

Solution Major Exam I

Phys. 310

Q1.

$$(a) \quad \Psi(x,0) = \frac{A}{\sqrt{12}} \phi_1(x) + \frac{1}{\sqrt{6}} \phi_2(x) + \frac{1}{\sqrt{3}} \phi_3(x) + \frac{1}{2} \phi_4(x)$$

$$\int_{-\infty}^{+\infty} \Psi^*(x,0) \Psi(x,0) dx = \frac{A^2}{12} \int \phi_1^* \phi_1 dx + \frac{1}{6} \int \phi_2^* \phi_2 dx + \frac{1}{3} \int \phi_3^* \phi_3 dx + \frac{1}{4} \int \phi_4^* \phi_4 dx$$

$$1 = \frac{A^2}{12} + \frac{1}{6} + \frac{1}{3} + \frac{1}{4} \quad \text{Since } \int \phi_m^* \phi_n dx = \delta_{m,n}$$

$$\Rightarrow \frac{A^2}{12} = \frac{1}{4} \Rightarrow A^2 = 3 \Rightarrow A = \sqrt{3}$$

$$\text{Thus } \Psi(x,0) = \frac{1}{2} \phi_1(x) + \frac{1}{\sqrt{6}} \phi_2(x) + \frac{1}{\sqrt{3}} \phi_3(x) + \frac{1}{2} \phi_4(x)$$

$$(b) \quad \Psi(x,t) = \frac{1}{2} e^{-iE_1 t/\hbar} \phi_1(x) + \frac{1}{\sqrt{6}} e^{-iE_2 t/\hbar} \phi_2(x) + \frac{1}{\sqrt{3}} e^{-iE_3 t/\hbar} \phi_3(x) + \frac{1}{2} e^{-iE_4 t/\hbar} \phi_4(x)$$

$$\Psi(x,t) = e^{-i\omega t/2} \left[e^{-i\omega t} \phi_1(x) \frac{1}{2} + \frac{1}{\sqrt{6}} e^{-2i\omega t} \phi_2(x) + \frac{1}{\sqrt{3}} e^{-3i\omega t} \phi_3(x) + \frac{1}{2} e^{-4i\omega t} \phi_4(x) \right]$$

$$\Psi(x,t) = e^{-i\frac{3}{2}\omega t} \left[\frac{1}{2} \phi_1(x) + \frac{1}{\sqrt{6}} e^{-i\omega t} \phi_2(x) + \frac{1}{\sqrt{3}} e^{-i2\omega t} \phi_3(x) + \frac{1}{2} e^{-i3\omega t} \phi_4(x) \right]$$

$$(c) \quad E = \int \Psi^*(x,t) H \Psi(x,t) dx = \int \Psi^*(x,0) H \Psi(x,0) dx$$

$$= \int \left(\frac{1}{2} \phi_1 + \frac{1}{\sqrt{6}} \phi_2 + \frac{1}{\sqrt{3}} \phi_3 + \frac{1}{2} \phi_4 \right) \left(\frac{E_1}{2} \phi_1 + \frac{E_2}{\sqrt{6}} \phi_2 + \frac{E_3}{\sqrt{3}} \phi_3 + \frac{E_4}{2} \phi_4 \right) dx$$

$$= \frac{E_1}{4} + \frac{E_2}{6} + \frac{E_3}{3} + \frac{E_4}{4} = \frac{1}{24} (6E_1 + 4E_2 + 8E_3 + 6E_4)$$

$$E = \frac{\hbar\omega}{24} \left[6\left(\frac{3}{2}\right) + 4\left(\frac{5}{2}\right) + 8\left(\frac{7}{2}\right) + 6\left(\frac{9}{2}\right) \right] = \frac{37}{12} \hbar\omega$$

Q2.

$$(a) \quad E = \int \Psi^*(x) H \Psi(x) dx = \int_{-\infty}^{+\infty} \Psi^*(x) \left(-\frac{\hbar^2}{2m} \frac{d^2}{dx^2} + V(x) \right) \Psi(x) dx$$

$$E = -\frac{\hbar^2}{2m} \int_{-\infty}^{+\infty} \Psi^*(x) \frac{d}{dx} \left(\frac{d\Psi}{dx} \right) dx + \int \Psi^*(x) V(x) \Psi(x) dx$$

$$E = -\frac{\hbar^2}{2m} \left\{ \Psi^*(x) \frac{d\Psi}{dx} \Big|_{-\infty}^{+\infty} - \int_{-\infty}^{+\infty} \left| \frac{d\Psi}{dx} \right|^2 dx \right\} + \int \Psi^*(x) V(x) \Psi(x) dx$$

$$E = \frac{\hbar^2}{2m} \int_{-\infty}^{+\infty} \left| \frac{d\Psi}{dx} \right|^2 dx + \int \Psi^*(x) V(x) \Psi(x) dx \geq V_{\min} \int \Psi^* \Psi dx = V_{\min}$$

Notice that the first quantity is > 0 .

(b)

$$[x, p] F(x) = \left(x \frac{\hbar}{i} \frac{d}{dx} - \frac{\hbar}{i} \frac{d}{dx} x \right) F(x) \quad ; \quad \forall F(x)$$

$$[x, p] F(x) = x \frac{\hbar}{i} \frac{dF}{dx} - \frac{\hbar}{i} F(x) - \frac{\hbar}{i} x \frac{dF}{dx} = -\frac{\hbar}{i} F(x) = i\hbar F(x)$$

$$\Rightarrow [x, p] = i\hbar \quad ; \quad p = \frac{\hbar}{i} \frac{d}{dx} \Rightarrow p^\nu = -\hbar^\nu \frac{d^\nu}{dx^\nu}$$

$$[x, \frac{d^\nu}{dx^\nu}] = -\frac{1}{\hbar^\nu} [x, p^\nu] = -\frac{1}{\hbar^\nu} \{ [x, p] p + p [x, p] \}$$

$$[x, \frac{d^\nu}{dx^\nu}] = -\frac{1}{\hbar} \{ i\hbar p + i\hbar p \} = -2ip/\hbar$$

$$[x^\nu, p^2] = x [x, p^2] + [x, p^2] x = 2i\hbar (x p + p x)$$

$$\text{from previous result } [x, p^\nu] = -\hbar^\nu [x, \frac{d^\nu}{dx^\nu}] = 2i\hbar p$$

$$\text{also using } [x, p] = i\hbar \Rightarrow xp = px + i\hbar \text{ or } px = xp - i\hbar$$

$$[x^\nu, p^\nu] = 2i\hbar (2xp - i\hbar) = 2i\hbar (2px + i\hbar)$$

$$(c) \quad F_n = [N, (a_+)^n]$$

$$n=1 \rightarrow F_1 = [N, a_+] = [a_+ a_-, a_+] = a_+ [a_-, a_+] + [a_+, a_-] a_+$$

$$F_1 = a_+$$

$$n=2 \rightarrow F_2 = [N, (a_+)^2] = a_+ [N, a_+] + [N, a_+] a_+$$

$$F_2 = (a_+)^2 + (a_+)^2 = 2(a_+)^2$$

$$\text{Guess } F_n = [N, (a_+)^n] = n(a_+)^n$$

assume it to be true for n , then

$$F_{n+1} = [N, (a_+)^{n+1}] = a_+ [N, (a_+)^n] + [N, a_+] (a_+)^n$$

$$= a_+ (n a_+^n) + (a_+)^n a_+$$

$$F_{n+1} = (n+1) a_+^{n+1}$$

Thus by induction we proved that $F_n = [N, a_+^n] = n a_+^n$

$$H(a_+^n \phi_0) = \hbar\omega \left(N + \frac{1}{2} \right) a_+^n \phi_0 = \frac{1}{2} \hbar\omega (a_+^n \phi_0) + \hbar\omega N (a_+^n \phi_0)$$

$$= \frac{1}{2} \hbar\omega (a_+^n \phi_0) + \{ [N, a_+^n] + a_+^n N \} \hbar\omega \phi_0$$

$$= \frac{1}{2} \hbar\omega (a_+^n \phi_0) + \{ n a_+^n + a_+^n 0 \} \hbar\omega \phi_0$$

$$H(a_+^n \phi_0) = \hbar\omega \left(n + \frac{1}{2} \right) a_+^n \phi_0 = E_n (a_+^n \phi_0)$$

$$\text{but } H \phi_n = E_n \phi_n \quad \text{no degeneracy} \Rightarrow a_+^n \phi_0 = A_n \phi_n$$

$$\text{From your formula sheet you know that } a_+^n \phi_0 = \sqrt{n!} \phi_n \quad \text{i.e. } A_n = \sqrt{n!}$$

Q3.

(a) $K = h\nu - W \geq 0 \Rightarrow \nu \geq \frac{W}{h} = \nu_0$

Thus photoelectric current will occur only if $\nu \geq \nu_0 = \frac{W}{h}$

(b) If Intensity is constant then $I = \frac{E}{tA} = \text{constant}$

But $E = n(h\nu)$; $n = \text{number of photon}$

$E = n(h\nu) = \text{const} \Rightarrow \text{as } \nu \uparrow \text{ then } n \downarrow$

Each photon gives one electron to photocurrent \Rightarrow Photocurrent decreases as $n \downarrow$

(c) $K = h\nu - W = \frac{hc}{\lambda} - W = e|V_s|$

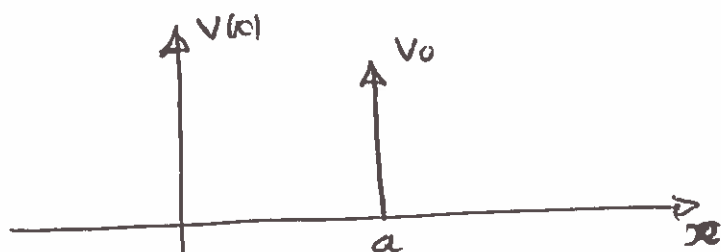
$\Rightarrow |V_s| = \frac{hc}{\lambda e} - \frac{W}{e} = \frac{2\pi\hbar c}{\lambda e} - \frac{W}{e} = \frac{(6.28)(197.3 \text{ eV}\cdot\text{nm})}{400 \text{ nm} \times e} - 1.5 \text{ V}$

$|V_s| = 3.1 \text{ V} - 1.5 \text{ V} = 1.6 \text{ V}$

(d) Interference experiment are the most important in exposing the wave like behavior. One such important in the history of Q.M. is Young double slit experiment for both electromagnetic and matter waves.

Q4.

$V(x) = V_0 \delta(x-a)$



(a) $E = \int \Psi^*(x) H \Psi(x) dx = \int_{-\infty}^{+\infty} \left(-\frac{\hbar^2}{2m} \frac{d^2 \Psi}{dx^2} + V_0 \delta(x-a) \right) \Psi(x) dx \quad \text{see (Q2.a)}$

$E = \frac{\hbar^2}{2m} \int \left| \frac{d\Psi}{dx} \right|^2 dx + V_0 |\Psi(a)|^2 > 0$

Since both quantities are definite positive.

(b) For $E > 0$ for both $x < a$ and $x > a$ we have

$\frac{d^2 \Psi}{dx^2} = -\frac{2mE}{\hbar^2} \Psi(x) = -k^2 \Psi(x) ; k = \frac{\sqrt{2mE}}{\hbar}$

Thus $\Psi(x) = \begin{cases} A e^{ik(x-a)} + B e^{-ik(x-a)} & \text{for } x < a \\ C e^{ik(x-a)} + D e^{-ik(x-a)} & \text{for } x > a \end{cases}$

the use of $(x-a)$ is for convenience and ease of calculation, you can avoid it.

Let us use $\psi(x) = \begin{cases} A e^{ikx} + B e^{-ikx} & \text{for } x < a \\ C e^{ikx} + D e^{-ikx} & \text{for } x > a \end{cases}$

then Schrodinger equation reads

$$\frac{d^2\psi}{dx^2} = -\frac{2mE}{\hbar^2} \psi(x) + \frac{2mV_0}{\hbar^2} \delta(x-a) \psi(x)$$

integrate this equation at the interface from $x = a-\epsilon$ to $x = a+\epsilon$

$$\int_{a-\epsilon}^{a+\epsilon} \frac{d^2\psi}{dx^2} dx = -\frac{2mE}{\hbar^2} \int_{a-\epsilon}^{a+\epsilon} dx \psi(x) + \frac{2mV_0}{\hbar^2} \int_{a-\epsilon}^{a+\epsilon} \delta(x-a) \psi(x) dx$$

$$\Rightarrow \left. \frac{d\psi}{dx} \right|_{a+\epsilon} - \left. \frac{d\psi}{dx} \right|_{a-\epsilon} = -\frac{2mE}{\hbar^2} \psi(a) (2\epsilon) + \frac{2mV_0}{\hbar^2} \psi(a)$$

limit as $\epsilon \rightarrow 0$ give

$$\left. \frac{d\psi}{dx} \right|_{a+} - \left. \frac{d\psi}{dx} \right|_{a-} = \frac{2mV_0}{\hbar^2} \psi(a)$$

(c) Using $\psi(a+) = \psi(a-)$; $\psi'(a+) - \psi'(a-) = \frac{2mV_0}{\hbar^2} \psi(a)$

gives

$$A e^{ika} + B e^{-ika} = C e^{ika} + D e^{-ika}$$

$$ik(C e^{ika} - D e^{-ika}) - ik(A e^{ika} - B e^{-ika}) = \frac{2mV_0}{\hbar^2} (A e^{ika} + B e^{-ika})$$

$$\Rightarrow \begin{cases} C e^{ika} + D e^{-ika} = A e^{ika} + B e^{-ika} \\ C e^{ika} - D e^{-ika} = A e^{ika} (1-2i\alpha) - B e^{-ika} (1+2i\alpha) \end{cases}$$

$$\alpha = \frac{mV_0}{\hbar^2 k}$$

$$\Rightarrow \begin{pmatrix} e^{ika} & e^{-ika} \\ e^{ika} & -e^{-ika} \end{pmatrix} \begin{pmatrix} C \\ D \end{pmatrix} = \begin{pmatrix} e^{ika} & e^{-ika} \\ e^{ika}(1-2i\alpha) & -e^{-ika}(1+2i\alpha) \end{pmatrix} \begin{pmatrix} A \\ B \end{pmatrix}$$

using $\begin{pmatrix} e^{ika} & e^{-ika} \\ e^{ika} & -e^{-ika} \end{pmatrix}^{-1} = -\frac{1}{2} \begin{pmatrix} -e^{-ika} & -e^{-ika} \\ -e^{ika} & e^{ika} \end{pmatrix} = \frac{1}{2} \begin{pmatrix} e^{-ika} & e^{-ika} \\ e^{ika} & -e^{ika} \end{pmatrix}$

$$\Rightarrow \begin{pmatrix} C \\ D \end{pmatrix} = \frac{1}{2} \begin{pmatrix} e^{-ika} & e^{-ika} \\ e^{ika} & -e^{ika} \end{pmatrix} \begin{pmatrix} e^{ika} & e^{-ika} \\ e^{ika}(1-2i\alpha) & -e^{-ika}(1+2i\alpha) \end{pmatrix} \begin{pmatrix} A \\ B \end{pmatrix}$$

$$\begin{pmatrix} C \\ D \end{pmatrix} = \frac{1}{2} \begin{pmatrix} 2 - 2i\alpha & e^{-2ika} (-2i\alpha) \\ e^{2ika} & 2 + 2i\alpha \end{pmatrix} \begin{pmatrix} A \\ B \end{pmatrix}$$

$$\Rightarrow \begin{pmatrix} C \\ D \end{pmatrix} = \begin{pmatrix} 1 - i\alpha & -i\alpha e^{-2ika} \\ i\alpha e^{2ika} & 1 + i\alpha \end{pmatrix} \begin{pmatrix} A \\ B \end{pmatrix} = M \begin{pmatrix} A \\ B \end{pmatrix}$$

$$M = \begin{pmatrix} 1 - i\alpha & -i\alpha e^{-2ika} \\ i\alpha e^{2ika} & 1 + i\alpha \end{pmatrix}; |M| = 1$$

(d) $D = 0$ for an incident wave from far left, that

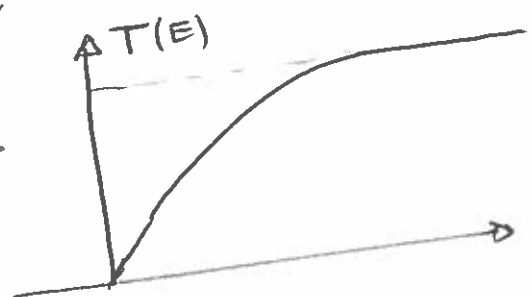
$$\begin{pmatrix} C \\ 0 \end{pmatrix} = \begin{pmatrix} M_{11} & M_{12} \\ M_{21} & M_{22} \end{pmatrix} \begin{pmatrix} A \\ B \end{pmatrix}$$

$$\Rightarrow \begin{aligned} C &= M_{11} A + M_{12} B \\ 0 &= M_{21} A + M_{22} B \end{aligned} \Rightarrow B = -\frac{M_{21}}{M_{22}} A$$

$$\text{then } C = A \left(M_{11} - \frac{M_{12} M_{21}}{M_{22}} \right) = A \frac{|M|}{M_{22}}$$

$$\Rightarrow T = \left| \frac{C}{A} \right|^2 = \frac{|M|^2}{|M_{22}|^2} = \frac{1}{|M_{22}|^2}$$

Since $|M| = 1$



$$T = \frac{1}{|M_{22}|^2} = \frac{1}{1 + \alpha^2} = \frac{1}{1 + \frac{(m v_0)^2}{\hbar^2 k^2}}$$

$$\alpha = \frac{m v_0}{\hbar k}; \quad k = \frac{2mE}{\hbar^2}$$

$$T(E) = \frac{1}{1 + \frac{(m v_0)^2}{\hbar^2 (2mE)}} = \frac{1}{1 + \frac{E_0}{E}} = \frac{E}{E + E_0}$$

$$E_0 = \frac{m v_0^2}{2\hbar^2}$$