Neo-Hookean Strain Energy Density Functions $$\operatorname{Updated}$$ on ${}_{\scriptscriptstyle -}$

Ida Ang

Contents

1	Definitions	2
2	Compressible Cases	2
3	Incompressible Cases	3
4	Summary	4

1 Definitions

Young's Modulus: E Poisson's Ratio: ν

Shear Modulus or first Lamé Parameter: μ

Second Lamé Parameter: λ

$$\mu = \frac{E}{2(1+\nu)} \quad \lambda = \frac{E\nu}{(1+\nu)(1-2\nu)} \tag{1.1}$$

Bulk Modulus: κ

$$\kappa = \lambda + \frac{2}{3}\mu\tag{1.2}$$

Deformation Gradient: **F**

Right Cauchy Green Tensor: $\mathbf{C} = \mathbf{F}^T \mathbf{F}$

1st Piola Kirchoff Stress: P

The third invariant of the deformation gradient

$$I_3 = J = \det \mathbf{F} \tag{1.3}$$

The first invariant of the right Cauchy Green Tensor:

$$I_1^{\mathbf{C}} = \operatorname{tr} \mathbf{C} \tag{1.4}$$

Lastly, we note a relationship between pressure and the bulk modulus

$$p = -\kappa I_3(I_3 - 1) \tag{1.5}$$

Proof of the derivative $\frac{\partial \operatorname{tr}(\mathbf{F}^{\mathbf{T}}\mathbf{F})}{\partial \mathbf{F}}$ in indicial:

$$\frac{\partial I_{1}^{\mathbf{C}}}{\partial \mathbf{F}} = \frac{\partial \operatorname{tr}(\mathbf{F}^{\mathbf{T}}\mathbf{F})}{\partial \mathbf{F}}$$

$$= \frac{\partial F_{kI}F_{kI}}{\partial F_{pQ}} \quad \text{Product Rule}$$

$$= \frac{\partial F_{kI}}{\partial F_{pQ}}F_{kI} + F_{kI}\frac{\partial F_{kI}}{\partial F_{pQ}}$$

$$= \delta_{kp}\delta_{IQ}F_{kI} + F_{kI}\delta_{kp}\delta_{IQ}$$

$$= F_{pQ} + F_{pQ} = 2F_{pQ}$$

Therefore, we have the following relationships:

$$\frac{\partial I_1^{\mathbf{C}}}{\partial \mathbf{F}} = 2\mathbf{F}
\frac{\partial I_3}{\partial \mathbf{F}} = I_3 \mathbf{F}^{-T}$$
(1.6)

2 Compressible Cases

There are several variations for the strain energy density function of a compressible neo-Hookean material

$$W(\mathbf{F}) = \frac{\mu}{2} (I_1^{\mathbf{C}} - 3 - 2 \ln I_3) + \frac{\kappa}{2} (I_3 - 1)^2$$

$$W(\mathbf{F}) = \frac{\mu}{2} (I_1^{\mathbf{C}} - 3 - 2 \ln I_3) + \frac{\lambda}{2} (\ln I_3)^2$$
(2.1)

We can calculate the first PK for the first expression using Eq. 1.6

$$\mathbf{P} = \frac{\partial W}{\partial \mathbf{F}} \quad \text{using chain rule}
= \frac{\partial W}{\partial I_{1}^{\mathbf{C}}} \frac{\partial I_{1}^{\mathbf{C}}}{\partial \mathbf{F}} + \frac{\partial W}{\partial I_{3}} \frac{\partial I_{3}}{\partial \mathbf{F}}
= \frac{\mu}{2} (2\mathbf{F}) + \left[\frac{\mu}{2} \left(-2 \frac{\partial \ln I_{3}}{\partial I_{3}} \right) + \frac{\kappa}{2} \frac{\partial (I_{3} - 1)^{2}}{\partial I_{3}} \right] I_{3} \mathbf{F}^{-T}
= \mu \mathbf{F} + \left[-\frac{\mu}{I_{3}} + \kappa (I_{3} - 1) \right] I_{3} \mathbf{F}^{-T}
= \mu \mathbf{F} - \mu \mathbf{F}^{-T} + \kappa I_{3} (I_{3} - 1) \mathbf{F}^{-T}
\mathbf{P} = \mu (\mathbf{F} - \mathbf{F}^{-T}) + \kappa I_{3} (I_{3} - 1) \mathbf{F}^{-T}$$
(2.2)

For the second expression

$$\mathbf{P} = \frac{\partial W}{\partial I_{1}^{\mathbf{C}}} \frac{\partial I_{1}^{\mathbf{C}}}{\partial \mathbf{F}} + \frac{\partial W}{\partial I_{3}} \frac{\partial I_{3}}{\partial \mathbf{F}}
= \frac{\mu}{2} (2\mathbf{F}) + \left[\frac{\mu}{2} \left(-2 \frac{\partial \ln I_{3}}{\partial I_{3}} \right) + \frac{\lambda}{2} \frac{\partial (\ln I_{3})^{2}}{\partial I_{3}} \right] I_{3} \mathbf{F}^{-T}
= \mu \mathbf{F} + \left[-\frac{\mu}{I_{3}} + \lambda \ln I_{3} \right] I_{3} \mathbf{F}^{-T}
= \mu \mathbf{F} - \mu \mathbf{F}^{-T} + \lambda I_{3} (\ln I_{3}) \mathbf{F}^{-T}
\mathbf{P} = \mu (\mathbf{F} - \mathbf{F}^{-T}) + \lambda I_{3} (\ln I_{3}) \mathbf{F}^{-T}$$
(2.3)

3 Incompressible Cases

For an incompressible material, $I_1 = \det(\mathbf{F}) = 1$; therefore, Eq. 3.1 can be modified,

$$W(\mathbf{F}) = \frac{\mu}{2} (I_1^{\mathbf{C}} - 3 - 2\ln(1)) + \frac{\kappa}{2} (1 - 1)^2$$
$$W(\mathbf{F}) = \frac{\mu}{2} (I_1^{\mathbf{C}} - 3 - 2\ln(1)) + \frac{\lambda}{2} (\ln(1))^2$$

Therefore, we can write

$$W(\mathbf{F}) = \frac{\mu}{2} (I_1^{\mathbf{C}} - 3) \tag{3.1}$$

To enforce incompressibility, we can utilize the Lagrange multiplier method where the lagrange multiplier, p, enforces the constraint that J=1:

$$W(\mathbf{F}) = \frac{\mu}{2} (I_1^{\mathbf{C}} - 3) + p(I_3 - 1)$$

$$W(\mathbf{F}) = \frac{\mu}{2} (I_1^{\mathbf{C}} - 3 - 2 \ln I_3) + p(I_3 - 1)$$
(3.2)

The 1st PK of Eq. 3.2 the first expression

$$\mathbf{P} = \frac{\partial W}{\partial I_1^{\mathbf{C}}} \frac{\partial I_1^{\mathbf{C}}}{\partial \mathbf{F}} + \frac{\partial W}{\partial I_3} \frac{\partial I_3}{\partial \mathbf{F}}$$
$$= \frac{\mu}{2} (2\mathbf{F}) + pI_3 \mathbf{F}^{-T}$$
$$\mathbf{P} = \mu \mathbf{F} + pI_3 \mathbf{F}^{-T}$$

The 1st PK of the second expression

$$\mathbf{P} = \frac{\partial W}{\partial I_{1}^{\mathbf{C}}} \frac{\partial I_{1}^{\mathbf{C}}}{\partial \mathbf{F}} + \frac{\partial W}{\partial I_{3}} \frac{\partial I_{3}}{\partial \mathbf{F}}$$

$$= \frac{\mu}{2} (2\mathbf{F}) + \left[\frac{\mu}{2} \left(-2 \frac{\partial \ln I_{3}}{\partial I_{3}} \right) + p \right] I_{3} \mathbf{F}^{-T}$$

$$= \mu \mathbf{F} + \left[-\frac{\mu}{I_{3}} + p \right] I_{3} \mathbf{F}^{-T}$$

$$= \mu \mathbf{F} - \mu \mathbf{F}^{-T} + p I_{3} \mathbf{F}^{-T}$$

$$\mathbf{P} = \mu (\mathbf{F} - \mathbf{F}^{-T}) + p I_{3} \mathbf{F}^{-T}$$
(3.3)

4 Summary

Compressible neo-Hookean, case 1

$$W(\mathbf{F}) = \frac{\mu}{2} (I_1^{\mathbf{C}} - 3 - 2 \ln I_3) + \frac{\kappa}{2} (I_3 - 1)^2$$

$$\mathbf{P} = \mu(\mathbf{F} - \mathbf{F}^{-T}) + \kappa I_3 (I_3 - 1) \mathbf{F}^{-T}$$
(4.1)

Compressible neo-Hookean, case 2

$$W(\mathbf{F}) = \frac{\mu}{2} (I_1^{\mathbf{C}} - 3 - 2 \ln I_3) + \frac{\lambda}{2} (\ln I_3)^2$$

$$\mathbf{P} = \mu(\mathbf{F} - \mathbf{F}^{-T}) + \lambda I_3 (\ln I_3) \mathbf{F}^{-T}$$

$$(4.2)$$

Incompressible neo-Hookean, case 1

$$W(\mathbf{F}) = \frac{\mu}{2} (I_1^{\mathbf{C}} - 3) + p(I_3 - 1)$$

$$\mathbf{P} = \mu \mathbf{F} + pI_3 \mathbf{F}^{-T}$$
(4.3)

Incompressible neo-Hookean, case 2

$$W(\mathbf{F}) = \frac{\mu}{2} (I_1^{\mathbf{C}} - 3 - 2 \ln I_3) + p(I_3 - 1)$$

$$\mathbf{P} = \mu(\mathbf{F} - \mathbf{F}^{-T}) + pI_3\mathbf{F}^{-T}$$
(4.4)