

PH110 : Waves and Electrodynamics

Tutorial 6

13. Find the potential in the infinite slot of Ex 3.3 if the boundary at $x=0$ consists of two metal stripes : one from $y=0$ to $y=a/2$ and held at constant potential V_0 and the other from $y=a/2$ to $y=a$ at potential $-V_0$.

$$\frac{\partial^2 V}{\partial x^2} + \frac{\partial^2 V}{\partial y^2} = 0$$

boundary conditions

$$V = 0 \text{ at } y = 0, a$$

$$V \rightarrow 0 \text{ as } x \rightarrow \infty$$

$$V = \begin{cases} V_0 & \text{at } x=0, 0 < y < a/2 \\ -V_0 & \text{at } x=0, a/2 \leq y < a \end{cases}$$

$$\text{Let } V = (A e^{kx} + B e^{-kx}) (C \sin ky + D \cos ky)$$

$$V = 0 \text{ at } y = 0$$

$$0 = (A e^{kx} + B e^{-kx}) (0 + D)$$

$$\therefore D = 0$$

$V = 0 \text{ at } y \rightarrow \infty$

$$0 = (A e^{kx} + B e^{-kx}) (C \sin k y)$$

$e^{kx} \rightarrow \infty \text{ as } x \rightarrow \infty$

$$\therefore A = 0$$

$$\therefore V(x, y) = C_1 e^{-kx} \sin k y \text{ sinky}$$

$V = V_0(y) \text{ at } x = 0$

$$V_0(y) = C \sin(k y)$$

$V = 0 \text{ at } y = a$

$$\therefore \sin(k a) = 0$$

$$ka = n\pi$$

$$k = \frac{n\pi}{a}$$

$$V(x, y) = C e^{-\frac{n\pi x}{a}} \sin\left(\frac{n\pi y}{a}\right)$$

$$V(x, y) = \sum_{n=1}^{\infty} C_n e^{-\frac{n\pi x}{a}} \sin\left(\frac{n\pi y}{a}\right)$$

$$V(0, y) = \sum_{n=1}^{\infty} C_n \sin\left(\frac{n\pi y}{a}\right) = V_0(y)$$

$$C_n = \frac{2V_0}{a} \left\{ \int_0^{a/2} \sin(n\pi y/a) dy - \int_{a/2}^a \sin(n\pi y/a) dy \right\}$$

$$C_n = \frac{2V_0}{a} \left\{ -\frac{\cos(n\pi y/a)}{(n\pi/a)} \Big|_0^{a/2} + \frac{\cos(n\pi y/a)}{(n\pi/a)} \Big|_{a/2}^a \right\}$$

$$C_n = \frac{2V_0}{\pi n} \left\{ -\cos\left(\frac{n\pi}{2}\right) + 1 + \cos(n\pi) - \cos\left(\frac{n\pi}{2}\right) \right\}$$

$$C_n = \frac{2V_0}{\pi n} \left\{ 1 + (-1)^n - 2\cos\left(\frac{n\pi}{2}\right) \right\}$$

$$1 + (-1)^n - 2\cos\left(\frac{n\pi}{2}\right) = \begin{cases} 0 & , (n-2)\%4 \neq 0 \\ 4 & , (n-2)\%4 = 0 \end{cases}$$

Therefore

$$C_n = \begin{cases} \frac{2V_0}{\pi n} & \text{if } n = 2, 6, 10, 14, \dots \\ 0 & \text{otherwise} \end{cases}$$

$$\therefore V(x,y) = \sum_{n=2,6,10,\dots} \frac{2V_0}{\pi n} \frac{e^{-n\pi x/a}}{n} \sin(n\pi y/a)$$

3-27. A sphere of radius R , centered at the origin, carries charge density

$$\rho(r, \theta) = k \frac{R}{r^2} (R - 2r) \sin\theta,$$

where k is a constant, and r, θ are usual spherical coordinates. Find approximate potential for points on z -axis far from sphere.

Monopole Term

$$\phi = \int \rho d\tau = kR \int \frac{1}{r^2} (R - 2r) \sin\theta r^2 \sin\theta dr d\theta d\phi$$

but the r integral

$$\int_0^R (R - 2r) dr = (Rr - r^2) \Big|_0^R = 0,$$

hence, $\phi = 0$.

Dipole Term

$$\int r \cos\theta \rho d\tau = kR \int r \cos\theta \cdot \frac{1}{r^2} \sin\theta (R - 2r) r^2 \sin\theta dr d\theta d\phi$$

but the θ integral

$$\int_0^\pi \sin^2\theta \cos\theta d\theta = \frac{\sin^3\theta}{3} \Big|_0^\pi = 0$$

so, dipole contribution is zero.

Quadrupole Term

$$\int r^2 \left(\frac{3}{2} \cos^2 \theta - \frac{1}{2} \right) \rho d\tau$$

$$= \frac{kR}{2} \int \int r^2 (3\cos^2 \theta - 1) \left[\frac{1}{r^2} (R-2r) \sin \theta \right] r^2 \sin \theta dr d\theta d\phi$$

r integral

$$\int_0^R r^2 (R-2r) dr = \left. \frac{r^3 R}{3} - \frac{r^4}{2} \right|_0^R = -\frac{R^4}{6}$$

θ -integral

$$\int_0^\pi ((3\cos^2 \theta - 1) \sin^2 \theta) d\theta = 2 \int_0^\pi \sin^2 \theta d\theta - 3 \int_0^\pi \sin^4 \theta d\theta$$

$$= 2 \left(\frac{\pi}{2} \right) - 3 \left(\frac{3\pi}{8} \right) = -\frac{\pi}{8}$$

ϕ integral

$$\int_0^{2\pi} d\phi = 2\pi$$

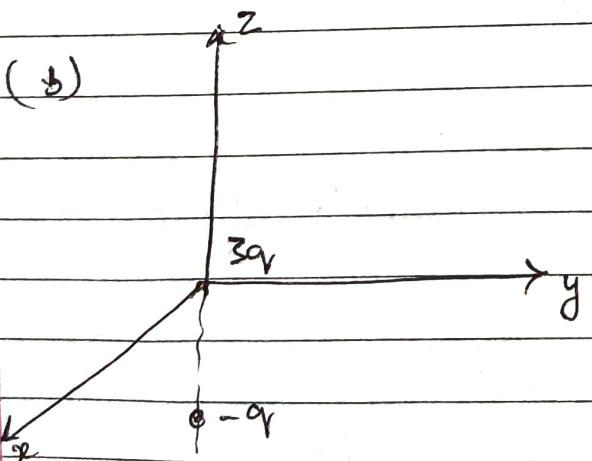
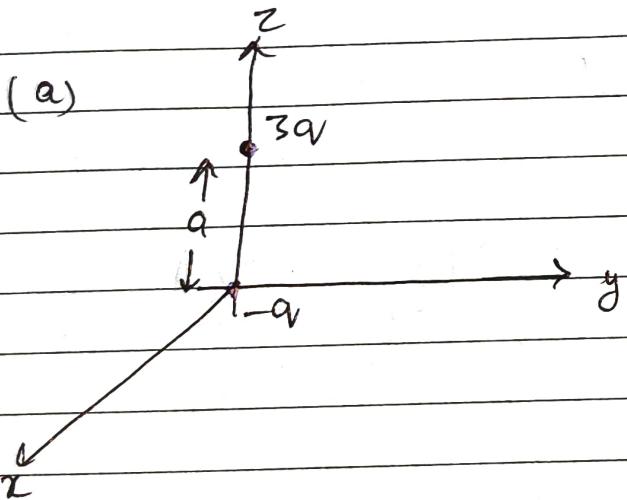
$$\therefore \text{the whole integral is} = \frac{1}{2} kR \left(-\frac{R^4}{6} \right) \left(-\frac{\pi}{8} \right) (2\pi)$$

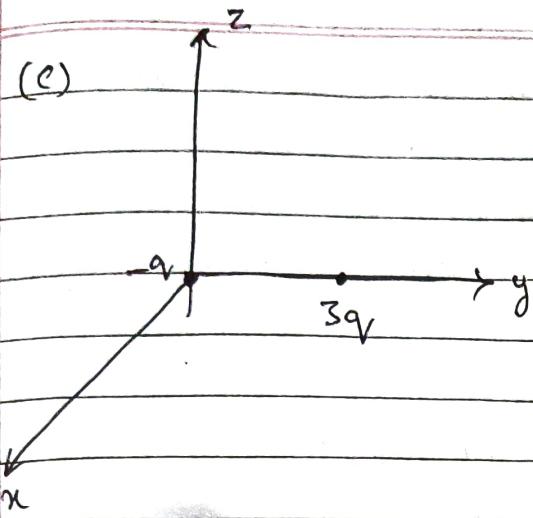
$$= \frac{k\pi^2 R^5}{48}$$

For point P on z-axis, the approximate potential is

$$V(z) \approx \frac{1}{4\pi\epsilon_0} \frac{k\pi^2 R^5}{48 z^3}$$

- (3) Two point charges $3q$ and $-q$, are separated by distance a . For each of the arrangements, find (i) the monopole moment, (ii) dipole moment
 (iii) approximate potential at large r .





(a) monopole moment $\Theta = 2q$

$$\text{dipole moment} = 3qa \hat{z}$$

$$V \approx \frac{1}{4\pi\epsilon_0} \left[\frac{\Theta}{r} + \frac{\mathbf{P} \cdot \hat{\mathbf{r}}}{r^2} \right] = \frac{1}{4\pi\epsilon_0} \left[\frac{2q}{r} + \frac{3qa \cos\theta}{r^2} \right]$$

(b) monopole moment $\Theta = 2q$

$$\text{dipole moment } \mathbf{p} = qa \hat{z}$$

$$V \approx \frac{1}{4\pi\epsilon_0} \left[\frac{2q}{r} + \frac{qa \cos\theta}{r^2} \right]$$

(c) monopole moment = $2q$

$$\text{dipole moment} = 3qa \hat{y}$$

$$V \approx \frac{1}{4\pi\epsilon_0} \left[\frac{2q}{r} + \frac{3qa \sin\theta \sin\phi}{r^2} \right]$$

(33)

A "par" dipole is situated at origin, pointing in z-direction.

(a) What is the force on a point charge q at $(x, 0, 0)$?

(b) What is the force on q at $(0, 0, a)$?

(c) How much work does it take to move q from $(a, 0, 0)$ to $(0, 0, a)$?

(a) This point is at $r=a$, $\theta=\pi$ and $\phi=0$. So

$$\vec{E} = \frac{p}{4\pi\epsilon_0 a^3} \hat{\theta}, \quad \theta = \frac{p}{4\pi\epsilon_0 a^3} (-\hat{z})$$

$$\therefore \vec{F} = q \vec{E} = -\frac{pq}{4\pi\epsilon_0 a^3} (\hat{z})$$

(b) $r=a$, $\theta=0$

$$\vec{E} = \frac{p}{4\pi\epsilon_0 a^3} (2\hat{r}) = \frac{2p}{4\pi\epsilon_0 a^3} \hat{z}$$

$$\therefore \vec{F} = q \vec{E} = \frac{2pq}{4\pi\epsilon_0 a^3} \hat{z}$$

$$\begin{aligned}
 (c) W &= qV [V_{(0,0,0)} - V_{(a,0,0)}] \\
 &= \frac{qV}{4\pi\epsilon_0 a^3} (\cos 0 - \cos \pi/2) \\
 &= \frac{\rho qV}{4\pi\epsilon_0 a^2}
 \end{aligned}$$

Chapter 4 Questions

2. According to Quantum Mechanics, the electron cloud for a hydrogen atom in ground state has a charge density

$$\rho(r) = \frac{q}{\pi a^3} e^{-2r/a}$$

where q is the charge of electron and a is Bohr's radius. Find the atomic polarizability of such an atom.

Let us first find the field at radius r using Gauss Law:

$$\int \mathbf{E} \cdot d\mathbf{a} = \frac{Q_{enc}}{\epsilon_0}$$

$$E = \frac{Q_{enc}}{4\pi\epsilon_0 r^2}$$

$$Q_{\text{enc}} = \int_0^r p \, dr$$

$$= \frac{4\pi q}{\pi a^3} \int_0^r e^{-2r/a} \bar{r}^2 \, dr$$

$$= \frac{4q}{a^3} \left[-\frac{a}{2} e^{-2r/a} \left(\bar{r}^2 + a\bar{r} + \frac{a^2}{2} \right) \right] \Big|_0^r$$

$$Q_{\text{enc}} = -\frac{2q}{a^2} \left[e^{-2r/a} \left(r^2 + ar + \frac{a^2}{2} \right) - \frac{a^2}{2} \right]$$

$$Q_{\text{enc}} = qr \left[1 - e^{-2r/a} \left(1 + \frac{2r}{a} + \frac{2r^2}{a^2} \right) \right]$$

So the field of electron cloud is

$$E_e = \frac{1}{4\pi\epsilon_0} \frac{qr}{r^2} \left[1 - e^{-2r/a} \left(1 + \frac{2r}{a} + \frac{2r^2}{a^2} \right) \right]$$

The proton will be shifted from $r=0$ to point d where $E_e = E$.

$$E = \frac{1}{4\pi\epsilon_0} \frac{qr}{d^2} \left[1 - e^{-2d/a} \left(1 + \frac{2d}{a} + \frac{2d^2}{a^2} \right) \right]$$

Expanding in powers of α/α

$$e^{-2d/a} = 1 - \left(\frac{2d}{a}\right) + \frac{1}{2} \left(\frac{2d}{a}\right)^2 - \frac{1}{3!} \left(\frac{2d}{a}\right)^3 \dots$$

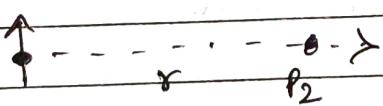
$$1 - e^{-2d/a} \left(1 + \frac{2d}{a} + \frac{2d^2}{a^2}\right) = 1 - \left(1 - \frac{2d}{a} + 2\left(\frac{d}{a}\right)^2 + \dots\right) \left(1 + \frac{2d}{a} + \frac{2d^2}{a^2}\right)$$

$$= \frac{4}{3} \left(\frac{d}{a}\right)^3 + \text{higher order terms}$$

$$\therefore E = \frac{q}{4\pi\epsilon_0 d^2} \left(\frac{4}{3} \frac{d^3}{a^3}\right)$$

$$\therefore \boxed{\alpha = 3\pi\epsilon_0 a^3}$$

5. In fig. p_1 and p_2 are (perfect) dipoles at distance r apart. What is the torque on p_1 due to p_2 ? What is the torque on p_2 due to p_1 ?



$$\text{Field of } p_1 \text{ at } p_2 (\theta = \pi/2) : E_1 = \frac{p_1}{4\pi\epsilon_0 r^3} \hat{\theta}$$

$$\text{Torque on } p_2 : N_2 = p_2 \times E_1$$

$$= p_2 E_1 \sin 90^\circ = p_2 E_1$$

$$= \frac{p_1 p_2}{4\pi\epsilon_0 r^3} \left(\begin{array}{l} \text{(points inside the} \\ \text{page)} \end{array}\right)$$

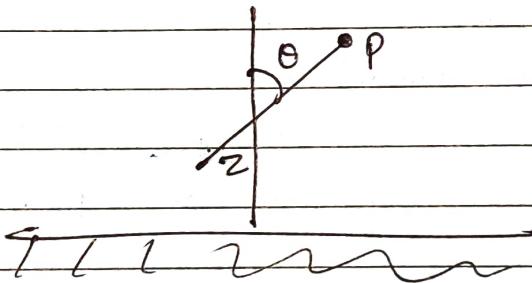
Field of p_2 at r_1 ($\theta = \pi$): $E_2 = \frac{p_2}{4\pi\epsilon_0 r^3} (-2\hat{r})$

(points to the right)

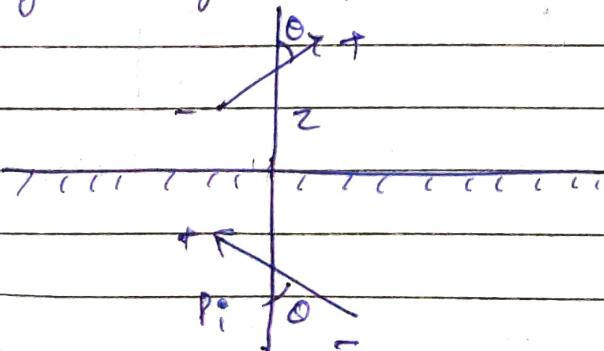
Torque on p_1 : $N_1 = p_1 \times E_2 = \frac{2p_1 p_2}{4\pi\epsilon_0 r^3}$

(points into the page)

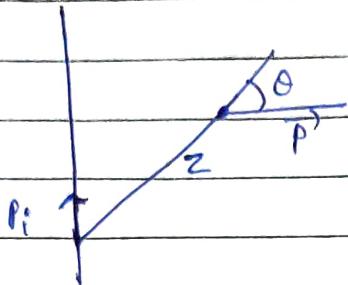
6. A perfect dipole p is situated at a distance r above an infinite grounded conducting plane. The dipole makes an angle θ with the perpendicular to the plane. Find the torque on p . If the dipole is free to rotate. In what direction will it come to rest?



Using image dipole as shown



redrawing placing P_i at origin



$$E_i = \frac{p}{4\pi\epsilon_0(2z)^3} (2\cos\theta \hat{r} + \sin\theta \hat{\theta})$$

$$p = p\cos\theta \hat{r} + p\sin\theta \hat{\theta}$$

$$N = p \times E_i$$

$$= \frac{p^2}{4\pi\epsilon_0(2z)^3} [(\cos\theta \hat{r} + \sin\theta \hat{\theta}) \times (2\cos\theta \hat{r} + \sin\theta \hat{\theta})]$$

$$= \frac{p^2}{4\pi\epsilon_0(2z)^3} [\cos\theta \sin\theta \hat{\phi} + 2\sin\theta \cos\theta (-\hat{\phi})]$$

$$N = \frac{p^2 \sin\theta \cos\theta}{4\pi\epsilon_0(2z)^3} (-\hat{\phi}) \quad (\text{out of the page})$$

$$\text{But } \sin\theta \cos\theta = \frac{\sin 2\theta}{2},$$

$$\text{so } N = \frac{p^2 \sin 2\theta}{4\pi\epsilon_0(2z)^3}$$

For $0 < \theta < \pi/2$, N tends to rotate p counterclockwise;
 for $\pi/2 < \theta < \pi$, N rotates p clockwise.

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Thus the stable configuration is perpendicular to surface, either in or out.