Conclusions & Outlook

3 This is the way the world ends

27

4 Not with a bang but a whimper.

- T.S. Eliot, The Hollow Men

One of the remaining unsolved mysteries in the universe is the presumed existence of dark matter. Various astrophysical observations indicate the presence of an unknown type of matter next to the known visible matter. As it does not interact electromagnetically, only very little is know about this new type of matter, and many theories exist to explain its nature and origin. These theories can be tested through a variety of techniques and experiments. Depending on the exact nature of the dark matter particles, they could also be produced at colliders in high-energy collisions between Standard Model particles. Currently, the largest particle accelerator in the world is the LHC at CERN, which provides proton-proton collisions with a record centre-of-mass energy of 13 TeV at a very high rate. In this thesis, two complementary searches for dark matter have been performed with data collected by the CMS detector located at one of the interaction points of this collider.

For both analyses, simplified model interpretations were considered. While dark matter does not interact electromagnetically, in these models it is assumed to interact with the ordinary matter through a new force. In the first analysis, the dark matter candidates are expected to interact very weakly with Standard Model particles and leave the detector undetected. They can however be detected when they recoil against another object in the event. In this case, the studied signature is one or more jets together with missing transverse energy. This search was already performed during Run 1, and has now been improved for Run 2. One of the main developments was the refinement of the prediction of events with a Z boson decaying to two neutrinos in association with one or more jets, which led to an improved sensitivity. While the first iteration of this analysis with Run 2 data sets less stringent limits on the scenario with a vector mediator compared to the Run 1 analysis, the new results including the discussed improvements and using data corresponding to 12.9 fb⁻¹ set stronger limits up to mediator masses of 1.95 TeV. Similar results are obtained for an axial-vector mediator. Scalar and pseudoscalar mediator masses are excluded up to 100 GeV and 430 GeV, respectively.

In the second analysis, a different simplified model is considered, where the dark matter particles are assumed to interact strongly with Standard Model particles, through a new, light mediator. The produced pair of stable, neutral dark matter candidates will then give a signature consisting of a pair of trackless jets. First, a feasibility study for this model and corresponding signature was performed and published. For this study, the analysis strategy was developed, and the achievable sensitivity was investigated. This model had never been tested at colliders before, and as the feasibility study showed promising results, this signature was now studied for the first time at CMS in this analysis, using a type of jets that is

90

thou should explain that this appear limit on the cross section drives or small medictor mers; so small that the pringlefield model is probably too simple, since hadron resonances can chapter ?

typically rejected by the jet identification criteria. While this search is very sensitive to new physics and the expected background was very efficiently reduced, some complications emerged as well. As an example, care was taken to account for the possibility of a wrongly chosen primary vertex, as this problem in the reconstruction can mimic signal-like events. A different issue was related to data taking and originated from the tracker APV pre-amplifier saturation problem, which occurred during the first half of the 2016 data taking period at CMS. Due to the nature of this issue, the data affected by this problem were rejected for the analysis, and data corresponding to a total of 16.1 fb⁻¹ were used. Moreover, the photon veto, which is crucial to reject background events coming from the production of a photon and a jet, was extended to take into account photon conversions as well. The outcome of this analysis yielded no observation of new physics, and all the considered dark matter masses, from 1 GeV up to 1 TeV, were excluded. Additionally, model-independent limits were derived as well, excluding production cross sections down to 0.18 fb.

Both results can be translated into limits on the dark matter-nucleon interaction cross section, in order to compare the results with other experiments. Depending on the dark matter mass, the monojet analysis can exclude cross sections between approximately 10^{-6} fb and $0.1\,\mathrm{pb}$, for the scalar or vector mediator case. In comparison, the trackless jets analysis excludes interaction cross sections of about $1-10\,\mathrm{mb}$, thus complementing the monojet results at higher cross sections.

In the future, the monojet analysis will be able to cover a larger part of phase space, going to higher mediator and dark matter masses. Projections can for example be made for the expected sensitivity at the high luminosity LHC, which is expected to deliver $3000~{\rm fb}^{-1}$ of integrated luminosity at a centre-of-mass energy of $14~{\rm TeV}$. The expected reach is shown in Figure 7.1 for an axial-vector and pseudoscalar mediator. The nominal scenario assumes that the level of control of the E_T^{miss} distribution will remain the same as in the current analysis. Two more scenarios are shown, where the systematic uncertainties are reduced by a factor 2 and 4.

For the pseudoscalar mediator, the systematic uncertainties are dominated by the uncertainty on lepton identification and isolation efficiencies used in the selection of events in the control regions at low missing transverse momentum. At high missing transverse momentum on the other hand, the systematic uncertainties are dominated by the statistical uncertainty. For the nominal scenario shown in the right plot of Figure 7.1, the systematic uncertainties are scaled by the luminosity at high E_T^{miss} and scaled according to the predictions for the uncertainty on the lepton identification and isolation efficiencies at low E_T^{miss} . Additionally, a scenario is shown where the systematic uncertainties are reduced by a factor 2, as well as a scenario where the systematic uncertainties are scaled by the luminosity over the full E_T^{miss} range.

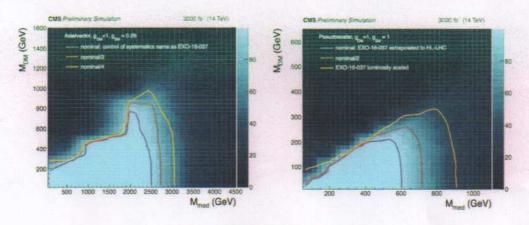


Figure 7.1: . Figures taken from [170].

The obtained prospects show that mediator masses up to at least 2.5 TeV and 600 GeV could be excluded in the axial-vector and pseudoscalar case, respectively. This corresponds to an increase in range of about 30%.

In the case of the trackless jets search, the entire considered dark matter mass range is already ex-

ma fraging

Ladd reant grans

fus com contrast

this with my

the comment

eto ove that this

is about

improvements

at high medicities

mans, which

translates to

comell Tells

33

37

- Cite: hep-gh/0604261

- cluded by the described results. In a next step, the mass range could be extended, or the search could be
- broadened by allowing more extra jets, or including missing transverse energy in the signature. Alterna-
- 3 tive signatures can also be obtained by assuming a different interaction cross section, such as emerging
- jets or a cluster of tracks in the muon systems. Aside from SIMPs, Hidden Valley models can also give rise to this kind of signature. In these models, the interaction with the hidden sector can for example
- happen through rare Higgs/boson decays. Additionally, trackless jets could also be produced by dark
- photons [171]. This model is now being studied as well, and the analysis can benefit from this first track-
- 8 less jets analysis, among other things concerning the issue of the tracker APV pre-amplifier saturation. In
- the considered scenario, the neutral dark photon is produced in association with an ordinary photon. The
- resulting signature is then composed of a photon and a trackless jet with energy deposits in the ECAL or

HCAL, or missing transverse energy, depending on the interaction strength.

I would repeat here explicitly that the cross section in the SIMP model is restricted to a mornow bound, alternise you lose the trackless get signature. This is why it was possible to exclude the model, and just up to some cross section