

Computational study of 2D Ising Model

Nishi Prabhat Hazarika

October 23, 2025

Objective

- ① To study how energy per spin, magnetization per spin, specific heat, and susceptibility change as a function of temperature.
- ② To study the effects of topology on the lattice behavior.
- ③ Extracting critical exponents via finite-size scaling and estimation of critical temperature.

The Ising Model – A Fundamental Framework in Statistical Physics

What Is It?

A mathematical model to study **magnetism** and **phase transitions** in materials.

- Spins $s_i = \pm 1$ on a lattice.
- Hamiltonian:

$$H = -J \sum_{\langle i,j \rangle} s_i s_j - h \sum_i s_i$$

- J : Interaction strength between neighbouring spins
- h : External magnetic field

Behavior

- Ferromagnetic ($J > 0$): Spins align
- Antiferromagnetic ($J < 0$): Spins anti-align
- **Phase Transition** in 2D
 - Above T_c : Disordered (paramagnetic)
 - Below T_c : Ordered (ferromagnetic)
- We take a 2D lattice of size $L \times L$ with random spins (± 1).

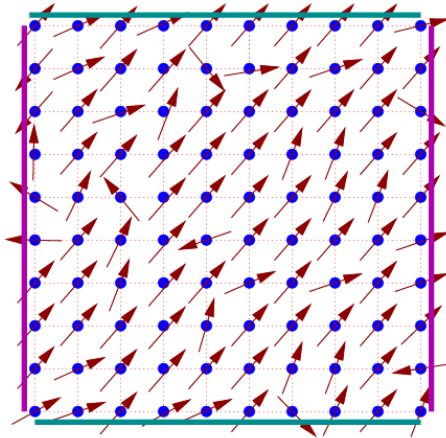


Figure: 2D square lattice with random spin orientation.

Thermodynamic quantities

The Hamiltonian for 2D ising model in absence of external magnetic field is

$$H = -J \sum_{\langle i,j \rangle} s_i s_j$$

① Energy :

$$E = -J \sum_{\langle i,j \rangle} s_i s_j$$

② Magnetization :

$$M = \frac{1}{N} \sum_i s_i$$

③ Specific heat:

$$C_v = \frac{1}{NT^2} (\langle E^2 \rangle - \langle E \rangle^2)$$

① Magnetic susceptibility :

$$\chi = \frac{1}{NT}(\langle M^2 \rangle - \langle M \rangle^2)$$

Metropolis Algorithm

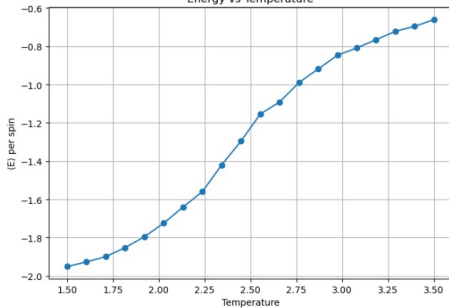
- 1 Pick a random spin: Choose the lattice site (i, j) at random.
- 2 Calculate the energy change ΔE if the spin is flipped:

$$\Delta E = E_{\text{new}} - E_{\text{old}} = -J(-s_{ij}) \sum_{\text{nn}} s_{\text{nn}} - \left[-J s_{ij} \sum_{\text{nn}} s_{\text{nn}} \right]$$

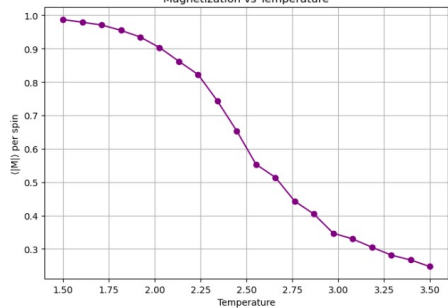
$$\Delta E = 2J s_{i,j} (s_{i+1,j} + s_{i-1,j} + s_{i,j+1} + s_{i,j-1})$$

- 3 Step 3: Decide whether to flip or not flip
 - If $\Delta E \leq 0$:
Flip the spin
 - If $\Delta E > 0$:
Flip with the probability $P = \exp(\frac{-\Delta E}{KT})$
- 4 Step 4: Repeat this for the desired Monte Carlo steps.

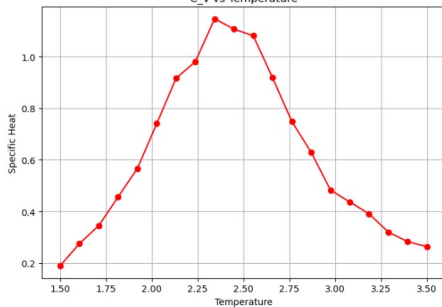
Energy vs Temperature



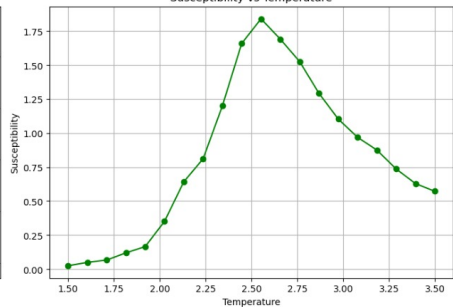
Magnetization vs Temperature



C_v vs Temperature



Susceptibility vs Temperature



Effects of topology on the lattice behaviour I

① Periodic boundary conditions:

- The lattice "wraps around" like a torus.
- The left edge is connected to the right edge, and the top is connected to the bottom.
- Mimics an infinite system by eliminating boundary effects.
- Every site has 4 neighbours, no matter its position.
- Most commonly used in simulations to approximate bulk behaviour.

② Open boundary conditions:

- The lattice edges are not connected: If a neighbour is outside the grid, it is ignored (or treated as zero).
- Models a finite, isolated patch of spins.

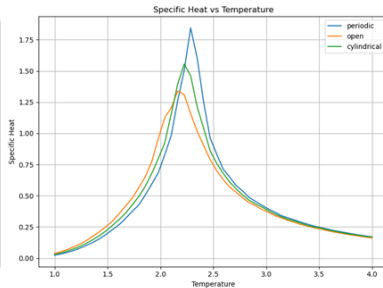
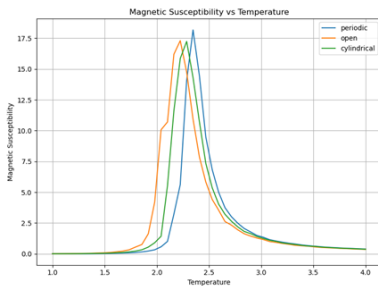
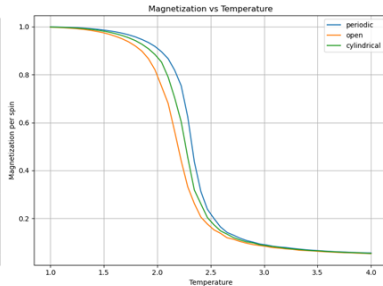
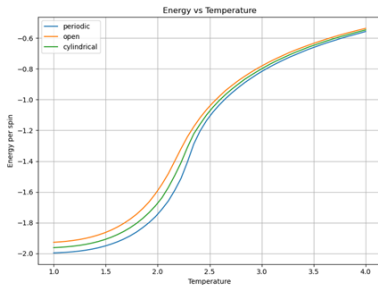
Effects of topology on the lattice behaviour II

- Edge spins have fewer neighbours, making alignment less favourable at the boundaries.
- Leads to edge effects: surface spins behave differently than interior ones.

③ Cylindrical boundary conditions:

- Like rolling the lattice into a cylinder: Wrap in one direction (usually vertical) and leave the other direction open (horizontal).
- Vertical neighbours use periodic wrapping, and horizontal neighbours are treated as open (only valid if they fall within bounds).
- Edge effects are intermediate between open and periodic.

Plots



Behaviour at different boundaries

Magnetization behaviour at Boundary:

- 1 Periodic: Minimizes edge effects, producing the clearest and sharpest phase transition.
- 2 More spin disorder near boundaries reduces net magnetization slightly and can smooth the transition.
- 3 Cylindrical: Behaves between open and periodic. Some boundary symmetry is preserved.

Behaviour at different boundaries

Energy behaviour at Boundary:

- 1 Periodic: Yield slightly lower ground-state energy due to maximized spin interactions.
- 2 Open: Edge effects slightly alter energy levels due to fewer neighbours.
- 3 Cylindrical: Intermediate behaviour between open and periodic boundary conditions.

Behaviour at different boundaries

Specific heat behaviour at Boundary:

- 1 Periodic: Exhibits the sharpest and highest peak, reflecting the most accurate bulk behaviour.
- 2 Open: Broader and lower peak due to increased surface effects and reduced interactions at edges.
- 3 Cylindrical: Intermediate behaviour, with one direction preserving periodicity.

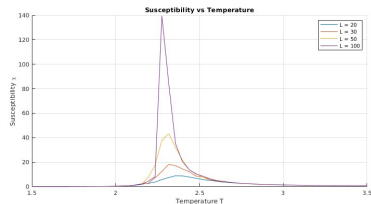
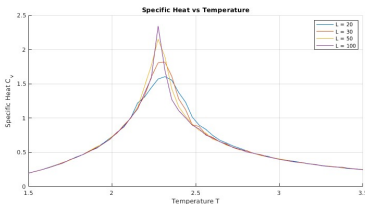
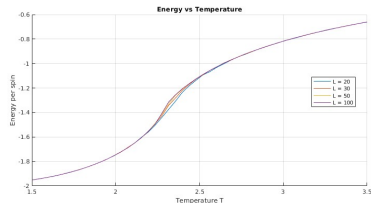
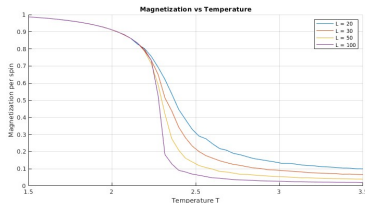
Magnetic susceptibility behaviour at Boundary:

- 1 Periodic: Produces a sharp, high peak—indicating strong magnetic response near the transition.
- 2 Open: Reduced overall susceptibility, and a smoother peak, due to less cooperative spin alignment at the edges.
- 3 Cylindrical: Intermediate behaviour between open and periodic cases.

Finite Size Scaling

- In theory, the phase transition occurs only in infinite systems, i.e. $L \rightarrow \infty$.
- We simulate Ising model on a finite lattice of size $L \times L$, so the transitions are smoothened out.
- **Finite Size Scaling** is a technique to extract infinite system behaviour from finite system simulations.

Plots of physical quantities at different L



The estimated critical temperature is $T_c = 2.3163$

Power Law Scaling and Critical Exponents

- $\zeta \sim |T - T_c|^{-\nu}$
- $M \sim |T - T_c|^\beta$
- $\chi \sim |T - T_c|^{-\gamma}$
- $C_v \sim |T - T_c|^{-\alpha}$

where, ν , β , γ , α are the critical exponents.

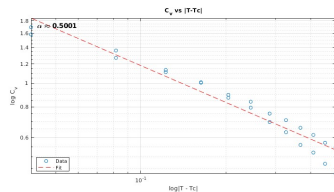
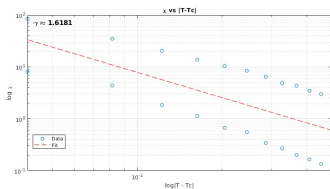
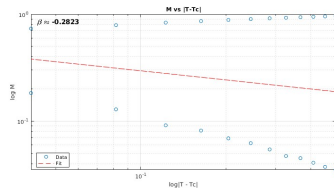
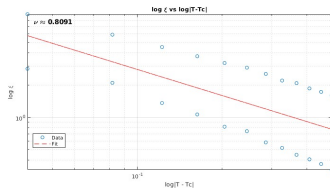
Critical exponents describe how physical quantities diverge or vanish as a system approaches a critical point.

At the critical point $T = T_c$, the system becomes scale invariant because the correlation length ζ , becomes infinite at T_c .

That means:

- No single length scale dominates
- Small and large clusters behave similarly

log-log plots to find critical exponents



Theoretical values: $\nu = 1$, $\beta = 0.125$, $\gamma = 1.75$, $\alpha = 0$

Here, there are two points for one $\log|T - T_c|$ value. This is because the value of T is approaching the critical temperature from both the sides.

Limitations of Power Law Fits in Simulations

- Power law behaviour is strictly valid exactly at T_c .
- Simulation is done in the region near T_c .
- Near T_c , correction to scaling terms are included.
- Singularity isn't pure power law scaling, but include logarithmic corrections.

How Finite-Size Scaling Overcomes Power-Law Limitations

① Exploiting Size Dependence

② Scaling Form:

$$Q(L, T) = L^p f(x), \quad (1)$$

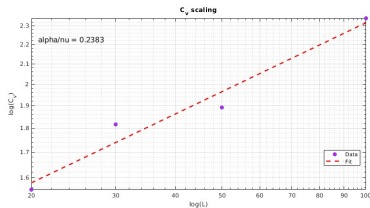
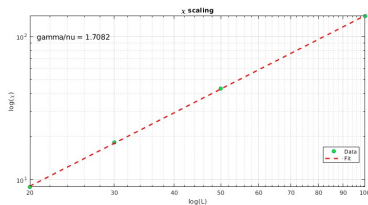
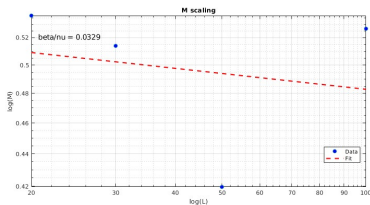
where, $Q(T, L)$ is the physical quantity of interest, p is a scaling dimension related to the critical exponent of Q , $f(x)$ is the scaling function, where the scaling variable $x = L^{\frac{1}{\nu}} |T - T_c|$.

③ Extracting Exponents from Scaling

The FSS hypothesis at $T = T_c$ states:

- ① Magnetization: $M \sim L^{-\beta/\nu}$
- ② Susceptibility: $\chi \sim L^{\gamma/\nu}$
- ③ Specific Heat: $C_v \sim L^{\alpha/\nu}$

Visualizing Finite-Size Scaling



Theoretical Critical Exponents (2D Ising):

$$\beta\nu = 1/8 = 0.125$$

$$\gamma\nu = 7/4 = 1.75$$

$$\alpha\nu = 0$$

Reference: Onsager solution

Conclusion

- Fitting simulation data using simple power laws near the critical temperature T_c can provide a rough estimate of critical behaviour.
- However, such fits are prone to inaccuracies due to:
 - Finite-size effects,
 - Corrections to scaling near T_c ,
 - Logarithmic or non-power-law deviations from ideal scaling.
- Finite-Size Scaling offers a more robust and accurate framework by:
 - Systematically analyzing dependence of observables on system size L ,
 - Exploiting universal scaling relations predicted by the theory of critical phenomena,
 - Allowing precise extraction of critical exponents.
- Therefore, FSS is essential for the reliable characterization of critical behaviour in finite systems.