# LECTURE 11: THE MEASUREMENT OF POWER SPECTRUM MULTIPOLES

### GONG-BO ZHAO

## 1. Measurement of power spectrum multipoles

(1) 
$$x = \chi \cos \delta \cos \alpha$$
$$y = \chi \cos \delta \sin \alpha$$
$$z = \chi \sin \delta$$

The Yamamoto estimator,

Date: April 17, 2019.

# North Celestial Pole celestial equator hour circle

South Celestial Pole

$$\hat{P}_{\ell}(k) = \frac{(2\ell+1)}{I} \int \frac{d\Omega_k}{4\pi} \left[ \int d\mathbf{r}_1 \int d\mathbf{r}_2 F\left(\mathbf{r}_1\right) F\left(\mathbf{r}_2\right) \times e^{i\mathbf{k}\cdot(\mathbf{r}_1-\mathbf{r}_2)} \mathcal{L}_{\ell}\left(\hat{\mathbf{k}}\cdot\hat{\mathbf{r}}_h\right) - P_{\ell}^{\text{noise}}(\mathbf{k}) \right],$$

$$\hat{P}_{\ell}^{\mathrm{Yama}}(k) = \frac{(2\ell+1)}{I} \int \frac{d\Omega_k}{4\pi} \left[ \int d\mathbf{r}_1 F\left(\mathbf{r}_1\right) e^{i\mathbf{k}\cdot\mathbf{r}_1} \times \int d\mathbf{r}_2 F\left(\mathbf{r}_2\right) e^{-i\mathbf{k}\cdot\mathbf{r}_2} \mathcal{L}_{\ell}\left(\hat{\mathbf{k}}\cdot\hat{\mathbf{r}}_2\right) - P_{\ell}^{\mathrm{noise}}(\mathbf{k}) \right],$$

which can not be evaluated using FFTs.

One solution is the following [2],

(2) 
$$A_n(\mathbf{k}) = \int d\mathbf{r} (\hat{\mathbf{k}} \cdot \hat{\mathbf{r}})^n F(\mathbf{r}) e^{i\mathbf{k} \cdot \mathbf{r}}$$

$$\hat{P}_{0}^{\text{Yama}}(k) = \frac{1}{I} \int \frac{d\Omega_{k}}{4\pi} \left[ A_{0}(\mathbf{k}) A_{0}^{*}(\mathbf{k}) \right] - P_{0}^{\text{noise}}$$

$$\hat{P}_{2}^{\text{Yama}}(k) = \frac{5}{2I} \int \frac{d\Omega_{k}}{4\pi} A_{0}(\mathbf{k}) \left[ 3A_{2}^{*}(\mathbf{k}) - A_{0}^{*}(\mathbf{k}) \right]$$

$$\hat{P}_{4}^{\text{Yama}}(k) = \frac{9}{8I} \int \frac{d\Omega_{k}}{4\pi} A_{0}(\mathbf{k}) \left[ 35A_{4}^{*}(\mathbf{k}) - 30A_{2}^{*}(\mathbf{k}) + 3A_{0}^{*}(\mathbf{k}) \right]$$

(4) 
$$\hat{\mathbf{k}} \cdot \hat{\mathbf{r}} = \frac{k_x r_x + k_y r_y + k_z r_z}{kr}$$

(5) 
$$A_{2}(\mathbf{k}) = \frac{1}{k^{2}} \left\{ k_{x}^{2} B_{xx}(\mathbf{k}) + k_{y}^{2} B_{yy}(\mathbf{k}) + k_{z}^{2} B_{zz}(\mathbf{k}) + 2 \left[ k_{x} k_{y} B_{xy}(\mathbf{k}) + k_{x} k_{z} B_{xz}(\mathbf{k}) + k_{y} k_{z} B_{yz}(\mathbf{k}) \right] \right\}$$

(6) 
$$B_{ij}(\mathbf{k}) \equiv \int d\mathbf{r} \frac{r_i r_j}{r^2} F(\mathbf{r}) e^{i\mathbf{k} \cdot \mathbf{r}}$$

(7) 
$$A_{4}(\mathbf{k}) = \frac{1}{k^{4}} \left\{ k_{x}^{4} C_{xxx} + k_{y}^{4} C_{yyy} + k_{z}^{4} C_{zzz} + 4 \left[ k_{x}^{3} k_{y} C_{xxy} + k_{x}^{3} k_{z} C_{xxz} + k_{y}^{3} k_{x} C_{yyx} + 4_{y}^{3} k_{z} C_{yyz} + k_{z}^{3} k_{z} C_{zzx} + k_{z}^{3} k_{y} C_{zzy} \right] + 6 \left[ k_{x}^{2} k_{y}^{2} C_{xyy} + k_{x}^{2} k_{z}^{2} C_{xzz} + k_{y}^{2} k_{z} C_{yzz} \right] + 12 k_{x} k_{y} k_{z} \left[ k_{x} C_{xyz} + k_{y} C_{yxz} + k_{z} C_{zxy} \right] \right\}$$

(8) 
$$C_{ijl}(\mathbf{k}) \equiv \int d\mathbf{r} \frac{r_i^2 r_j r_l}{r^4} F(\mathbf{r}) e^{i\mathbf{k}\cdot\mathbf{r}}$$

This works, but not efficient. A better way is using the addition theorem [3],

(9) 
$$\mathcal{L}_{\ell}\left(\hat{\mathbf{r}}_{1}\cdot\hat{\mathbf{r}}_{2}\right) = \frac{4\pi}{2\ell+1} \sum_{m=-\ell}^{\ell} Y_{\ell m}\left(\hat{\mathbf{r}}_{1}\right) Y_{\ell m}^{\star}\left(\hat{\mathbf{r}}_{2}\right)$$

which is a generalisation of the well-known formula below,

(10) 
$$\sin(x+y) = \sin(x)\cos(y) + \cos(x)\sin(y) \\ \cos(x+y) = \cos(x)\cos(y) - \sin(x)\sin(y)$$

Now,

(11) 
$$\widehat{P}_{\ell}(k) = \frac{2\ell + 1}{I} \int \frac{\mathrm{d}\Omega_k}{4\pi} F_0(\mathbf{k}) F_{\ell}(-\mathbf{k})$$

(12) 
$$F_{\ell}(\mathbf{k}) \equiv \int d\mathbf{r} F(\mathbf{r}) e^{i\mathbf{k}\cdot\mathbf{r}} \mathcal{L}_{\ell}(\hat{\mathbf{k}}\cdot\hat{\mathbf{r}})$$

$$= \frac{4\pi}{2\ell+1} \sum_{m=-\ell}^{\ell} Y_{\ell m}(\hat{\mathbf{k}}) \int d\mathbf{r} F(\mathbf{r}) Y_{\ell m}^{*}(\hat{\mathbf{r}}) e^{i\mathbf{k}\cdot\mathbf{r}}$$

where

(13) 
$$Y_{\ell m}(\theta, \phi) \equiv \begin{cases} \sqrt{\frac{2\ell+1}{2\pi} \frac{(\ell-m)!}{(\ell+m)!}} \mathcal{L}_{\ell}^{m}(\cos\theta) \cos m\phi & m > 0 \\ \sqrt{\frac{2\ell+1}{4\pi}} \mathcal{L}_{\ell}^{m}(\cos\theta) & m = 0 \\ \sqrt{\frac{2\ell+1}{2\pi} \frac{(\ell-|m|)!}{(\ell+|m|)!}} \mathcal{L}_{\ell}^{|m|}(\cos\theta) \sin|m|\phi & m < 0 \end{cases}$$

### References

- [1] D. Bianchi, H. Gil-Marn, R. Ruggeri and W. J. Percival, Mon. Not. Roy. Astron. Soc. **453**, no. 1, L11 (2015) doi:10.1093/mnrasl/slv090 [arXiv:1505.05341 [astro-ph.CO]].
- [2] N. Hand, Y. Li, Z. Slepian and U. Seljak, JCAP 1707, no. 07, 002 (2017) doi:10.1088/1475-7516/2017/07/002 [arXiv:1704.02357 [astro-ph.CO]].