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## PHYS422 HW.1

# 28/1/2023

#### Problem 1): $\psi_1 = A e^{i k_1 x} + b e^{-i k_1 x}$ :

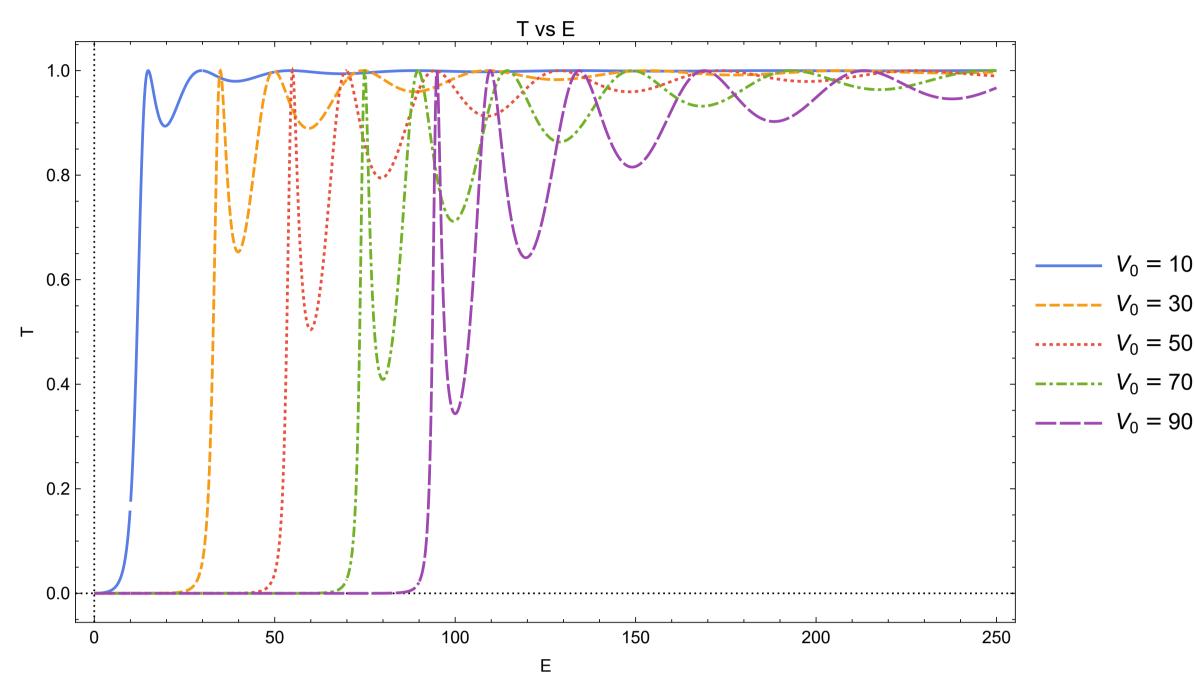
$$\psi_{2} = d e^{-i k_{2} x} + c e^{i k_{2} x};$$

$$\psi_{3} = f e^{i k_{3} x};$$

$$fSol = f /. Solve[{ \psi_{1} = \psi_{2} /. x \to 0, D[\psi_{1}, x] = D[\psi_{2}, x] /. x \to 0, \psi_{3} = \psi_{2} /. x \to a, D[\psi_{3}, x] = D[\psi_{2}, x] /. x \to a}, {b, c, d, f}][[1];$$

$$T = \frac{1}{A^2} \; \text{FullSimplify@} \left( \text{ComplexExpand@Abs[fSol]} \; /. \; k_1 \rightarrow k_3 \; /. \; k_2 \rightarrow \; \sqrt{2 \, m} \; \frac{(E - V_0)}{\hbar^2} \; /. \; k_3 \rightarrow \; \sqrt{2 \, m} \; \frac{E}{\hbar^2} \; \right)^2 \; /. \; a \rightarrow 1 \; /. \; \hbar \rightarrow 1 \; /. \; m \rightarrow 1$$
 
$$\frac{4 \; \text{E} \; (E - V_0)}{4 \; \text{E}^2 - 4 \; \text{E} \; V_0 \; + \; \text{Sin} \left[ \sqrt{2} \; \sqrt{E - V_0} \; \right]^2 \; V_0^2}$$

plt = Table [T,  $\{V_0, 10, 100, 20\}$ ]; Plot[plt, {E, 0, 250}, PlotRange → All, PlotTheme → {"Scientific", "BoldColor", "Monochrome"}, FrameLabel → {"E", "T"}, PlotLegends  $\rightarrow$  Table[TraditionalForm[ $V_0 = t$ ], {t, 10, 100, 20}], ImageSize  $\rightarrow$  Large, PlotLabel  $\rightarrow$  "T vs E"]



We can notice that the Transmission coefficient is almost zero when  $E < V_0$ . Moreover, when  $E > V_0$ , the transmission is not always 1, it oscillates closer to 1 as E gets bigger.

#### Problem 3): We first solve Schrodinger's equation for two regions, the first is when 0<x<a. The second is when x>a:

 $\psi 1 = \left( \psi[x] / . \text{ Assuming} \left[ m > 0 \&\& E > 0 \&\& V_0 > 0, \text{ DSolve} \left[ \left\{ \frac{-\hbar^2}{2 m} D[\psi[x], \{x, 2\}] - V_0 \psi[x] = E \psi[x], \psi[0] = 0 \right\}, \psi[x], x \right] \right] \right) [1] / . C[1] \rightarrow A;$ 

Assuming 
$$\left[m > 0 \&\& E > 0 \&\& V_0 > 0 \&\& a > 0$$
, FullSimplify  $\left[\psi 1\right] / \frac{\sqrt{2 (E + V_0) m}}{\hbar} \rightarrow \alpha\right] / / TraditionalForm$ 

 $-2 i A \sin(\alpha x)$  $\psi 2 = \left[ \psi[x] /. DSolve \left[ \frac{-\hbar^2}{2m} D[\psi[x], \{x, 2\}] = E \psi[x], \psi[x], x \right] \right] [[1]];$ 

Assuming 
$$\left[\mathbf{m} > 0 \&\& \mathbf{E} > 0 \&\& \mathbf{V}_0 > 0 \&\& \mathbf{a} > 0$$
, FullSimplify  $\left[\psi 2\right] / \frac{\sqrt{2 \mathbf{E} \mathbf{m}}}{\hbar} \rightarrow \mathbf{k}\right] / / \text{TraditionalForm}$ 

$$c_1 \cos(k x) + c_2 \sin(k x)$$

 $c_1\cos(k\,x) + c_2\sin(k\,x)$  $\alpha = \frac{\sqrt{2 (E + V_0) m}}{\pi}$ ;  $k = \frac{\sqrt{2 E m}}{\pi}$ . Now we apply the continuity of the wavefunction and its derivative at the boundary x=a. The solutions of

 $\left\{c_1 \to -2 i A \left(\sin(a \alpha) \cos(a k) - \sqrt{\frac{E + V_0}{E}} \cos(a \alpha) \sin(a k)\right), c_2 \to -2 i A \left(\sin(a \alpha) \sin(a k) + \sqrt{\frac{E + V_0}{E}} \cos(a \alpha) \cos(a k)\right)\right\}$ 

Sol = Solve[{
$$\psi$$
2 ==  $\psi$ 1, D[ $\psi$ 1, x] == D[ $\psi$ 2, x]} /. x  $\rightarrow$  a, {C[1], C[2]}];

Assuming[ $m > 0 \&\& E > 0 \&\& V_0 > 0 \&\& a > 0$ , FullSimplify[Sol] /.  $\frac{\sqrt{2(E + V_0) m}}{\hbar} \rightarrow \alpha$  /.  $\frac{\sqrt{2Em}}{\hbar} \rightarrow k$ ][1] // TraditionalForm

### Problem 7):

$$\langle x \rangle = \sqrt{\frac{h}{4\pi m \omega}} \langle 0 \mid a + a^{\dagger} \mid 0 \rangle = 0$$

the coefficients are in terms of A:

$$\langle x^{2} \rangle = \frac{h}{4 \pi m \omega} \langle \mathbf{0} \mid (a + a^{\dagger}) * (a + a^{\dagger}) \mid \mathbf{0} \rangle = \frac{h}{4 \pi m \omega} \langle \mathbf{0} \mid a a + a^{\dagger} a^{\dagger} + a a^{\dagger} + a^{\dagger} a \mid \mathbf{0} \rangle = \frac{h}{4 \pi m \omega} \langle \mathbf{0} \mid a a^{\dagger} \mid \mathbf{0} \rangle = \frac{h}{4 \pi m \omega} \langle \mathbf{0} \mid a a^{\dagger} \mid \mathbf{0} \rangle = \frac{h}{4 \pi m \omega} \langle \mathbf{0} \mid a a^{\dagger} \mid \mathbf{0} \rangle = \frac{h}{4 \pi m \omega} \langle \mathbf{0} \mid a a^{\dagger} \mid \mathbf{0} \rangle = \frac{h}{4 \pi m \omega} \langle \mathbf{0} \mid a a^{\dagger} \mid \mathbf{0} \rangle = \frac{h}{4 \pi m \omega} \langle \mathbf{0} \mid a a^{\dagger} \mid \mathbf{0} \rangle = \frac{h}{4 \pi m \omega} \langle \mathbf{0} \mid a a^{\dagger} \mid \mathbf{0} \rangle = \frac{h}{4 \pi m \omega} \langle \mathbf{0} \mid a a^{\dagger} \mid \mathbf{0} \rangle = \frac{h}{4 \pi m \omega} \langle \mathbf{0} \mid a a^{\dagger} \mid \mathbf{0} \rangle = \frac{h}{4 \pi m \omega} \langle \mathbf{0} \mid a a^{\dagger} \mid \mathbf{0} \rangle = \frac{h}{4 \pi m \omega} \langle \mathbf{0} \mid a a^{\dagger} \mid \mathbf{0} \rangle = \frac{h}{4 \pi m \omega} \langle \mathbf{0} \mid a a^{\dagger} \mid \mathbf{0} \rangle = \frac{h}{4 \pi m \omega} \langle \mathbf{0} \mid a a^{\dagger} \mid \mathbf{0} \rangle = \frac{h}{4 \pi m \omega} \langle \mathbf{0} \mid a a^{\dagger} \mid \mathbf{0} \rangle = \frac{h}{4 \pi m \omega} \langle \mathbf{0} \mid a a^{\dagger} \mid \mathbf{0} \rangle = \frac{h}{4 \pi m \omega} \langle \mathbf{0} \mid a a^{\dagger} \mid \mathbf{0} \rangle = \frac{h}{4 \pi m \omega} \langle \mathbf{0} \mid a a^{\dagger} \mid \mathbf{0} \rangle = \frac{h}{4 \pi m \omega} \langle \mathbf{0} \mid a a^{\dagger} \mid \mathbf{0} \rangle = \frac{h}{4 \pi m \omega} \langle \mathbf{0} \mid a a^{\dagger} \mid \mathbf{0} \rangle = \frac{h}{4 \pi m \omega} \langle \mathbf{0} \mid a a^{\dagger} \mid \mathbf{0} \rangle = \frac{h}{4 \pi m \omega} \langle \mathbf{0} \mid a a^{\dagger} \mid \mathbf{0} \rangle = \frac{h}{4 \pi m \omega} \langle \mathbf{0} \mid a a^{\dagger} \mid \mathbf{0} \rangle = \frac{h}{4 \pi m \omega} \langle \mathbf{0} \mid a a^{\dagger} \mid \mathbf{0} \rangle = \frac{h}{4 \pi m \omega} \langle \mathbf{0} \mid a a^{\dagger} \mid \mathbf{0} \rangle = \frac{h}{4 \pi m \omega} \langle \mathbf{0} \mid a a^{\dagger} \mid \mathbf{0} \rangle = \frac{h}{4 \pi m \omega} \langle \mathbf{0} \mid a a^{\dagger} \mid \mathbf{0} \rangle = \frac{h}{4 \pi m \omega} \langle \mathbf{0} \mid a a^{\dagger} \mid \mathbf{0} \rangle = \frac{h}{4 \pi m \omega} \langle \mathbf{0} \mid a a^{\dagger} \mid \mathbf{0} \rangle = \frac{h}{4 \pi m \omega} \langle \mathbf{0} \mid a a^{\dagger} \mid \mathbf{0} \rangle = \frac{h}{4 \pi m \omega} \langle \mathbf{0} \mid a a^{\dagger} \mid \mathbf{0} \rangle = \frac{h}{4 \pi m \omega} \langle \mathbf{0} \mid a a^{\dagger} \mid \mathbf{0} \rangle = \frac{h}{4 \pi m \omega} \langle \mathbf{0} \mid a a^{\dagger} \mid a a^{$$

$$\Delta x = \left[ \left\langle x^{2} \right\rangle - \left\langle x \right\rangle^{2} \right]^{\frac{1}{2}} = \sqrt{\frac{h}{4 \pi m \omega}}$$

 $\langle p_x \rangle = i \sqrt{\frac{h \, m \, \omega}{A \, \pi}} \, \left\langle \mathbf{0} \mid a - a^{\dagger} \mid \mathbf{0} \right\rangle = \mathbf{0}$ 

$$\left\langle p_{x}^{2}\right\rangle =-\frac{h\,m\,\omega}{4\,\pi}\,\left\langle \mathbf{0}\mid\,\left(a-a^{\dagger}\right)\,\star\,\left(a-a^{\dagger}\right)\,\mid\,\mathbf{0}\right\rangle =-\frac{h\,m\,\omega}{4\,\pi}\,\left\langle \mathbf{0}\mid\,a\,a+a^{\dagger}\,a^{\dagger}-a\,a^{\dagger}-a^{\dagger}\,a\mid\,\mathbf{0}\right\rangle =\frac{h\,m\,\omega}{4\,\pi}\,\left\langle \mathbf{0}\mid\,a\,a^{\dagger}\mid\,\mathbf{0}\right\rangle =\frac{h\,m\,\omega}{4\,\pi}$$

$$\triangle p_{x} = \left[ \left\langle p_{x}^{2} \right\rangle - \left\langle p_{x} \right\rangle^{2} \right]^{\frac{1}{2}} = \sqrt{\frac{h \, m \, \omega}{4 \, \pi}}$$

$$\triangle x \triangle p_x = \sqrt{\frac{h\,m\,\omega}{4\,\pi}}\,\,\sqrt{\frac{h}{4\,\pi\,m\,\omega}}\,= \frac{h}{4\,\pi}$$
 This is a "minimum-uncertainty" wave packet because it has the minimum  $\triangle x$  possible and the

minimum  $\triangle p_x$  possible. This restriction is described in the uncertainty principle  $\triangle x \triangle p_x \ge \frac{n}{\sqrt{\pi}}$ .

### a) "f" means /=3. Therefore $j = / \pm \frac{1}{2} \Longrightarrow j = \frac{5}{2}$ or $j = \frac{7}{2}$

For "f" state, for:

Problem 15):

b) for 
$$j = \frac{5}{2} \Longrightarrow m_j = \pm \frac{5}{2}, \pm \frac{3}{2}, \pm \frac{1}{2}$$

for 
$$j = \frac{7}{2} \Longrightarrow m_j = \pm \frac{7}{2}, \pm \frac{5}{2}, \pm \frac{3}{2}, \pm \frac{1}{2}$$

c) The total number of  $m_i$  states = 2j + 1

 $j = \frac{5}{2}$  the total number of states is:  $2 \times \frac{5}{2} + 1 = 6$ 

 $j = \frac{7}{2}$  the total number of states is:  $2 \times \frac{7}{2} + 1 = 8$ 

d) If we use  $m_l$  and  $m_s$  we will count the states for all possible values of j so the total number of states =  $2 \times (2 / + 1)$ For "f" state, the total number of states is:  $2 \times (3 \times 2 + 1) = 14$