini](#)
[examples\Supra_4.dat](#)
[examples\Supra_4.cal](#)

[examples\Supra_5.ini](#)
[examples\Supra_5.dat](#)
[examples\Supra_5.cal](#)

RFQ cell

| Mnemonic | Parameter | Definition |
|----------|-----------------------|---------------------------|
| RFQ_CELL | <i>V</i> | Effective gap voltage (V) |
| | <i>R_o</i> | Vane radius (mm) |
| | <i>A₁₀</i> | Acceleration parameter |
| | <i>m</i> | Modulation |
| | <i>L</i> | Length (mm) |

| | | |
|--|------------|-------------------------------|
| | θ_s | RF phase (deg) |
| | Type | |
| | Tc | Transverse curvature (mm) |
| | Aoi | Transverse focusing parameter |

[RFQ cell matrix](#)

See following examples:

[examples\RFQ_1.ini](#)

[examples\RFQ_1.dat](#)

[examples\RFQ_1.cal](#)

[examples\RFQ_2.ini](#)

[examples\RFQ_2.dat](#)

[examples\RFQ_2.cal](#)

Diagnostic elements

| Diagnostic type | Mnemonic | Parameter | Definition |
|--|------------------------|--|---|
| Current measurement | DIAG_CURRENT | N Ib | Diagnostic number Wanted beam current (mA) |
| Positions measurement | DIAG_POSITION | N X Y dm | Diagnostic number Wanted X beam position (mm), if $X < 1e50$ Wanted Y beam position (mm), if $Y < 1e50$ Diagnostic Accuracy (mm) |
| Divergences measurement | DIAG_DIVERGENCE | N X' Y' dm | Diagnostic number Wanted X rms beam divergence (mrad) Wanted Y rms beam divergence (mrad) Diagnostic Accuracy (mrad) |
| Size measurement | DIAG_SIZE | N Sx Sy ΔP Dm $d\Delta P$ fo | Diagnostic number Wanted X rms beam size (mm) Wanted Y rms beam size (mm) Wanted rms Phase spread (°) Size Accuracy (mm) Phase spread Accuracy (°) ΔP Low-pass filter frequency (MHz) |
| Divergence measurement (*) if S<0 Twiss parameter alpha will set <0 | DIAG_SIZEP | N $Sx' (*)$ $Sy' (*)$ ΔW Dm dW | Diagnostic number Wanted X' rms beam divergence (mrad) Wanted Y' rms beam divergence (mrad) Wanted rms Energy spread (MeV) Divergence Accuracy (mrad) Energy spread Accuracy (%) |
| Delta size Measurement | DIAG_DSIZE | N $x_n - y_n$ dm | Diagnostic number Wanted x-y rms beam delta size (mm) dx-y size Accuracy (mm) |
| Delta size Measurement 2 (At least 2 are needed) | DIAG_DSIZE2 | N $x_n - x_{n-1}$ $y_n - y_{n-1}$ dm | Diagnostic number Wanted x rms beam delta size (mm) Wanted y rms beam delta size (mm) dx & dy size Accuracy (mm) |
| Delta phase spread Measurement 2 (At least 2 are Needed) | DIAG_DSIZE3 | N $\Delta P_n - \Delta P_{n-1}$ $d\Delta P$ fo | Diagnostic number Wanted rms delta phase spread (°) Phase spread accuracy (°) $d\Delta P$ Low-pass filter frequency (MHz) |
| Phase measurement | DIAG_PHASE | N Θ | Diagnostic number Wanted centroid Phase (°) |

| | | | |
|--|------------------------|--|--|
| Beam energy - Perfect linac energy measurement | DIAG_DENERGY | N W | Diagnostic number Wanted Energy (MeV) |
| Beam phase - Perfect linac phase measurement | DIAG_DPHASE | N Θ | Diagnostic number Wanted centroid Phase (°) |
| Luminosity | DIAG_LUMINOSITY | N Lu Dlu | Diagnostic number Wanted luminosity (mm ⁻²) Luminosity accuracy (mm ⁻²) |
| Waist setting | DIAG_WAIST | N f_x f_y d_{xy} | Diagnostic number XX' waist asked (for f_x not equal to 0) YY' waist asked (for f_y not equal to 0) Transverse waist Accuracy |
| Achromat setting | DIAG_ACHROMAT | N n $f1$ $f2$ | Diagnostic number First element number If =1 set achromatic position If =1 set achromatic angle |
| Emittance setting (rms values) | DIAG_EMIT | N Exx' Eyy' E_{pw} | Diagnostic number Wanted emittance if greater than 0 Wanted emittance if greater than 0 Wanted emittance if greater than 0 |
| Emittance setting (99% value) | DIAG_EMIT_99 | N Exx' Eyy' E_{pw} | Diagnostic number Wanted emittance if greater than 0 Wanted emittance if greater than 0 Wanted emittance if greater than 0 |
| Transfer matrix setting | DIAG_SET_MATRIX | N $N^{°ele}$ $Row(i)$ $Column(j)$ k M_{ij} | Diagnostic number Transfer matrix form N°ele to diag. position Row transfer matrix term Column transfer matrix term Corrector coefficient Wanted transfer matrix term value |
| Beam Twiss parameters setting | DIAG_TWISS | $\alpha_{xx'}$ $\beta_{xx'}$ $\alpha_{yy'}$ $\beta_{yy'}$ $\alpha_{zdp/p}$ $\beta_{zdp/p'}$ | Wanted alp_X-X' Wanted bet_X-X' (mm/mrad) Wanted alp_Y-Y' Wanted bet_Y-Y' (mm/mrad) Wanted alp_Z-dp/p Wanted bet_Z-dp/p (mm/mrad) |

TRACE_WIN is able to verify if a beam line contains enough diagnostic elements to control the beam and correct the errors coming from the input beam or from the different element errors. In order to put one diagnostics in the line use the elements “*DIAG_...*” followed by the diagnostic number and the parameters which have to be imposed at the diagnostic location (current, position, size or emittance). To choose the different elements, which have to be adjusted to control the beam, use the [adjust commands](#) followed by the number of the associated diagnostic.

There are two ways of using these diagnostics. The first one is for example to see if your design is able to control the beam position at a given location when you input beam is not at the centre or if a few elements, like steerers induce misalignment. The second way is to use diagnostics associated with a error study in order to see if you scheme of misalignment correction for example is good or not.

The diagnostics adjustments are independent process, which occurs after the matching process. If any diagnostic element is present any adjustment is started. Finally, at the end of a line design process, only diagnostic elements should be used to tuned it.

Warning : They are treated in the order of the diagnostic number.

[Adjust and diagnostic examples](#)

New (23-12-2009): The diagnostic influence can be weighted by putting a correcting factor between brackets like following:

DIAG_EMIT(1e-4) 1 0.1 0.1 0.2 , here this diagnostic emittance criteria is reduced by a factor 10000.

By default, without bracket, the correction is set to 1.

Copy following examples:

[examples\Supra_2.ini](#)

[examples\Supra_2.dat](#)

[examples\Supra_2.cal](#)

Funneling gap

| | Mnemonic | Parameter | Definition |
|---------------|-------------------|-------------------------------------|---|
| Funneling Gap | FUNNEL_GAP | E_oTL θ_s R p | Effective gap voltage (V) RF phase (deg) (absolute or relative) Aperture (mm) 1: θ_s is absolute phase, 0: θ_s is relative |
| Frame change | CHFRAME | X_o X'_o | (mm) (deg) |

[Funneling gap matrix](#)

Alpha magnet

| Mnemonic | Parameter | Definition |
|---------------------|----------------------------------|---|
| ALPHA_MAGNET | θ K R $plan$ | Entrance angle (°) (T/m) Aperture (mm) 0 (x)/1 (y) |

[Alpha magnet matrix](#)

Electrostatic Acceleration

| Mnemonic | Parameter | Definition |
|-----------------------|-------------------|---|
| ELECTROSTA_ACC | V_o | Voltage (V) |
| | L K R | Length (mm) Transverse defocal (eV/ mm ²) Aperture (mm) |

[Electrostatic Acceleration matrix](#)

Field Map

| | | |
|--|--|--|
| | | |
|--|--|--|

| Mnemonic | Parameter | Definition |
|------------------|------------------------------|----------------------------------|
| FIELD_MAP | <i>geom</i> | Field map type |
| | <i>L</i> | Field map length (mm) |
| | <i>θ_i</i> | RF input field phase (°) |
| | <i>R</i> | Aperture (mm) |
| | <i>kb</i> | Magnetic field intensity factor |
| | <i>ke</i> | Electric field intensity factor |
| | <i>Ki</i> | Space charge compensation factor |
| | <i>Ka</i> | Aperture flag |
| | <i>FielName</i> | File name without extension |

File name is the field map name without extension and without path if files is en the data file (*.dat) path, otherwise put the full file name with path without extension

Equations du transport dans un champ électromagnétique

“geom” parameter defines the field map type as following:

Unit figure : Static electric field
Tens figure : Static magnetic field
Hundreds figure : RF electric field
Thousands figure : RF magnetic field

0 - No field
1 - 1D : Fz(z)
2 - Not available
3 - Not available
4 - 2D cylindrical static or RF electric field : Ez(r,z), Er(r,z) and B θ (r,z) for RF
5 - 2D cylindrical static or RF magnetic field : Bz(r,z), Br(r,z) and E θ (r,z) for RF
6 - 2D Cartesian field: Fx(x,y), Fy(x,y)
7 - 3D Cartesian field: Fx(x,y,z), Fy(x,y,z), Fz(x,y,z)
8 - 3D cylindrical field : Fr(r, θ , z), F θ (r, θ ,z), Fz(r, θ ,z) Not implemented yet

θ_i : RF phase when the generator particle enters the cavity.
The RF field is E=E0.cos(θ), B=-B0.sin(θ)

FileName: is the file name without extension where are localized the field maps (.bsz, .bsr, .edx ...), the space charge compensation map (.scc) and the beam pipe radius evolution with z (.ouv). See [the file formats description](#)

Field in 1D: The field can be described in 1D, Field according to z position

- 1 file contains the static electric field : filename.esz
- 1 file contains the static magnetic field : filename.bsz
- 1 file contains the RF electric field : filename.edz
- 1 file contains the RF magnetic field : filename.bdz

Field in 2D: The fields can be described in 2D Cartesian (invariant through translation on z axis) or in cylindrical (invariant through rotation around the z axis) coordinates.

In Cartesian coordinates (x,y)

- 2 files contain the static electric field : filename.esx, filename.esy
- 2 files contain the static magnetic field : filename.bsx, filename.bsy
- 2 files contain the RF electric field : filename.edx, filename.edy
- 2 files contain the RF magnetic field : filename.bdx, filename.bdy

In cylindrical coordinates (r,z, θ)

- 2 files contain the static electric field : filename.esr, filename.esz
- 2 files contain the static magnetic field : filename.bsr, filename.bsz
- 3 files contain the RF fields (TM) : filename.edr, filename.edz, filename.bdq
- 3 files contain the RF fields (TE) : filename.bdr, filename.bdz, filename.edq

Field in 3D:

The fields can be described in 3D either Cartesian or cylindrical frame.

In Cartesian coordinates (x,y,z)

- 3 files contain the static electric field : filename.esx, filename.esy, filename.esz
- 3 files contain the static magnetic field : filename.bsx, filename.bsy, filename.bsz
- 3 files contain the RF electric field : filename.edx, filename.edy, filename.edz
- 3 files contain the RF magnetic field : filename.bdx, filename.bdy, filename.bdz

In cylindrical coordinates (r, θ , z)

- 3 files contain the static electric field : filename.esr, filename.esq, filename.esz
- 3 files contain the static magnetic field : filename.bsr, filename.bsq, filename.bsz
- 3 files contain the RF electric field : filename.edr, filename.edq, filename.edz

- 3 files contain the RF magnetic field : filename.bdr, filename.bdq, filename.bdz

ki: If it is greater than 0, a space charge compensation map or current map has to be created in “FileName.scc” and ki is the normlization factor.

[Space charge compensation file syntax](#)

ka: If it is equal to 1, a beam pipe radius map has to be created in “FileName.ouv”.

[Aperture map file syntax](#)

Losses consideration:

Particle is lost in FIELD_MAP when it comes out of aperture defined in element syntax or field_map frame size defined in field_map files. This behavior is true except when aperture flag is set to 1 in order to use beam pipe radius file map. Now, if its flag is set to 2, no pipe file is required, but particle will not be lost when it comes out field map frame. That allows superpose small size field map inside bigger one.

All these files have to be located in the directory of the main data file.

Copy following example:

[examples\LBE_1.ini](#)

[examples\LBE_1.dat](#)

[examples\LBE_1.cal](#)

[examples\solenoid1.bsz](#)

[examples\solenoid1.ouv](#)

[examples\solenoid2.bsz](#)

[examples\solenoid2.ouv](#)

Multipole Field Map

| Mnemonic | Parameter | Definition |
|-----------|--------------|---|
| MULTIPOLE | <i>Order</i> | Multipole order |
| | <i>L</i> | Field map length (mm) |
| | <i>step</i> | Number of step along x & y direction |
| | <i>B</i> | Magnetic field on pole (T) |
| | <i>R</i> | Aperture (mm) |
| | <i>Lsol</i> | Physical length (mm) of solenoid (<i>Order=0</i>) |
| | <i>Zstep</i> | Step number for solenoid case (<i>Order=0</i>) |

Multipole element generates a static magnetic field map file where the step along x & y direction is defined as 2.R/step. Warning, 2.R/step must be much lower than the beam size. But, the field map size depends on *step* parameter.

And the step along z direction (beam) is defined by TraceWin calculation step.

Order = 0: (Special mode) for Solenoid field map (Br(r,z) & Bz(r,z))

Order = 1: for dipole

Order = 2: for quadrupole

Order = 3: for sextupole

Order = 4: for octupole

...

Funneling gap ($EoTL$, ϕ_s)

$EoTL$ (V) is the maximum energy gain, ϕ_s (°) is the synchronous phase.

$$\beta_z = \sqrt{\frac{1 - \frac{1}{\gamma_i^2}}{1 + x'^2 + y'^2}}$$

$$K = |q| E_0 TL \cdot \cos \varphi$$

$$x' = x + \frac{K}{\beta_z^2 \sqrt{(\gamma_i \cdot mc^2)^2 + \frac{K^2}{\beta_z^2}}}$$

Where x' being the horizontal beam centroid slope.

$$R_{xx} = R_{yy} = R_{zz} = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \quad R_{xz} = \begin{bmatrix} 0 & 0 \\ \frac{2\pi |q| E_0 TL \cdot \sin \varphi_s}{\gamma_i \beta_z^3 \cdot \lambda \cdot mc^2} & 0 \end{bmatrix}$$

Drift (Δs)

Δs (mm) is the drift length:

The non null 2x2 transfer sub-matrixes are:

$$R_{xx} = R_{yy} = \begin{bmatrix} 1 & \Delta s \\ 0 & 1 \end{bmatrix}, \quad \text{and} \quad R_{zz} = \begin{bmatrix} 1 & \frac{\Delta s}{\gamma^2} \\ 0 & 1 \end{bmatrix}.$$

Thin lens (f_x, f_y)

f_x, f_y are focal length in meter

The non null 2x2 transfer sub-matrixes are:

$$R_{xx} = \begin{bmatrix} 1 & 0 \\ \frac{-1}{f_x} & 1 \end{bmatrix}, \quad R_{yy} = \begin{bmatrix} 1 & 0 \\ \frac{-1}{f_y} & 1 \end{bmatrix}, \quad \text{and} \quad R_{zz} = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}.$$

Beam rotation ($\theta_{xy}, \theta_{xz}, \theta_{zy}, dx, dy, d_{xp}, d_{yp}$)

θ_{xy} are the beam rotation XY angle in degree

The ellipsoid can be brought upright by rotations $-\theta_{xy}$ accomplished by applying the transfer matrixes:

$$R_{xy} = \left[\begin{array}{cc|cc|cc} \cos(\theta_{xy}) & 0 & \sin(\theta_{xy}) & 0 & 0 & 0 \\ 0 & \cos(\theta_{xy}) & 0 & \sin(\theta_{xy}) & 0 & 0 \\ \hline -\sin(\theta_{xy}) & 0 & \cos(\theta_{xy}) & 0 & 0 & 0 \\ 0 & -\sin(\theta_{xy}) & 0 & \cos(\theta_{xy}) & 0 & 0 \\ \hline 0 & 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 \end{array} \right],$$

The rotation is then applied:

$$[\sigma] = R_{xy} \cdot [\sigma] \cdot R_{xy}^T,$$

Quadrupole (Δs , G)

Δs (mm) is the quadrupole length, G (T/m) its gradient.

For electrostatic quadrupole $G = \frac{V_0}{\beta_c R^2}$, where V_0 is the voltage between electrode, β_c beam speed and R is the half distance between electrode.

Lets use $k = \sqrt{\frac{G}{B\rho}}$, with $B\rho = \frac{m_0 c \beta \gamma}{q}$, the magnetic rigidity of the particle.

If $q \cdot G$ is positive, the quadrupole is focusing in the horizontal direction, else it's defocusing.

The non null 2x2 transfer sub-matrixes are:

$$R_{zz} = \begin{bmatrix} 1 & \frac{\Delta s}{\gamma^2} \\ 0 & 1 \end{bmatrix}.$$

In the longitudinal direction, we have:

In the transverse direction, two possibilities:

□ Focusing quadrupole (in horizontal direction, $q \cdot G > 0$)

$$R_{xx} = \begin{bmatrix} \cos(k\Delta s) & \frac{\sin(k\Delta s)}{k} \\ -k \sin(k\Delta s) & \cos(k\Delta s) \end{bmatrix}, \quad R_{yy} = \begin{bmatrix} \cosh(k\Delta s) & \frac{\sinh(k\Delta s)}{k} \\ k \sinh(k\Delta s) & \cosh(k\Delta s) \end{bmatrix}.$$

□ Defocusing quadrupole (in horizontal direction, $q \cdot G < 0$)

$$R_{xx} = \begin{bmatrix} \cosh(k\Delta s) & \frac{\sinh(k\Delta s)}{k} \\ k \sinh(k\Delta s) & \cosh(k\Delta s) \end{bmatrix}, \quad R_{yy} = \begin{bmatrix} \cos(k\Delta s) & \frac{\sin(k\Delta s)}{k} \\ -k \sin(k\Delta s) & \cos(k\Delta s) \end{bmatrix}.$$

Bunched cavity or thin gap ($E_0 TL$, ϕ_s , p , β_s , T_s , kT 's, $k^2 T$'s, kS ', $k^2 S$ '")

$E_0 TL$ (eV) is the maximum energy gain, ϕ_s (°) is the synchronous phase.

The reduced energy change in the gap is:

$$\gamma_s = \gamma_e + \frac{|q| E_0 TL \cos(\phi_s)}{mc^2}.$$

The phase shift in the gap is:

$$\Delta\phi = \frac{q E_0 TL \cdot \sin(\phi_s)}{mc^2 \cdot \gamma^2 \cdot \beta} \left(\frac{kT'}{T} \right)$$

If $\beta_s = 0$: $\Delta\phi = 0$ else

The changes in the normalized momentum caused by the gap are given by:

$$k_{xy} = \frac{-q\pi E_0 TL \sin(\varphi_s)}{mc^2 \beta^2 \cdot \gamma^2 \lambda}, \quad \text{With } \bar{\gamma} = \frac{\gamma_e + \gamma_s}{2} \text{ and } \bar{\beta} = \frac{\beta_e + \beta_s}{2}.$$

$$k_z = \frac{2q\pi E_0 TL \sin(\varphi_s)}{mc^2 \beta^2 \lambda}.$$

The non null 2x2 transfer sub-matrixes are:

$$R_{xx} = \begin{bmatrix} k_x \cdot C & 0 \\ k_{xy} & (\beta\gamma)_e \\ (\beta\gamma)_s & (\beta\gamma)_s \end{bmatrix}, \quad R_{yy} = \begin{bmatrix} k_y C & 0 \\ k_{xy} & (\beta\gamma)_e \\ (\beta\gamma)_s & (\beta\gamma)_s \end{bmatrix}, \quad R_{zz} = \begin{bmatrix} 1 & 0 \\ k_z & (\beta\gamma)_e \\ (\beta\gamma)_s & (\beta\gamma)_s \end{bmatrix}.$$

$$\text{- If } \beta_s = 0: \quad k_x = k_y = 1 - \frac{q E_0 TL \cos(\varphi_s)}{2mc^2 \beta^2 \cdot \bar{\gamma}}$$

$$\text{- If } \beta_s \neq 0: \quad k_x = \frac{-q E_0 T L \cos(\varphi_s)}{2mc^2 \cdot \gamma^3 \cdot \beta^2} \left(\bar{\gamma}^2 + \frac{kT'}{T} \right), \quad k_y = \frac{-q E_0 T L \cos(\varphi_s)}{2mc^2 \cdot \gamma^3 \cdot \beta^2} \left(\bar{\gamma}^2 - \frac{kT'}{T} \right)$$

(*) See [Transit time factor definition](#)

C: is a coefficient allowing to keep the matrix determinant equal to:

$$\frac{(\beta\gamma)_e}{(\beta\gamma)_s} : C = \sqrt{\frac{(\beta\gamma)_e}{(\beta\gamma)_s} \frac{1}{k_x k_y}}$$

Solenoid (Δs , B)

Δs (mm) is the solenoid length, B (T) its axis magnetic field.

Lets use $k = \frac{B}{2B\rho}$, with $B\rho = \frac{m_0 c \beta \gamma}{q}$, the magnetic rigidity of the particle.

The non null 2x2 transfer sub-matrixes are:

$$R_{xx} = R_{yy} = \begin{bmatrix} \cos^2(k\Delta s) & \frac{\sin(k\Delta s) \cos(k\Delta s)}{k} \\ -k \sin(k\Delta s) \cos(k\Delta s) & \cos^2(k\Delta s) \end{bmatrix},$$

$$R_{xy} = -R_{yx} = \begin{bmatrix} \sin(k\Delta s) \cos(k\Delta s) & \frac{\sin^2(k\Delta s)}{k} \\ -k \sin^2(k\Delta s) & \sin(k\Delta s) \cos(k\Delta s) \end{bmatrix},$$

$$R_{zz} = \begin{bmatrix} 1 & \frac{\Delta s}{\gamma^2} \\ 0 & 1 \end{bmatrix}.$$

Bending magnet ($\Delta\alpha$, $|\rho|$, n , ouv , HV)

$\Delta\alpha$ ($^\circ$) is the rotation angle, ρ (mm) is the curvature radius, n is the field index, $ouv(mm)$ the aperture and HV means: horizontal (=0) or vertical bend (=1).

A positive bend ($\alpha > 0$) bends the particles to the right in the horizontal plane, regardless of the sign of the charge on the particle. A negative α bends particles to the left.

$$h = \frac{1}{|\rho|} \frac{\Delta\alpha}{|\Delta\alpha|}, \quad k_x = \sqrt{(1-n) \cdot h^2}, \quad k_y = \sqrt{n \cdot h^2}, \quad \Delta s = |\rho| \Delta\alpha.$$

The non null 2x2 transfer sub-matrixes are:

$$R_{xx} = \begin{bmatrix} \cos(k_x \Delta s) & \frac{\sin(k_x \Delta s)}{k_x} \\ -k_x \sin(k_x \Delta s) & \cos(k_x \Delta s) \end{bmatrix}, \quad R_{yy} = \begin{bmatrix} \cos(k_y \Delta s) & \frac{\sin(k_y \Delta s)}{k_y} \\ -k_y \sin(k_y \Delta s) & \cos(k_y \Delta s) \end{bmatrix},$$

$$R_{zz} = \begin{bmatrix} 1 & \frac{-h^2(k_x \Delta s \beta^2 - \sin(k_x \Delta s))}{k_x^3} + \frac{\Delta s}{\gamma^2} \left(1 - \frac{h^2}{k_x^2}\right) \\ 0 & 1 \end{bmatrix},$$

$$R_{zx} = \begin{bmatrix} \frac{-h \sin(k_x \Delta s)}{k_x} & \frac{-h(1 - \cos k_x \Delta s)}{k_x^2} \\ 0 & 0 \end{bmatrix}, \quad R_{xz} = \begin{bmatrix} 0 & \frac{h(1 - \cos k_x \Delta s)}{k_x^2} \\ 0 & \frac{h \sin(k_x \Delta s)}{k_x} \end{bmatrix}.$$

Edge angle on bending magnet (β , $|\rho|$, g , $K1$, $K2$, ouv , HV)

β ($^\circ$) is the edge angle, ρ (mm) is the curvature radius in the bending magnet, g (mm) is the gap between the poles of the bending magnet, $K1$ and $K2$ are used in a development for the fringe-field correction. If they are equal to zero, $K1 = 0.45$ and $K2 = 2.8$. Set small values to cancel fringe-field correction.

ouv (mm) is the aperture and HV means: horizontal (=0) or vertical bend (=1).

The edge angle is treated as a thin lens. Ψ is the fringe-field correction.

$$\Psi = K_1 \frac{g}{|\rho|} \left(\frac{1 + \sin^2(\beta)}{\cos(\beta)} \right) \cdot \left(1 - K_1 K_2 \frac{g}{|\rho|} \tan(\beta) \right).$$

The non null 2x2 transfer sub-matrixes are:

$$R_{xx} = \begin{bmatrix} 1 & 0 \\ \frac{\tan(\beta)}{|\rho|} & 1 \end{bmatrix}, \quad R_{yy} = \begin{bmatrix} 1 & 0 \\ -\frac{\tan(\beta - \Psi)}{|\rho|} & 1 \end{bmatrix} \quad \text{And} \quad R_{zz} = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}.$$

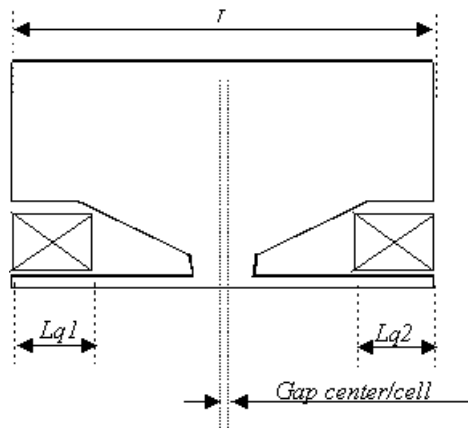
Thin steering magnet (BL_x , BL_y , r)

$$x' = x + \frac{qBL_y \cdot c}{mc^2 \beta \gamma} \quad \text{And} \quad y' = y - \frac{qBL_x \cdot c}{mc^2 \beta \gamma}$$

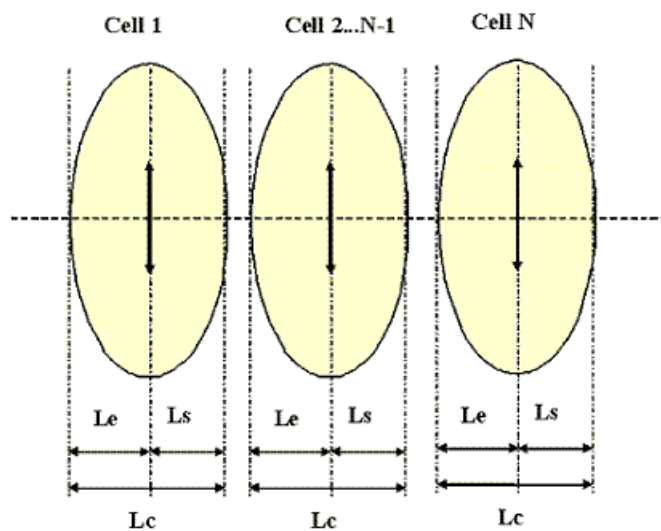
Where x' and y' being respectively the horizontal and vertical beam centroid slope.

DTL cell (L, Lq1, Lq2, g, B1', B2', EoTL, θ_s , r, p, β_s , Ts, kT's, k²T"s)

The dimension $g(mm)$ is defined as: $gap\ position = \frac{L}{2} - g - Lq_1$.



Cavity multi-gap (m, N, β_g , EoT, θ_s , r, p, kEoTi, kEoTo, Dzi, Dzo, β_s , Ts, kT's, k²T"i, Ts, kT'i, k²T"i, To, kT'o, k²T"o)



If $m=0$:

For cell1 : $Le = \frac{1}{2} \beta \lambda + dzi$, $Ls = \frac{1}{2} \beta \lambda$, $Lc = \beta \lambda$, $EoT(1) = EoT (1+kEoTi).(Ti/Ts)$

For cell2 to N-1 : $L_e = L_s = \frac{1}{2} \beta \lambda$, $L_c = \beta \lambda$, $E_o T(2..N-1) = E_o T$
 For cellN : $L_e = \frac{1}{2} \beta \lambda$, $L_s = \frac{1}{2} \beta \lambda + dzo$, $L_c = \beta \lambda$, $E_o T(N) = E_o T.(1+kE_o T_o).(T_o/T_s)$
 If m=1:
 For cell1 : $L_e = 1/4 \beta \lambda + dzi$, $L_s = 1/4 \beta \lambda$, $L_c = \frac{1}{2} \beta \lambda$, $E_o T(1) = E_o T.(1+kE_o T_i).(T_i/T_s)$
 For cell2 to N-1 : $L_e = L_s = 1/4 \beta \lambda$, $L_c = \frac{1}{2} \beta \lambda$, $E_o T(2..N-1) = E_o T.(1+kE_o T_i).(T_i/T_s)$
 For cellN : $L_e = 1/4 \beta \lambda$, $L_s = 1/4 \beta \lambda + dzo$, $L_c = \frac{1}{2} \beta \lambda$, $E_o T(N) = E_o T.(1+kE_o T_o).(T_o/T_s)$
 If m=2:
 For cell1 : $L_e = 1/4 \beta \lambda + dzi$, $L_s = \frac{1}{2} \beta \lambda$, $L_c = \frac{3}{4} \beta \lambda$, $E_o T(1) = E_o T.(1+kE_o T_i).(T_i/T_s)$
 For cell2 to N-1 : $L_e = L_s = \frac{1}{2} \beta \lambda$, $L_c = \beta \lambda$, $E_o T(2..N-1) = E_o T.(1+kE_o T_i).(T_i/T_s)$
 For cellN : $L_e = \frac{1}{2} \beta \lambda$, $L_s = 1/2 \beta \lambda + dzo$, $L_c = \frac{3}{4} \beta \lambda$, $E_o T(N) = E_o T.(1+kE_o T_o).(T_o/T_s)$
For all cases : $E_o T_L = E_o T(x).(L_s + L_e)$

Electrostatic Acceleration (Vo, Δs, K)

V_o (V) is the voltage, Δs (mm) is the step length, L (mm) is the element length, and K (eV/m²) is the transverse defocalisation contribution.

The reduced energy change in the gap is:

$$\gamma_s = \gamma_e + \frac{|q| \frac{V_o}{L} \Delta_s}{mc^2}.$$

Lets use:

$$K^* = \frac{q K \Delta_s}{m_0 c^2 \beta^2 \gamma} \quad \text{and} \quad \delta = \sqrt{\frac{(\beta \gamma)_e}{(\beta \gamma)_s}} \quad \text{with} \quad \bar{\gamma} = \frac{\gamma_e + \gamma_s}{2} \quad \text{and} \quad \bar{\beta} = \frac{\beta_e + \beta_s}{2}.$$

The non null 2x2 transfer sub-matrixes are:

$$R_{xx} = R_{yy} = \begin{bmatrix} 1 - K^* \frac{\Delta_s}{2} & \frac{\Delta_s}{2} \left(1 - K^* \frac{\Delta_s}{2} + \delta \right) \\ - K^* \delta & - K^* \delta \frac{\Delta_s}{2} + \delta^2 \end{bmatrix}$$

$$R_{zz} = \begin{bmatrix} 1 & \Delta_s \\ 0 & \delta^2 \end{bmatrix}$$

Sinus cavity or CCL (L, N, EoT, θs)

L (mm) is the cavity length, N is the number of cells, $E_o T$ (eV/m) is the mean electric field of the cavity, θ_s (°) is the phase of the synchronous particle at the entrance of the cavity (relative to the R.F. phase).

□Fields

Let's assume the electric field on the axis:

$$E_z(z, t) = E_0 \sin(\omega t + \varphi_0) \sin\left(\frac{K}{\beta_c} z\right), \quad \text{With} \quad \beta_c = \frac{2L}{N\lambda}.$$

The transverse electric field component can be deduced from Maxwell equations with a first order expansion:

$$E_r(z, t, r) = -\frac{r}{2} \frac{dE_z}{dz} = -\frac{K E_0}{2 \beta_c} \sin(\omega t + \varphi_0) \cos\left(\frac{K}{\beta_c} z\right) \cdot r,$$

$$\Rightarrow \begin{cases} E_x(z, t, x) = -\frac{KE_0}{2\beta_c} \sin(\omega t + \varphi_0) \cos\left(\frac{K}{\beta_c} z\right) \cdot x \\ E_y(z, t, y) = -\frac{KE_0}{2\beta_c} \sin(\omega t + \varphi_0) \cos\left(\frac{K}{\beta_c} z\right) \cdot y \end{cases}$$

The same way, the transverse magnetic field component can be deduced:

$$B_\theta(z, t, r) = \frac{r}{2c^2} \frac{dE_z}{dt} = \frac{KE_0}{2c} \cos(\omega t + \varphi_0) \sin\left(\frac{K}{\beta_c} z\right) \cdot r,$$

$$\Rightarrow \begin{cases} B_x(z, t, y) = -\frac{KE_0}{2c} \cos(\omega t + \varphi_0) \sin\left(\frac{K}{\beta_c} z\right) \cdot y \\ B_y(z, t, x) = \frac{KE_0}{2c} \cos(\omega t + \varphi_0) \sin\left(\frac{K}{\beta_c} z\right) \cdot x \end{cases}$$

With $K = \frac{2\pi}{\lambda} = \frac{\omega}{c}$, $\omega = 2\pi f$, f the R.F frequency and c the speed of light in vacuum.

□ Longitudinal motion

Let's p_s be the synchronous particle momentum at a given z , its evolution is given by:

$$\dot{p}_s = q E_z(z, t_s) = q E_{zs}.$$

The particle momentum is: $p = \sqrt{p_x^2 + p_y^2 + p_z^2} = p_z \sqrt{1 + x'^2 + y'^2}$,

$$\begin{cases} \dot{p}_x = q(E_x - v_z B_y + v_y B_z) \\ \dot{p}_y = q(E_y + v_z B_x - v_x B_z) \\ \dot{p}_z = q(E_z + v_x B_y - v_y B_x) \end{cases}$$

With

The evolution with time of the particle momentum p is given by:

$$\dot{p} = \frac{p_z}{p} (\dot{p}_z + x' \dot{p}_x + y' \dot{p}_y)$$

Let's δ be the momentum of the particle relative to that of the synchronous particle:

$$\delta = \frac{p - p_s}{p_s}.$$

We have: $p = (1 + \delta)p_s = p_z \sqrt{1 + x'^2 + y'^2}$.

The evolution of δ is given by the equation:

$$\dot{\delta} = \frac{1}{p_s^2} (p_s \dot{p} - \dot{p}_s p),$$

$$\dot{\delta} = \frac{1}{p_s^2} \left(\frac{p_s p_z}{p} (\dot{p}_z + x' \dot{p}_x + y' \dot{p}_y) - \dot{p}_s p \right),$$

$$\dot{\delta} = \frac{q}{p_s} \left[(1 + x'^2 + y'^2)^{1/2} \cdot (E_z + x' E_x + y' E_y) - (1 + \delta) E_{zs} \right]$$

$$E_z = E_{zs} + \frac{dE_z}{d\varphi} \Big|_{\varphi=\varphi_s} \cdot \phi, \quad \phi = \varphi - \varphi_s \quad \text{and} \quad \frac{d}{dz} = \frac{1}{\beta_s c} \cdot \frac{d}{dt},$$

We finally find at first order:

$$\frac{d\delta}{dz} = \delta' = \frac{q E_0}{\gamma_s \beta_s^2 m c^2} \sin\left(\frac{Kz}{\beta_c}\right) [\cos(\omega t_s + \varphi_0) \cdot \phi - \sin(\omega t_s + \varphi_0) \cdot \delta]$$

Both focusing and damping effects can be observed.

1. Thin lens approximation

$$\phi = -\frac{K}{\beta_s} \cdot \tilde{x}$$

Using: β_s , the matrix transport can then be written over a small step dz :

$$\begin{pmatrix} \tilde{\mathcal{E}} \\ \tilde{\delta} \end{pmatrix}_{z+dz/2} = \begin{pmatrix} 1 & dz/2\gamma_I^{*2} \\ 0 & 1 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ K_1 & K_2 \end{pmatrix} \begin{pmatrix} 1 & dz/2\gamma_O^{*2} \\ 0 & 1 \end{pmatrix} \cdot \begin{pmatrix} \tilde{\mathcal{E}} \\ \tilde{\delta} \end{pmatrix}_{z-dz/2} = M_z \cdot \begin{pmatrix} \tilde{\mathcal{E}} \\ \tilde{\delta} \end{pmatrix}_{z-dz/2},$$

With:

γ_s and γ_e being the synchronous normalized energy before and after the "gap",

$$K_1 = -\frac{q E_0 K}{\gamma_s \beta_s^3 m c^2} \sin\left(\frac{Kz}{\beta_c}\right) \cos(\omega_s + \varphi_0) \quad \text{And}$$

$$K_2 = 1 - \frac{q E_0}{\gamma_s \beta_s^2 m c^2} \sin\left(\frac{Kz}{\beta_c}\right) \sin(\omega_s + \varphi_0).$$

□ Transverse motion

The evolution of x' with time is driven by the equation:

$$\dot{x}' = \frac{d}{dt} \left(\frac{p_x}{p_z} \right) = \frac{1}{p_z^2} (\dot{p}_x p_z - \dot{p}_z p_x)$$

$$\dot{x}' = \frac{q}{p_s} \cdot \frac{(1 + x'^2 + y'^2)^{1/2}}{1 + \delta} \cdot (\mathcal{E}_x - v_z B_y - x'(\mathcal{E}_z + v_x B_y - v_y B_x))$$

At first order, we finally have:

$$x'' = -\frac{q E_0}{\gamma_s \beta_s^2 m c^2} \left[\frac{K}{2} \left(\frac{1}{\beta_c} \sin(\omega_s + \varphi_0) \cos\left(\frac{Kz}{\beta_c}\right) + \beta_s \cos(\omega_s + \varphi_0) \sin\left(\frac{Kz}{\beta_c}\right) \right) \cdot x \right. \\ \left. + \sin(\omega_s + \varphi_0) \sin\left(\frac{Kz}{\beta_c}\right) \cdot x' \right]$$

Both focusing and damping effects can be observed.

2. Thin lens approximation

The matrix transport over a small step dz can then be written:

$$\begin{pmatrix} x \\ x' \end{pmatrix}_{z+dz/2} = \begin{pmatrix} 1 & dz/2 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ k_1 & k_2 \end{pmatrix} \begin{pmatrix} 1 & dz/2 \\ 0 & 1 \end{pmatrix} \cdot \begin{pmatrix} x \\ x' \end{pmatrix}_{z-dz/2} = M_x \cdot \begin{pmatrix} x \\ x' \end{pmatrix}_{z-dz/2},$$

$$\text{With: } k_1 = -\frac{q E_0 K}{2 \gamma_s \beta_s^2 m c^2} \left[\frac{1}{\beta_c} \sin(\omega_s + \varphi_0) \cos\left(\frac{Kz}{\beta_c}\right) + \beta_s \cos(\omega_s + \varphi_0) \sin\left(\frac{Kz}{\beta_c}\right) \right]$$

$$\text{And } k_2 = K_2$$

3. Thick lens approximation

The differential equation can be written:

$$x'' + 2 \cdot \alpha \cdot x' + \beta \cdot x = 0,$$

$$\text{With: } \beta = \frac{q E_0 K}{2 \gamma_s \beta_s^2 m c^2} \left[\frac{1}{\beta_c} \sin(\omega_s + \varphi_0) \cos\left(\frac{Kz}{\beta_c}\right) + \beta_s \cos(\omega_s + \varphi_0) \sin\left(\frac{Kz}{\beta_c}\right) \right]$$

$$\text{And: } \alpha = \frac{q E_0}{2 \gamma_s \beta_s^2 m c^2} \cdot \sin(\omega_s + \varphi_0) \cdot \sin\left(\frac{Kz}{\beta_c}\right),$$

$$\text{And: } \omega = \sqrt{\beta - \alpha^2},$$

The solution of this differential equation gives:

$$M_x = e^{-\alpha dz} \cdot \begin{pmatrix} \cos(\omega dz) + \frac{\alpha}{\omega} \sin(\omega dz) & \frac{\sin(\omega dz)}{\omega} \\ -\frac{\beta \sin(\omega dz)}{\omega} & \cos(\omega dz) - \frac{\alpha}{\omega} \sin(\omega dz) \end{pmatrix} \quad \text{If } \omega \text{ is real,}$$

And

$$M_x = e^{-\alpha dz} \cdot \begin{pmatrix} ch(|\omega| dz) + \frac{\alpha}{|\omega|} \cdot sh(|\omega| dz) & \frac{sh(|\omega| dz)}{|\omega|} \\ \frac{\beta sh(|\omega| dz)}{|\omega|} & ch(|\omega| dz) - \frac{\alpha}{|\omega|} \cdot sh(|\omega| dz) \end{pmatrix}, \text{ If } j \cdot \omega \text{ is real.}$$

□Transport through a sin-like cavity

$$dz = \frac{\beta_c \lambda}{2n}$$

The N_c -cells cavity is divided in $n \cdot N_c$ steps of length:

As input, we have:

$$\gamma_0 = \gamma_{in}; \quad t_s = \frac{dz}{2\beta_{in}c}; \quad z_s = \frac{dz}{2}; \quad \begin{pmatrix} x \\ x' \end{pmatrix}_0, \quad \begin{pmatrix} \delta \\ \delta \end{pmatrix}_0: \text{ Particle co-ordinates at cavity input.}$$

Then, we loop until reaching the end of the cavity:

For i from 0 to $n \cdot N_c - 1$ do

$$\begin{aligned} \gamma_O^* &= \gamma_{i+1}^* = \gamma_i + \frac{qE_0}{mc^2} \sin(\omega_s + \varphi_0) \sin\left(\frac{K}{\beta_c} \cdot z_s\right) \cdot dz, \quad \beta_{i+1} = \sqrt{1 - \gamma_{i+1}^{-2}}, \\ \gamma_i &= \frac{\gamma_{i+1} + \gamma_i}{2}, \quad \beta_s = \sqrt{1 - \gamma_s^{-2}}, \\ \begin{pmatrix} x \\ x' \end{pmatrix}_{i+1} &= M_x \cdot \begin{pmatrix} x \\ x' \end{pmatrix}_i, \quad \begin{pmatrix} \delta \\ \delta \end{pmatrix}_{i+1} = M_z \cdot \begin{pmatrix} \delta \\ \delta \end{pmatrix}_i, \\ t_s &= t_s + \frac{dz}{\beta_{i+1}c}, \quad z_s = z_s + dz. \end{aligned}$$

RFQ cell (V, ro, A10, m, L, θs, Type ,ρ, A01)

L (mm) is the RFQ cell length, V (V) is the mean voltage of the cell, θ_s (°) is the phase of the synchronous particle, r_o is..., m is ... and type...

Cell type:

- ±2: Accelerating cell.
- ±3: Front-end cell.
- ±4: Transcell.

The sign of type being.....

$$W_{i+1} = W_i + |q|dzE_z \text{ And } \overline{W} = \frac{1}{2}(W_i + W_{i+1})$$

$$\Phi_{i+1} = \Phi_i + dz \frac{2\pi}{\beta\lambda}$$

$$E_z = \frac{\pi A_{10} V}{2L} \sin\left(\frac{\pi}{L} z\right) \sin(\omega_s + \varphi_0)$$

□Thin lens approximation (Longitudinal)

$\omega = 2\pi f$, f the R.F frequency and c the speed of light in vacuum.

The matrix transport can then be written over a small step dz :

$$\begin{pmatrix} \delta \\ \delta \end{pmatrix}_{z+dz/2} = \begin{pmatrix} 1 & dz/2\gamma_i^{*2} \\ 0 & 1 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ K_1 & K_2 \end{pmatrix} \begin{pmatrix} 1 & dz/2\gamma_O^{*2} \\ 0 & 1 \end{pmatrix} \cdot \begin{pmatrix} \delta \\ \delta \end{pmatrix}_{z-dz/2} = M_z \cdot \begin{pmatrix} \delta \\ \delta \end{pmatrix}_{z-dz/2},$$

With:

γ_s and γ_e being the synchronous normalized energy before and after the "gap",

$$K_1 = -\frac{|q|d_z A_{10} V}{\gamma_s \beta_s^2 2mc^2} \left(\frac{\pi}{L}\right)^2 \cdot C_3 \cos(\omega_s + \varphi_0)$$

$$K_2 = 1 - \frac{|q|d_z A_{10} V}{\gamma_s \beta_s^2 2mc^2} \left(\frac{\pi}{L}\right) \cdot C_3 \cos(\omega_s + \varphi_0)$$

C_3 depend of the cell *type*.

$$\begin{aligned} \blacksquare \pm 2 \text{ or } \pm 3: \quad C_3 &= \sin\left(\frac{\pi}{L} z\right) \\ \blacksquare \pm 4: \quad C_3 &= \frac{1}{2} \sin\left(\frac{\pi}{L} z\right) \end{aligned}$$

□Thin lens approximation (Transverse)

The matrix transport over a small step dz can then be written:

$$\begin{pmatrix} x \\ x' \end{pmatrix}_{z+dz/2} = \begin{pmatrix} 1 & dz/2 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} k_{x1} & 0 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} 1 & dz/2 \\ 0 & 1 \end{pmatrix} \cdot \begin{pmatrix} x \\ x' \end{pmatrix}_{z-dz/2} = M_x \cdot \begin{pmatrix} x \\ x' \end{pmatrix}_{z-dz/2},$$

$$\begin{pmatrix} y \\ y' \end{pmatrix}_{z+dz/2} = \begin{pmatrix} 1 & dz/2 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} k_{y1} & 0 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} 1 & dz/2 \\ 0 & 1 \end{pmatrix} \cdot \begin{pmatrix} y \\ y' \end{pmatrix}_{z-dz/2} = M_y \cdot \begin{pmatrix} y \\ y' \end{pmatrix}_{z-dz/2}$$

With:
$$k_{x1} = -\frac{|q|d_z}{\gamma_s \beta_s^2 2mc^2} \cos(\omega_s + \varphi_0) \left[S \frac{V}{r_0^2} A_{01} C_1 - \left(\frac{\pi}{L}\right)^2 \frac{A_{10} V}{4} C_2 \right]$$

$$k_{y1} = -\frac{|q|d_z}{\gamma_s \beta_s^2 2mc^2} \cos(\omega_s + \varphi_0) \left[-S \frac{V}{r_0^2} A_{01} C_1 - \left(\frac{\pi}{L}\right)^2 \frac{A_{10} V}{4} C_2 \right]$$

And $k_{x2} = k_{y2} = K_2$.

$$\begin{pmatrix} x \\ x' \end{pmatrix}_{i+1} = M_x \cdot \begin{pmatrix} x \\ x' \end{pmatrix}_i \text{ And } \begin{pmatrix} y \\ y' \end{pmatrix}_{i+1} = M_y \cdot \begin{pmatrix} y \\ y' \end{pmatrix}_i$$

C_1 , C_2 and S depend of the cell *type*.

$$\begin{aligned} \blacksquare \pm 2: \quad C_1 &= 1, \quad C_2 = \sin\left(\frac{\pi}{L} z\right), \quad S = -\text{sign}(\text{type}) \\ \blacksquare \pm 3: \quad C_1 &= \frac{1}{4} \left(3 \cos\left(\frac{1}{2} \frac{\pi}{L} z - \frac{\pi}{2}\right) + \cos\left(3 \cdot \left(\frac{1}{2} \frac{\pi}{L} z - \frac{\pi}{2}\right)\right) \right), \quad C_2 = 0, \quad S = -\text{sign}(\text{type}[n+1]) \\ \blacksquare -3: \quad C_1 &= \frac{3}{4} \left(\cos\left(\frac{1}{2} \frac{\pi}{L} z\right) - \frac{1}{3} \cos\left(\frac{3}{2} \frac{\pi}{L} z\right) \right), \quad C_2 = 0, \quad S = -\text{sign}(\text{type}[n-1]) \\ \blacksquare \pm 4: \quad C_1 &= 1, \quad C_2 = \frac{1}{2} \left(\cos\left(\frac{\pi}{L} z\right) + 1 \right), \quad S = -\text{sign}(\text{type}[n+1]) \\ \blacksquare -4: \quad C_1 &= 1, \quad C_2 = -\frac{1}{2} \left(\cos\left(\frac{\pi}{L} z\right) - 1 \right), \quad S = -\text{sign}(\text{type}[n-1]) \end{aligned}$$

With $\text{type}[n+1]$ being the type from the next cell and $\text{type}[n-1]$ the type from preceding cell.

□Transport through a rfq cell

The rfq cell is divided in N steps of length: $dz = \frac{L}{N}$.

As input, we have:

$$\gamma_0 = \gamma_{in}; \quad t_s = \frac{dz}{2\beta_{in} c}; \quad z_s = \frac{dz}{2}; \quad \begin{pmatrix} x \\ x' \end{pmatrix}_0, \quad \begin{pmatrix} \tilde{x} \\ \tilde{\delta} \end{pmatrix}_0: \text{ Particle co-ordinates at cavity input.}$$

Then, we loop until reaching the end of the cavity:

For i from 0 to N-1 do

$$\begin{aligned}
 &= \gamma_O^*, \gamma_O^* = \gamma_{i+1} = \gamma_i + \frac{|q|\pi^4 V}{2 L m c^2} \sin(\omega t_s + \varphi_0) \sin\left(\frac{\pi}{L} \cdot z_s\right) \cdot dz, \beta_{i+1} = \sqrt{1 - \gamma_{i+1}^{-2}}, \\
 &= \frac{\gamma_{i+1} + \gamma_i}{2}, \beta_s = \sqrt{1 - \gamma_s^{-2}}, \\
 &\begin{pmatrix} x \\ x' \end{pmatrix}_{i+1} = M_x \cdot \begin{pmatrix} x \\ x' \end{pmatrix}_i, \begin{pmatrix} \delta x \\ \delta \end{pmatrix}_{i+1} = M_z \cdot \begin{pmatrix} \delta x \\ \delta \end{pmatrix}_i, \\
 &= t_s + \frac{dz}{\beta_{i+1} c}, z_s = z_s + dz.
 \end{aligned}$$

Transfer matrix of alpha magnet

The following paper come from a CEA report named:

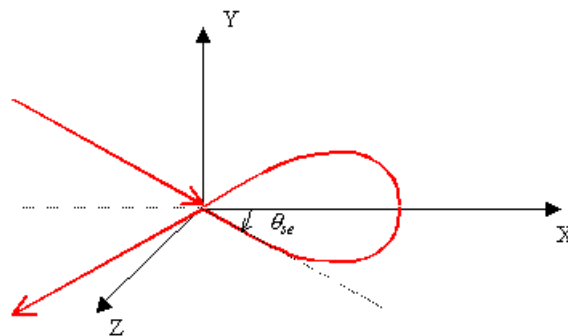
“Transfer matrix of a constant gradient alpha magnet for ELSA extension”

Ref: DSM/DAPNIA/SACM/2002/13

In this paper, the motion equation of a particle in such a device is linearised to get the transfer matrix. The matrix has been validated with a step-by-step integration of the motion of a particle in the magnetic field.

Magnetic field and fixed coordinate system

The trajectory of the synchronous particle (ideal trajectory) is, by definition, in the plan (X, Y). The frame origin is the point where the particle enters the magnet. The X direction is the direction of the magnetic field gradient. At any position, the main trajectory makes an angle θ_s with the X axis. The entrance angle for negative particle θ_{se} is: -40.71° . With this particular angle, the synchronous particle exits the magnet at the same position as the entering particle (and with the opposite angle).



In this frame, the magnetic field is:

$$\vec{B} = \begin{pmatrix} B_x = k \cdot Z \\ B_y = 0 \\ B_z = k \cdot X \end{pmatrix}$$

Trajectory of the synchronous particle

The synchronous particle with charge q moves in the (X, Y) plan.

Its motion equations with time t are:

$$\begin{cases} \frac{dp_x}{dt} = q \cdot v_y \cdot B_z \\ \frac{dp_y}{dt} = -q \cdot v_x \cdot B_z \end{cases}$$

p is the particle momentum, v its velocity,

giving:

$$\begin{cases} \frac{dp_x}{ds} = q \cdot \frac{p_y}{p_0} \cdot k \cdot X \\ \frac{dp_y}{ds} = -q \cdot \frac{p_x}{p_0} \cdot k \cdot X \\ \frac{dX}{ds} = \frac{p_x}{p_0} \\ \theta_s = \arcsin\left(\frac{p_y}{p_0}\right) = \arccos\left(\frac{p_x}{p_0}\right) \end{cases}$$

$p_0 = \beta_0 \gamma_0 m_0 c$ is the particle momentum modulus,

s is the curvilign abscissa, with $ds = v_0 \cdot dt$.

These equations have to be solved using as initial conditions:

$$\begin{cases} p_x = p_0 \cos \theta_{se} \\ p_y = p_0 \sin \theta_{se} \\ X = 0 \\ \theta_s = \theta_{se} \end{cases}$$

The maximum *penetration of the particle* X_M in the magnet can be calculated the following way:

$$\frac{dp_y}{dX} = -q \cdot k \cdot X \quad \Rightarrow \quad p_y = p_{y0} - \frac{1}{2} \cdot q \cdot k \cdot X^2.$$

At maximum penetration, one has: $p_y = p_0$, giving:

$$X_M = \sqrt{\frac{2 \cdot (p_{y0} - p_0)}{q \cdot k}}.$$

The length L of the trajectory is obtained from the integration of:

$$dt = \frac{\gamma_0 m_0 \cdot dX}{p_x}.$$

One obtains:

$$L = 2 \cdot p_0 \cdot \int_0^{X_M} \left(p_{x0}^2 + q \cdot k \cdot p_{y0} \cdot X^2 - 0.25 \cdot q^2 k^2 \cdot X^4 \right)^{\frac{1}{2}} \cdot dX.$$

This integral cannot be solved analytically but can be easily calculated numerically.

The moving coordinate system

In the transfer matrix formalism, a beam particles is referenced, at a given curvilign abscissa s , to the synchronous particle in a 6D phase-space, with a 6-coordinates vector:

$$\vec{r} = \begin{pmatrix} x \\ x' \\ y \\ y' \\ \varphi \\ \delta \end{pmatrix},$$

with:

- x is the particle transverse position in the deviation plan,

- x' is the particle transverse slope in the x direction. $x' = \frac{p_x}{p_s} = \frac{dx}{ds}$, p_x the x-component of the particle momentum, and p_s the s-component of the particle momentum,

- $$\varphi = 2\pi \cdot f_{RF} \cdot (t - t^{(S)}),$$

- $$\delta = \frac{p - p^{(s)}}{p^{(s)}}.$$

In the matrix formalism, the particle vector change from a point (1) to a point (2) is given by :

$$(\vec{r})_{(2)} = [T] \cdot (\vec{r})_{(1)},$$

$$T_{i,j} = \frac{\partial r_i)_{(2)}}{\partial r_j)_{(1)}}.$$

The matrix coefficients can be calculated by varying the input particle coordinate along each direction independently and looking at the output coordinates.

Matrix third and fourth columns

Figure 1 illustrates the geometry of the problem in a Cartesian coordinate system with origin $O(0,0)$. The vertical axis is Y and the horizontal axis is X . A vertical line segment OE is shown on the Y -axis, with E at $(0, \rho_0 + x - \rho)$. A point A is located on the Y -axis at $(0, \rho_0)$. A point B is located on the Y -axis at $(0, \rho_0 + x)$. A point C is located on a red curve at $(\rho_0 \sin(\theta_0), \rho_0 \cos(\theta_0))$. A point D is located on a blue curve L at (X_D, Y_D) . A line segment CD is drawn, and its projection onto the Y -axis is x_f . A line segment BD is drawn, and its projection onto the Y -axis is x'_f . The angle θ_0 is shown between the Y -axis and the line segment OD .

D is at the intersection of (O, C) and the circle with centre E and radius ρ . Its coordinates satisfy:

$$\begin{cases} X_D = \tan(\theta_0) \cdot Y_D \\ X_D^2 + (Y_D - (\rho_0 + x - \rho))^2 = \rho^2 \end{cases}$$

Giving:

$$\left(1 + \tan^2 \theta_0\right) \cdot Y_D^2 - 2 \cdot (x - d\rho) \cdot Y_D + (\rho_0 + x) \cdot (x - 2 \cdot d\rho - \rho_0) = 0,$$

with: $\rho = \rho_0 + d\rho = \rho_0 \cdot \left(1 + \frac{d\rho}{\rho_0}\right)$.

As: $\rho = \frac{p}{q \cdot B}$, one has: $\frac{d\rho}{\rho_0} = -\frac{dB}{B_0} = \sin \theta_s \cdot \frac{x}{X_s}$.

The solution of (X) is obtained giving a reduced discriminator:

$$\Delta' = (x - d\rho)^2 - \frac{(\rho_0 + x) \cdot (x - d\rho - \rho_0)}{\cos^2 \theta_0},$$

giving:

$$\Delta' = \frac{\rho_0^2}{\cos^2 \theta_0} \cdot \left[\left(\frac{x}{\rho_0} - \frac{d\rho}{\rho_0} \right)^2 \cdot \cos^2 \theta_0 - \left(1 + \frac{x}{\rho_0} \right) \cdot \left(\frac{x}{\rho_0} - 2 \cdot \frac{d\rho}{\rho_0} - 1 \right) \right]$$

A first order development in $\frac{x}{\rho_0} \ll 1$ and $\frac{d\rho}{\rho_0} \ll 1$ gives:

$$\Delta' = \frac{\rho_0^2}{\cos^2 \theta_0} \cdot \left[1 + 2 \cdot \frac{d\rho}{\rho_0} \right].$$

This gives the coordinates of D:

$$\begin{cases} X_D = \rho_0 \cdot \sin \theta_0 \cdot \left(1 + \cos \theta_0 \cdot \frac{x}{\rho_0} + (1 - \cos \theta_0) \cdot \frac{d\rho}{\rho_0} \right) \\ Y_D = \rho_0 \cdot \cos \theta_0 \cdot \left(1 + \cos \theta_0 \cdot \frac{x}{\rho_0} + (1 - \cos \theta_0) \cdot \frac{d\rho}{\rho_0} \right) \end{cases}$$

The final position of the particle in the moving frame is then :

$$x_f = \frac{Y_D - Y_C}{\cos \theta_0} = \left(\cos \theta_0 + (1 - \cos \theta_0) \cdot \frac{\rho_0 \cdot \sin \theta_s}{X_s} \right) \cdot x = T_{11} \cdot x$$

Calculus of $T_{2,1}$

$$x'_f = \tan(\theta_0 - \theta) = \frac{\tan \theta_0 - \tan \theta}{1 + \tan \theta_0 \cdot \tan \theta},$$

with: $\tan \theta = \frac{X_D - X_F}{Y_D - Y_F}$.

At first order, one has:

$$\tan \theta = \tan \theta_0 \cdot \left(1 + \frac{x - d\rho}{\rho_0 \cos \theta_0} \right),$$

giving :

$$x'_f = -\sin \theta_0 \cdot \left(1 - \frac{\rho_0 \cdot \sin \theta_s}{X_s} \right) \cdot \frac{x}{\rho_0} = T_{2,1} \cdot x$$

Calculus of $T_{5,1}$

$$\varphi = \frac{2\pi \cdot f_{RF}}{\beta_0 c} (L - \rho_0 \theta_0) = \frac{2\pi \cdot f_{RF}}{\beta_0 c} \rho_0 \theta_0 \cdot \left(\frac{d\rho}{\rho_0} + \frac{d\theta}{\theta_0} \right).$$

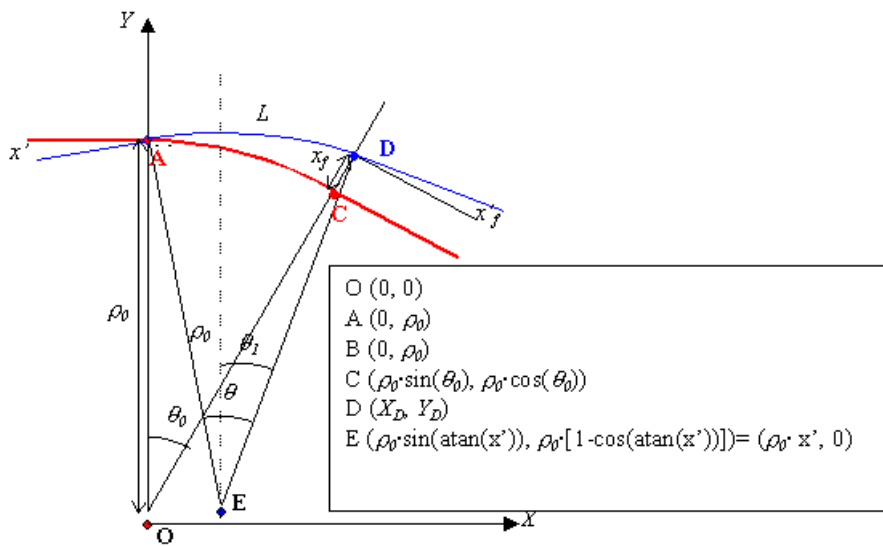
with: $d\theta = \cos^2 \theta_0 \cdot d(\tan \theta) = \cos^2 \theta_0 \cdot (\tan \theta - \tan \theta_0) = \sin \theta_0 \cdot \left(\frac{x - d\rho}{\rho_0} \right)$.

One obtains finally:

$$\varphi = \frac{2\pi \cdot f_{RF}}{\beta_0 c} \cdot \left(\sin \theta_0 + (\theta_0 - \sin \theta_0) \cdot \frac{\rho_0 \sin \theta_s}{X_s} \right) \cdot x = T_{5,1} \cdot x$$

The other terms ($T_{3,1}$, $T_{4,1}$, $T_{6,1}$) are equal to zero.

Matrix second column: variation with x'



Calculus of $T_{1,2}$

D is at the intersection of (O, C) and the circle with centre E and radius ρ . Its coordinates satisfy:

$$\begin{cases} X_D = \tan \theta_0 \cdot Y_D \\ (X_D - \rho_0 \cdot x')^2 + Y_D^2 = \rho_0^2 \end{cases}$$

Giving:

$$(1 + \tan^2 \theta_0) \cdot Y_D^2 - 2 \cdot \rho_0 \cdot x' \cdot \tan \theta_0 \cdot Y_D - \rho_0^2 \cdot (1 - x'^2) = 0$$

The solution of (X) is obtained giving a reduced discriminator:

$$\Delta' = \rho_0^2 \cdot \tan^2 \theta_0 \cdot x'^2 + \frac{\rho_0^2 \cdot (1 - x'^2)}{\cos^2 \theta_0}$$

A first order development in $x' \ll 1$ gives:

$$\Delta' = \frac{\rho_0^2}{\cos^2 \theta_0}$$

This gives the coordinates of D:

$$\begin{cases} X_D = \rho_0 \cdot \sin \theta_0 \cdot (1 + \sin \theta_0 \cdot x') \\ Y_D = \rho_0 \cdot \cos \theta_0 \cdot (1 + \sin \theta_0 \cdot x') \end{cases}$$

The final position of the particle in the moving frame is then :

$$x_f' = \frac{Y_D - Y_C}{\cos \theta_0} = \rho_0 \cdot \sin \theta_0 \cdot x' = T_{1,2} \cdot x'$$

Calculus of $T_{2,2}$

$$x_f' = \tan(\theta_0 - \theta_1) = \frac{\tan \theta_0 - \tan \theta_1}{1 + \tan \theta_0 \cdot \tan \theta_1}$$

with: $\tan \theta_1 = \frac{X_D - X_E}{Y_D - Y_E}$.

At first order, one has:

$$\tan \theta = \frac{X_D - X_E}{Y_D - Y_E},$$

giving :

$$x_f' = \cos \theta_0 \cdot x' = T_{2,2} \cdot x'$$

Calculus of $T_{5,2}$

$$\varphi = \frac{2\pi \cdot f_{RF}}{\beta_0 c} (L - \rho_0 \theta_0) = \frac{2\pi \cdot f_{RF}}{\beta_0 c} \rho_0 \cdot d\theta.$$

with : $d\theta = \cos^2 \theta_0 \cdot d(\tan \theta) = \cos^2 \theta_0 \cdot (\tan \theta - \tan \theta_0),$

as: $\tan \theta = \tan(\theta_1 + \theta_0) = \tan \theta_0 \cdot \left(1 + \frac{1 - \cos \theta_0}{\sin \theta_0 \cdot \cos \theta_0} \cdot x'\right),$

one has: $d\theta = (1 - \cos \theta_0) \cdot x'.$

One obtains finally:

$$\varphi = \frac{2\pi \cdot f_{RF}}{\beta_0 c} \cdot \rho_0 \cdot (1 - \cos \theta_0) \cdot x' = T_{5,2} \cdot x'.$$

The other terms ($T_{3,2}, T_{4,2}, T_{6,2}$) are equal to zero.

Matrix sixth column: variation with δ

$$x'_f = \tan(\theta_0 - \theta) = \frac{\tan \theta_0 - \tan \theta}{1 + \tan \theta_0 \cdot \tan \theta}$$

wih: $\tan \theta = \frac{X_D - X_E}{Y_D - Y_E}$.

At first order, one has:

$$\tan \theta = \tan \theta_0 \cdot \left(1 - \frac{d\rho}{\rho_0 \cos \theta_0} \right),$$

giving :

$$x'_f = \sin \theta_0 \cdot \delta = T_{5,5} \cdot \delta$$

Calculus of $T_{5,6}$

$$\varphi = \frac{2\pi \cdot f_{RF}}{\beta_0 c} (L - \rho_0 \theta_0) = \frac{2\pi \cdot f_{RF}}{\beta_0 c} \rho_0 \theta_0 \cdot \left(\frac{d\rho}{\rho_0} + \frac{d\theta}{\theta_0} \right).$$

with: $d\theta = \cos^2 \theta_0 \cdot d(\tan \theta) = \cos^2 \theta_0 \cdot (\tan \theta - \tan \theta_0) = -\sin \theta_0 \cdot \frac{d\rho}{\rho_0}$,

One obtains finally:

$$\varphi = \frac{2\pi \cdot f_{RF}}{\beta_0 c} \cdot \rho_0 \theta_0 \cdot \left(1 - \frac{\sin \theta_0}{\theta_0} \right) \cdot \delta = T_{5,6} \cdot \delta$$

The terms ($T_{3,6}$, $T_{4,6}$) are equal to zero, the term $T_{6,6}$ is equal to 1.

Matrix fifth column: variation with φ

The output position, slope, energy do not depend on the input phase φ .

The terms ($T_{1,5}$, $T_{2,5}$, $T_{3,5}$, $T_{4,5}$, $T_{6,5}$) are equal to zero, the term $T_{5,5}$ is equal to 1.

Matrix third and fourth columns: motion in y (or Z)

The equation of motion along Z direction is :

$$\frac{dp_z}{dt} = -q \cdot \frac{p_y}{\gamma m} \cdot B_x = -q \cdot \frac{p_y}{\gamma m} \cdot k \cdot Z$$

For the matrix calculation, one uses a first order development of the force, giving:

$$\frac{dp_z}{dt} = -q \cdot \frac{p_0 \cdot \sin \theta_0}{\gamma m} \cdot k \cdot Z$$

This equation is the classical one in a quadrupole with gradient: $k \cdot \sin \theta_0$.

The associated matrix coefficients are:

$$T_{3,3} = T_{4,4} = \cos(\sqrt{K} \rho_0 \theta_0),$$

$$T_{3,4} = \frac{\sin(\sqrt{K} \rho_0 \theta_0)}{\sqrt{K}},$$

$$T_{4,3} = -\sqrt{K} \cdot \sin(\sqrt{K} \rho_0 \theta_0),$$

with: $K = \frac{k \cdot \sin \theta_s}{p_0}$.

The other coefficients are equal to 0.

Alpha magnet matrix

The final matrix of a fraction of a alpha magnet (on which, X_s and θ_s are kept almost constant) :

$$\begin{pmatrix} \cos \theta_0 + (1 - \cos \theta_0) \cdot \frac{\rho_0 \cdot \sin \theta_s}{X_s} & \rho_0 \cdot \sin \theta_s & 0 & 0 & 0 & \rho_0 \cdot (1 - \cos \theta_0) \\ -\frac{\sin \theta_0}{\rho_0} \cdot \left(1 - \frac{\rho_0 \cdot \sin \theta_s}{X_s}\right) & \cos \theta_0 & 0 & 0 & 0 & \sin \theta_0 \\ 0 & 0 & \cos(\sqrt{K} \rho_0 \theta_0) & \frac{\sin(\sqrt{K} \rho_0 \theta_0)}{\sqrt{K}} & 0 & 0 \\ 0 & 0 & -\sqrt{K} \cdot \sin(\sqrt{K} \rho_0 \theta_0) & \cos(\sqrt{K} \rho_0 \theta_0) & 0 & 0 \\ K_{\#} \cdot \left(\sin \theta_0 + (\theta_0 - \sin \theta_0) \cdot \frac{\rho_0 \sin \theta_s}{X_s}\right) & K_{\#} \cdot \rho_0 \cdot (1 - \cos \theta_0) & 0 & 0 & 1 & K_{\#} \cdot \rho_0 \cdot (\theta_0 - \sin \theta_0) \\ 0 & 0 & 0 & 0 & 0 & 1 \end{pmatrix}$$

with: $K = \frac{k \cdot \sin \theta_s}{X_s}$,

and: $K_{\#} = \frac{2\pi \cdot f_{RF}}{\beta_0 c}$.

The matrix of the full element is a product of all matrixes for varying X_s and θ_s .

[Explanation about the way to obtain this matrix](#)

Equations du transport dans un champ électromagnétique

Le but de ce rapport est d'écrire les équations exactes du transport d'une particule dans un champ électromagnétique (RF ou statique). Elles seront utilisées dans les codes de calculs multiparticulaires (par exemple Partran). Les équations sont ensuite linéarisées afin d'être introduite dans les codes d'enveloppes (par exemple, TraceWIN).

- Equations générales de la dynamique**

La variations de la quantité de mouvement \vec{p} d'une particule de charge q et de masse γm dans un champ électromagnétique est :

$$\frac{d\vec{p}}{dt} = q \cdot \left(\vec{E} + \frac{\vec{p}}{\gamma m} \times \vec{B} \right)$$

En coordonnées cartésiennes, l'équation se projette :

$$\begin{cases} \dot{p}_x = \frac{dp_x}{dt} = q \cdot \left(E_x + \frac{p_z}{\gamma m} \cdot (y' \cdot B_z - B_y) \right) \\ \dot{p}_y = \frac{dp_y}{dt} = q \cdot \left(E_y + \frac{p_z}{\gamma m} \cdot (B_x - x' \cdot B_z) \right) \\ \dot{p}_z = \frac{dp_z}{dt} = q \cdot \left(E_z + \frac{p_z}{\gamma m} \cdot (x' \cdot B_y - y' \cdot B_x) \right) \end{cases}$$

avec :

$$\frac{dx}{ds} = x' = \frac{p_x}{p_z} \quad \text{et} \quad \frac{dy}{ds} = y' = \frac{p_y}{p_z}$$

Nous avons alors :

$$\frac{dx'}{ds} = \frac{d(p_x/p_z)}{dt} \cdot \frac{dt}{ds} = \frac{1}{\beta_z c} \cdot \frac{\dot{p}_x - p_z \cdot x'}{p_z}$$

Ce qui donne :

$$\frac{dx'}{ds} = \frac{1}{\gamma \beta_z^2 (mc^2/q)} \cdot \left(E_x - x' \cdot E_z + \beta_z c \cdot (x' y' \cdot B_x - (1 + x'^2) \cdot B_y + y' \cdot B_z) \right)$$

De même nous avons :

$$\frac{dy'}{ds} = \frac{1}{\gamma \beta_z^2 (mc^2/q)} \cdot \left(E_y - y' \cdot E_z + \beta_z c \cdot ((1 + y'^2) \cdot B_x - x' y' \cdot B_y - x' \cdot B_z) \right)$$

En ce qui concerne la dynamique longitudinale, les variables utilisées sont généralement soit l'énergie cinétique T soit la quantité de mouvement p de la particule. La variation de ces grandeurs est déduite de la relation :

$$p^2 = p_x^2 + p_y^2 + p_z^2 = p_z^2 \cdot (1 + x'^2 + y'^2) = \frac{(E_0 + T)^2 - E_0^2}{c^2}$$

Donnant :

$$\frac{dp}{ds} = \frac{\frac{dp_z}{ds} \cdot (1 + x'^2 + y'^2) + p_z \cdot \left(x' \cdot \frac{dx'}{ds} + y' \cdot \frac{dy'}{ds} \right)}{\sqrt{1 + x'^2 + y'^2}} \quad p^2 = p_x^2 + p_y^2 + p_z^2 = p_z^2 \cdot (1 + x'^2 + y'^2) = \frac{(E_0 + T)^2 - E_0^2}{c^2}$$

soit :

$$\frac{dp}{ds} = \frac{q}{\beta \cdot c} \cdot (E_z + x' \cdot E_x + y' \cdot E_y)$$

et :

$$\frac{dT}{ds} = q \cdot (E_z + x' \cdot E_x + y' \cdot E_y)$$

Dans le code PIC, ces équations doivent être intégrées pas à pas.

• Linéarisation

Dans le cadre de l'hypothèse de paraxialité, la linéarisation des équations nous donne :

$$\begin{aligned} \frac{dx'}{ds} &= \frac{1}{\gamma \beta_z^2 (mc^2/q)} \cdot (E_x - x' \cdot E_z + \beta_z c \cdot (-B_y + y' \cdot B_z)) \\ \frac{dy'}{ds} &= \frac{1}{\gamma \beta_z^2 (mc^2/q)} \cdot (E_y - y' \cdot E_z + \beta_z c \cdot (B_x - x' \cdot B_z)) \\ \frac{dp}{ds} &= \frac{1}{\beta \cdot c} \cdot (E_z + x' \cdot E_x + y' \cdot E_y) \end{aligned}$$

Dans TraceWIN, les variables longitudinales sont :

$$z = -\frac{\varphi}{2\pi} \cdot \beta \lambda \quad \text{et} \quad \delta = \frac{p - p_s}{p_s}$$

p_s est la quantité de mouvement de la particule synchrone.

L'évolution de δ est alors :

$$\frac{d\delta}{ds} = \frac{1}{p_s} \cdot \left(\frac{dp}{ds} - (1 + \delta) \cdot \frac{dp_s}{ds} \right) = \frac{1}{\beta \cdot c} \cdot (E_z + x' \cdot E_x + y' \cdot E_y)$$

Les champs peuvent être modélisés par :

$$E_x = E_{x0} + \frac{\partial E_x}{\partial x} \cdot x + \frac{\partial E_x}{\partial y} \cdot y + \frac{\partial E_x}{\partial z} \cdot z \dots$$

$$B_x = B_{x0} + \frac{\partial B_x}{\partial x} \cdot x + \frac{\partial B_x}{\partial y} \cdot y + \frac{\partial B_x}{\partial z} \cdot z \dots$$

Les équations du mouvement linéarisées deviennent alors :

$$\frac{dx'}{ds} = \frac{q}{\gamma_s \beta_s^2 mc^2} \cdot \left((E_{x0} - \beta_s c \cdot B_{y0}) - E_{x0} \cdot x' + \beta_s c \cdot B_{x0} \cdot y' - [(2 - \beta_s^2) \cdot E_{x0} + \beta_s c \cdot B_{y0}] \cdot \delta \right)$$

$$\frac{dy'}{ds} = \frac{q}{\gamma_s \beta_s^2 mc^2} \cdot \left((E_{y0} + \beta_s c \cdot B_{x0}) - \beta_s c \cdot B_{x0} \cdot x' - E_{y0} \cdot y' - [(2 - \beta_s^2) \cdot E_{y0} - \beta_s c \cdot B_{x0}] \cdot \delta \right)$$

$$\frac{d\delta}{ds} = \frac{q}{\gamma_s \beta_s^2 mc^2} \cdot \left(\frac{\partial E_x}{\partial x} \cdot x + \frac{\partial E_x}{\partial y} \cdot y + \frac{\partial E_x}{\partial z} \cdot z + E_{x0} \cdot x' + E_{y0} \cdot y' - (2 - \beta_s^2) \cdot E_{x0} \cdot \delta \right)$$

Nous avons utilisé ici :

$$d(\gamma \beta^2)^{-1} = -\frac{2 - \beta_s^2}{\gamma_s \beta_s^2} \cdot \delta$$

$$d\beta = \frac{\beta_s}{\gamma_s^2} \cdot \delta$$

$$d(\beta^{-1}) = -\frac{\delta}{\beta_s \cdot \gamma_s^2}$$

• Quelques éléments de transport

• Éléments à symétrie de révolution

De nombreux éléments (solénoïdes, lentilles de Einzel, cavités RF (DTL, CCL, supra-elliptique,...)) présentent une symétrie de révolution autour de l'axe du faisceau. Dans ce cas, les champs sont représentés dans le référentiel cylindrique : (r, θ, z) .

On a alors :

$$\begin{pmatrix} V_x(r \cdot \cos \theta, r \cdot \sin \theta) \\ V_y(r \cdot \cos \theta, r \cdot \sin \theta) \end{pmatrix} = \begin{pmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{pmatrix} \cdot \begin{pmatrix} V_r(r, \theta) \\ V_\theta(r, \theta) \end{pmatrix}$$

• Solénoïde magnétique

Dans un solénoïde, on a, au premier ordre :

$$\vec{E} = \vec{0},$$

$$B_\theta = 0,$$

$$B_r = \frac{\partial B_z}{\partial r} \cdot r,$$

$$B_z = B_{z0}.$$

Les équations linéarisées du mouvement deviennent alors :

$$\frac{dx'}{ds} = \frac{q}{\gamma_s \beta_s^2 mc^2} \cdot \left(\beta_s c \cdot B_{z0} \cdot y' - \beta_s c \cdot \frac{\partial B_z}{\partial r} \cdot y \right)$$

$$\frac{dy'}{ds} = \frac{q}{\gamma_s \beta_s^2 mc^2} \cdot \left(-\beta_s c \cdot B_{z0} \cdot x' + \beta_s c \cdot \frac{\partial B_z}{\partial r} \cdot x \right)$$

• Lentille électrostatique de Einzel

Dans une lentille électrostatique de Einzel, on a, au premier ordre :

$$\vec{B} = \vec{0},$$

$$E_\theta = 0,$$

$$E_r = \frac{\partial E_z}{\partial r} \cdot r,$$

$$E_x = E_{x0}.$$

Les équations linéarisées du mouvement deviennent alors :

$$\frac{dx'}{ds} = \frac{q}{\gamma_s \beta_s^2 mc^2} \cdot \left(\frac{\partial E_x}{\partial r} \cdot x - E_{x0} \cdot x' \right)$$

$$\frac{dy'}{ds} = \frac{q}{\gamma_s \beta_s^2 mc^2} \cdot \left(\frac{\partial E_y}{\partial r} \cdot y - E_{x0} \cdot y' \right)$$

$$\frac{d\delta}{ds} = -\frac{q}{\gamma_s \beta_s^2 mc^2} \cdot (2 - \beta_s^2) \cdot E_{x0} \cdot \delta$$

A ces équations, il faut ajouter la variation de l'énergie de la particule synchrone :

$$\frac{dT_s}{ds} = q \cdot E_{x0}.$$

- Cavités accélératrice RF

Dans une cavité accélératrice RF, on a :

$$E_x(r, s, t) = E_{x0}(s) \cdot \left(1 - \frac{1}{4} \cdot \left(\frac{1}{E_{x0}(s)} \cdot \frac{d^2 E_{x0}(s)}{ds^2} + \frac{\omega^2}{c^2} \right) \cdot r^2 \right) \cdot \cos(\omega \cdot t + \varphi),$$

$$E_y(r, s, t) = -\frac{1}{2} \cdot \frac{dE_{x0}(s)}{ds} \cdot \left(1 - \frac{1}{8} \cdot \left(\frac{1}{E_{x0}(s)} \cdot \frac{d^2 E_{x0}(s)}{ds^2} + \frac{\omega^2}{c^2} \right) \cdot r^2 \right) \cdot r \cdot \cos(\omega \cdot t + \varphi),$$

$$B_\theta(r, s, t) = -\frac{\omega}{2 \cdot c^2} \cdot E_{x0}(s) \cdot \left(1 - \frac{1}{8} \cdot \left(\frac{1}{E_{x0}(s)} \cdot \frac{d^2 E_{x0}(s)}{ds^2} + \frac{\omega^2}{c^2} \right) \cdot r^2 \right) \cdot r \cdot \sin(\omega \cdot t + \varphi).$$

Soit au premier ordre :

$$E_x(r, s, t) = E_{x0}(s) \cdot \cos(\omega \cdot t + \varphi),$$

$$E_y(r, s, t) = -\frac{1}{2} \cdot \frac{dE_{x0}(s)}{ds} \cdot r \cdot \cos(\omega \cdot t + \varphi),$$

$$B_\theta(r, s, t) = -\frac{\omega}{2 \cdot c^2} \cdot E_{x0}(s) \cdot r \cdot \sin(\omega \cdot t + \varphi).$$

Soient, en coordonnées cartésiennes :

$$E_x(r, s, t) = -\frac{1}{2} \cdot \frac{dE_{x0}(s)}{ds} \cdot x \cdot \cos(\omega \cdot t + \varphi) = \frac{\partial E_x}{\partial x} \cdot x,$$

$$E_y(r, s, t) = -\frac{1}{2} \cdot \frac{dE_{x0}(s)}{ds} \cdot y \cdot \cos(\omega \cdot t + \varphi) = \frac{\partial E_y}{\partial y} \cdot y,$$

$$B_x(r, s, t) = \frac{\omega}{2 \cdot c^2} \cdot E_{x0}(s) \cdot y \cdot \sin(\omega \cdot t + \varphi) = \frac{\partial B_x}{\partial y} \cdot y,$$

$$B_y(r, s, t) = -\frac{\omega}{2 \cdot c^2} \cdot E_{x0}(s) \cdot x \cdot \sin(\omega \cdot t + \varphi) = \frac{\partial B_y}{\partial x} \cdot x.$$

Les équations linéarisées du mouvement deviennent alors :

$$\frac{dx'}{ds} = \frac{q}{\gamma_s \beta_s^2 mc^2} \cdot \left(-E_{x0} \cdot \cos \varphi \cdot x' + \left(\frac{\partial E_x}{\partial x} - \beta_s c \cdot \frac{\partial B_y}{\partial x} \right) \cdot x \right),$$

$$\frac{dy'}{ds} = \frac{q}{\gamma_s \beta_s^2 mc^2} \cdot \left(-E_{x0} \cdot \cos \varphi \cdot y' + \left(\frac{\partial E_y}{\partial y} + \beta_s c \cdot \frac{\partial B_x}{\partial y} \right) \cdot y \right),$$

$$\frac{d\delta}{ds} = \frac{q}{\gamma_s \beta_s^2 mc^2} \cdot \left(\frac{\partial E_x}{\partial z} \cdot z - (2 - \beta_s^2) \cdot E_{x0} \cdot \cos \varphi \cdot \delta \right).$$

- **Eléments sans symétrie particulière**

- **Quadripôle magnétique**

Dans un quadripôle magnétique :

$$\vec{E} = \vec{0}, \quad B_x = G \cdot y \quad \text{et} \quad B_y = G \cdot x$$

On obtient alors :

$$\frac{dx'}{ds} = -\frac{q}{\gamma_s \beta_s^2 mc^2} \cdot \beta_s c \cdot G \cdot x,$$

$$\frac{dy'}{ds} = \frac{q}{\gamma_s \beta_s^2 mc^2} \cdot \beta_s c \cdot G \cdot y.$$

- **Quadripole électrique**

Dans un quadripôle électrique :

$$\vec{B} = \vec{0}, \quad E_x = -G \cdot x \quad \text{et} \quad E_y = G \cdot y$$

On obtient alors :

$$\frac{dx'}{ds} = \frac{q}{\gamma_s \beta_s^2 mc^2} \cdot G \cdot x$$

$$\frac{dy'}{ds} = \frac{q}{\gamma_s \beta_s^2 mc^2} \cdot G \cdot y$$

Change frequency

FREQ *f(Mhz)*

Change the R.F. frequency of the following structure, the beam frequency is not affected.

Begin of lattice

LATTICE *n1, n2*

Define the periodic focusing lattices, *n1* is the number of element per lattice, *n2* is the number of lattice per macro-lattice (generally 1). You don't have to repeat this command; all following elements are considered as part of the lattice structure until reaching the commands [LATTICE_END](#) or [END](#).

Copy following examples:

[examples\DTL_1.ini](#)

[examples\DTL_1.dat](#)

[examples\DTL_1.cal](#)

End of lattice

LATTICE_END

End the periodic focusing lattices.

Structure file end

END

A structure file description must finish by the word “END”.

Set phase advance

SET_ADV k_{0t}

Set the transverse phase advance with no current to k_{0t}

[See the phase advance definition](#)

To set a transverse phase advance law

The periodic focusing lattices structure has to be defined by using the command [LATTICE](#) and [LATTICE_END](#), in order to indicate the number of lattice and the number of element per lattice.

The zero-current transverse phase advance law can be imposed in two ways:

When the option “*Use phase advance definition*” of “*Main*” page is checked, TRACE_WIN imposes the zero-current transverse phase advance law describe in the [Sigma0 file](#). This file is editable in the “*Main*” page

If the option “*Use phase advance definition*” is unchecked that means you have to use the commands [SET_ADV](#) in you data file in order to describe your phase advance law.

Phase advance example

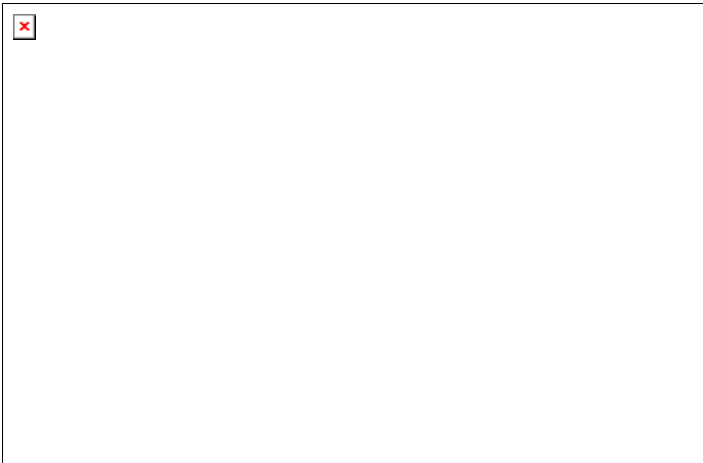
FODO lattices in a DTL tank

```
;Tank 1
LATTICE 2 1
SET_ADV 30
DTL_CEL 87.5277 28 28 0.00548673 71.3977 -71.6524 74699.4 -50 10 0 0.102751 0.723252 -0.475485 -0.194532
DTL_CEL 87.9518 28 28 0.0080072 -71.6524 71.9632 76721.8 -49.5 10 0 0.103248 0.726628 -0.470582 -0.195991
DTL_CEL 88.3919 28 28 0.0112292 71.9632 -72.2903 79282.7 -49 10 0 0.103764 0.734598 -0.458632 -0.197949
DTL_CEL 88.8487 28 28 0.0144758 -72.2903 72.5931 81864.3 -48.5 10 0 0.104299 0.742159 -0.447189 -0.199371
DTL_CEL 89.3222 28 28 0.0177645 72.5931 -72.9101 84484.9 -48 10 0 0.104853 0.74948 -0.436027 -0.200384
DTL_CEL 89.8133 28 28 0.0212224 -72.9101 73.1977 87251.7 -47.5 10 0 0.105428 0.757488 -0.423767 -0.201319
DTL_CEL 90.3217 28 28 0.0245936 73.1977 -73.4992 89953.5 -47 10 0 0.106024 0.764331 -0.413161 -0.201532
DTL_CEL 90.8475 28 28 0.0280032 -73.4992 73.7703 92694.2 -46.5 10 0 0.10664 0.770931 -0.402823 -0.201301
DTL_CEL 91.3909 28 28 0.0314415 73.7703 -74.0554 95477.4 -46 10 0 0.107276 0.777317 -0.392802 -0.200986
SET_ADV 60
DTL_CEL 91.952 28 28 0.0349097 -74.0554 74.3058 98302.2 -45.5 10 0 0.107934 0.783481 -0.38306 -0.200411
DTL_CEL 92.5298 28 28 0.0379774 74.3058 -74.5691 100888 -45 10 0 0.108611 0.787243 -0.377176 -0.2003
DTL_CEL 93.124 28 28 0.0410704 -74.5691 74.8056 103520 -44.5 10 0 0.109308 0.790922 -0.371402 -0.200107
DTL_CEL 93.7349 28 28 0.0442596 74.8056 -75.056 106233 -44 10 0 0.110024 0.794763 -0.365289 -0.199539
DTL_CEL 94.3626 28 28 0.0473999 -75.056 75.2835 108961 -43.5 10 0 0.11076 0.798277 -0.35973 -0.199184
DTL_CEL 95.0078 28 28 0.0507833 75.2835 -75.5255 111880 -43 10 0 0.111516 0.802715 -0.352626 -0.198381
DTL_CEL 95.6703 28 28 0.0539705 -75.5255 75.7413 114710 -42.5 10 0 0.112293 0.80606 -0.347306 -0.197902
DTL_CEL 96.3499 28 28 0.0571744 75.7413 -75.9704 117592 -42 10 0 0.11309 0.809328 -0.342081 -0.197341
DTL_CEL 97.0469 28 28 0.0603973 -75.9704 76.1784 120527 -41.5 10 0 0.113907 0.812524 -0.336963 -0.196748
DTL_CEL 97.7612 28 28 0.0636044 76.1784 -76.3982 123492 -41 10 0 0.114744 0.815493 -0.332338 -0.196611
SET_ADV 40
DTL_CEL 98.4937 28 28 0.0670821 -76.3982 76.5998 126672 -40.5 10 0 0.115603 0.819426 -0.325941 -0.195469
DTL_CEL 99.244 28 28 0.0703571 76.5998 -76.8124 129777 -40 10 0 0.116483 0.822424 -0.321067 -0.19464
DTL_CEL 100.013 28 28 0.0738459 -76.8124 77.0048 133061 -39.5 10 0 0.117384 0.826103 -0.315053 -0.193443
DTL_CEL 100.8 28 28 0.0771374 77.0048 -77.209 136276 -39 10 0 0.118307 0.828909 -0.31049 -0.192635
DTL_CEL 101.604 28 28 0.0804404 -77.209 77.395 139551 -38.5 10 0 0.11925 0.831652 -0.306016 -0.191796
DTL_CEL 102.428 28 28 0.0839794 77.395 -77.5934 143023 -38 10 0 0.120216 0.835125 -0.30028 -0.190411
DTL_CEL 103.271 28 28 0.0875319 -77.5934 77.7738 146561 -37.5 10 0 0.121204 0.838521 -0.29462 -0.188871
DTL_CEL 104.132 28 28 0.0908581 77.7738 -77.9644 150025 -37 10 0 0.122213 0.841052 -0.290457 -0.187938
LATTICE_END
END
```

“*LATTICE 2 1*” defines the periodic focusing lattices, 2 is the number of element per lattice (2 DTL cell), 1 is the number of lattice per macro-lattice (generally 1). And TRACE_WIN imposes a phase advance linear continuity between each “SET_ADV” command. This linear continuity can be per lattice or per meter according to the option “*Linear phase advance per meter*” of the “*Main*” page. This option is very useful, if you have lattice length discontinuity in order to keep continuity in the phase advance law per meter.

TRACE_WIN calculates the quadrupole gradients to obtain the asked zero-current phase advance law, using the [lattice transfer matrix](#).

Warning: it is not always possible to find the required quadrupole gradients.



Copy following examples:

[examples\DTL_1.ini](#)

[examples\DTL_1.dat](#)

[examples\DTL_1.cal](#)

-

-

Matching commands

MATCH_FAM_GRAD f_n, n
MATCH_FAM_FIELD f_n, n
MATCH_FAM_PHASE f_n, n
MATCH_FAM_LFOC f_n, n
MATCH_FAM_LENGTH f_n, n

f_n is the family or section number and n is the matching element number. When one of these commands precedes an element, it is used (modified) to match the section f_n . For example, if two independent cavities are needed to match the beam at the entrance of a section two commands “*MATCH_FAM_FIELD*” or “*MATCH_FAM_PHASE*” or “*MATCH_FAM_LFOC*”, with different matching element numbers n , have to be placed before these cavity entries. But if these two cavities must have the same field, they must have the same matching element number n . If you need for the matching five different quadrupoles and 2 different cavities, then the last command “*MATCH_FAM...*” must have $n = 7$. All combinations are possible. You can use the number of quadrupole, the number of cavity and the number of quadrupole or cavity coupled, you want. The coupled elements are not necessary consecutive. Two examples are presented below. These commands can also be used to match the beam at the entrance of the linac. The only constraint about the family numbers f_n is to have a different number per matching family, you can use for example 5,4,8,1...If you set $n=0$ in each comment of a matching, TRACE_WIN sets automatically a different number for each command, but no coupling are possible between elements.

| | |
|-------------------|--|
| MATCH_FAM_GRAD: | the quadrupole gradient is adjusted. |
| MATCH_FAM_FIELD: | the field is adjusted. |
| MATCH_FAM_PHASE: | the synchronous phase is adjusted. |
| MATCH_FAM_LFOC: | the longitudinal focalization is adjusted by moving the field and the synchronous phase and the energy gain is kept. |
| MATCH_FAM_LENGTH: | the element length is adjusted. |

Matching Way

TRACE_WIN is able to match the beam at the entrance of the linac or between the different sections. In these two cases the criterion for a good matching is either keep the longitudinal and transverse phase advances as smooth as possible or have at the input and the output lattice the same [Twiss parameters](#) (taking into account the [acceleration parameters](#)).

TRACE_WIN is also able to impose Twiss parameters at a position in the linac. In this case you have to insert a [SET_TWISS](#) command behind the matching commands.

These three processes are named “*Optimization*”, and can be stop by using the menu “*Stop*”.

Optimization

To make the optimization of the beam at the entrance of the linac or between the different sections, TRACE_WIN calculates the phase advance on N lattices, using the option “*Nbr of phase advance period to smooth*” of “*Matching*” page “. The choice of the criterion (Twiss parameter or smooth phase advances) depends of this option. If the number of lattice is too small (below 6) or if the number of lattice to optimize is set to 0, the [Twiss criterion](#) is used. You can watch the optimization process by plotting the phase advance chart, “*Beam*” and starting the “*Synch.*” option in the chart. The results of the optimizations can be found in the [Results file](#). The optimization

automatically stops when the criterion reaches “**Max. Number of iteration**” defined in the “**Matching**” page, but it can be stopped before in the menu “**Stop/Optimize**”. During a matching some other criterions can be included in order to control the beam size, the beam separation, the beam transverse position or the beam emittance growth...(See [matching commands](#)) In this case all the different criterions are added.

☐ Phase advance criterion:



$$\text{With } d^2 \sigma_{i,j} = \frac{\sigma_{i+1,j} + \sigma_{i-1,j} - 2 \sigma_{i,j}}{\sigma_{i,j}}$$

With $M=2$ for DC beam (x,y) or $M=3$ for bunched beam (x,y,z), and N is the number of lattice to optimize. $\sigma_{i,\varphi}$ is in °/m.

See also the [phase advance](#) definition and setting

Matching section or family

If the option “**Matching with family & Twiss commands**” of page “**Matching**” is checked, TRACE_WIN changes some quads and cavities strength or element lengths pointed by “**MATCH_FAM...**” commands, in order to match the beam between two sections. The syntax is not always very easy to use because it has been defined in order to be able to represent most of the cases. Some examples are shown below to help you and if you meet some difficulties yet, send an Email with your data file attached to the authors.

[Matching example 1](#)

[Matching example 2](#)

[Matching example 3](#)

[Matching example 4](#)

[Matching example 5](#)

Copy following examples:

[examples\Supra_1.ini](#)

[examples\Supra_1.dat](#)

[examples\Supra_1.cal](#)

[examples\DTL_2.ini](#)

[examples\DTL_2.dat](#)

[examples\DTL_2.cal](#)

Matching example 3

Matching line calculation between a RFQ and a 352 MHz DTL. Two 704MHz buncher cavities and 4 quadrupoles are used in order to match the beam to the structure. A command [SET_SIZE_MAX](#) has been include in order to try to reduce the beam size. To help the matching optimization two drift lengths are adjusted by using the command “**MATCH_FAM_LENGTH**”.

```

, ***** RFQ-DTL Matching line *****
,
DRIFT 1e-05 10
FREQ 704
SET_SIZE_MAX 0.025 25 4.5 4.5 25 1
DRIFT 20 10
MATCH_FAM_GRAD 1 0
QUAD 56 55.9938 10
MATCH_FAM_LENGTH 1 0
DRIFT 58.85 10
MATCH_FAM_FIELD 1 0
GAP 251000 -90 10 0 0 0 0 0 0
DRIFT 58.85 10
MATCH_FAM_GRAD 1 0
QUAD 56 -56.1374 10
MATCH_FAM_LENGTH 1 0
DRIFT 50.0 10
MATCH_FAM_GRAD 1 0
QUAD 56 62.7152 10
DRIFT 58.85 10
MATCH_FAM_FIELD 1 0
GAP 282000 -90 10 0 0 0 0 0 0
DRIFT 58.85 10
MATCH_FAM_GRAD 1 0
QUAD 56 -56.5712 10
DRIFT 70 10
FREQ 352
QUAD 28 71.3977 10

, ***** DTL *****
,
LATTICE 2 1
SET_ADV 46
DTL_CEL 87.5277 28 28 0.00548673 71.3977 -71.6524 74699.4 -50 10 0 0.102751 0.723252 -0.475485 -0.194532
DTL_CEL 87.9518 28 28 0.0080072 -71.6524 71.9632 76721.8 -49.5 10 0 0.103248 0.726628 -0.470582 -0.195991
DTL_CEL 88.3919 28 28 0.0112292 71.9632 -72.2903 79282.7 -49 10 0 0.103764 0.734598 -0.458632 -0.197949
DTL_CEL 88.8487 28 28 0.0144758 -72.2903 72.5931 81864.3 -48.5 10 0 0.104299 0.742159 -0.447189 -0.99371
DTL_CEL 89.3222 28 28 0.0177645 72.5931 -72.9101 84484.9 -48 10 0 0.104853 0.74948 -0.436027 -0.200384

```

```
DTL_CEL 89.8133 28 28 0.0212224 -72.9101 73.1977 87251.7 -47.5 10 0 0.105428 0.757488 -0.423767 -0.201319
DTL_CEL 90.3217 28 28 0.0245936 73.1977 -73.4992 89953.5 -47 10 0 0.106024 0.764331 -0.413161 -0.201532
...
..
.
```

The calculation result is the following:

The first line result contains the 4 quadrupole gradients in T/m, the second is the cavity field corrections and the last the drift lengths in mm.

```
Matching_Between_Section_0_to_1
56.3399 -56.1173 62.8806 -56.4843
0.934963 0.991466
0.00834559 0.000887792 0.0294733 0.0874924
```

Matching example 4

Matching line calculation between a two super conducting cavity families by adjusting the synchronous phases of the last cavities of the first family and the first cavities of the second family. Four quadrupoles are also adjusted in order to match the beam.

```
.....
DRIFT 325 100
QUAD 400 10 100
DRIFT 400 100
QUAD 400 -10 100
DRIFT 950 100
NCELLS 1 5 0.6579000 1.11349e+07 -12.331 100 0 0.1801715 0.2934260 -10.0848222 14.7245834 ...
DRIFT 475 100
NCELLS 1 5 0.6579000 1.11616e+07 -10.584 100 0 0.1801715 0.2934260 -10.1057623 14.7737516 ...
DRIFT 475 100
NCELLS 1 5 0.6579000 1.11873e+07 -8.890 100 0 0.1801715 0.2934260 -10.1259198 14.8211363 ...
DRIFT 625 100
SET_ADV 70
DRIFT 325 100
MATCH_FAM_GRAD 8 1
QUAD 400 10 100
DRIFT 400 100
MATCH_FAM_GRAD 8 2
QUAD 400 -10 100
DRIFT 950 100
MATCH_FAM_PHASE 8 3
NCELLS 1 5 0.6579000 1.12121e+07 -7.248 100 0 0.1801715 0.2934260 -10.1453288 14.8668118 ...
DRIFT 475 100
MATCH_FAM_PHASE 8 3
NCELLS 1 5 0.6579000 1.12361e+07 -5.657 100 0 0.1801715 0.2934260 -10.1640222 14.9108503 ...
DRIFT 475 100
MATCH_FAM_PHASE 8 3
NCELLS 1 5 0.6579000 1.12593e+07 -4.114 100 0 0.1801715 0.2934260 -10.1820316 14.9533216 ...
DRIFT 625 100
LATTICE_END
DRIFT 1187.35 100

; Second superconducting family
LATTICE 13 1
SET_ADV 85
DRIFT 325 100
MATCH_FAM_GRAD 8 4
QUAD 500 10 100
DRIFT 400 100
MATCH_FAM_GRAD 8 5
QUAD 500 -10 100
DRIFT 975 100
MATCH_FAM_PHASE 8 6
NCELLS 1 5 0.8458000 1.14453e+07 -94.435 100 0 0.1560511 0.2270898 -10.0611253 13.2916014 ...
DRIFT 525 100
MATCH_FAM_PHASE 8 6
NCELLS 1 5 0.8458000 1.14816e+07 -92.773 100 0 0.1560511 0.2270898 -10.0883587 13.3415267 ...
DRIFT 525 100
MATCH_FAM_PHASE 8 6
NCELLS 1 5 0.8458000 1.15177e+07 -91.116 100 0 0.1560511 0.2270898 -10.1153599 13.3910179 ...
DRIFT 525 100
MATCH_FAM_PHASE 8 6
NCELLS 1 5 0.8458000 1.15536e+07 -89.464 100 0 0.1560511 0.2270898 -10.1421056 13.4400346 ...
DRIFT 650 100
DRIFT 325 100
QUAD 500 10 100
DRIFT 400 100
QUAD 500 -10 100
DRIFT 975 100
...
```

Matching example 5

TRACE_WIN sets Twiss parameter to the entrance of a RFQ. To impose Twiss parameter in your structure by adjusting elements you have to use the same “MATCH_FAM...” commands, but a [SET_TWISS](#) command has been put in front of the element where you want to impose Twiss parameters at its output. The first parameter of [SET_TWISS](#) and “MATCH_FAM...” commands is the family or the section number.

```
...
DRIFT 50.1 100
MATCH_FAM_FIELD 1 0
GAP 154335 -90 100 0 0
DRIFT 50.1 100
MATCH_FAM_LENGTH 1 0
DRIFT 60 100
MATCH_FAM_GRAD 1 0
QUAD 60 16.9089 100
MATCH_FAM_LENGTH 1 0
DRIFT 40 100
DRIFT 20 100
MATCH_FAM_GRAD 1 0
QUAD 60 -16.9089 100
DRIFT 50.1 100
MATCH_FAM_FIELD 1 0
GAP 154335 -90 100 0 0
DRIFT 50.1 100
DRIFT 20 100
MATCH_FAM_GRAD 1 0
QUAD 60 16.9089 100
MATCH_FAM_LENGTH 1 0
DRIFT 40 100
DRIFT 20 100
DIAPHRAGME 2.7 2.9 1
FREQ 352.21
DRIFT 0.00001 10
RFQ_CELL 120000 4.96362 0 1 109.954 -60 3
RFQ_CELL 120000 4.96362 0.114946 1.13305 27.4884 -60 4
SET_TWISS 1 -1.5440 0.2462 0.7924 0.1898 -0.0754 0.9523 0 0 0 0 0
RFQ_CELL 120000 4.96362 0.114946 1.13305 27.4884 -60 2
RFQ_CELL 120000 4.96362 0.115356 1.13354 27.5261 -60 -2
RFQ_CELL 120000 4.96362 0.115799 1.13406 27.567 -59.9802 2
RFQ_CELL 120000 4.96362 0.116304 1.13465 27.611 -59.9208 -2
RFQ_CELL 120000 4.96362 0.116874 1.13532 27.6552 -59.822 2
RFQ_CELL 120000 4.96362 0.117509 1.13609 27.6995 -59.6843 -2
RFQ_CELL 120000 4.96362 0.118213 1.13692 27.7441 -59.5081 2
RFQ_CELL 120000 4.96362 0.118987 1.13786 27.789 -59.2943 -2
RFQ_CELL 120000 4.96362 0.119833 1.13888 27.8343 -59.0435 2
RFQ_CELL 120000 4.96362 0.120755 1.13999 27.8799 -58.7569 -2
...
..
.
```

Matching example 1

4 quadrupoles and 2 cavities are used to match the beam between section 1 and 2. The longitudinal matching is done by a cavity field modification. This is the simplest case.

| | |
|--------------------------------|---|
| LATTICE 7 1 | «section 1» |
| DRIFT 100 100 | |
| MATCH_FAM_GRAD 1 1 | |
| QUAD 150 10 100 | «first quadrupole used for matching » |
| DRIFT 178.56 100 | |
| MATCH_FAM_GRAD 1 2 | |
| QUAD 150 -10 100 | «second quadrupole used for matching» |
| DRIFT 488.64 100 | |
| MATCH_FAM_FIELD 1 3 | |
| GAP_M 2.94577e+06 -30.5122 100 | «first cavity used for matching, the field is adjusted» |
| DRIFT 722.08 100 | |

| | |
|--------------------------------|--|
| LATTICE 9 1 | «section 2» |
| DRIFT 150 100 | |
| MATCH_FAM_GRAD 1 4 | |
| QUAD 350 10 100 | «third quadrupole used for matching» |
| DRIFT 100 100 | |
| MATCH_FAM_GRAD 1 5 | |
| QUAD 350 -10 100 | «fourth quadrupole used for matching» |
| DRIFT 935.126 100 | |
| MATCH_FAM_FIELD 1 6 | |
| GAP_M 3.04084e+06 -30.5122 100 | «second cavity used for matching, the field is adjusted» |
| DRIFT 1385.81 100 | |
| GAP_M 3.07847e+06 -30.5122 100 | |
| DRIFT 729.063 100 | |

Matching example 2

4 quadrupoles and 4 cavities are used for the matching between section 1 and 2. The longitudinal matching is done by a cavity phase adjustment. Here, the cavities are coupled 2 by 2.

| | |
|--------------------------------|--|
| LATTICE 8 1 | «section 1» |
| ... | |
| DRIFT 100 100 | |
| MATCH_FAM_GRAD 1 1 | |
| QUAD 150 10 100 | «first quadrupole used for matching » |
| DRIFT 178.56 100 | |
| MATCH_FAM_GRAD 1 2 | |
| QUAD 150 -10 100 | «second quadrupole used for matching» |
| DRIFT 488.64 100 | |
| MATCH_FAM_PHASE 1 3 | |
| GAP_M 2.94577e+06 -30.5122 100 | «first cavity used for matching, the phase is adjusted. $n=3$ ». |
| MATCH_FAM_PHASE 1 3 | |
| GAP_M 2.94577e+06 -30.5122 100 | «second cavity used for matching. It will conserve the same phase as the first cavity. $n=3$ » |
| DRIFT 722.08 100 | |
| LATTICE 10 1 | «section 2» |
| DRIFT 150 100 | |
| MATCH_FAM_GRAD 1 4 | |
| QUAD 350 10 100 | «third quadrupole used for matching» |
| DRIFT 100 100 | |
| MATCH_FAM_GRAD 1 5 | |
| QUAD 350 -10 100 | «fourth quadrupole used for matching» |
| DRIFT 935.126 100 | |
| MATCH_FAM_PHASE 1 6 | |
| GAP_M 3.04084e+06 -30.5122 100 | «third cavity used for matching, the phase is adjusted. $n=6$ ». |
| DRIFT 1385.81 100 | |
| MATCH_FAM_PHASE 1 6 | |
| GAP_M 3.04084e+06 -30.5122 100 | «fourth cavity used for matching, It will conserve the same phase as the third cavity. $n=6$ » |
| DRIFT 729.063 100 | |

Set Twiss parameters

SET_TWISS f_n , α_x , $\beta_x(mm/mrad)$, α_y , $\beta_y(mm/mrad)$, α_z , $\beta_z(mm/mrad)$, $k_{ax}, k_{bx}, k_{ay}, k_{by}, k_{az}, k_{bz}$

The forces of elements pointed with the matching commands "MATCH_FAM_XXX" are adjusted to impose the Twiss parameters given by the "SET_TWISS" command at the output element following the command. More than one command "SET_TWISS" can be used in the same optimization. In this case, use the same f_n parameter.

□ Twiss criterion:

$$vcr = \frac{1}{2M} \sum_{j=1}^M \left(\sqrt{\frac{R_j + \sqrt{R_j^2 - 4}}{2}} \right) - 1$$

With $R_j = \beta_j \gamma_{j0} + \beta_{j0} \gamma_j - 2 \alpha_j \alpha_{j0}$

f_n is the section or family number.

$\beta_j, \gamma_j, \alpha_j$ being the beam Twiss parameters of the space phase xx', yy', zz' .

The 6 following parameters “ k ” are optional, set one to “1” allows to not taking account of the corresponding Twiss parameters

Write “SET_TWISS f_n ” corresponds to write “SET_TWISS $f_n 0 0 0 0 0 1 1 1 1 1 1$ ”. That disables the SET_TWISS command and allows different other optimisations by inserting another commands like for example “SET_SIZE”.

Copy following example:

[examples\LBE_1.ini](#)
[examples\LBE_1.dat](#)
[examples\LBE_1.cal](#)
[examples\solenoid1.bsz](#)
[examples\solenoid1.ouv](#)
[examples\solenoid2.bsz](#)
[examples\solenoid2.ouv](#)

Set beam centroid position

SET_POSITION $k(m^{-1}), x(mm), x'(mrad), y(mm), y'(mrad)$

x, x', y, y' are the centroid beam positions imposed at the point where this command appears, k is used in the criterion calculation.

$$\square \text{ Set position criterion: } vcr = \frac{k}{2} \cdot \sqrt{(x - x_0)^2 + (y - y_0)^2 + (x' - x'_0)^2 + (y' - y'_0)^2}$$

With $x(mm), x'(mrad), y(mm), y'(mrad)$ being the beam centroid transverse positions at the place where the command “SET_POSITION” appears and $x_0(mm), x'_0(mrad), y_0(mm), y'_0(mrad)$ being the beam imposed positions.

Set achromatic line

SET_ACHROMAT $k f1 f2$

Located after a deviation for instance this command allows to make achromatic the preceding line. k is used to balance the criterion calculation

($f1=1$ means set achromatic position).

($f2=1$ means set achromatic angle).

($f1=f2$ means set achromatic position and angle).

Example:

```
DRIFT 300 100
START_ACHROMAT
EDGE 22.5 500 100 0.45 2.8 50 0
BEND 45 500 0 50 0
EDGE 22.5 500 100 0.45 2.8 50 0
DRIFT 200 100
MATCH_FAM_GRAD 4
QUAD 200 2
DRIFT 200 100
EDGE 22.5 500 100 0.45 2.8 50 0
BEND 45 500 0 50 0
EDGE 22.5 500 100 0.45 2.8 50 0
DRIFT 107 100
SET_ACHROMAT 1.0 1 1
DRIFT 100 100
SET_TWISS 4
DRIFT 100 100
end
```

The command [SET_TWISS](#) is always obligatory to allow the matching of the gradupole. And the option “*Matching with family & Twiss commands*” of page “*Matching*” has to be checked

Copy following example:

[examples\Achromat_2.ini](#)

[examples\Achromat_2.dat](#)
[examples\Achromat_2.cal](#)

[examples\Achromat_1.ini](#)
[examples\Achromat_1.dat](#)

Set maximum beam max

SET_SIZE_MAX $k, N, x(mm), y(mm), \varphi(^{\circ})/z(mm), k2 (0/1)$

x, y, φ are the imposed beam size max in the N elements following this command, N, k and $k2$ is used in the criterion calculation.

$$vcr = \frac{k}{M} \cdot \left(\exp\left(-\left(\frac{x_0}{x}\right)^4\right) + \exp\left(-\left(\frac{y_0}{y}\right)^4\right) + \exp\left(-\left(\frac{\varphi_0}{\varphi}\right)^4\right) \right)$$

□ Set size max criterion:

With $x(mm), y(mm), \varphi(^{\circ})$ being the beam maximum sizes in the N elements following the command “SET_SIZE_MAX”, and $x_0(mm), y_0(mm), \varphi_0(^{\circ})$ being the beam imposed sizes. If φ is lower than 0 the longitudinal size used is $z(mm)$. If one of these last parameters is set to 0, no optimization is done on this size and M is reduced by one. If $k2=0$, the transverses sizes are calculated without taking account the beam centroid position. If one of the parameters x, y or φ is equal to 0, no optimization is done on this size and M is reduced by one. The sizes are the effective beam sizes (rms*sqrt(5) for bunched beam or rms*sqrt(4) for CW beam)

Set beam size

SET_SIZE $k, x(mm), y(mm), \varphi(^{\circ})/z(mm), k2 (0/1)$

x, y, φ are the imposed beam size in the output element following this command, N, k is used in the criterion calculation. If $k2=0$, the transverses sizes are calculated without taking account the beam centroid position.

$$vcr = \frac{k}{M} \cdot \sqrt{\left(\frac{x - x_0}{x_0}\right)^2 + \left(\frac{y - y_0}{y_0}\right)^2 + \left(\frac{\varphi - \varphi_0}{\varphi_0}\right)^2}$$

□ Set size criterion:

With $x(mm), y(mm), \varphi(^{\circ})$ being the beam sizes in the output elements following the command “SET_SIZE”, and $x_0(mm), y_0(mm), \varphi_0(^{\circ})$ being the beam imposed sizes. If φ is lower than 0 the longitudinal size used is $z(mm)$. If one of these last parameters is equal to 0, no optimization is done on this size and M is reduced by one. The sizes are the effective beam sizes (rms*sqrt(5) for bunched beam or rms*sqrt(4) for CW beam)

Set beam separation

SET_SEPARATION k, S_x, S_y

S_x, S_y are the centroid beam positions divided by the imposed beam size max at the point where this command appears, k is used in the criterion calculation.

$$vcr3 = k \cdot \sqrt{(S_{x_0} - S_x)^2 + (S_{y_0} - S_y)^2}$$

□ Set separation criterion:

Minimize emittance growth

MIN_EMIT_GROW *k, N, ex, ey, ez, f*

k and *N* are used in the criterion calculation.

ex, *ey*, *ez* and *f* are optional: if *ex* equal to 1 the criterion doesn't take into account ϵ_x beam emittance. That is the same way for *ey* and *ez*. *f* parameter define if the criterion calculation is relative or not (see formula below).

□ The emittance growth criterion:

$$\nu_{CR} = k \cdot \sum_N \sqrt{\left(\frac{\epsilon_x - \epsilon_{x0}}{\epsilon_{x0}}\right)^2 + \left(\frac{\epsilon_y - \epsilon_{y0}}{\epsilon_{y0}}\right)^2 + \left(\frac{\epsilon_z - \epsilon_{z0}}{\epsilon_{z0}}\right)^2}$$

□ If (*f*=0)

$$\nu_{CR} = k \cdot \left(|\epsilon_x| + |\epsilon_y| + |\epsilon_z| \right) \text{ at the command position}$$

□ If (*f*=1)

With ϵ_{x0} , ϵ_{y0} , ϵ_{z0} being the beam emittances where this command appears, and ϵ_x , ϵ_y , ϵ_z being the beam emittances after *N* elements.

Magnetic steerer**STEERER *B_x(T), B_y(T), B_{max}***

The magnetic steerer is inserted in the element (quadrupole) placed just after the instruction “Steerer” (keeping the same length).

$$x' = x' - \frac{q \cdot \Delta s \cdot B_y c}{mc^2 \beta \gamma} \quad \text{And} \quad y' = y' + \frac{q \cdot \Delta s \cdot B_x c}{mc^2 \beta \gamma}$$

Where *x'* and *y'* being respectively the horizontal and vertical beam centroid slope.

Used in diagnostic optimization, *B_{max}* is the maximum limit for *B_x* and *B_y*, if it's greater than zero.

Chopper**CHOPPER *N, U(V), D(mm), C(mm), p(0/1)***

The chopper is inserted in the *N* elements placed just after the instruction “Chopper” (keeping the same length). *U* is the voltage between axis and plates and *C* is the chopper transverse position. $\pm D$ is the distance between axis and plates

$$\text{If } p=0 \quad x' = x' + \frac{q \cdot \Delta s \cdot U}{mc^2 \beta^2 \gamma \cdot D} \quad \text{and if } p=1 \quad y' = y' + \frac{q \cdot \Delta s \cdot U}{mc^2 \beta^2 \gamma \cdot D}$$

Where *x'* and *y'* are respectively the horizontal and vertical beam centroid slope.

See following example:

[examples\Chopper_1.ini](#)

[examples\chopper_1.dat](#)

Input beam errors

- **ERROR_BEAM_STAT** *r(0/1), dx(mm), dy(mm), dφ(°), dxp(mrad), dyp(mrad), de(MeV), dEx(%), dEy(%), dEz(%), mx(%), my(%), mz(%), dib(mA)*
- **ERROR_BEAM_DYN**

Four kind of error can be set:

- ▣ **Beam displacement:** (dx, dy, dp, dxp, dyp, de) The beam input position is not centered.
- ▣ **Emittance growth:** (dEx, dEy, dEz) The input beam emittance is increased by a percentage.
- ▣ **Beam mismatch:** (mx, my, mz) The input beam is mismatched by a percentage. A 20 % mismatch in x plane means α_x and β_x are multiplied by 1.2².
- ▣ **Beam Current error:** (dib) Allows to study the effect of the input beam current variation.

This command concerns only the input beam. If r is equal to 0, the errors are constant and equal to each value of the command line. If r is equal to 1, the errors are distributed randomly; each value of the command line is then the maximum error. The errors are uniformly distributed between +/- the maximum error.

[NCP CPL STAT DYN meaning](#)

Quadrupole errors

- **ERROR_QUAD_NCPL_STAT** $N, r(0/1), dx(mm), dy(mm), \varphi_x(^{\circ}), \varphi_y(^{\circ}), \varphi_z(^{\circ}), dg(\%), dz(mm)$
- **ERROR_QUAD_NCPL_DYN**
- **ERROR_QUAD_CPL_STAT**
- **ERROR_QUAD_CPL_DYN**

dx and dy being respectively the horizontal and vertical magnetic element displacement. dz being the longitudinal shift. $\varphi_x, \varphi_y, \varphi_z$ being respectively the quadrupole rotation around x, y, z-axis and dg being the gradient amplitude error. These errors are applied in the N elements following this command, excepted, if a new error command appears. If r is equal to 0, the errors are constant and equal to each value of the command line. If r is equal to 1, the errors are distributed randomly, each value of the command line is then the maximum error. The errors are uniformly distributed between +/- the maximum error.

This error set affects:

- Quadrupole
- Solenoid
- Quadrupole of DTL
- Field map if defined as static magnetic field

[NCP CPL STAT DYN meaning](#)

Copy following examples:

[examples\Supra_4.ini](#)
[examples\Supra_4.dat](#)
[examples\Supra_4.cal](#)

[examples\Supra_5.ini](#)
[examples\Supra_5.dat](#)
[examples\Supra_5.cal](#)

Cavity errors

- **ERROR_CAV_NCPL_STAT** $N, r(0/1), dx(mm), dy(mm), \varphi_x(^{\circ}), \varphi_y(^{\circ}), E(\%), \varphi(^{\circ}), dz(mm)$
- **ERROR_CAV_NCPL_DYN**
- **ERROR_CAV_CPL_STAT**
- **ERROR_CAV_CPL_DYN**

dx and dy being respectively the horizontal and vertical electric element displacement. dz being the longitudinal shift, φ_x, φ_y being respectively the cavity rotation around x, y-axis. E being the field amplitude error. φ being the field phase error. These errors are applied in the N elements following this command, excepted, if a new error command appears. If r is equal to 0, the errors are constant and equal to each value of the command line. If r is equal to 1, the errors are distributed randomly, each value of the command line is then the maximum error. The errors are uniformly distributed between +/- the maximum error.

This error set affects:

- Bunched cavity
- Cavity multi-gap
- DTL ([Warning the DTL are only concerns by the field errors](#))
- Sinus cavity
- Field map if not defined as static magnetic field

[NCP CPL STAT DYN meaning](#)

-

Bend errors

- **ERROR_BEND_NCPL_STAT** $N, r(0/1), dx(mm), dy(mm), \varphi_x(^{\circ}), \varphi_y(^{\circ}), \varphi_z(^{\circ}), dg(\%), dz(mm)$
- **ERROR_BEND_NCPL_DYN**
- **ERROR_BEND_CPL_STAT**
- **ERROR_BEND_CPL_DYN**

dx and dy being respectively the horizontal and vertical magnetic element displacement. dz being the longitudinal shift. $\varphi_x, \varphi_y, \varphi_z$ being respectively the bend rotation around x, y, z -axis and dg being the magnetic field amplitude error. These errors are applied in the N elements following this command, excepted, if a new error command appears. If r is equal to 0, the errors are constant and equal to each value of the command line. If r is equal to 1, the errors are distributed randomly, each value of the command line is then the maximum error. The errors are uniformly distributed between +/- the maximum error.

This error set affects only magnet bend and edge

Valid only if the BEND element is not cut in several parts.

Warning, the “ERROR_BEND_CPL_X” commands concern only the magnetic field amplitude errors, all other errors (shift, rotation,...) are not couple.

[NCP CPL STAT DYN meaning](#)

-

See detail about [bend error treatment](#).

Adjust commands

-

- **ADJUST** $N, v, n, min, max, first_step$

Red parameters are optional.

This command has to be associated with a [Diagnostic elements](#) number N . v , an integer, is the v -th variable to be adjust in the next element list of variable. As an exemple if $v=2$ and the next element is a quad then the gradient will be adjusted. n (if it's different from zero) allows to link ADJUST commands each other, therefore, two ADJUST commands with the same n will give the same value to the v variable point. v lower than zero allows to replace both command “ADJUST_DPHASE” and “ADJUST_DFIELD” when v points to phase or field variable.

New in version 2.2.0.16 : *first_step* parameter has been added in order to choose the first step of optimization.

New in version 2.2.1.9

The n parameter allow to link two variables: Now if you set n parameter to opposite sign you'll get opposite variation for the linked variables. For exemple two drifts coupled with ($n=1$ & $n=-1$) will keep the same total sum length.

5 consecutives ADJUST command can be set.

(See [Adjust and diagnostic examples](#))

- **ADJUST_STEERER** *N, max, first_step*
- **ADJUST_STEERER_BX** *N, max, first_step*
- **ADJUST_STEERER_BY** *N, max, first_step*

These commands have to be associated with [Diagnostic elements](#) number *N* and should be placed before a [STEERER](#) command. It allows to adjust the horizontal and/or vertical magnetic field steering.

Particular case of steerer adjustment: When the number of steerers corresponds to the number of BMP, no optimisation is performed and the resolution of the system is directly made by a matrix inversion

New in version 2.2.0.20

Input beam parameters can be adjusted in order to fit with a set of diagnostic values at different positions of a line (Example adjusted beam emittances and Twiss parameters at the input of a simulated line in order to obtain the beam sizes measured at different positions with real beam)

- **ADJUST_BEAM_TWISS** *N, AlpX_flag, betX_flag, AlpY_flag, BetY_flag, AlpZ_flag, betZ_flag*
If flag is set to 1 the selected Twiss parameter will be adjusted, if you want to have alpX=AlpY set alpX_flag=1 and alpY_flag=2, same way for Bet_flag)
- **ADJUST_BEAM_EMIT** *N, Ex_flag, Ey_flag, Ez_flag*
If flag is set to 1 the selected emittance will be adjusted, if you want to have Ex=Ey set Ex_flag=1 and Ey_flag=2)
- **ADJUST_BEAM_CENTROID** *N, X_flag, Xp_flag, Y_flag, Yp_flag, Z_flag, Zp_flag*
If flag is set to 1 the selected beam centroid parameter will be adjusted, if you want to have X=Y set X_flag=1 and Y_flag=2, same way for Xp_flag)
- **ADJUST_BEAM_CURRENT** *N, I_flag*

All “*ADJUST_BEAM_XXX*” commands must be located in front of the first element. Some elements can be also adjusted in the same matching process.

Adjust and diagnostic examples

Example 1

Beam alignment from a RFQ to a DTL. Two [ADJUST](#) commands are associated with two [STEERER](#) and with two [Diagnostic elements](#).

```

DRIFT 0.00001 6.5
DRIFT 56.0 6.5
ADJUST_STEERER 1                                «Adjust the steerer inside the following quadrupole »
STEERER -.01 0.01
QUAD_ERROR 1 0 0.1 0.1 0.3 0.3 0.5 0           «Include 0.1mm misalignment and 0.3° rotation errors in the quadrupole »
QUAD 56.0 55.0 6.5
DRIFT 2.15 6.5
DRIFT 58.85 6.5
GAP_M 180000 -90 6.5
DRIFT 58.85 6.5
ADJUST_STEERER 1                                «Adjust the steerer inside the following quadrupole »
STEERER .01 0.01
QUAD 56.0 -55.0 6.5
DRIFT 82.9 6.5
DIAG_POSITION 1 0 0                             « Position monitor with x and y position imposed to 0 »
DRIFT 20 6.5
QUAD 56.0 55.0 6.5
DRIFT 58.85 6.5
GAP_M 210000 -90 6.5
DRIFT 38.85 6.5
QUAD 56.0 -55.0 6.5
DRIFT 54.3 6.5
DIAG_POSITION 1 0 0                             « Position monitor with x and y position imposed to 0 »
DRIFT 20.0 6.5
...

```

You can see the result in the result file. This result is not used in a statistical error study where an adjustment is calculated for each

random error distribution.

Diagnostic_1

-0.0661744 -0.0663664 0.0554 0.0458

Example 2

Triplet adjusting in order to match a beam size

```
ADJUST 1 2 1           « Quadrupole 1 and 3 are linked, the point variable '2' is the gradient »
QUAD 152 2.14 52
DRIFT 119 52
ADJUST 1 1 0 200 300  « Here, the quadrupole length is adjusted, but the length will keep between (200 and 300) »
QUAD 281 -2.13 52
DRIFT 119 52
ADJUST 1 2 1
QUAD 152 2.0 52
DRIFT 3190 100
APERTURE 2 2 0
DIAG_SIZE 1 2 2
DRIFT 3590 100
DIAG_SIZE 1 4 4
DRIFT 100 100
...
```

Example 3

Set achromatic line part from element 3 to achromatic diagnostic position

```
DRIFT 30 100
DRIFT 30 100
DRIFT 30 100
DRIFT 30 100
EDGE 22.5 500 100 0.45 2.8 50 0
BEND 45 500 0 50 0
EDGE 22.5 500 100 0.45 2.8 50 0
DRIFT 200 100
ADJUST 11 2 1
QUAD 100 -200 20
DRIFT 200 100
ADJUST 11 2 1
QUAD 100 200 20
DRIFT 200 100
ADJUST 11 2 1
QUAD 100 -200 20
DRIFT 200 100
EDGE 22.5 500 100 0.45 2.8 50 0
BEND 45 500 0 50 0
EDGE 22.5 500 100 0.45 2.8 50 0
DRIFT 107 100
DIAG_ACHROMAT 11 3 1 1
DRIFT 100 100
DRIFT 100 100
End
```

Example 4

Find the inout beam Twiss parameters and emittances given the rms sizes measured at the end of the line.

```
ADJUST_BEAM_CURRENT 99 1
ADJUST_BEAM_EMIT 99 1 2 0 ; Ex=Ey
ADJUST_BEAM_TWISS 99 1 1 2 2 0 0 ; Rond beam
DRIFT 0 100
DRIFT 200 100
DRIFT 200 100
QUAD 100 -200 20
DRIFT 200 100
ADJUST 11 2 1
QUAD 100 200 20
DRIFT 200 100
DIAG_SIZE 99 2.5. 3.6
DRIFT 200 100
DIAG_SIZE 99 4.2. 1.6
DRIFT 0 100
END
```

RFQ errors

- **ERROR_RFQ_CEL_NCPL_STAT** *N*, *r(0/1)*, *dR(mm)*, *d (mm)*, *E(%)*, *φ(°)*, *TEpe(mm)*, *TEpa(mm)*, *DEpe(mm)*, *DEpa(mm)*, *DELong(mm)* *TSVerti(mm)*, *TSHori(mm)*, *DSVerti(mm)*, *DSHori(mm)*, *DSLLong(m)*

- **ERROR_RFQ_CEL_NCPL_DYN** (not implemented yet)

dR is the error for the longitudinal profile, d is the error for the transverse curvature of the electrode, E is the voltage amplitude error. φ being the field phase error (uneffective).

For the transverse plane, TEpe is the perpendicular tilt error by electrode, TEpa is the parallel tilt error by electrode, DEpe is the perpendicular displacement error by electrode, DEpa is the parallel displacement error by electrode, DELong is the a longitudinal displacement error by electrode.

For each segment (block of four electrodes defined with the RFQ_GAP command), TSVerti is the vertical tilt error, TSHori is the horizontal tilt error, DSVerti is the vertical displacement error, DSHori is the horizontal displacement error and DSLong is the longitudinal displacement error.

These errors are applied in the N elements following this command, excepted, if a new error command appears. If r is equal to 0, the errors are constant and equal to each value of the command line. If r is equal to 1, the errors are distributed randomly, each value of the command line is then the maximum error. RFQ is a special element in Trace_Win and all these errors are no effect in a envelope calculation except the phase and filed errors.

[NCP CPL STAT DYN meaning](#)

Plot distribution

PLOT_DST

Indicate to the multiparticle code to store beam distribution (in the *.[plt file](#)) at the command location

Gas pressure

GAS C, N, P

Only used in multiparticle simulation, this command set the gas pressure parameter until a new command *GAS*

C is the cross section at 1 MeV (m²), *N* is the atomic number and *P* is the pressure (hPa).

You have to select ‘*gas stripping*’ or ‘*gas scattering*’ in the ‘*Multiparticle*’ page.

The cross section values are used for gas scattering calculation.

Several “gas” commands can be set simultaneously (Max. 9).

Minimize envelope variation

MIN_ENV_VARIATION k, N

Minimize the variation of the maximum and minimum transverse beam envelope in the quadrupoles included in *N* elements following the command. For bunched beam, the phase-spread variation is also minimizing in the accelerating elements. *k* is the pondering factor of this criterion.

Minimize phase variation

MIN_PHASE_VARIATION k, N, θ

During a matching procedure, where some synchronous phases of cavities are adjusted in order to match a beam, it’s often useful to limit the maximum phase amplitude variation.

k is used in the criterion calculation. θ is the max angle variation wanted during a matching, *N* is the number of elements where this

command is applied but only cavities are concerned.

NCPL CPL STAT DYN meaning

NCPL: means: No Coupled. The errors are individually applied on each element

CPL: means: Coupled. The errors are coupled on the N elements. In other words, a rotation error corresponds to an N elements block rotation.

STAT: means: Static. The effect of these errors can be detected and corrected with appropriate diagnostic and correctors. For example, beam position measurement coupled with steerers can compensate the quadrupole or cavities misalignments. Correction strategy should be known to be able to estimate their impact on beam dynamics

DYN: means: Dynamic. The effect of these errors cannot be measured and then corrected. Fortunately, they have usually lower amplitude than static errors. They are, for example, the vibrations of the elements or the RF field control errors (in phase or amplitude). The knowledge of the correction scheme is not needed to study their statistic impact. They are responsible of orbit oscillations around the corrected orbit (this notion of orbit is also extended in the longitudinal motion).

To read a particle file

The particle file from the page “**Main**” allows to define the input particle file, but if you want to force the envelope or multiparticle calculation to load a new particle file characteristics (Twiss parameters, current, emittances, centroid) at a given position in the linac, you have to include the command below followed by the full name of the [particle file](#).

READ_DST full_particle_file_name

Warning: This command is read at its position and not like other commands at the end of the following element.

Set RFQ vane geometry

- **RFQ_GEOM type N dz**

Put this command just before the first [RFQ cell](#).

According to *type* parameter:

If *type* = 0: Toutatis generates vane geometry file with 50steps/cell. It's the default case.

If *type* = 1: Toutatis reads vane geometry file

Example: **RFQ_GEOM 1 c:\my_project\rfq\My_rfq.vane**

If *type* = 2: and $N = 0$: Toutatis generates vane geometry file with a cell step = dz

Example: **RFQ_GEOM 2 0 0.001** for a 1 mm step

If *type* = 2: and $N \neq 0$: Toutatis generates vane geometry file with N step/cell

Example: **RFQ_GEOM 2 20**

For statistical error study case: (for X cases)

Example: **RFQ_GEOM 1 c:\my_project\rfq\My_rfq.vane**

Toutatis, for each case, will look for a RFQ geometry vane file called:

c:\my_project\rfq\My_rfq000001.vane for first run

c:\my_project\rfq\My_rfq000002.vane for second one

...

..

.

c:\my_project\rfq\My_rfq00000X.vane for last one

Set RFQ electrode type

- **FOUR_RODS**

This option generates electrode profile with varying transverse curvature. By default the transverse curvature is proportional to Ro (four vane type). Put this command just before the first [RFQ cell](#) element.

- **TWOTERMS**

This command allows to generate a longitudinal profile (way to modulate) which is governed by the classical 2 terms potential. In case this option is used, only m and a are taken into account, Ro is recomputed. The sinus modulation is the default in case nothing is specified. Put this command just before the first [RFQ cell](#).

Set RFQ coupling gap

RFQ_GAP $Lp(m)$, $Lg(m)$, $Sl(m)$, $St(m)$

This command defines in a RFQ structure a resonant coupling gap. The position of the command in the data file doesn't matter, if you respect 2 rules: place it before a [RFQ cell](#) element, and if you need more than one coupling gap you must avoid to put the new command before the same RFQ cell element. Lp is the longitudinal position of the center of the gap. Lg is gap width. Sl is the half-ellipse size in the beam direction and St is the half-ellipse height in the perpendicular beam direction.

To read a multiparticle output file

If you choose to run multiparticle code (Partran or Toutatis), you can avoid to run a linac part or even all the linac which has been already computed. Then, you have to indicate to TraceWin, which part doesn't have to be run, by insert at the beginning of this part the command "**READ_OUT**" followed by the full name of the [Partran or Toutatis output file](#)

READ_OUT *full_output_multiparticle_file_name*

Warning: This command is read at its position and not like other commands at the end of the following element.

Shift

SHIFT $dx(mm)$, $dy(mm)$

Move the following element

PARTRAN step calculation

PARTRAN_STEP *step1, step2*

Step1 is the new step of calculation per metre, *step2* is the new step of space-charge calculation per meter until reaching a new “*PARTRAN_STEP*” command. The default *step1* and *step2* value is put in the “*Multiparticle*” page, see “*Partran step of calcul*”. These two step concern only DRIFT and FIELD_MAP elements, all other elements are treated in 2 steps.

Magnetic or electric static field

FIELD $B_x(T)$, $B_y(T)$, $B_z(T)$ $E_x(V/m)$ $E_y(V/m)$ $E_z(V/m)$

Add to the following elements a magnetic or electric force until a new command “*FIELD*” canceling the preceding. In PARTRAN simulation, only the field command of the first element is considering.

$$\begin{aligned} x' &= x' - \frac{q \cdot \Delta s \cdot B_y c}{mc^2 \beta \gamma} & y' &= y' + \frac{q \cdot \Delta s \cdot B_x c}{mc^2 \beta \gamma} \\ x' &= x' + \frac{q \cdot \Delta s \cdot E_y}{mc^2 \beta^2 \gamma} & y' &= y' + \frac{q \cdot \Delta s \cdot E_x}{mc^2 \beta^2 \gamma} & \gamma &= \gamma + \frac{q E_z}{mc^2} \end{aligned}$$

Where x' , y' and γ being respectively the horizontal and vertical beam centroid slope and the reduced energy

Set beam phase advance

SET_BEAM_PHASE_ADV k , N , $\varphi_x(^{\circ})$, $\varphi_y(^{\circ})$, $\varphi_z(^{\circ})$

φ_x , φ_y , φ_z are the imposed beam phase advance in space charge during N elements from command. k is used in the criterion calculation.

$$\square \text{ Set beam adv criterion: } vcr = \frac{k}{M} \cdot \sqrt{\left(\frac{\varphi_x - \varphi_{x0}}{\varphi_{x0}}\right)^2 + \left(\frac{\varphi_y - \varphi_{y0}}{\varphi_{y0}}\right)^2 + \left(\frac{\varphi_z - \varphi_{z0}}{\varphi_{z0}}\right)^2}$$

With $\varphi_x(^{\circ})$, $\varphi_y(^{\circ})$, $\varphi_z(^{\circ})$ being the beam phase advances in the output elements following the command “SET_BEAM_PHASE_ADV”, and $\varphi_{x0}(^{\circ})$, $\varphi_{y0}(^{\circ})$, $\varphi_{z0}(^{\circ})$ being the imposed beam phase advances. If one of these last parameters is equal to 0, no optimization is done on this phase and M is reduced by one.

Set beam energy and phase

SET_BEAM_E0_P0 k , ΔE , $\Delta \varphi$, k_e , k_p

ΔE and $\Delta \varphi$ are the imposed delta beam energy and phase.

$$\square \text{ Set beam energy and phase criterion: } vcr = k \cdot \sqrt{k_e (\Delta \varphi - \Delta \varphi_0)^2 + k_p (\Delta E - \Delta E_0)^2}$$

With ΔE_0 being the difference between the beam energy and the linac energy (close to 0 in the ideal case without error). and $\Delta \varphi_0$ being the difference between the beam phase and the linac phase (close to 0 in the ideal case without error)
 ke , kp allow to take into account or no one of both.

Set synchronous phase

SET_SYNC_PHASE

Allows fixing a synchronous phase into an accelerating element where the phase parameter is defined as the input phase. Concerns the following element: (So, replace the input phase value by the wanted synchronous phase value)

FIELD_MAP
 RFQ_CEL
 CAVSIN
 NCELLS

Minimize field variation

MIN_FIELD_VARIATION k , N , ke

During a matching procedure, where some electric fields of cavities are adjusted in order to match a beam, it's often useful to limit the maximum field amplitude variation.

k is used in the criterion calculation. ke is the max field variation (initial field amplitude divided by new field amplitude) wanted during a matching, N is the number of elements where this command is applied but only cavities are concerned.

Duplicate elements

REPEAT_ELE k , n

Allows to duplicate the following n elements k times. Caution: The associated commands are also repeated except: 'LATTICE', 'LATTICE_END' and 'SET_ADV' one. A 'REPEAT_ELE' command cannot include another one.

Example: 100 times a period.

```
DRIFT 5 30
DRIFT 5 30
REPEAT_ELE 100 4
LATTICE 4 1
SET_ADV 20
QUAD 100 1530 0
DRIFT 100 30
QUAD 100 -15 30 0
DRIFT 100 30
LATTICE_END
DRIFT 5 30
DRIFT 5 30
end
```

Example: Repeat 10 times a quadrupole.

```
DRIFT 5 30
REPEAT_ELE 10 1
QUAD 10 15 30 0
DRIFT 5 30
```

END

Example: Cut in 10 steps a dipole with Edge and Fringe-field.

```
EDGE 20 600 40 0 0 25 0
BEND -90 600 0 20 0
EDGE 20 600 40 0 0 25 0
```

Become,

```
EDGE 20 600 40 0 0 25 0
BEND -9 600 0 20 0
EDGE 0 600 40 1e-12 1e-12 25 0
REPEAT_ELE 8 3
EDGE 0 600 40 1e-12 1e-12 25 0
BEND -9 600 0 20 0
EDGE 0 600 40 1e-12 1e-12 25 0
EDGE 0 600 40 1e-12 1e-12 25 0
BEND -9 600 0 20 0
EDGE 20 600 40 0 0 25 0
```

Superpose field maps

SUPERPOSE_MAP *shift (mm)*

A FIELD_MAP element already allows to superpose 4 field map types (electrostatic, electrodynamic, magnetic or electric field map). With the “SUPERPOSE_MAP command, different FIELD_MAP elements can be superposed and shifted. There are no limitations excepted than only one RF field map is allowed among all field maps.

If you want to use field map aperture or current in case of superposed field map, you have to include a empty field map element at the position 0 including aperture and/or current field map.

Example: quadrupole inside a solenoid.

```
DRIFT 5 30
SUPERPOSE_MAP 400
FIELD_MAP 70 100 0 42 -0.3 0 0 0 qpole_field_map_file
SUPERPOSE_MAP 0
FIELD_MAP 70 1000 0 100 -1.3 0 0 0 solenoid_fiedl_map_file
DRIFT 5 30
End
```

Set beam phase error

SET_BEAM_PHASE_ERROR *phase (°)*

This command located before a cavity allows to set to specific value the beam phase error (typically to zero). That means the phase errors coming from static errors can be canceled. The beam phase errors coming from dynamic errors are not canceled by this command. This command works only when option “*Match using diagnostics*” is selected in “*Matching*” page.

General description

Usual mathematics formulas

$$\gamma = \frac{W}{mc^2} + 1 \qquad \gamma^{-2} = 1 - \beta^2 \qquad \lambda = \frac{c}{f} \qquad \frac{dp}{p} = \frac{\gamma}{\gamma + 1} \frac{dW}{W}$$

γ and β are the usual relativistic parameters, λ and f are respectively the free-space wavelength and the frequency of the RF. W is the kinetic energy, mc^2 the mass and p the momentum of a particle.

Description

If all the forces are considered linear, the phase-space coordinates of a particle at location s_2 can be deduced from those at the location s_1 along a transport system, by a single matrix multiplication:

$$\vec{x}(s_2) = R \cdot \vec{x}(s_1),$$

$$\vec{x} = \begin{pmatrix} x \\ x' = dx/ds \\ y \\ y' = dy/ds \\ z \\ \delta = \Delta p / p_s \end{pmatrix}.$$

Where $\vec{x}(s)$ is a vector representing the particle position in the phase-space:

x , y and z being respectively the horizontal, vertical and longitudinal position of the particle in the bunch (relative to a synchronous particle). p_s is the synchronous particle momentum, and $\delta = (p - p_s)/p_s$, with p being the particle momentum.

R is the 6x6 transfer matrix between s_1 and s_2 . In TRACE_WIN this matrix is partitioned into 2x2 matrices to simplify and accelerate the calculations.

$$R = \begin{bmatrix} a_{00} & a_{01} & a_{02} & a_{03} & a_{04} & a_{05} \\ a_{10} & a_{11} & a_{12} & a_{13} & a_{14} & a_{15} \\ a_{20} & a_{21} & a_{22} & a_{23} & a_{24} & a_{25} \\ a_{30} & a_{31} & a_{32} & a_{33} & a_{34} & a_{35} \\ a_{40} & a_{41} & a_{42} & a_{43} & a_{44} & a_{45} \\ a_{50} & a_{51} & a_{52} & a_{53} & a_{54} & a_{55} \end{bmatrix} = \begin{bmatrix} R_{xx} & R_{xy} & R_{xz} \\ R_{yx} & R_{yy} & R_{yz} \\ R_{zx} & R_{zy} & R_{zz} \end{bmatrix}.$$

In order to be able to use this formalism, the [space charge](#) force is considered as linear. To calculate the space-charge effect, the real beam is replaced by an equivalent uniform beam having identical rms properties (sizes and emittances). The total emittance of the equivalent uniform bunched beam in each phase plane is then 5 times the rms emittance, and its envelope size is $\sqrt{5}$ times its rms size. For a continuous beam these factors are 4 for the emittance and $2(= \sqrt{4})$ for envelope.

Twiss parameters

Lets define $\langle w \rangle$ as the mean value of the w particle property over the beam at location s .

$$\begin{pmatrix} \bar{x} = \langle x \rangle \\ \bar{x}' = \langle x' \rangle \\ \bar{y} = \langle y \rangle \\ \bar{y}' = \langle y' \rangle \\ \bar{z} = \langle z \rangle \\ \mathbf{0} \end{pmatrix}$$

The beam phase-space position is defined by:

The beam rms sizes are defined by: $\sigma_w = \sqrt{\langle (w - \bar{w})^2 \rangle}$, with w used for x , x' , y , y' , z or δ .

The beam rms correlation: $\overline{wv} = \langle (w - \overline{w}) \cdot (v - \overline{v}) \rangle$, with w and v used for x, x', y, y', z or δ .

The beam rms unnormalized emittances: $\tilde{\epsilon}_w = \sqrt{\overline{w^2 w'^2} - \langle w w' \rangle^2}$ with w used for x, y and z.

Note: In the preceding definition, z' is defined as $z' = (v_z - v_{zs})/v_{zs}$, v_z and v_{zs} being the longitudinal velocities of respectively the particle and the synchronous particle. Most of the time, δ is used rather than z' . In that last case, the emittance is defined by the [conversion](#) from ϵ_z to $\epsilon\delta$.

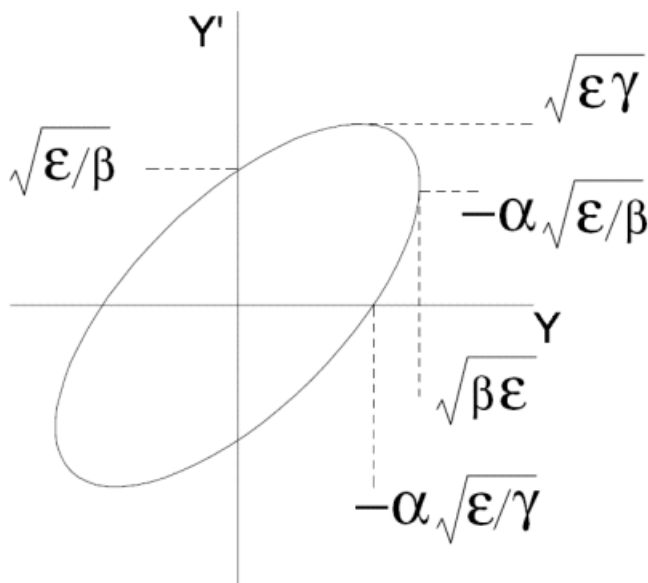
In case of linear forces, the beam can be represented in sub phase planes by ellipses whose equation can be written:

$$\gamma_w w^2 + 2\alpha_w w w' + \beta_w w'^2 = \epsilon_w, \text{ where :}$$

ϵ_w is the unnormalized beam effective emittance (which is the full emittance of a homogenous beam) define as 5 times the rms-emittance for bunched beam and 4 times that of a continuous beam.

$\alpha_w = -\overline{w w'} / \tilde{\epsilon}_w$, $\beta_w = \overline{w'^2} / \tilde{\epsilon}_w$, and $\gamma_w = \overline{w'^2} / \tilde{\epsilon}_w$ are the beam Twiss parameters satisfying the relationship:

$$\beta_w \gamma_w - \alpha_w^2 = 1.$$



The beam can be represented by a matrix, called the σ -matrix defined as:

$$[\sigma] = 5 \cdot \begin{bmatrix} \overline{x^2} & \overline{xx'} & \overline{xy} & \overline{xy'} & \overline{xz} & \overline{x\delta} \\ \overline{xx'} & \overline{x'^2} & \overline{x'y} & \overline{x'y'} & \overline{x'z} & \overline{x'\delta} \\ \overline{xy} & \overline{x'y} & \overline{y^2} & \overline{yy'} & \overline{yz} & \overline{y\delta} \\ \overline{xy'} & \overline{x'y'} & \overline{yy'} & \overline{y'^2} & \overline{y'z} & \overline{y'\delta} \\ \overline{xz} & \overline{x'z} & \overline{yz} & \overline{y'z} & \overline{z^2} & \overline{z\delta} \\ \overline{x\delta} & \overline{x'\delta} & \overline{y\delta} & \overline{y'\delta} & \overline{z\delta} & \overline{\delta^2} \end{bmatrix}.$$

The evolution of the σ -matrix along the line from s_1 to s_2 can be calculated with the transfer matrix R :

$$[\sigma](s_2) = R \cdot [\sigma](s_1) \cdot R^T,$$

Where R^T is the transpose of R and $[\sigma]$ is the beam σ -matrix. Like with the transfer matrixes the σ -matrixes can be partitioned into 2x2 matrixes:

$$[\sigma] = \begin{bmatrix} \sigma_{xx} & \sigma_{xy} & \sigma_{xz} \\ \sigma_{yx} & \sigma_{yy} & \sigma_{yz} \\ \sigma_{zx} & \sigma_{zy} & \sigma_{zz} \end{bmatrix}.$$

The elements are divided into small steps, whose transfer matrixes are used to transport the beam σ -matrix. The space-charge effect is applied at each step.

Definition of the matched beam

The 2x2 extracted σ -matrix can be written in terms of Twiss parameters.

$$\sigma_{ww} = \begin{bmatrix} \beta_w \varepsilon_w & -\alpha_w \varepsilon_w \\ -\alpha_w \varepsilon_w & \gamma_w \varepsilon_w \end{bmatrix}.$$

We observe that $\det(\sigma_{ww}) = (\beta_w \gamma_w - \alpha_w^2) \cdot \varepsilon_w^2 = \varepsilon_w^2$

$$R = \begin{bmatrix} R_{xx} & R_{xy} & R_{xz} \\ R_{yx} & R_{yy} & R_{yz} \\ R_{zx} & R_{zy} & R_{zz} \end{bmatrix}.$$

Let's R be the transfer matrix of a lattice of a periodic structure:

The 2x2 extracted diagonal matrix can be written:

$$R_{ww} = \begin{bmatrix} \cos(\sigma_{w0}) + \alpha_{w0} \sin(\sigma_{w0}) & \beta_{w0} \sin(\sigma_{w0}) \\ -\gamma_{w0} \sin(\sigma_{w0}) & \cos(\sigma_{w0}) - \alpha_{w0} \sin(\sigma_{w0}) \end{bmatrix},$$

Where α_{w0} , β_{w0} , and γ_{w0} are the Twiss parameters of the beam matched to the lattice, and σ_{w0} is the zero-current phase advance per lattice in the $[w-w']$ phase plane.

Twiss parameters and acceleration

In case of acceleration, the determinant of the transfer matrix is not equal to 1, and the matrix cannot be written with the Twiss parameters, as defined before. To extract the Twiss parameters of the matched beam from the transfer matrix when there is acceleration, we use the matrix defined as below:

$$R' = \begin{bmatrix} 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & \beta_o \gamma_o & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & \beta_o \gamma_o & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & 0 & \beta_o \gamma_o \end{bmatrix} \cdot R \cdot \begin{bmatrix} 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 1/\beta_i \gamma_i & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1/\beta_i \gamma_i & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1/\beta_i \gamma_i \end{bmatrix},$$

Where β_i and γ_i are the relativistic parameters at the input and β_o and γ_o are the relativistic parameters at the output of the lattice. Now, the R' matrix determinant equals 1.

The Twiss parameters of the matched beam under acceleration conditions can then be deduced from the R' matrix:

$$\begin{aligned} \sigma_{x0} &= \cos^{-1} \left(\frac{r'_{00} + r'_{11}}{2} \right), & \sigma_{y0} &= \cos^{-1} \left(\frac{r'_{22} + r'_{33}}{2} \right), & \sigma_{z0} &= \cos^{-1} \left(\frac{r'_{44} + r'_{55}}{2} \right), \\ \alpha_{x0} &= \frac{(r'_{00} - r'_{11})}{2 \sin(\sigma_{x0})}, & \alpha_{y0} &= \frac{(r'_{22} - r'_{33})}{2 \sin(\sigma_{y0})}, & \alpha_{z0} &= \frac{(r'_{44} - r'_{55})}{2 \sin(\sigma_{z0})}, \\ \beta_{x0} &= \beta_s \gamma_s \frac{r'_{01}}{\sin(\sigma_{x0})}, & \beta_{y0} &= \beta_s \gamma_s \frac{r'_{23}}{\sin(\sigma_{y0})}, & \beta_{z0} &= \beta_s \gamma_s^3 \frac{r'_{45}}{\sin(\sigma_{z0})}, \\ \gamma_{x0} &= \frac{-r'_{10}}{\beta_s \gamma_s \sin(\sigma_{x0})}, & \gamma_{y0} &= \frac{-r'_{32}}{\beta_s \gamma_s \sin(\sigma_{y0})}, & \gamma_{z0} &= \frac{-r'_{54}}{\beta_s \gamma_s^3 \sin(\sigma_{z0})}, \end{aligned}$$

Where, r'_{ij} is a R' matrix coefficient (i^{th} row, j^{th} column) and α_{wo} , β_{wo} , and γ_{wo} are output Twiss parameters of the matched beam in the $[w-w']$ phase plane.

Conversions between $[z-z']$, $[z-\delta]$ and $[\Delta\phi-\Delta W]$ phase planes

β and γ being the beam reduced velocity and energy, λ the RF wavelength in vacuum, mc^2 the particle rest energy, we have in the paraxial approximation conditions the following relationship between parameters:

$$\Delta\phi = -\frac{360^\circ}{\beta\lambda} \cdot z$$

But if the beam shows a divergence this relation becomes:

$$\Delta\phi = -\frac{360^\circ}{\beta\lambda} \cdot \sqrt{1 + \frac{(\Delta x')^2}{4} + \frac{(\Delta y')^2}{4}} z$$

$\Delta\phi$ and z being the RF phase and the position of a beam particle relative to the synchronous one.

$$\Delta W = \beta^2 \gamma^3 mc^2 \cdot z' = \beta^2 \gamma mc^2 \cdot \delta$$

In the general case with a synchronous particle and a generator particle these relation become:

$$\delta = \frac{\Delta W \gamma_G + (\gamma_G - 1)(m_{sG}^2 - m_{cG}^2)}{m_{cG}^2 (\gamma_G^2 - 1)} \quad \delta = z' \gamma_G^2 + \frac{(m_{sG}^2 - m_{cG}^2)}{m_{cG}^2}$$

ΔW , z' and δ being the energy, velocity and momentum of a beam particle relative to the synchronous one.

$$\varepsilon_w = \frac{360^\circ \cdot mc^2}{\lambda} \cdot \varepsilon_{zn} = \frac{360^\circ \cdot mc^2}{\lambda} \cdot \beta \gamma^3 \cdot \varepsilon_z = \frac{360^\circ \cdot mc^2}{\lambda} \cdot \beta \gamma \cdot \varepsilon_{z\delta}$$

ε_w and $\varepsilon_{z\delta}$ being the normalized longitudinal emittances, ε_z and ε_{zn} being the unnormalized longitudinal emittances of the beam in respectively the $[z-z']$ and the $[z-\delta]$ phase planes.

$$\beta_w = \frac{360^\circ}{mc^2 \gamma^3 \beta^3 \lambda} \cdot \beta_z = \frac{360^\circ}{mc^2 \gamma \beta^3 \lambda} \cdot \beta_{z\delta}$$

β_w , β_z and $\beta_{z\delta}$ are the β -Twiss parameters of the beam in respectively the $[\Delta\phi-\Delta W]$, $[z-z']$ and $[z-\delta]$ phase planes.

$$\alpha_w = \alpha_z = \alpha_{z\delta}$$

α_w , α_z and $\alpha_{z\delta}$ are the α -Twiss parameters of the beam in respectively the $[\Delta\phi-\Delta W]$, $[z-z']$ and $[z-\delta]$ phase planes.

$$\gamma_w = \frac{mc^2 \gamma^3 \beta^3 \lambda}{360^\circ} \cdot \gamma_z = \frac{mc^2 \gamma \beta^3 \lambda}{360^\circ} \cdot \gamma_{z\delta}$$

γ_w , γ_z and $\gamma_{z\delta}$ are the γ -Twiss parameters of the beam in respectively the $[\Delta\phi-\Delta W]$, $[z-z']$ and $[z-\delta]$ phase planes.

Normalization of the emittance

$$\varepsilon_{tn} = \beta \gamma \cdot \varepsilon_t$$

$$\varepsilon_{zn} = \beta \gamma^3 \cdot \varepsilon_z = \beta \gamma \cdot \varepsilon_{z\delta}$$

Transit time factor definition

These following definitions are used in all accelerating element.

T : is the usual time factor transit given by SUPERFISH.

T' : is the T' from SUPERFISH time -2π

T'' : is the T'' from SUPERFISH time $-4\pi^2$

$$T(\beta) = T_s + kT'_s \cdot (1 - \kappa) + k^2 T''_s \cdot \frac{(1 - \kappa)^2}{2},$$

$$kT'(\beta) = kT'_s + k^2 T''_s \cdot (1 - \kappa) \dots \quad \text{With } \kappa = \frac{k_s}{k} = \frac{\beta}{\beta_s}.$$

The electric field is corrected according to $T(\beta)$.

Coordinate transformations are given in Wangler's book page 202:

Phase advance definition

TRACE_WIN calculates the particle phase advance in two ways. The first one is the extraction of the phase advance from the transfer matrix of the lattice (μ or σ_o). The second one is done by the beta function integration along a lattice giving the phase advance with (σ) or without σ_o space charge:

$$\sigma_x = \int_0^L \frac{dl}{\beta(l)} \quad \text{with } L \text{ is the lattice length and } \beta = \frac{\langle x^2 \rangle}{\epsilon_x}.$$

x is the beam RMS or effective size and ϵ_x the unnormalized RMS or effective emittance.

The first phase advance type can be plotted from the "**Phase advance**"->"**Structure**".of "**Chart**" page And the second one can be plotted from the "**Phase advance**"->"**Beam**". of "**Chart**" page

Residual orbit

Use in the error studies in order to know the beam gravity evolution, It's defined like below:

$$\begin{aligned} \bar{x} &= \frac{1}{N} \sum_1^N x & x_{rms} &= \sqrt{\frac{1}{N} \sum_1^N x^2 - \bar{x}^2} \\ \bar{y} &= \frac{1}{N} \sum_1^N y & y_{rms} &= \sqrt{\frac{1}{N} \sum_1^N y^2 - \bar{y}^2} \\ \bar{r} &= \sqrt{\bar{x}^2 + \bar{y}^2} & r_{rms} &= \sqrt{\frac{1}{N} \sum_1^N r^2 - \bar{r}^2} \end{aligned}$$

x & y are the beam gravity position and N is the number of run or linac.

Halo definition

If the motion is uncoupled between phase planes, the following quantities are kinematic invariants of motion:

$$I_2^i \equiv \langle q_i^2 \rangle \langle p_i^2 \rangle - \langle q_i p_i \rangle^2$$

$$I_4^i \equiv \langle q_i^4 \rangle \langle p_i^4 \rangle + 3 \langle q_i^2 p_i^2 \rangle^2 - 4 \langle q_i p_i^3 \rangle \langle q_i^3 p_i \rangle$$

Then we define the halo intensity parameter, H_i , in the i th phase plane, as

$$H_i \equiv \frac{\sqrt{3I_4^i}}{2I_2^i} - 2 = \frac{\sqrt{3 \langle q_i^4 \rangle \langle p_i^4 \rangle + 9 \langle q_i^2 p_i^2 \rangle^2 - 12 \langle q_i p_i^3 \rangle \langle q_i^3 p_i \rangle}}{2 \langle q_i^2 \rangle \langle p_i^2 \rangle - 2 \langle q_i p_i \rangle^2}$$

From ref: “*PHYSICAL REVIEW SPECIAL TOPICS - ACCELERATORS AND BEAMS, VOLUME 5, 124202 (2002)*“, “*Beam halo definitions based upon moments of the particle distribution (C. K. Allen and T. P. Wangler)*“

Space charge

In the dynamics the space charge is calculated at each step. The number of calculation and obviously the calculation time depend on the parameter “*Step of calculation per $\beta\lambda$* ”. The space-charge calculation as it is done in TRACE_WIN is presented in this paragraph.

In a free space, the motion equation of a particle, feeling only space-charge force, can be written along one direction in the **beam frame R^*** :

$$\frac{d^2 w^*}{dt^{*2}} = F_w^*,$$

F_w^* being the w component of the space-charge force (divided by the particle rest mass).

All the quantities with a star (*) are expressed in R^* , all quantities without star are expressed in the **laboratory frame R** . We have, from the Lorentz transform:

$$dt^* = dt/\gamma, \quad x^* = x, \quad y^* = y, \quad z^* = \gamma \cdot z,$$

And:

$$ds = \beta \cdot dt.$$

The derivation with s , the beam longitudinal position used as independent variable gives:

$$\frac{d^2 w^*}{ds^2} = F_w^* = \frac{F_w^*}{(\gamma\beta)^2}.$$

In the longitudinal direction, we have:

$$\frac{d\delta}{ds} = \gamma^2 \frac{d^2 z}{ds^2} = \gamma \frac{d^2 z^*}{ds^2} = \gamma \frac{F_w^*}{(\gamma\beta)^2} = \gamma F_w^*.$$

DC beam

When the beam is continuous, there is no space-charge force acting along the longitudinal direction. We have:

$$F_x = \frac{2K}{a_x(a_x + a_y)} \cdot x, \quad F_y = \frac{2K}{a_y(a_x + a_y)} \cdot y,$$

$$K = \frac{|q| \cdot I}{\pi \epsilon_0 m c^3 \beta^3 \gamma^3}$$

With β , the beam generalized perveance.

a_x, a_y , are the semiaxes of the homogeneous ellipse (2 times the rms beam sizes), I is the average beam current, ϵ_0 is the vacuum permittivity.

Bunched beam

The w -component of the space-charge force acting on the particle in the bunch frame is:

$$F_w^* = \frac{K_w}{a_x^* a_y^* a_z^*} \cdot W^*$$

Giving:

$$F_w = \frac{K_w}{a_x^* a_y^* a_z^*} \cdot \frac{1}{(\gamma\beta)^2} \cdot W^*$$

With: w for x, y or z ,

$$K_w = \frac{3 \cdot |q| \cdot I \lambda h}{4\pi\epsilon_0 mc} \cdot f_w(a_x^*, a_y^*, a_z^*)$$

$$f_w(a_x, a_y, a_z) = \int_0^\infty \frac{ds}{(a_w^2 + s) \sqrt{(a_x^2 + s)(a_y^2 + s)(a_z^2 + s)}}$$

And

is the form factor such as:

$$\forall a_x, a_y \text{ and } a_z, \quad f_x + f_y + f_z = 1.$$

a_x, a_y, a_z , are the semiaxes in the laboratory frame of the homogeneous ellipsoid ($\sqrt{5}$ times the rms beam sizes), $a_x^* = a_x$,

$a_y^* = a_y$ and $a_z^* = \gamma a_z$ are the beam semiaxes along longitudinal direction in the beam frame, I is the average beam current, h is the harmonic number (ratio between the RF frequency and the bunch frequency), ϵ_0 is the vacuum permittivity. Note that the longitudinal beam size in the beam frame is equal to γ times the one in the laboratory frame.

Form factor calculation

The form factor integral calculation depends on the ratios a_z^* / a_x^* and a_z^* / a_y^* . If they are lower than 12 the integral is calculated by the Gauss method with a very good precision. If they are greater than 12 (which happens obviously when the beam is ultra-relativistic due to the Lorentz transformation) an expended development is used and slightly reducing the result precision.

Space-charge application

Frame change

The space charge impulse should be applied in the beam frame. Before any application, the beam σ -matrix should be written from the laboratory frame to the beam frame by making the transformation:

$$[\sigma] = R_\gamma \cdot [\sigma] \cdot R_\gamma^T,$$

$$R_\gamma = \begin{pmatrix} 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & \gamma & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 \end{pmatrix}.$$

With:

After the application of space-charge impulse (taking into account the beam coupling), the σ -matrix should be written back to the laboratory frame:

$$[\sigma] = R_{\gamma^{-1}} \cdot [\sigma] \cdot R_{\gamma^{-1}}^T,$$

$$R_{\gamma^{-1}} = \begin{pmatrix} 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & \gamma^{-1} & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 \end{pmatrix}$$

With:

Beam coupling or tilted.

Due to the elements (magnetic coil or dipole) or initial conditions, the beam ellipsoid in $[x-y-z]$ space can be tilted. In this case, the beam ellipsoid (in the beam frame) must first be transformed to a coordinate system in which it is upright before applying the space-charge impulses.

If the ellipsoid is tilted in the $[x-y]$ plane, the angle between the x-axis and the axis of the elliptical projection on the $[x-y]$ plane is:

$$\theta_{xy} = \frac{1}{2} \tan^{-1} \left(\frac{2\sigma_{02}}{\sigma_{22} - \sigma_{00}} \right) \text{ With } \sigma_{ij}, \sigma \text{-matrix elements.}$$

$$\text{If the ellipsoid is tilted in the } [x-z] \text{ plane: } \theta_{xz} = \frac{1}{2} \tan^{-1} \left(\frac{2\sigma_{04}}{\sigma_{44} - \sigma_{00}} \right).$$

$$\text{If the ellipsoid is tilted in the } [y-z] \text{ plane: } \theta_{yz} = \frac{1}{2} \tan^{-1} \left(\frac{2\sigma_{24}}{\sigma_{44} - \sigma_{22}} \right).$$

The ellipsoid can be brought upright by rotations of angles $-\theta_{xy}$, $-\theta_{xz}$ and $-\theta_{yz}$ accomplished by applying the transfer matrices:

$$R_{xy} = \begin{bmatrix} \cos(\theta_{xy}) & 0 & \sin(\theta_{xy}) & 0 & 0 & 0 \\ 0 & \cos(\theta_{xy}) & 0 & \sin(\theta_{xy}) & 0 & 0 \\ -\sin(\theta_{xy}) & 0 & \cos(\theta_{xy}) & 0 & 0 & 0 \\ 0 & -\sin(\theta_{xy}) & 0 & \cos(\theta_{xy}) & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 \end{bmatrix},$$

R_{xz} and R_{yz} can be obtained the same way.

The rotation is then applied:

$$[\sigma] = R_{yz} \cdot R_{xz} \cdot R_{xy} \cdot [\sigma] \cdot R_{xy}^T \cdot R_{xz}^T \cdot R_{yz}^T,$$

$$[\sigma] = Rot \cdot [\sigma] \cdot Rot^T.$$

When the ellipse is upright, the space-charge impulses can be applied. The three reverse rotations can then be applied:

$$[\sigma] = R_{xy}^T \cdot R_{xz}^T \cdot R_{yz}^T \cdot [\sigma] \cdot R_{yz} \cdot R_{xz} \cdot R_{xy},$$

$$[\sigma] = Rot^T \cdot [\sigma] \cdot Rot.$$

Space-charge impulse

The space-charge transfer matrix R_{ce} applies on a distance Δs (the calculation step):

$$R_{ce} = \begin{bmatrix} 1 & 0 & 0 & 0 & 0 & 0 \\ F_x \Delta s & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & F_y \Delta s & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & \gamma F_z \Delta s & 1 \end{bmatrix}.$$

The space-charge impulse is applied in the bunch frame, where the beam ellipse is upright:

$$[\sigma] = R_{cs} \cdot [\sigma] \cdot R_{cs}^T.$$

Finally, the total space charge effect is given by:

$$[\sigma] = R_{y-1} \cdot Rot^T \cdot R_{cs} \cdot Rot \cdot R_y \cdot [\sigma] \cdot R_y^T \cdot Rot^T \cdot R_{cs}^T \cdot Rot \cdot R_{y-1}^T.$$

Data file

The data file (“*.dat”) contains the list of elements and commands. It must be ended by the command “END”. The elements and the commands syntax are described in the rubric “Element definitions” and “Command definitions”. Comments line begins by the character ‘;’.

A name for each element can be specified with 8 characters max, (See example below).

Result files are automatically created at the first use of the data file. At the end of a run, TRACE_WIN creates another data file with the same name but located in the calculation directory, which contains the new element list. With, for example, the quadrupole values calculated to have the wanted phase advance law. If calculation directory is the same than the data file directory, the name of the new data file start with “new_...”

Warning:

Each command concerns the following element, par example “SET_TWISS” will impose some Twiss parameters at the output of the following element.

Two identical commands cannot be consecutive.

Example 1:

```

; *****
DRIFT 1e-08 100
SPACE_CHARGE_COMP 0.7
DRIFT 350 100
DRIFT 60 100
DRIFT 192 100
MATCH_FAM_GRAD 1 1
ADJUST 1 2 1 0 0
SOLENOID 410 0.25 100
DRIFT 100 100
MATCH_FAM_GRAD 1 2
STEERER 0 0 100 0
ADJUST_STEERER 2
ADJUST 1 2 2 0 0
QUAD 200 0.18 100 0
DRIFT 150 100
; *****

```

Example 2:

```

; *****
DR1 : DRIFT 1e-08 100
SP1 : SPACE_CHARGE_COMP 0.7
DR2 : DRIFT 350 100
DR3 : DRIFT 60 100
SOL 1 : SOLENOID 410 0.25 100
DR5 : DRIFT 100 100
QPF1 : QUAD 200 0.18 100 0
DR6 : DRIFT 150 100
; *****

```

Init project file

The init file “project_name.ini” contains all the TraceWin project parameters. It can be loaded, saved, copied by using the TraceWin menu.

Results file

Created by TRACE_WIN, his name is “*Data_file_name.cal*” and it is located in the data file directory and contains the results of the already done matching calculations, to avoid redundant calculations. See the following example.

```
Twiss_parameters_of_matched_beam
0.3167415265 0.1850852302 0.5246751875
-0.0938830920 0.0822867115 -0.0875778140

Matching_Between_Section_1_to_2
-8.11004 8.16711 -8.19803 8.21871
-3.8207 -2.8207 -4.1476 -1.1476
0.00834559 0.000887792
BEAM_FAM_69_0.DST
```

The three first lines are written after a matching beam calculation. The second line contain the Twiss parameters β_{xx} , β_{yy} , β_{zz} , and the last α_{xx} , α_{yy} , α_{zz} .

The five following lines are written after a matching calculation, it contains the result of a matching between two sections. The first line contains the quadrupole gradients witch have been adjusted (“MATCH_FAM_GRAD” command), the second line is either the phase shift or the field factor correction or both, of the accelerator elements witch have been adjusted (“MATCH_FAM_PHASE”, “MATCH_FAM_FIELD” or “MATCH_FAM_LFOC” command). The third line correspond to the element length witch have been adjusted (“MATCH_FAM_LENGTH” command) and the last is the name of a beam distribution file (located in the file data path), witch is save when the matching family calculated is done with the option “*With beam from Partran*”. All of these lines are optional and depend from the “MATCH_FAM...”command in your data file.

For more details, see the [matching commands](#) and its examples. You can also force the optimization process of calculation with starting values by the following commands

```
Init_Matching_Between_Section_1_to_2
-8.11004 8.16711 -8.19803 8.21871
-2.8207 -2.8207 -3.1476 -3.1476
```

To put in comment a result, simply add the char “;” as a first character.

This file contains also all the diagnostic results like the following example. For more details, see the [adjust commands](#) and its examples. You can also force the optimization process of diagnostic calculation with starting values by using “*Init_*” syntax

```
Diagnostic_10
10.7992 -10.5893 5.81701

Init_Diagnostic_10
10.7992 -10.5893 5.81701
```

For all these result, an extension “_PAR” is added when the result comes from a multiparticle optimization ‘*With Partran*’ is checked

Sigma0 file

Created by TRACE_WIN or No, his name is “*.sig”. It’s located in the data file directory and contains the transverse phase advances law with no current, one value per lattice. See following example:

```
60
60
61
62
...
..
.
```

Magnetic or electric Field map

The field map file syntax is the following in the ASCII format:

**Fz are in MV/m for electric field or in T for magnetic field.
The dimensions are in metre.**

- Dimension 1 :

nz zmax

Norm

for k=0 to nz

Fz(k.zmax/nz)

Return

- Dimension 2 :

nz zmax

nr rmax

Norm

for k=0 to nz

for i=0 to nr

Fz(k.z_{max}/nz, i.r_{max}/nr)

Return

or

nx xmin xmax

ny ymin ymax

Norm

for k=0 to ny

for i=0 to nx

Fz(k.x_{max}/nx, i.y_{max}/ny)

Return

- Dimension 3 :

nz zmax

nx xmin xmax

ny ymin ymax

Norm

for k=0 to nz

for j=0 to ny

for i=0 to nx

Fz(k.z_{max}/nz, y_{min}+j.(y_{max}-y_{min})/ny, x_{min}+i.(x_{max}-x_{min})/nx)

Return

The field map file syntax is the following in the BINARY format:

- Dimension 1 :

nz (integer 4 bytes) zmax (double 8 bytes)

Norm (double 8 bytes)

for k=0 to nz

Fz(k.zmax/nz) (float 4 bytes)

- Dimension 2 :

nz (integer 4 bytes) zmax (double 8 bytes)

nr (integer 4 bytes) rmax (double 8 bytes)

Norm (double 8 bytes)

for k=0 to nz

for i=0 to nr

Fz(k.z_{max}/nz, i.r_{max}/nr) (float 4 bytes)

- Dimension 3 :

nz (integer 4 bytes) zmax (double 8 bytes)

nx (integer 4 bytes) xmin (double 8 bytes) xmax (double 8 bytes)

ny (integer 4 bytes) ymin (double 8 bytes) ymax (double 8 bytes)

Norm (double 8 bytes)

for k=0 to nz

for j=0 to ny

for i=0 to nx

$Fz(k \cdot z_{\max}/nz, y_{\min} + j \cdot (y_{\max} - y_{\min})/ny, x_{\min} + i \cdot (x_{\max} - x_{\min})/nx)$ (float 4 bytes)

Warning: The lattice has to be regular.

The normalization factor is equal to k_e/Norm or k_b/Norm .

**Fz are in MV/m for electric field or in T for magnetic field.
The dimensions are in metre.**

Current or space charge compensation map

“FileMapName.scc”

A flag in “FIELD_MAP” element syntax allow to include it.

The space charge compensation or current file syntax is like following:

- Space charge compensation according to Z format:

0 N

for i=0 to N-1

$Z_i \text{ Scc}_i$

- Current evolution according to Z file format:

1 N

for i=0 to N-1

$Z_i I_i$

- Z_i is the position (m)

- Scc_i is the space charge compensation at the Z_i position, (1 for 100%)

- I_i is the current (mA) at the Z_i position

Partran and Trace_Win codes make an interpolation in between this figure.

Aperture map

“FileMapName.ouv”

A flag in [FIELD_MAP](#) element syntax allow to include it.

For the field map elements, sometime we need to define a beam pipe radius geometry according to z axis. The file syntax is the following:

Warning in case of superposed field map these aperture map have to be defined in the first [FIELD_MAP](#) element and have to get a length equivalent to all field_map.

- Aperture according to Z format:

N

for i=0 to N-1

$Z_i \text{ Ouv}_i$

- Z_i is the position (m)

- Ouv_i is the aperture radius(m) at Z_i .

The first location Zi has to be 0.

Input files for multiparticle programs

At the end of a calculation TRACE_WIN creates the input files for multiparticle programs, PARTRAN (*Data_file_name.par*), and TOUTATIS (*toutatis.dat*).

Particle distribution

Dist_Error_Env.dat

Contain the beam distribution at the end of each element after an envelope calculation. This file is created if “*nbr of particles*” is greater than 10 and “*Use aperture element*” of “*Main*” page is selected. During an error study the condition “*nbr of particles*” is sufficient. You can visualize these results in the “*Error*” page by setting the “*Distribution file*” and using the right buttons

Dist_Error_PAR.dat

Contain the beam distribution at the end of each element after a multiparticle calculation. This file is created either by PARTRAN or TOUTATIS. You can visualize these results in the “*Error*” page by setting the “*Distribution file*” and using the right buttons

- N: Number of linac = 1
- Element number
- Element aperture (cm)
- Element aperture (cm)
- $\sum_{1}^N \sqrt{x^2 + y^2}$ (cm)
- $\sum_{1}^N x^2 + y^2$ (cm²)
- $\sum_{1}^N x$ (cm)
- $\sum_{1}^N x^2$ (cm²)
- $\sum_{1}^N y$ (cm)
- $\sum_{1}^N y^2$ (cm²)
- 100 integers corresponding to the particle distribution along the aperture divided in 100 steps.
- 100 Doubles corresponding to the power distribution along the aperture divided in 100 steps.
- $\sum_{1}^N particle_lost$
- $\sum_{1}^N particle_lost^2$
- Max particle lost
- Min particle lost
- $\sum_{1}^N power_lost$ (w)

- $\sum_1^N power_lost^2$ (w)
- Max power lost (w)
- Min power lost (w)

In case of statistical error study 2 new files are created name **Dist_Error_tot_ENV.dat** and **Dist_Error_tot_PAR.dat** witch contain the sum of the 2 preceeding files (**N>1**).

If the “*Nbr of Step*” parameter of the tab-sheet “*Error*” is bigger than 1 the name of the 2 files become for example for 5 steps

“Dist_Error_Tot_Env_0.2000.dat“ for 20%

“Dist_Error_Tot_Env_0.4000.dat“ for 40%,

...

“Dist_Error_Tot_Env_1.0000.dat“ for 100%,

Steerer strength file

The file “*Steerer_Values.txt*” is created after diagnostic position calculation. Its syntax is number of diagnostic follows by all the steerer strengths associated in T.m (Plane X and Y).

In case of statistical error study a new file named “*Steerer_Values_Tot.txt*” is written including all the steerer strength of the whole linac simulated.

Magnetic streeping file

The file “*MAGSTRIPI.LOS*” is created only in multiparticle mode if option “*Magnetic streeping*” is selected in “*Option*” of multiparticle codes. You can directly exploit these results using “*Stripping*” button in tab-sheet “*Graphs*”. It contains the probability losses due to Lorentz magnetic streeping. The syntax is the following. (One line per element):

Number of element, Position (m), Energy (MeV), Losses probability

In case of statistical error study a new file named “*MAGSTRIPI_TOT.LOS*” is written including the probability sum of the whole linacs simulated. You can directly exploit these results using “***Stripping losses probability results***” button in “***Errors***” page. The syntax is the following (One line per element):

Number of element, Number of linac simulated, Sum of losses probability

The probability corresponding to Sum of Magnetic streeping probability divided by Number of linac simulated.

Gas streeping file

The file “*GASSTRIPI.LOS*” is created only in multiparticle mode if option “*Gas streeping*” is selected in “*Option*” of multiparticle codes and if command [Gas pressure](#) is included in the data file. It contains the probability losses due to Gas streeping. The syntax is the following. (One line per element):

Number of element, Position (m), Energy (MeV), Losses probability

In case of statistical error study a new file named “*GASSTRIPI_TOT.LOS*” is written including the probability sum of the whole linacs simulated. You can directly exploit these results using “***Stripping losses probability results***” button in “***Errors***” page. The syntax is the following (One line per element):

Number of element, Number of linac simulated, Sum of losses probability

The probability corresponding to Sum of Gas stripping probability divided by Number of linac simulated.

Particle losses distribution

The following file are created only while statistical error study of k linacs

Dist_Error_Tot_Env.loss

Contain the number of losses at the end of each element after an envelope calculation. This file is created if “*nbr of particles*” is greater than 10. You can visualize these results in the “*Error*” page by setting the “*Distribution file*” and using the right buttons

Dist_Error_Tot_PAR.loss

Contain the number of losses at the end of each element after a multiparticle calculation. This file is created either by PARTRAN. You can visualize these results in the “*Error*” page by setting the “*Distribution file*” and using the right buttons

- Number of elements
- Number of particles

- Linac 1 (Losses element 1, Lost power element 1, losses element 2, Lost power element 2,... losses element N, Lost power element N)
- Linac 2 (Losses element 1, Lost power element 1, losses element 2, Lost power element 2,... losses element N, Lost power element N)
-
- Linac k (Losses element 1, Lost power element 1, losses element 2, Lost power element 2,... losses element N, Lost power element N)

Input & Output particle distribution

These following files are created while multiparticle simulation

part_dtl1.dst: Binary file containing the output beam distribution at the end of the linac.

part_rfql.dst: Binary file containing the beam distribution at the entrance of the linac.

A .dst file use a binary format. It contains informations of a beam at a given longitudinal position: number of particles, beam current, repetition frequency and rest mass as well as the 6D particles coordinates. The format is the following:

```
2×CHAR(125,100)+INT(Np)+DOUBLE(Ib(A))+DOUBLE(freq(MHz))+CHAR(0)+
Np×[6×DOUBLE(x(cm),x'(rad),y(cm),y'(rad),phi(rad),Energie(MeV))]+
DOUBLE(mc2(MeV))
```

Comments:

CHAR is 1 byte long,

INT is 4 bytes long,

DOUBLE is 8 bytes long.

Np is the number of particles,

Ib is the beam current,

freq is the bunch frequency,

mc2 is the particle rest mass.

dtl1.plt: Binary file containing the beam distribution at the end of each element.

A .plt file use a binary format. It contains information of a beam at many longitudinal positions: longitudinal position, number of particles, beam current, repetition frequency and rest mass as well as the 6D particles coordinates. The format is the following:

```
2×CHAR(125,100)+INT(Ne)+INT(Np)+DOUBLE(Ib(A))+DOUBLE(freq(MHz))+
DOUBLE(mc2(MeV))+
Ne×[CHAR(0)+INT(Nelp)+DOUBLE(Zgen)+DOUBLE(phase0(deg))+DOUBLE(wgen(MeV))+
Np×[7×FLOAT(x(cm),x'(rad),y(cm),y'(rad),phi(rad),Energie(MeV),Loss)] ]
```

Comments:

CHAR is 1 byte long,

INT is 4 bytes long,

FLOAT is a Real 4 bytes long.

DOUBLE is a Real 8 bytes long.

Ne is the number of different positions,

Np is the number of particles,

Ib is the beam current,

freq is the bunch frequency,

mc2 is the particle rest mass,

Nelp is the longitudinal element position,

Zgen is the longitudinal position in meter,

Phase0 & wgen are the phase and energy references of the beam,

Compressions format



Error file result

The final results can be found on the calculation directory. Named “*Error_study_Name_TRA.txt*” when the result comes from an envelope calculation and “*Error_study_Name_PAR.txt*” when it is a PARTRAN study. The format of these files is the same for each kind of error. For each step of calculation one line of 11 parameters is written with the following format.

- Step of error (0->1)
- 1-(Nbr of particles)/(Nbr of particles, reference case)
- (Emittance rms xx',yy',zz')/ (Reference rms emittance xx',yy',zz')-1
- X beam center (m)
- Y beam center (m)
- X' beam center (rad)
- Y' beam center (rad)
- Energy beam center (keV)
- Phase beam center (deg)
- RMS Beam X size (m)
- RMS Beam Y size (m)
- RMS Beam X' size (m)
- RMS Beam Y' size (m)
- RMS Beam Energy size (keV)
- RMS Beam Phase size (deg)
- Halo parameter xx'
- Halo parameter yy'
- Halo parameter zz'

Halo definition.

All these values are relative to the output beam without errors.

In case of statistical study, where each step of calculation contains several runs, the format becomes:

- Step of error (0->1)
- AVERAGE(1-(Nbr of particles)/(Nbr of particles, reference case))
- (AVERAGE(Emittance rms xx' , yy' , zz'))/ (Reference rms emittance xx' , yy' , zz')-1
- RMS(X beam center (m))
- RMS(Y beam center (m))
- RMS(X' beam center (rad))
- RMS(Y' beam center (rad))
- RMS(Energy beam center (keV))
- RMS(Phase beam center (deg))
- AVERAGE(RMS Beam X size (m))
- AVERAGE(RMS Beam Y size (m))
- AVERAGE(RMS Beam X' size (m))
- AVERAGE(RMS Beam Y' size (m))
- AVERAGE(RMS Beam Energy size (keV))
- AVERAGE(RMS Beam Phase size (deg))
- AVERAGE(Halo parameter xx')
- AVERAGE(Halo parameter yy')
- AVERAGE(Halo parameter zz')
-

And a file call “Error_study_Name_TRA_tot.txt” or “Error_study_Name_PAR_tot.txt” is written containing all run results.

Partran and Toutatis output

The final multiparticle results contain one line by element output, the first line being the input beam parameters. The format is like following.

- Element number
- Element position (m)
- Relativistic parameters: (γ -1)
- Centroid position: x(mm), y(mm), Phase($^{\circ}$), x' (mrad), y' (mrad), W(MeV)
- RMS_SIZE(x(mm), y(mm), Phase(deg))
- RMS (xx' (mm.mrad), yy' (mm), Phase.Energy(deg.MeV))
- Normalized rms emit: xx' (mm.mrad), yy' (mm.mrad), **PW (Deg.MeV).**
- Halo parameters: ($H_{xx'}$, $H_{yy'}$, $H_{z.dp/p}$)
- Number of particles
- Phase advance with space charge (deg): σ_x , σ_y , σ_z .
- Emittance at 99%: $\epsilon_{xx'}$, $\epsilon_{yy'}$, $\epsilon_{z.dp/p}$
- Beam currant (mA) used for space charge calculation
- Aperture (mm)
- Normalized 4D transverse emittance $E_{xx',yy'}$, (mm.mrad)²
- Normalized rms emit (mm.mrad): $\epsilon_{rr'}$.
- Phase advance with space charge (deg): σ_r
- Lost power (w)
- Maximum excursion particle : Xmax(mm), Ymax(mm),
- Normalized long. rms emit: $\epsilon_{z.dp/p}$ (mm.mrad) [replace PW(Deg.Mev)]
- Dispersion: Dh (mm), Dv(mm)
- Derivative dispersion: Dh' (mrad), Dv'(mrad)

Since TraceWin version 2.1.0.0 Longitudinal rms emittance PW is set to zero and has been replaced by $\epsilon_{z.dp/p}$

Bend error treatment


Traitement des erreurs dans les dipôles dans les codes PARTRAN & TRACEWIN *Nicolas Pichoff, Didier URIOT*

| | |
|--|---|
| <p>$DY :$</p> <p>$y_e \leftarrow DY$ $y_s \leftarrow y_e + DY$</p> | <div style="border: 1px solid black; height: 250px; margin-top: 10px;"> <div style="position: absolute; top: 5px; left: 5px; width: 20px; height: 20px; border: 1px solid black; text-align: center; line-height: 20px; color: red;">×</div> </div> |
|--|---|

| | |
|---|---|
| <p>$DX :$</p> <p>DRIFT ($DX \cdot \sin(\theta/2)$) à l'entrée</p> <div style="border: 1px solid black; height: 40px; margin-top: 5px;"> <div style="position: absolute; top: 5px; left: 5px; width: 20px; height: 20px; border: 1px solid black; text-align: center; line-height: 20px; color: red;">×</div> </div> <p>DRIFT ($DX \cdot \sin(\theta/2)$) à la sortie</p> | <div style="border: 1px solid black; height: 250px; margin-top: 10px;"> <div style="position: absolute; top: 5px; left: 5px; width: 20px; height: 20px; border: 1px solid black; text-align: center; line-height: 20px; color: red;">×</div> </div> |
|---|---|

DZ :

DRIFT $(DZ \cdot \cos(\theta/2))$ à l'entrée



DRIFT ($-DZ \cdot \cos(\theta/2)$) à la sortie

-

$$M_{\theta_x} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \cos \theta_x & \sin \theta_x \\ 0 & -\sin \theta_x & \cos \theta_x \end{pmatrix} \approx \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 - \frac{1}{2}\theta_x^2 & \theta_x \\ 0 & -\theta_x & 1 - \frac{1}{2}\theta_x^2 \end{pmatrix}$$

$$M_{\theta_y} = \begin{pmatrix} \cos \theta_y & 0 & -\sin \theta_y \\ 0 & 1 & 0 \\ \sin \theta_y & 0 & \cos \theta_y \end{pmatrix} \approx \begin{pmatrix} 1 - \frac{1}{2} \theta_y^2 & 0 & -\theta_y \\ 0 & 1 & 0 \\ \theta_y & 0 & 1 - \frac{1}{2} \theta_y^2 \end{pmatrix}$$

$$M_{\theta_z} = \begin{pmatrix} \cos \theta_z & \sin \theta_z & 0 \\ -\sin \theta_z & \cos \theta_z & 0 \\ 0 & 0 & 1 \end{pmatrix} \approx \begin{pmatrix} 1 - \frac{1}{2} \theta_z^2 & \theta_z & 0 \\ -\theta_z & 1 - \frac{1}{2} \theta_z^2 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

Nous avons donné aussi le développement au deuxième ordre de la matrice (attention, dans ce cas, le déterminant est non nul). La matrice exprimant la rotation totale du dipôle est le produit des trois matrices de rotation. Cependant, ce produit n'est pas commutatif.

Par contre, en supposant que les angles de rotation sont petits et en restant au deuxième ordre, on peut alors trouver une matrice simplifiée pour la rotation des éléments :

$$M_{\theta} \approx \begin{pmatrix} 1 - \frac{1}{3} \cdot (\theta_x^2 + \theta_y^2 + \theta_z^2) & \theta_z & -\theta_y \\ -\theta_z & 1 - \frac{1}{3} \cdot (\theta_x^2 + \theta_y^2 + \theta_z^2) & \theta_x \\ \theta_y & -\theta_x & 1 - \frac{1}{3} \cdot (\theta_x^2 + \theta_y^2 + \theta_z^2) \end{pmatrix}$$

C'est cette matrice que l'on utilisera dans PARTRAN. Elle est évidemment approchée, mais très certainement moins que la connaissance de l'amplitude des erreurs. Remarquons tout de même que son déterminant est non nul (à l'ordre 4). Le choix de la diagonale n'est pas unique. C'est cependant celui qui minimise le déterminant pour des angles de rotation égaux dans toutes les directions.

La quantité de mouvement \vec{P} dans le repère $(\vec{i}_e, \vec{j}_e, \vec{k}_e)$ s'exprime alors \vec{P}^* dans le repère $(\vec{i}_e^*, \vec{j}_e^*, \vec{k}_e^*)$:

$$\vec{P}^* = M_{-\theta} \cdot \vec{P}. \quad (1)$$

Une particule de coordonnées $(x_e, y_e, 0)$ dans le repère $(\vec{i}_e, \vec{j}_e, \vec{k}_e)$ centré sur E possède alors les coordonnées (x_e^*, y_e^*, z_e^*) dans le repère $(\vec{i}_e^*, \vec{j}_e^*, \vec{k}_e^*)$ centré sur E^* telles que :

$$\begin{pmatrix} x_e^* \\ y_e^* \\ z_e^* \end{pmatrix} = T_{R \rightarrow R_e} \cdot \left((Id - M_{-\theta}) \cdot \begin{pmatrix} 0 \\ 0 \\ -l \end{pmatrix} + M_{-\theta} \cdot T_{R_e \rightarrow R} \cdot \begin{pmatrix} x_e \\ y_e \\ 0 \end{pmatrix} \right)$$

l est la demi-longueur de la corde à la trajectoire principale de l'aimant.

On a : $l = \rho \cdot \sin \varphi / 2$.

soit :

$$\begin{pmatrix} x_e^* \\ y_e^* \\ z_e^* \end{pmatrix} = \begin{pmatrix} (1 - \varepsilon) \cdot x_e - (\theta_z \cos \varphi / 2 + \theta_x \sin \varphi / 2) \cdot y_e - (\theta_y \cos \varphi / 2 - \varepsilon \cos \varphi / 2) \cdot l \\ (1 - \varepsilon) \cdot y_e + (\theta_z \cos \varphi / 2 + \theta_x \sin \varphi / 2) \cdot x_e + \theta_x \cdot l \\ -\theta_y \cdot x_e - (\theta_z \sin \varphi / 2 - \theta_x \cos \varphi / 2) \cdot y_e - (\theta_y \sin \varphi / 2 + \varepsilon \cos \varphi / 2) \cdot l \end{pmatrix} \quad (2)$$

$$\text{avec : } \varepsilon = \frac{1}{3} \cdot (\theta_x^2 + \theta_y^2 + \theta_z^2)$$

La procédure en entrée est la suivante :

Une particule entre dans le dipôle avec les coordonnées $(x_e, x'_e, y_e, y'_e, phi_e, E_e)$ dans le référentiel $(\vec{i}_e, \vec{j}_e, \vec{k}_e)$.

Sa quantité de mouvement réduite $\vec{P}_e = (x'_e \cdot p_{ze}, y'_e \cdot p_{ze}, p_{ze})$ dans $(\vec{i}_e, \vec{j}_e, \vec{k}_e)$ est calculée avec

$$p_{ze} = \sqrt{\frac{\gamma^2 - 1}{1 + x_e'^2 + y_e'^2}}.$$

Elle est ensuite transformée en \vec{p}_e^* avec (1), puis on prend : $x_e'^* = \frac{p_{xe}^*}{p_{ze}^*}, \dots$

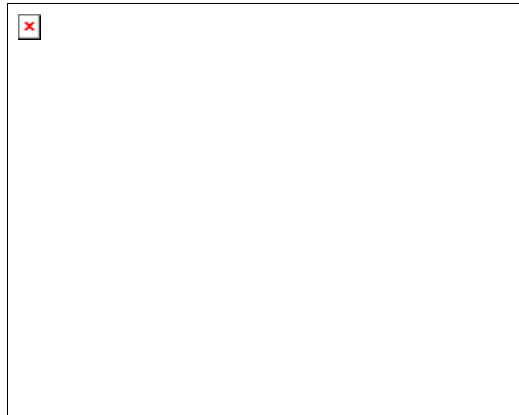
Sa position dans le repère $(\vec{i}_e^*, \vec{j}_e^*, \vec{k}_e^*)$ est obtenue avec (2).

Un drift de longueur $-Z_e^*$ doit être appliqué avant la matrice du dipôle.


• Transport dans le dipôle

Le coin d'entrée est traité en lentille mince "à la PARMILA" en prenant en compte l'influence du champ de fuite magnétique.

La matrice du dipôle telle qu'elle était utilisée avant est remplacée par un transport non linéaire dans le référentiel $(\vec{i}_e^*, \vec{j}_e^*, \vec{k}_e^*)$. Le calcul est détaillé ci-dessous :



Les coordonnées des points sont données dans le repère (X,Y) centré sur O.

 $E(0, \rho_0 + x_e), S_0(\rho_0 \cos \theta_0, \rho_0 \sin \theta_0)$

La trajectoire de la particule est un cercle de rayon de courbure ρ :

$$\rho = \frac{p}{qB} = \rho_0 \cdot \frac{1 + \delta p/p_0}{1 + \delta B/B_0}$$

Recherchons les coordonnées du centre de ce cercle C :

$$\Rightarrow \begin{cases} x_e' = \frac{X_C}{Y_E - Y_C} \\ (X_E - X_C)^2 + (Y_E - Y_C)^2 = \rho^2 \\ \begin{cases} X_C = x_e' \cdot \frac{\rho}{\sqrt{1 + x_e'^2}} \\ Y_C = Y_E - \frac{\rho}{\sqrt{1 + x_e'^2}} \end{cases} \end{cases}$$

Recherchons les coordonnées de la position de sortie S de la particule :

$$\Rightarrow \begin{cases} X_S = Y_S \cdot \tan \theta_0 \\ (X_S - X_C)^2 + (Y_S - Y_C)^2 = \rho^2 \\ \begin{cases} X_S = \sin \theta_0 \cdot \left[X_C \sin \theta_0 + Y_C \cos \theta_0 + \sqrt{\rho^2 - (X_C \cos \theta_0 - Y_C \sin \theta_0)^2} \right] \\ Y_S = \cos \theta_0 \cdot \left[X_C \sin \theta_0 + Y_C \cos \theta_0 + \sqrt{\rho^2 - (X_C \cos \theta_0 - Y_C \sin \theta_0)^2} \right] \end{cases} \end{cases}$$

On en déduit sa position de sortie $x_s = SS_0$ de la particule dans le référentiel lié à la trajectoire de référence :

$$x_s = X_C \sin \theta_0 + Y_C \cos \theta_0 + \sqrt{\rho^2 - (X_C \cos \theta_0 - Y_C \sin \theta_0)^2} - \rho_0$$

Calculons l'angle de la particule en sortie dans le référentiel lié à la trajectoire de référence $x_s' = \arctan(\theta_0 - \theta_1)$:

$$\tan \theta_1 = \frac{X_S - X_C}{Y_S - Y_C},$$

$$\Rightarrow x'_s = \frac{(Y_S - Y_C) \cdot \tan \theta_0 - (X_S - X_C)}{(X_S - X_C) \cdot \tan \theta_0 + (Y_S - Y_C)}$$

Calculons la longueur de la trajectoire : $l = \rho \cdot \theta$

$$\tan \theta = \frac{(X_S - X_C) + (Y_S - Y_C) \cdot x'_e}{(Y_S - Y_C) - (X_S - X_C) \cdot x'_e}$$

$$\Rightarrow l = \rho \cdot \arctan \left(\frac{(X_S - X_C) + (Y_S - Y_C) \cdot x'_e}{(Y_S - Y_C) - (X_S - X_C) \cdot x'_e} \right).$$

Dans le plan vertical, l'aimant peut-être considéré comme un espace de glissement :

$$\begin{cases} y'_s = y'_e \\ y_s = y_e + y'_e \cdot l. \end{cases}$$

• Sortie du dipôle

En sortie du dipôle, la particule est à la position $(x_s^*, x_s'^*, y_s^*, y_s'^*, phi_s^*, E_s^*)$.

En sortie, on a :

$$\begin{pmatrix} \vec{i}^* \\ \vec{j}^* \\ \vec{k}^* \end{pmatrix} = \begin{pmatrix} \cos \theta / 2 & 0 & -\sin \theta / 2 \\ 0 & 1 & 0 \\ \sin \theta / 2 & 0 & \cos \theta / 2 \end{pmatrix} \cdot \begin{pmatrix} \vec{i}_s^* \\ \vec{j}_s^* \\ \vec{k}_s^* \end{pmatrix} = T_{r_s^* \rightarrow R^*} \cdot \begin{pmatrix} \vec{i}_s^* \\ \vec{j}_s^* \\ \vec{k}_s^* \end{pmatrix}$$

Le vecteur \vec{P}^* dans le repère $(\vec{i}_s^*, \vec{j}_s^*, \vec{k}_s^*)$ s'exprime alors \vec{P} dans le repère $(\vec{i}_s, \vec{j}_s, \vec{k}_s)$:

$$\vec{P} = M_\theta \cdot \vec{P}^*. \quad (3)$$

Une particule de coordonnées $(x_s^*, y_s^*, 0)$ dans le repère $(\vec{i}_s^*, \vec{j}_s^*, \vec{k}_s^*)$ centré sur S^* possède alors les coordonnées (x_s, y_s, z_s) dans le repère $(\vec{i}_s, \vec{j}_s, \vec{k}_s)$ centré sur S telles que :

$$\begin{pmatrix} x_s \\ y_s \\ z_s \end{pmatrix} = T_{R \rightarrow r_s^*} \cdot \left((Id - M_\theta) \cdot \begin{pmatrix} 0 \\ 0 \\ l \end{pmatrix} + M_\theta \cdot T_{r_s^* \rightarrow R} \cdot \begin{pmatrix} x_s^* \\ y_s^* \\ 0 \end{pmatrix} \right) \quad (4)$$

soit :

$$\begin{pmatrix} x_s \\ y_s \\ z_s \end{pmatrix} = \begin{pmatrix} (1 - \varepsilon) \cdot x_s^* + (\theta_z \cos \varphi / 2 - \theta_x \sin \varphi / 2) \cdot y_s^* - (\theta_y \cos \varphi / 2 - \varepsilon \cos \varphi / 2) \cdot l \\ (1 - \varepsilon) \cdot y_s^* - (\theta_z \cos \varphi / 2 - \theta_x \sin \varphi / 2) \cdot x_s^* + \theta_x \cdot l \\ \theta_y \cdot x_s^* - (\theta_z \sin \varphi / 2 + \theta_x \cos \varphi / 2) \cdot y_s^* + (\theta_y \sin \varphi / 2 + \varepsilon \cos \varphi / 2) \cdot l \end{pmatrix} \quad (2)$$

$$\text{avec : } \varepsilon = \frac{1}{3} \cdot (\theta_x^2 + \theta_y^2 + \theta_z^2)$$

La procédure en sortie est la suivante :

Une particule sort du dipôle avec les coordonnées $(x_s^*, x_s'^*, y_s^*, y_s'^*, phi_s^*, E_s^*)$ dans le référentiel $(\vec{i}_s^*, \vec{j}_s^*, \vec{k}_s^*)$.

Sa quantité de mouvement réduite $\vec{P}_s^* = (x_s'^* \cdot p_{zs}^*, y_s'^* \cdot p_{zs}^*, p_{zs}^*)$ dans $(\vec{i}_s^*, \vec{j}_s^*, \vec{k}_s^*)$ est calculée avec

$$p_{zs}^* = \sqrt{\frac{\gamma^2 - 1}{1 + x_s'^{*2} + y_s'^{*2}}}.$$

Elle est ensuite transformée en \vec{P}_s avec (3), puis on prend : $x'_s = \frac{p_{xs}}{p_{zs}}, \dots$

Sa position dans le repère $(\vec{i}_s, \vec{j}_s, \vec{k}_s)$ est obtenue avec (2).

Un drift de longueur $-z_s$ doit être appliqué.

TraceWin News

22-02-2009 : Fixed Bug, The multiparticle transport in gap element was wrong with charge greater then one. (Edgar Sargsyan / CERN)

02-19-2009 : New: A new command, RFQ_GEOM has been added (See Help File -> commands -> set RFQ vane geometry)

10-09-2008 :TraceWin v2 see New in [Tracewin v2](#)

02-06-2008 : Fixed Bug, the NCELLS mod 2 did work correctly and the Help file was not correct too. (G. Bellodi, CERN)

01-04-2008 : Fixed Bug, the bucket plot was wrong in acceleration with RF harmonique. (M. Eshraqi, CERN)

11-03-2008 : New: A new command, SET_BEAM_PHASE_ERROR has been added (See Help File) .(Giulia Bellodi, CERN)

11-03-2008 : New: A new command, TWOTERMS, allows to select the longitudinall profil in RFQ (See Help file).

14-02-2008 : Fixed Bug, the NCELLS element didn't work in Partran when these optional parameters was not set. (Giulia Bellodi, CERN)

13-02-2008 : New: In multiparticle box, the number of space-charge calculation per meter or beta_Lambda has been added, previously it was beta_Lambda divided by 2, but in few case that was too low. The command "PARTRAN_STEP" still allows to change this parameter along the structure.

10-02-2008 : New: Halo parameters are now calculated in Toutatis simulations (see "multiparticle options Box").

10-02-2008 : New: RFQ errors are now divided in segment errors and vane errors (See ERROR_RFQ_CEL_NCPL_STAT command).

10-01-2008 : Fixed Bug, The particles exclude from emittance calculation was not (E-Esyn)/Esyn>Elimt as indicated in multiparticle options box, but (E-Esync)>Elimt.

11-01-2008 : New: The first order dispersion evolution visible in envelope chart is plotted in pink collar and the bigger order dispersion are plotted in green.

10-01-2008 : Fixed Bug, The dispersion values were wrong with input beam including dp/p error.

30-11-2007 : New: The possibility to include several cores for each computer has been include in the remote computer box.

05-10-2007 : New: Two options in error studies have been added.

- The first one : "keep all results file" allows to keep the output distribution file (*.dst) and distribution file (Dist_error_PAR.dat") for each run of a statistical error study.

- The second one : "Input or Ouput beam as reference" allows to select as reference in emittance growthes and lost calculations the input or the ouput beam of the reference linac without error.

By default it was the ouput beam. In case of "Input beam" no run without error is needed before a error study.

- 18-09-2007 : New:** In envelope plots, the dispersion value is now available.
- 17-09-2007 : Fixed Bug,** The error level printed was wrong during error calculations. (Rita Paparella, INFN)
- 10-09-2007 : New:** A new TraceWin Server able to works under Windows and Linux on remote computer is available (See TraceWinServer Help Topic).
- 25-04-2007 : Fixed Bug,** The MIN_FIELD_VARIATION command did work correctly. (A.. Ponton, IPN Orsay)
- 05-04-2007 : New:** A low-pass filter has been include in DIAG_SIZE and DIAG_DSIZE3 on the phase spread measurement.
- 21-03-2007 : Fixed Bug,** The secondary beam transport in bend was wrong in enveloppe when the particule charge of the second beam was different from the main beam. (J.L. Biarrot, IPN Orsay)
- 23-02-2007 : New:** A better memory management allows to get up the number of particle to 20.000.000.
- 21-02-2007 : New:** A new type of scrapper is available (See APERTURE element) allowing to select a beam fraction.
- 11-02-2007 : New:** TraceWin can work on 64bits windows computer (but, It's still a 32 bits code).
- 17-01-2007 : New:** Now, even if all particles are lost, the outputs are still available.
- 11-01-2007 : Fixed Bug,** The steerer didn't work in FIELD_MAP element. (J.L. Biarrot, IPN Orsay)
- 20-12-2006 : New:** A new command "SUPERPOSE_MAP" allows to superpose different FIELD_MAP elements and shift them. (See Help File).
- 22-11-2006 : Fixed Bug,** TraceWin version "1.4.1.67 to 68" doesn't work correctly in DTL_CEL element. (Edgar Sargsyan / CERN)
- 22-11-2006 : Fixed Bug,** TraceWin version "1.4.1.67" doesn't work in Edge element with negative charge.(Rita Paparella, INFN)
- 13-11-2006 : Fixed Bug,** TraceWin version "1.4.1.66" doesn't support toutatis.
- 22-10-2006 : New:** A new command allowing to easely duplicate elment or elements block has been included (See REPEAT_ELE command Help).
- 21-10-2006 : New:** The number of elements was limited to 5000. This limitation has been remove.
- 06-10-2006 : New:** A new option in DRIFT element has been added in order to define rectangular aperture, see DRIFT element help (C. prostinar, Rutherford Lab.)
- 04-10-2006 : New:** A 4 order solenoid field map has been added, see MUTIPOLE element help.
- 18-09-2006 : New:** A peperpot element has been added, see APERTURE element help.
- 12-09-2006 : New:** For Linux Wine emulation use command "wine Trace_Win.exe wine". And add dll files "Riched20.dll" and "Riched32.dll"
- 22-07-2006 : New:** XP design interface. The main menu "Theme" allows to swich from NT theme to XP theme.
- 21-07-2006 : Fixed Bug,** That sometime froze the Envelope plots during optimization.
- 20-07-2006 : Fixed Bug,** That sometime froze the TraceWin Servor under XP (twserver).
- 15-07-2006 : New:** New compiler used.
- 14-06-2006 : Fixed Bug,** The losses (W) calculation don't taking into account charge of particle. (J.L. Biarrot, IPN Orsay)
- 12-05-2006 : Fixed Bug,** The fringe-field correction applied in the edge angle output in bend was wrong in multiparticle mode if the dipole was cut in several part. (J.L. Biarrot, IPN Orsay)
- 11-05-2006 : New:** A new diagnostic element "DIAG_ACHROMAT", see Help file.
- 11-05-2006 : New:** 2 new commands "ADJUST_STERRER_BX" and "ADJUST_STERRER_BY", see Help file.
- 15-03-2006 : New:** A name (8 caracteres max) can be attributed to each element. The syntax is "Name : Type_element".
ex: "QP12 : QUAD 10.0 100. 100."
- 15-03-2006 : New:** TraceWin Translates the data file to BETA code syntax.
- 07-03-2006 : New:** A new diagnostic has been introduce to minimize the beam luminosity: DIAG_LUMINOSITY (See Help file).
- 30-01-2006 : New:** The "ADJUST" command syntax has been changed (See Help file).
- 10-01-2006 : New:** The elipe and envelope charts (multiparticle or envelope simulations) can be saved or copied in vector format (EMF file, windows meta file format), that gives much better quality output.
- 10-01-2006 : Fixed Bug,** The fringe-field correction applied in the edge angle output in bend was wrong in multiparticle mode. (J.L. Biarrot, IPN Orsay)
- 01-12-2005 : New:** A new optimization algorithm has been implemented, you can select it by the "Optimization Options" boxe.

- 28-11-2005 : New:** The documentation has been revised (See Help file).
- 24-11-2005 : New:** A new diagnostic element had been implemented (Diag_Waist) allowing to impose 'Alp' twiss parameter to zero (See Help file).
- 23-11-2005 : New:** The losses in bend are now estimated with the gap of magnets defined in edge element and aperture defined in bend element.
- 22-11-2005 : Fixed Bug,** The project loading for specific project did not works. (J.L. Biarrot, IPN Orsay)
- 17-11-2005 : Fixed Bug,** In output charts the gradation of colours were badly saved in BMP or JPG format.
- 14-09-2005 : New:** A synoptic allowing to see bend deviation and edge orientation is now available in tab-set-graph (button "synoptic").
- 03-09-2005 : New:** All Tranfert matrix term evolutions can be plotted (See "Transfer Matrix" button in "Graphs" tab-sheet).
- 03-11-2005 : Fixed Bug,** The "SET_ACHROMAT" command has been improved, see Help file.
- 01-11-2005 : Fixed Bug,** The "APERTURE" element plot was wrong in Y plane . (J.L. Biarrot, IPN Orsay)
- 18-10-2005 : Fixed Bug,** In Bend calculation (Robin Ferdinand / Asclepios).
- 01-09-2005 : New:** The multiparticle input beam distribution type selection and generator has been changed (see multiparticle option).
- 11-09-2005 : Fixed Bug,** the data file could be partialy erase when it was located in remote disk, some protections have been inserted in order to avoid that. (Rainer Cee/GSI)
- 01-09-2005 : New:** The beam matrix is now visible (In 'Beam parameter' & 'Twiss parameter') and can be changed in the 'Beam parameter' in order to add some extra special matrix terms.
- 01-07-2005 : New:** The TraceWin Server has been corrected in order to spare the Windows XP firewall.
- 03-09-2005 : Fixed Bug,** In multiparticle simulation when beam current is lower than 1e-5 mA (Robin Ferdinand / Asclepios).
- 06-07-2005 : Fixed Bug,** The project manager failed in a new installation in a new machine.
- 01-07-2005 : New:** A new command has been developped: "SET_SYNC_PHASE" (see help file).
- 26-06-2005 : New:** A new element has been developped: Beam Rotation (see help file) (Robin Ferdinand / Asclepios).
- 22-06-2005 : New:** A new element has been developped: Thin lens (see help file).
- 06-06-2005 : Fixed Bug,** The Gaussup space-charge module was failed. (Maud Baylac / LPSC-Grenoble).
- 06-06-2005 : Fixed Bug,** 'SET_TWISS' command did not work with CW beam.
- 31-05-2005 : New:** A new tool allowing to visualize the field maps has been implemented.
- 26-05-2005 : New:** "The files needed for TraceWin Error Help Email are now automaticly included.
- 24-05-2005 : Fixed Bug,** The phase advance wihtout current (structure) charts were sometime wrong. (Edgar Sargsyan / CERN)
- 12-05-2005 : Fixed Bug,** FIELD_MAP option 4 did not work correctly in multiparticle simulation.
- 26-04-2005 : Fixed Bug,** TraceWin failed in first run case. (Andeas Sauer)
- 26-04-2005 : New:** 2 new parameters in 'options' -> 'multiparticle code options' allows to exclude or not, particles in emittance calculations according to phase max and dE/E max.
- 22-04-2005 : Fixed Bug,** Electrostatic accelerating lens did not take into account the beam charge.
- 17-04-2005 : New:** Some longitudinal diagnostics have been implemented and some modifications have been made on the associated adjust commands (See Help file).
- 16-04-2005 : New:** The optimization parameters are summarize in a new box "Optimization options" in "Options".
- 10-04-2005 : Fixed Bug,** which deletes without reason the project manager configuration.
- 06-04-2005 : New:** Bend with field index has been implemented in partran (N. Pichoff). Warning, at present time only first order is considered. (R.Ferdinand / Ganil)
- 05-04-2005 : New:** Some new data have been include in the error file result, "Statistic_Errors_1_PAR.txt" (rms sizes, halo parameters), see "Error file result" in Help. Some dedicated buttons have been implemented in Tab-Set "Errors" in order to exploit these new data . (Edgar Sargsyan / CERN & E. Froidefond / LPSC-Grenoble).
- 25-03-2005 : Fixed Bug,** In few cases, a "EzeroDivide" error occurred for continue beam. (Ferdinand Robin / Ganil)
- 17-03-2005 : New:** All the dissipated powers calculated in TraceWin takes in account the energy of the particles lost rather than the beam synchronous energy.

15-03-2005 : New: The number of particle in TraceWin & Partran are now no limited par the memory system. If the number of particle is too big compare to available memory the beam will be skipped in few parts.

15-03-2005 : Fixed Bug, load a imported project wasn't imported the good mass. (Maud Baylac / LPSC-Grenoble).

07-03-2005 : Fixed Bug, in Button "Transfer Matrix" in tab-sheet "Graphs" which makes it sometime unstable.

16-02-2005 : New: A new option in multiparticle code allowing to increase the number of space charge calculation in a DTL cel has been implemented (See mulitparticle options).

13-02-2005 : New: A new type of statistique errors has been include: The input beam errors (See button 'Error Selection' and command "ERROR_BEAM_STAT" & "ERROR_BEAM_DYN").

04-02-2005 : Fixed Bug, in DTL cell with error, partran and TraceWin wasn't agree.(Maud Baylac / LPSC-Grenoble)

01-02-2005 : Fixed Bug, in thin GAP (without beta parameter) a bug had been fixed in the longitudinal plan.(Edgar Sargsyan / CERN)

21-01-2005 : New: Two new phase space coordinates have been developed: "theta" and "thetap". They only available in multiparticle output. (N. Pichoff / CEA).

20-01-2005 : New: A new file containing random errors list are generated in calculation path (E. Froidefond / LPSC-Grenoble).

18-01-2005 : New: Some usual examples have been added in the "Example" path of project manager.

14-01-2005 : Fixed Bug, which deletes without reason the data file.

10-01-2005 : New: The documentation has been updated.

06-01-2005 : Fixed Bug, in output chart, the axis color managment was wrong.

03-01-2005 : New: A new parameter "Duty cycle" has been implemented, it is used in power losses calculation.

15-11-2004 : New: A new command has been implemented "SET_ACHROMAT" see help file.

15-12-2004 : Fixed Bug, in Partran: The longitudinal coefficient for distribution 36 was wrong. (Edgar Sargsyan / CERN)

07-11-2004 : New: A new element has been implemented: MULTIPOLE element. It is based on a field map generation (Concetta Ronsivalle / Frascati).

30-11-2004 : Fixed Bug, in magnetic bend treatment according to the main or second beam (Rainer Cee/GSI)

25-10-2004 : Fixed Bug, in the synchronous phase calculation RF FIELD_MAP element.

22-10-2004 : Fixed Bug, in the longitudinal losses calculation according |E-Eo| and |Ph-Ph0| criterion (E. Froidefond / LPSC-Grenoble).

24-09-2004 : Fixed Bug, in mismatching error studies, the results was dependent of the number of step. (J.L. Biarrotte / IPN).

26-08-2004 : New: Some datas are been added in the statistical error result file in order to obtain X and Y residual orbit and rms, max and min losses results.

16-08-2004 : Fixed Bug, in the funnel_gap element (N. Pichoff / CEA).

22-06-2004 : Fixed Bug, in the compress level mode (Andreas Sauer Frankfort).

07-05-2004 : New: The diagnostic errors can be disabled in tab-set "option".

06-05-2004 : Fixed Bug, Zp, phase, energy amplitude in longitudinal plots were wrong, with in accelerating mode.

05-05-2004 : Fixed Bug, Z sign in longitudinal plot was wrong.

04-05-2004 : New: A new RFQ error is been implemented (longitudinal displacement error by segment).

28-24-2004 : Fixed Bug, Bend with negative angle failed. (Marta Novati/ INFN).

03-04-2004 : New: Commands "ERROR_BEND_NCPL_DYN" and "ERROR_BEND_NCPL_STAT" in ordre to apply magnetic errors on the dipole (See Help file).

01-04-2004 : New: Error treatment concerning dipole has been include in TraceWin and Partran.

23-03-2004 : New: Element quadrupole electrostatic (QUAD_ELE), see Help.

18-03-2004 : New: Diagnostic elements, see Help.

03-03-2004 : New: The coeff: sqrt(5) between rms and total size can be changed by using the scale button of envelope or distribution output charts.

03-01-2004 : New: Flag to enable or not Diagnostic computation has been included in "Option" Tab-Set.

01-29-2004 : New: Button allows to extract beam parameters form the dst file.

01-29-2004 : Fixed Bug, Output field amplitude with longitudinal errors was wrong. (J.L. Biarrotte / IPN).

- 01-23-2003** : **Warning**: In the version 1.15.0.1 all the errors selected have been deselected for compatibility reason.
- 01-23-2003** : **New**: Command to minimize field variation during matching "MIN_FIELD_VARIATION".
- 01-18-2004** : **Fixed Bug**, Dynamic errors was 2 times applied for adjacent quadrupoles.
- 01-04-2004** : **Fixed Bug**, in the chopper element, about the particle charge. (P. Bertrand / GANIL).
- 12-15-2003** : **New**: TOUTATIS code generates plt file like PARTRAN (See options of mutiparticle codes).
- 11-27-2003** : **Fixed Bug**, in the beam phase and energy centroid extracted from PARTRAN output file. (J.L. Biarrotte / IPN).
- 11-24-2003** : **Fixed Bug**, in the "PARTRAN_STEP" command where the step did not depend of the element length. (N. Pichoff / CEA).
- 11-22-2003** : **Fixed Bug**, in TraceWin Server. It did not work under XP.
- 11-21-2003** : **New**: A new file containing all the steerer values in the statistical error studies: "Steerer_Values_Tot_xxx.txt". The "Steerer_Values.txt" contains the last run.
- 11-16-2003** : **Fixed Bug**, in the FIELD_MAP element in the aperture map reading.
- 11-16-2003** : **Fixed Bug**, in the particles file creating procedure in DC beam. (N. Pichoff / CEA).
- 11-14-2003** : **New** possibility to save in the particle output distribution charts the particles in an ascii file. Use "Save" button, then selects "txt" as file type.
- 11-12-2003** : **Fixed Bug**, in the phase space chart menus. (P. Bertrand / GANIL).
- 11-10-2003** : **New** , The new command "SET_BEAM_E0_P0" allows to match the beam energy and phase.
- 11-10-2003** : **New** , The possibility to match the beam after the errors are applied, has been implemented (MATCH_FAM_FIELD -x and MATCH_FAM_FIELD -x) use a negative number for the matching number.
- 11-03-2003** : **Fixed Bug**, in the "SET_TWISS" command: Match only the Alpha parameter using flag system doesn't work when the beta parameter was equal to 0 (E. Froidefond / LPSC-Grenoble).
- 10-29-2003** : **New** , The emittance values from particle file are not extracted when the file are opened.
- 10-27-2003** : **New** , A separation between the scale of the envelope charts and phase space charts has been implemented
- 10-09-2003** : **New** possibility to select a compress format for the multiparticle output file (Only the last PARTRAN version or the Library version are compatible, otherwise use 0%)
- 10-09-2003** : **New** the "MONET" and "SCDYN" codes are removed, because not used.
- 10-02-2003** : **Fixed Bug**, the "jpg" pictures cannot be used as background in the output charts (R.Duperrier CEA-France).
- 10-02-2003** : **Fixed Bug**, the value of "nbr of lattice to optimize" saved was always 0 (J. L. Biarrotte IPN-France).
- 10-02-2003** : **New** matching result format has been implemented in the result file ("Matching_Between_Section_#_to_#+1" -> "Matching_#")
- 07-16-2003** : **Fixed Bug**, in few case the help file could stop a calculation.
- 07-15-2003** : **New** command: "SET_BEAM_PHASE_ADV allows to set beam phase advances in space charge, (see help).
- 07-15-2003** : **New** way to use "SET_SIZE" and "SET_SIZE_MAX" command, (see help).
- 07-15-2003** : **New** : Some more information during optimisation are now plotting.
- 07-15-2003** : **New** possibility to use more than 1 "SET_TWISS" command in the same optimisation.
- 07-11-2003** : **New** space charge routine, "CE_CYL" (for CEA-DAM developments) has been implemented .
- 07-10-2003** : **New** possibility to select the number of step, when you save in data file format the envelope behaviour in the envelope charts.
- 07-10-2003** : **Fixed Bug**, in static magnetic field map reading (type 0070).
- 07-10-2003** : **Fixed Bug**, in steerer command where the sign of the charge was ignored.
- 06-25-2003** : **New** possibility to obtain some help about element or command syntax directly in the edit tab-set by using right mouse button.
- 06-24-2003** : **Fixed Bug**, in envelope and distribution charts the last button did not work in Windows NT4 (J-L. Coacolo IPN-France).
- 06-19-2003** : **New** possibility to save in the particle output distribution charts the selected particles in a "dst" file. Use "Save" button, then selects "dst" as file type.
- 06-19-2003** : **New** possibility to fix an electrostatic or magnetostatic field like magnetic field of earth by using the new command "FIELD Bx(T) By(T) Bz(T) Ex(V/m) Ey(V/m) Ez(V/m)". For the moment, Bz has no influence.

06-16-2003 : **New** possibility to obtain all type of output phase spaces, Example (Z-X or Phase-X).

06-04-2003 : **New** possibility to taking into account the aperture in field map when the option "use aperture element" is checked.

06-02-2003 : **Fixed Bug**, in space charge compensation interpolation in field map element.

05-26-2003 : **New** command: "SHIFT dx dy" allows to move a element in the x and y direction.

05-26-2003 : **New** possibility to load more than one field map in memory in the same time (see tabsheet: "Options"), in order to accelerate computing.

05-20-2003 : **Fixed Bug**, emittance growth in field map *.edz due to approximation in phase synchronous calculation.

05-19-2003 : **New** element pour Partran developers named "PARTRAN_ELE N° lg(m) par1 par2 ... par19".

05-18-2003 : **New** possibility to save in data file format the envelope behaviour in the envelope charts (uses save button).

New in TraceWin V2

The project manager has been abandoned. Now, TraceWin use a usual project file system (*.ini). Our old project file can still be read by the new TraceWin which will convert them. Be careful, your own particles defined in old project file will be lost.

Each TraceWin instance manages a copy of the input data files, that means you can now perfectly use several TraceWin(s) in the same time with the same projects file and data files without conflict between each over.

The data file and result file are not automatically saved in TraceWin v2.

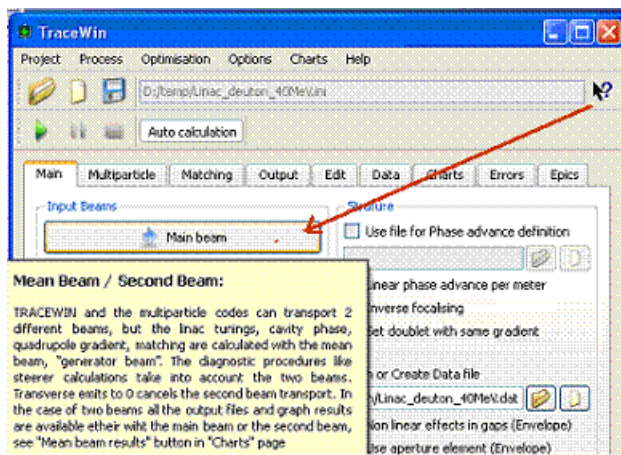
The help has been upgraded. A big difference about it is; the “*ways of using*” old help file part it is directly included in the GUI. It’s vey important to have a look in the new short “*Ways of using*” help part to be able to use it.

General description

TRACE_WIN’s program is organized in 8 pages and 3 toolbars. You can find more details about these pages below.

| | |
|--------------------|--|
| Menu | : Shows the 10 last current opened projects. |
| First ToolBar | : Open save or create a project file (configuration file, *.ini), show the current open project. |
| Second ToolBar | : To launch the process, to have a break or stop it and set “auto_calculation”. |
| Thirs ToolBar | : Visible only during matching, to stop or visualise criteria variation. |
| Main page | : To set input beam parameters, structure options and calculation options. |
| Multiparticle page | : To configure multiparticle code options. |
| Matching page | : To configure beam matching options |
| Output page | : To visualize the calculation stages |
| Edit page | : To modify or visualize the main input and output files. |
| Data page | : To visualize the elements and commands list from the data file |
| Charts page | : To visualize the results with plots. |
| Errors page | : To parameterize the error study and visualize results. |
| Epics page | : For EPICS virtual machine configuration |

Each input or widget item of TraceWin GUI owns explanatory text. The default way for users to view the help is to move the focus to the relevant widget and press Shift+F1. The help text appears immediately. A second ways is to use the help button, see following picture.

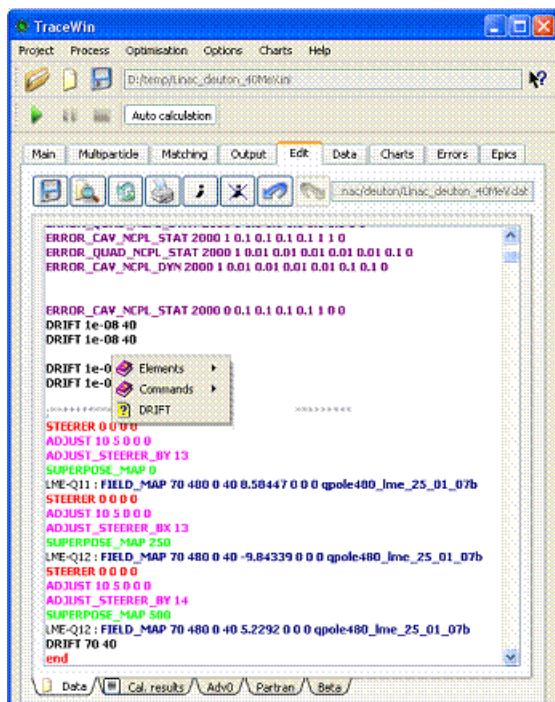


TRACE_WIN process is organized in several stages. The stages can automatically run one behind the other or not ("**Auto calculation**" button of ToolBar). Some of them can be disabled according to some options and commands.

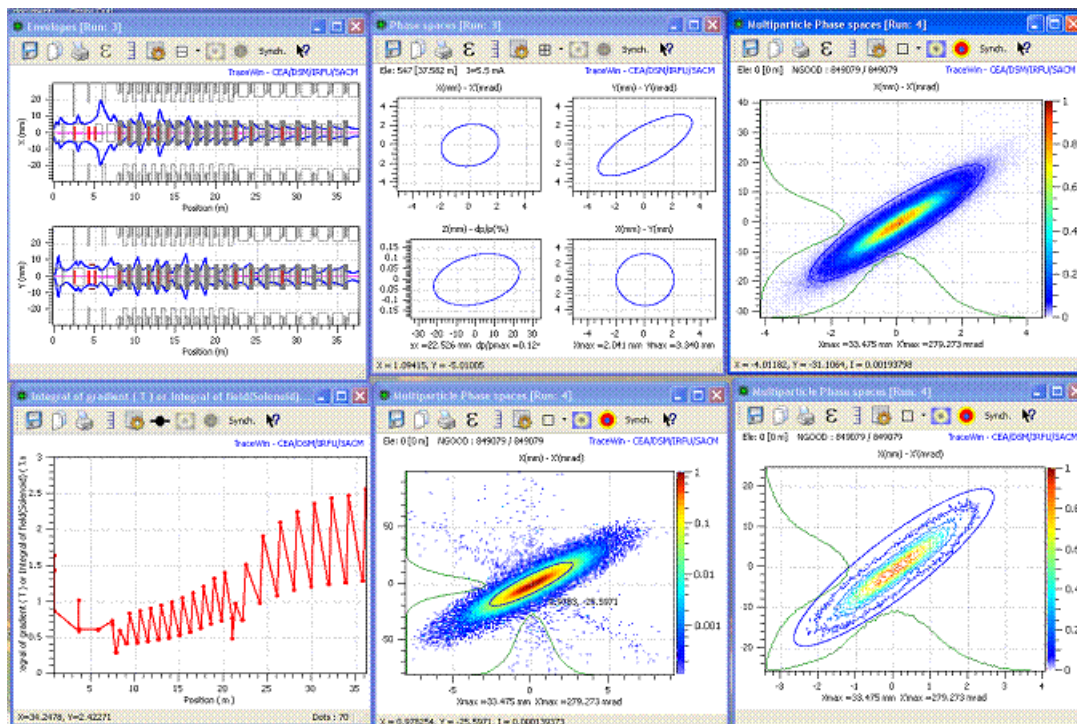
| | | | Results saved | Affected or used beam | |
|----|---|---|---------------|-----------------------|-----|
| | Different stages | Needed Condition(s) | | Main | Sec |
| 1 | Read input data file and set tab in " Data " page | Data file | | | |
| 2 | Transport of the reference particle | | | X | |
| 3 | Set Phase advance law (set Quadrupole, Solenoid, Field map strengths) | SET_ADV commands in data file or " Use file for phase adv definition " checked LATTICE commands in data file | | X | |
| 4 | Read particle files | Particle file defined in " Main " page | | X | X |
| 5 | Calculate the input matched beam | "Calculate match beam" checked in " Matching " page | | X | X |
| 6 | Set quads or cavity strengthes to match the beam through the different linac sections or set Twiss parameters (In order of theirs positions) | "Matching with family" checked in page " Matching " MATCH_FAM commands and (LATTICE or SET_TWISS command) in data file | X | X | |
| 7 | Diagnostics (Example: Steerers) calculations (In the order of theirs numbers) | "Match with diagnostics" checked in page " Matching " Diagnostic elements and Adjust commands in data file | X | X | X |
| 8 | Input beam distribution (*.dst) is adjusted in order to fit the input beam defined in " Main " page or to fit the input matched beam calculated in (5) | "Calculate match beam" -> "With partran" checked in " Matching " page "Use particle file" in "Input distribution type" in " Multipart " page Particle file defined in " Main " page | | | |
| 9 | Repetition of the preceding stages (5)(6)(7) using mutliparticle code | On or several options " With partran " checked in " Match " page | X | X | X |
| 10 | Random errors generator initialised | "Reinitialize random generator" checked in " Main " page. | | | |
| 11 | Apply Static errors | "Include error defined in..." checked in " Main " page " ..Data No " in " Main " page set to errors defined in " Errors setup " of page " Error " " ERROR_xxx_STAT_xxxx " commands in data file | | X | X |
| 12 | Diagnostics (Example: Steerers) calculations | "Match wiht diagnostics" in " Match " page Diagnostic elements and ADJUST commands in data file | | X | X |
| 13 | Repetition of the preceding stages (10 to 13) using mutliparticle code | "Launch Partran" checked in " Multiparticle " page | | X | X |
| 14 | Apply dynamic errors | "Include error defined in..." checked in " Main " page " ..Data No " in " Main " page set to errors defined in " Errors setup " of page " Error " " ERROR_xxx_DYN_xxxx " commands in data file | | X | X |
| 15 | Calculates the transport line envelope | Always | | X | X |
| | Losses and beam parameters variations | " Nbr of particles " in " Main " page greater | | | |

| | | | | | |
|----|--|--|--|---|---|
| | estimated in envelope transport | then 10 "Use aperture element" checked in "Main" page | | X | X |
| 16 | Write new data file in " <i>calculation directory</i> " | Always | | | |
| 17 | Make the Error studies (envelope). N linacs, Loop with stage (10,11,12,14,15) | " <i>Study Envelope</i> " checked in " <i>Errors</i> " page and error selection done | | X | X |
| 18 | Make the Error studies (Particles). N linacs, Loop with stage (10,11,12,13,14,15,19) | " <i>Study Multiparticle</i> " checked in " <i>Errors</i> " page and error selection done | | X | X |
| 19 | Write input files of multiparticle codes PARTRAN and TOUTATIS and launch them | "Launch Partran or Toutatis" checked in "Multiparticle" page | | X | X |

Somme help about element or command in the data file editor can be obtained by using the right mouse button. If you don't release it, the element number is shown before.



Different kind of plot, **explore and test** all the top buttons to well understand all available options.



Zoom with left button mouse.

Zoom back with double click.

Move, only after zooming, with right button mouse.

Save, in several picture format, pdf, ps, data ASCII file.

Copy in png format.

Plot options: Color, size, font, dots type... are configurable.

Synch.: All chart can be synchronize with current process.

Envelope plot:

Element info can be obtain by right button menu on a element, the envelope types (X , Y , X' , Y' , Z , Z' , $Phase$, $Energy$, Z , dp/p , ...) can be selected by using this menu.

The plot area contains at most 6 plots.

Phase space plot:

- The selection of the **plot phase-space coordinates** is obtained by clicking right on the chart. In this case, one can either select one of the 8 proposed 2D phase-space or chose the variable for the horizontal and vertical axis.
- On can “**Select**” or “**Unselect**” the visible particles on the chart area. The sizes, colors of the selected or non selected particles are configurable by using the option button (very convenient to localize in 6D phase-space some particles and study their behavior in a line).
- The plot area contains at most 6 plots.
- **The Emittance button “ε”** calculates statistic information about plotted distribution. The code does the same calculation in the total number of phase-space plotted whatever they are.

$$\text{Emit [rms]} = \sqrt{\langle (x - \langle x \rangle)^2 \rangle \cdot \langle (y - \langle y \rangle)^2 \rangle - \langle (x - \langle x \rangle) \cdot (y - \langle y \rangle) \rangle^2}$$

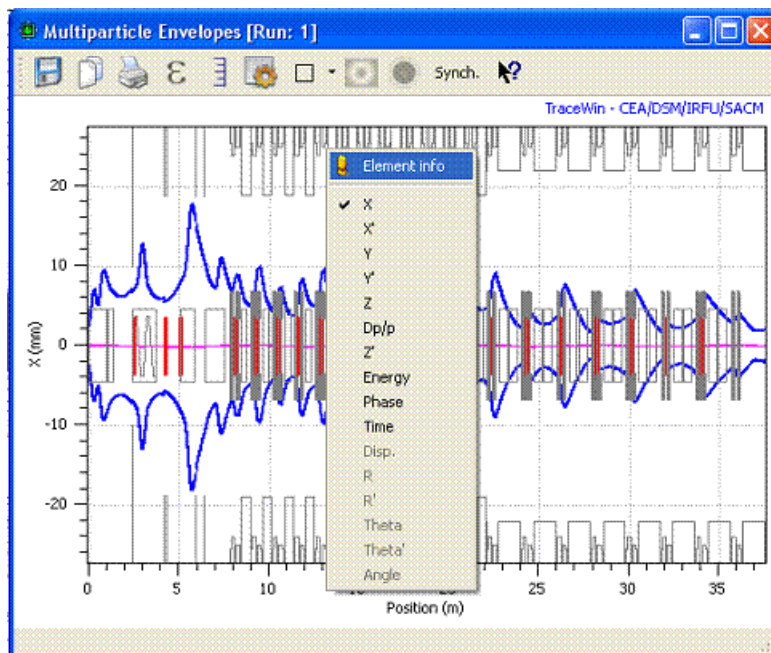
$$\text{Bet} = \langle (x - \langle x \rangle)^2 \rangle / \text{Emit [rms]}$$

$$\text{Alp} = - \langle (x - \langle x \rangle) \cdot (y - \langle y \rangle) \rangle / \text{Emit [rms]}$$

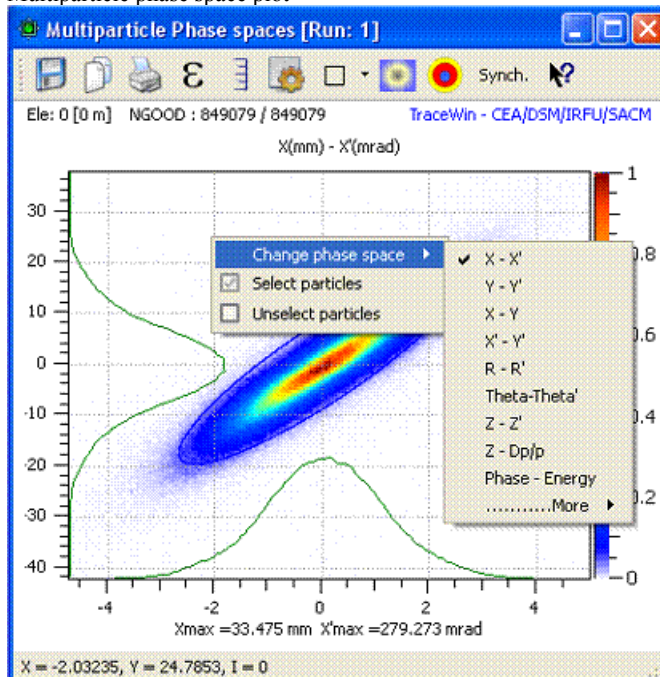
Emit [xx%] gives the ellipsis surface divided by π homothetic to the rms ellipsis containing xx% of the beam particles. The xx is either the calculated fraction of particles inside an homothetic ellipses whose area is N times the rms ellipses, or a given fraction of the beam. The plotted ellipses are those last ellipses. Position of the c.o.g. and size of the beam can be also calculated. Finally, one can plot 2 graphs, the first showing the evolution of the number of particles outside a given emittance (scaled to the rms emittance calculated above), the second showing the evolution of the number of particles outside a given size (scaled to the rms size calculated above). The last button (σ beam) shows 6x6 the beam matrix.

The emittances are calculated according to “**Energy and Phase limit**” defined in “**Multiparticle**” page.

Envelope plot example



Multiparticle phase space plot



Errors study

Copy following examples:

[examples\Supra_4.ini](#)
[examples\Supra_4.dat](#)
[examples\Supra_4.cal](#)

[examples\Supra_5.ini](#)
[examples\Supra_5.dat](#)
[examples\Supra_5.cal](#)