

PH110: Waves and Electromagnetics

Tutorial 7

Q.9 A dipole p is at a distance r from a point charge q , and oriented so that p makes an angle θ with a vector r from q to p .

- (a) What is the force on p ?
- (b) What is the force on q ?

(a) $F = (p \cdot \nabla) E$

$$E = \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \hat{r} = \frac{q}{4\pi\epsilon_0} \frac{x\hat{x} + y\hat{y} + z\hat{z}}{(x^2 + y^2 + z^2)^{3/2}}$$

$$F_x = \left(p_x \frac{\partial}{\partial x} + p_y \frac{\partial}{\partial y} + p_z \frac{\partial}{\partial z} \right) \frac{q}{4\pi\epsilon_0} \frac{x}{(x^2 + y^2 + z^2)^{3/2}}$$

$$F_x = \frac{q}{4\pi\epsilon_0} \left\{ p_x \left[\frac{1}{(x^2 + y^2 + z^2)^{3/2}} - \frac{3x \cdot 2x}{2} \frac{1}{(x^2 + y^2 + z^2)^{5/2}} \right] \right\}$$

$$+ \frac{q}{4\pi\epsilon_0} \left\{ p_y \left[-\frac{3x \cdot 2y}{2} \frac{1}{(x^2 + y^2 + z^2)^{5/2}} \right] + p_z \left[-\frac{3x \cdot 2z}{2} \frac{1}{(x^2 + y^2 + z^2)^{5/2}} \right] \right\}$$

$$F_x = \frac{q}{4\pi\epsilon_0} \left[\frac{p_z}{r^3} - \frac{3x}{r^5} (p_x x + p_y y + p_z z) \right]$$

$$F_x = \frac{q}{4\pi\epsilon_0} \left[\frac{p}{r^3} - \frac{3r(p \cdot r)}{r^5} \right]_x$$

$$\therefore F = \frac{q}{4\pi\epsilon_0 r^3} [p - 3(p \cdot \hat{r})\hat{r}]$$

$$(b) E = \frac{1}{4\pi\epsilon_0} \frac{1}{r^3} \left\{ 3(p(-\hat{r}))(-\hat{r}) - p \right\}$$

$$= \frac{1}{4\pi\epsilon_0} \frac{1}{r^3} [3(p \cdot \hat{r})\hat{r} - p]$$

The sign indicates that \mathbf{r} points towards S.

$$\therefore F = qE$$

$$\therefore F = \frac{q}{4\pi\epsilon_0 r^3} [3(p \cdot \hat{r})\hat{r} - p]$$

The forces are equal and opposite, as expected from Newton's Third Law.

Q.11 A short cylinder of radius a and length L , carries a 'frozen in' uniform polarization P , parallel to its axis. Find the bound charge and sketch the electric field.

(i) for $L \gg a$

(ii) for $L \ll a$

(iii) for $L \approx a$

$$\rho_b = 0$$

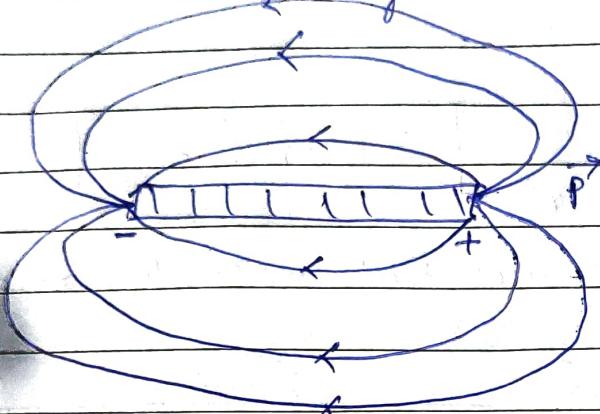
$$\sigma_b = \pm P$$

(plus sign taken at one end -

the one P points towards : minus sign at the other end - the one P points away from).

(i) $L \gg a$

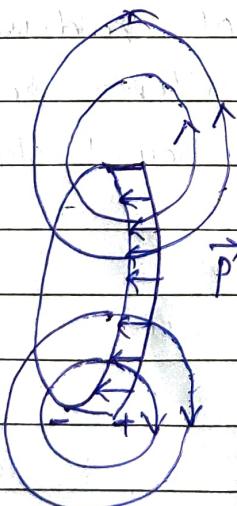
Then the end looks like point charges, and the whole thing is like a physical dipole, of length L and charge $P\pi a^2$



Like a dipole

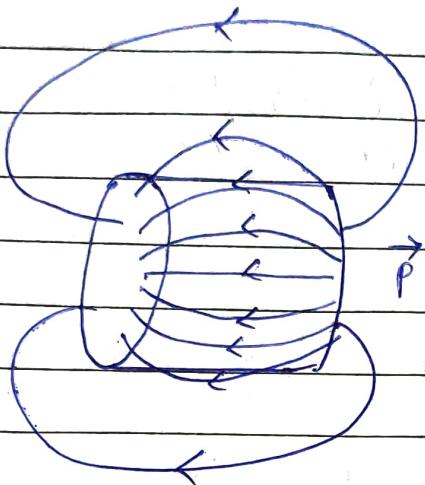
(ii) $L \ll a$

Then it is like a circular plate capacitor: field is nearly uniform inside; nonuniform "fringing" field at the edges



Like a parallel plate capacitor

(iii) $L \approx a$



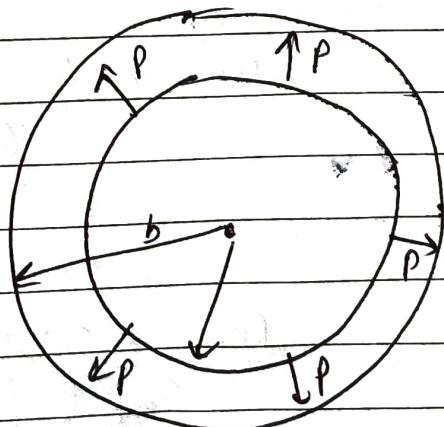
Q.15 A thick spherical shell (inner radius a , outer radius b) is made of dielectric material with a 'frozen in' polarisation

$$P(r) = \frac{k}{r} \hat{r}$$

where k is a constant and r is the distance from centre (there is no free charge in the problem).

Find the electric field in all different regions by two methods:

- (a) Locate all the bound charge, and use Gauss's law to calculate the field it produces



- (b) use eq 4.23 to find D , and then get E from eq. 4.21



$$P_b = -\nabla \cdot P$$

$$P_b = -\frac{1}{r^2} \frac{d}{dr} \left(r^2 \frac{k}{r} \right)$$

$$P_b = -\frac{k}{r^2}$$

$$\therefore \sigma_b = P \cdot \hat{n}$$

$$= \begin{cases} + P \cdot \hat{r} = k/b & (\text{at } br=b) \\ - P \cdot \hat{r} = -k/a & (\text{at } r=a) \end{cases}$$

Gauss's Law \Rightarrow

$$E = \frac{1}{4\pi\epsilon_0} \frac{Q_{enc}}{r^2} \hat{r}$$

For $r < a$

$Q_{enc} = 0$ (No charge enclosed)

$E = 0$ (Field is zero)

For $r > b$,

$$Q_{enc} = 0$$

$$\therefore E = 0$$

For $a < r < b$

$$Q_{enc} = \left(-\frac{k}{a}\right)(4\pi a^2) + \int_a^r \left(-\frac{k}{r^2}\right) 4\pi r^2 dr$$

$$Q_{enc} = -4\pi ka - 4\pi k(r-a)$$

$$Q_{enc} = -4\pi kr$$

$$\therefore \mathbf{E} = -\left(\frac{\kappa}{\epsilon_0 r}\right) \hat{r}$$

$$(b) \oint \mathbf{D} \cdot d\mathbf{a} = Q_{\text{enc}} = 0$$

$$\mathbf{D} = 0 \text{ everywhere}$$

$$\mathbf{D} = \epsilon_0 \mathbf{E} + \mathbf{P} = 0$$

$$\therefore \mathbf{E} = \left(-\frac{1}{\epsilon_0}\right) \mathbf{P}$$

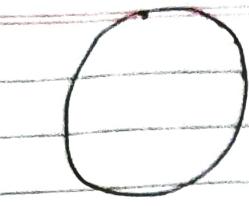
Q.16 Suppose the field inside a large piece of dielectric is ϵ_0 , so that the electric displacement is $\mathbf{D}_0 = \epsilon_0 \mathbf{E}_0 + \mathbf{P}$

(a) Now a small spherical cavity (Fig. 4.19a) is hollowed out of the material. Find the field at the centre of the cavity in terms of ϵ_0 and \mathbf{P} . Also, find the displacement at the centre of the cavity in terms of \mathbf{D}_0 and \mathbf{P} . Assume the polarisation is 'frozen in' so it doesn't change when the cavity is excavated.

(b) Do the same for a long needle-shaped cavity running parallel to \mathbf{P} (Fig. 4.19 b)

(c) Do the same for a thin wafer-shaped cavity perpendicular to \mathbf{P} (Fig. 4.19 c)

Assume that the cavities are small enough that \mathbf{P} , ϵ_0 and \mathbf{D}_0 are essentially uniform.



4.19 (a)



4.19(b)



4.19(c)

- (a) It is same as E_0 minus the field at the centre of a sphere with uniform polarization P which is equal to $-P/3\epsilon_0$.

$$\therefore E = E_0 + \frac{P}{3\epsilon_0}$$

$$\therefore D = \epsilon_0 E$$

$$D = \epsilon_0 E_0 + \frac{P}{3}$$

$$\therefore D = D_0 - \frac{2P}{3}$$

- (b) Same as E_0 minus the field of \pm charges at the two ends of the needle - but these are small and far away, so

$$E = E_0$$

$$D = \epsilon_0 E = \epsilon_0 E_0$$

$$D = D_0 - P$$

- (c) Same as E minus the field of a parallel plate capacitor with upper plate at $\sigma = P$, which is equal to $-\frac{P}{\epsilon_0}$.

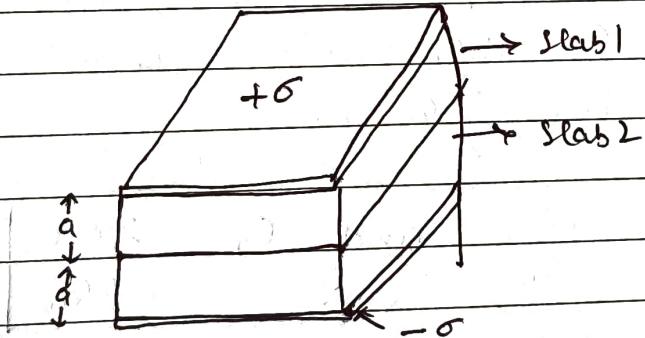
$$\therefore E = \epsilon_0 + \frac{P}{\epsilon_0}$$

$$\therefore D = \epsilon_0 E \\ = \epsilon_0 \epsilon_0 + P$$

$$\text{So, } D = D_0$$

Q.18 The space between the plates of a parallel-plate capacitor (Fig 4-24) is filled with two slabs of linear dielectric material. Each slab has thickness a , so the total distance between the plates is $2a$.

Slab 1 has dielectric constant of 2 and slab 2 has $k = 1.5$. The free charge density on the top plate is σ and on the bottom plate is $-\sigma$.



- (a) Find the electric displacement D in each slab.
- (b) Find the electric field E in each slab.
- (c) Find the polarisation P in each slab.
- (d) Find the potential difference between plates.
- (e) Find the location and amount of all bound charge.
- (f) Now that you know all the charge, recalculate the field in each slab, and confirm your answer to (b).



Applying $\int D \cdot da = \sigma_{\text{face}}$ to the Gaussian surface

$$DA = \sigma A$$

$$\therefore D = \sigma \quad (\text{as } D = 0 \text{ inside the metal plate})$$

This is true in both the slabs: D points down.

$$(b) D = \sigma E$$

$$E_1 = \frac{\sigma}{\epsilon_1} \quad (\text{in slab 1})$$

$$E_2 = \frac{\sigma}{\epsilon_2} \quad (\text{in slab 2})$$

$$\text{But, } \epsilon = \epsilon_0 \epsilon_r$$

$$\text{so, } \epsilon_1 = 2\epsilon_0$$

$$\epsilon_2 = \frac{3}{2}\epsilon_0$$

$$\text{So, } E_1 = \frac{\sigma}{2\epsilon_0} \quad \text{and} \quad E_2 = \frac{2\sigma}{3\epsilon_0}$$

$$(c) P = \epsilon_0 \chi_e E, \text{ so}$$

$$P = \frac{\epsilon_0 \chi_e d}{\epsilon_0 \epsilon_r} = \left(\frac{\chi_e}{\epsilon_r} \right) \sigma$$

$$\therefore \chi_e = \epsilon_r - 1$$

$$\therefore P = \left(1 - \frac{1}{\epsilon_r} \right) \sigma$$

$$\therefore P_1 = \frac{\sigma}{2} \quad \text{and} \quad P_2 = \frac{\sigma}{3}$$

(d) $V = E_1 a + E_2 a$

$$V = \frac{\sigma a}{2\epsilon_0} + \frac{2\sigma a}{3\epsilon_0}$$

$$V = \frac{7\sigma a}{6\epsilon_0}$$

(e) $P_b = 0$

$$\sigma_b = \begin{cases} +P_1 & \text{at bottom of slab 1} = \sigma/2 \\ -P_1 & \text{at top of slab 1} = -\sigma/2 \end{cases}$$

$$\sigma_b = \begin{cases} +P_2 & \text{at bottom of slab 2} = \sigma/3 \\ -P_2 & \text{at top of slab 2} = -\sigma/3 \end{cases}$$

(f) In slab 1,

$$\text{total surface charge above} = \sigma - \frac{\sigma}{2} = \frac{\sigma}{2}$$

$$\begin{aligned} \text{total surface charge below} &= \frac{\sigma}{2} - \frac{\sigma}{3} + \frac{\sigma}{3} - \sigma \\ &= -\frac{\sigma}{2} \end{aligned}$$

$$\therefore E_1 = \frac{\sigma}{2\epsilon_0}$$

In slab 2,

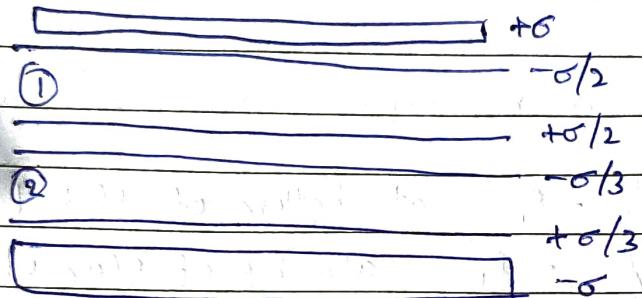
$$\text{total surface charge above} = \sigma - \frac{\sigma}{2} + \frac{\sigma}{2} - \frac{\sigma}{3} = \frac{2\sigma}{3}$$

$$\text{total surface charge below} = \frac{\sigma}{2} - \sigma = -\frac{2\sigma}{3}$$

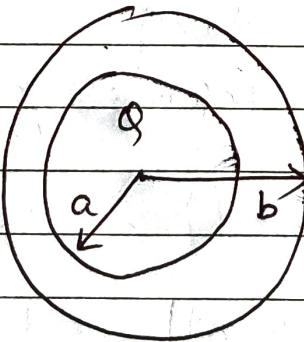
L:

$E_2 = \frac{2\sigma}{3\epsilon_0}$

E_1 and E_2 match with the result of (b)



Q.26 A spherical conductor, of radius a , carries a charge Q . It is surrounded by linear dielectric material of susceptibility χ_e , out of the radius b . Find the energy of this configuration.



$$D = \begin{cases} 0 & , r < a \\ \frac{Q}{4\pi r^2} \hat{r} & , r > a \end{cases}$$

$$E = \begin{cases} 0 & , r < a \\ \frac{Q}{4\pi\epsilon_0 r^2} \hat{r} & , a \leq r < b \\ \frac{Q}{4\pi\epsilon_0 r^2} \hat{r} & , r > b \end{cases}$$

$$W = \frac{1}{2} \int D \cdot E dV$$

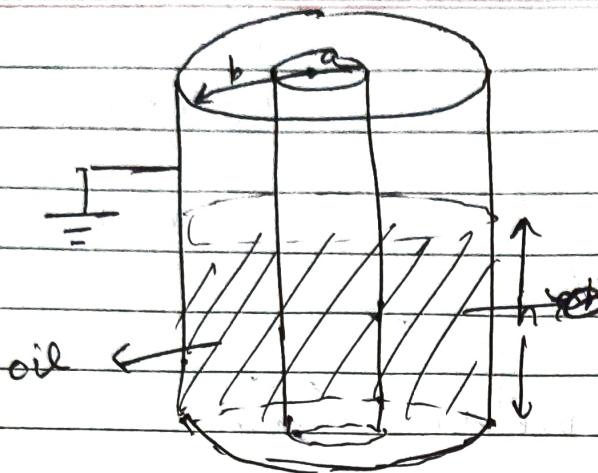
$$W = \frac{1}{2} \frac{Q^2 \cdot 4\pi}{(4\pi)^2} \left\{ \frac{1}{\epsilon} \int_a^b \frac{1}{r^2} \frac{1}{r^2} r^2 dr + \frac{1}{\epsilon_0} \int_b^\infty \frac{1}{r^2} dr \right\}$$

$$= \frac{Q^2}{8\pi} \left\{ \frac{1}{\epsilon} \left(-\frac{1}{r} \right) \Big|_a^b + \frac{1}{\epsilon_0} \left(-\frac{1}{r} \right) \Big|_b^\infty \right\}$$

$$= \frac{Q^2}{8\pi\epsilon_0} \left\{ (1 + \chi_e) \left(\frac{1}{a} - \frac{1}{b} \right) + \frac{1}{b} \right\}$$

$$\boxed{2. W = \frac{Q^2}{8\pi\epsilon_0 (1 + \chi_e)} \left(\frac{1}{a} + \frac{\chi_e}{b} \right)}$$

Q.28 . Two long co-axial cylindrical metal tubes (inner radius a , outer radius b) stand vertically in a tank of dielectric oil (susceptibility χ_e , mass density ρ). The inner one is maintained at potential V , and the outer one is grounded. To what height h does the oil rise in the space between the tubes?



Let us find the capacitance as a function of h .

$$\text{air part: } E = \frac{2\lambda}{4\pi\epsilon_0 S} \quad \left. \begin{array}{l} \\ V = \frac{2\lambda}{4\pi\epsilon_0} \ln\left(\frac{b}{a}\right) \end{array} \right\} \quad \frac{\lambda}{\epsilon_0} = \frac{\lambda'}{\epsilon}$$

$$\text{oil part: } D = \frac{2\lambda'}{4\pi\epsilon} \quad \left. \begin{array}{l} E = \frac{2\lambda'}{4\pi\epsilon_0 S} \\ V = \frac{2\lambda'}{4\pi\epsilon} \ln\left(\frac{b}{a}\right) \end{array} \right\} \quad \lambda' = \frac{\epsilon}{\epsilon_0} \lambda \quad \lambda' = \epsilon_r \lambda$$

$$V = \frac{2\lambda'}{4\pi\epsilon} \ln\left(\frac{b}{a}\right)$$

$$\therefore Q = \lambda' h + \lambda(l-h)$$

$$\therefore Q = \epsilon_r \lambda h - \lambda h + \lambda l$$

$$Q = \lambda [(\epsilon_r - 1)h + l]$$

$$Q = \lambda [x_e h + l] \quad ; \quad (l \text{ is total height})$$

$$c = \frac{Q}{V} = \frac{\lambda (x_e h + l)}{2\pi \ln(b/a)} (4\pi\epsilon_0)$$

$$\therefore c = \frac{2\pi\epsilon_0 (x_e h + l)}{\ln(b/a)}$$

Net upward force,

$$F = \frac{1}{2} V^2 \frac{dc}{dh} = \frac{1}{2} V^2 \frac{2\pi\epsilon_0 x_e}{\ln(b/a)}$$

The gravitational force,

$$F = mg$$

$$F = \rho\pi (b^2 - a^2) gh$$

$$\therefore F_{\text{upward}} = F_{\text{downward}}$$

$$\frac{V^2 \pi \epsilon_0 x_e}{\ln(b/a)} = \rho\pi (b^2 - a^2) gh$$

$$h = \boxed{\frac{V^2 \epsilon_0 x_e}{\rho (b^2 - a^2) g \ln(b/a)}}$$