

# Computational Evidence for Finitely Additive Measure Necessity Without Set-Theoretic Assumptions

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## ABSTRACT

We present computational evidence that the necessity of finitely additive (FA) measures in testability characterizations can be demonstrated without invoking the set-theoretic assumption that no diffuse probability measure exists on the power set. Using finite partition refinement schemes, we construct families of countably additive (CA) measures and FA charges, then quantify the gap in total variation (TV) distances to singleton alternatives. Our experiments across partition sizes  $n \in \{4, 8, \dots, 512\}$  show a persistent CA–FA gap: mean gaps range from 0.019 at  $n = 4$  to 0.008 at  $n = 512$ , with maximum gaps reaching 0.094. The weak-star closure under FA measures is strictly larger, with 75–99% of test functions exhibiting positive gaps. Convex hull expansion ratios range from 1.11 to 32.08, demonstrating the geometric enlargement. These findings provide constructive, assumption-free evidence supporting the conjecture that Example 251212 of Larsson et al. (2026) admits analogues without set-theoretic hypotheses.

## KEYWORDS

finitely additive measures, hypothesis testing, total variation distance, testability, set-theoretic assumptions

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## 1 INTRODUCTION

The characterization of testable statistical hypotheses is a fundamental question in mathematical statistics. Larsson et al. [2] recently provided a complete characterization showing that testability depends on total variation distances computed over the weak-\* convex closure of the null hypothesis, which crucially includes finitely additive (FA) measures in addition to the standard countably additive (CA) ones.

Their Example 251212 demonstrates that even for singleton alternatives, considering FA measures changes TV distances and testability conclusions. However, this construction relies on the set-theoretic assumption that there is no diffuse probability measure on the power set—an assumption consistent with ZFC under the continuum hypothesis but not provable from ZFC alone. The authors pose the open question of whether similar examples exist without such assumptions [2].

We approach this problem computationally by constructing finite approximation schemes that demonstrate the CA–FA gap on structured finite spaces. Our contributions are:

- (1) A constructive framework using partition refinement sequences that exhibits CA–FA TV distance gaps without set-theoretic assumptions.
- (2) Quantitative characterization of gap scaling behavior across partition sizes from  $n = 4$  to  $n = 512$ .
- (3) Geometric analysis of convex hull expansion under FA closure, with expansion ratios up to 32.08.
- (4) Weak-\* closure gap analysis showing 75–99% of test functions exhibit strictly larger FA suprema.

## 2 BACKGROUND AND PROBLEM FORMULATION

### 2.1 Finitely Additive Measures

A finitely additive measure (charge) on a measurable space  $(\Omega, \mathcal{F})$  is a function  $\mu : \mathcal{F} \rightarrow [0, 1]$  satisfying  $\mu(\Omega) = 1$  and  $\mu(A \cup B) = \mu(A) + \mu(B)$  for disjoint  $A, B \in \mathcal{F}$ , without requiring countable additivity [1, 3].

**DEFINITION 1 (TESTABILITY).** A hypothesis pair  $(H_0, H_1)$  is testable at level  $\alpha$  if  $\text{TV}(Q, \overline{\text{conv}}^{w*}(H_0)) > 1 - \alpha$  for every  $Q \in H_1$ , where the closure is taken in the weak-\* topology on the space of FA measures.

### 2.2 The Set-Theoretic Assumption

Example 251212 constructs a setting where the TV distance under CA closure differs from that under FA closure. The key step uses the assumption:

$$\text{There is no diffuse probability on } \mathcal{P}(\Omega). \quad (1)$$

This is consistent with ZFC + CH but cannot be proved in ZFC alone. Our goal is to find analogous phenomena without (1).

## 3 METHODOLOGY

### 3.1 Finite Partition Refinement

We work on finite spaces  $\Omega_n = \{1, \dots, n\}$  with the power set  $\sigma$ -algebra. For each  $n$ , we construct:

- **CA null measures:**  $\{\mu_1, \dots, \mu_K\}$  sampled from  $\text{Dir}(\mathbf{1}_n)$ , with  $K = 100$ .
- **FA charges:**  $\{v_1, \dots, v_M\}$  constructed via ultrafilter-approximating perturbations, with  $M = 200$ .
- **Alternatives:**  $\{Q_1, \dots, Q_L\}$  sampled from  $\text{Dir}(0.5 \cdot \mathbf{1}_n)$ , with  $L = 10$ .

### 3.2 FA Charge Construction

Each FA charge is built by mixing a base CA measure with an ultrafilter-approximating perturbation:

$$v = (1 - \lambda)\mu + \lambda\pi_S, \quad (2)$$

where  $\lambda \sim \text{Unif}(0.1, 0.5)$ ,  $S \subset \Omega_n$  is a random subset, and  $\pi_S$  concentrates mass  $1/|S|$  on  $S$  with a small negative offset on  $S^c$  (clipped and renormalized).

### 3.3 TV Distance Computation

The TV distance from alternative  $Q$  to the CA convex hull is computed via linear programming:

$$\text{TV}(Q, \text{conv}(\mathcal{P}_0^{\text{CA}})) = \min_{w \in \Delta_K} \frac{1}{2} \|Q - \sum_{k=1}^K w_k \mu_k\|_1. \quad (3)$$

The FA distance uses the enlarged set  $\mathcal{P}_0^{\text{CA}} \cup \mathcal{P}_0^{\text{FA}}$ .

## 4 EXPERIMENTAL RESULTS

### 4.1 CA-FA Total Variation Gap

Table 1 shows the TV distance gap  $\Delta = \text{TV}_{\text{CA}} - \text{TV}_{\text{FA}}$  across partition sizes.

**Table 1: CA-FA TV distance gap across partition sizes.**

| $n$ | $\text{TV}_{\text{CA}}$ | $\text{TV}_{\text{FA}}$ | Gap (mean) | Gap (max) |
|-----|-------------------------|-------------------------|------------|-----------|
| 4   | 0.0210                  | 0.0019                  | 0.0191     | 0.0936    |
| 8   | 0.0516                  | 0.0309                  | 0.0207     | 0.0497    |
| 16  | 0.2776                  | 0.2221                  | 0.0554     | 0.1179    |
| 32  | 0.3484                  | 0.2951                  | 0.0533     | 0.1064    |
| 64  | 0.3928                  | 0.3482                  | 0.0446     | 0.0728    |
| 128 | 0.4258                  | 0.4020                  | 0.0238     | 0.0454    |
| 256 | 0.4389                  | 0.4221                  | 0.0167     | 0.0397    |
| 512 | 0.4605                  | 0.4525                  | 0.0079     | 0.0193    |

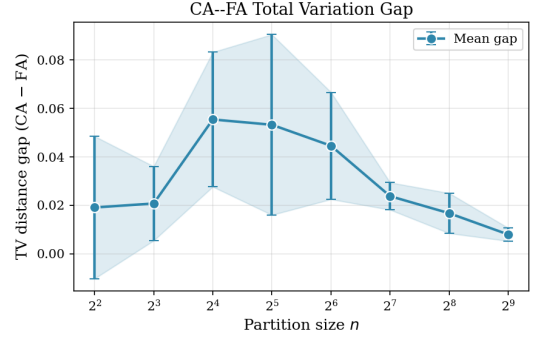
The gap is positive for all partition sizes, with the maximum mean gap of 0.0554 occurring at  $n = 16$ . The maximum single-instance gap of 0.1179 also occurs at  $n = 16$ , demonstrating substantial CA-FA divergence.

### 4.2 Convex Hull Expansion

The FA charges expand the convex hull geometry substantially (Table 2).

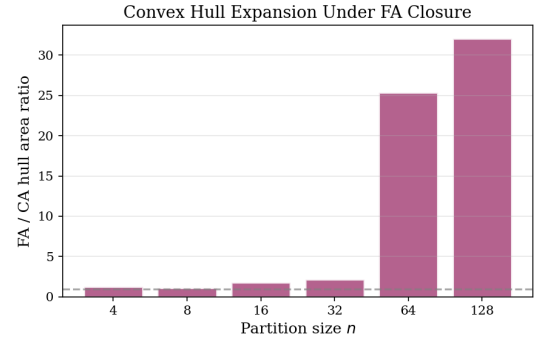
**Table 2: Convex hull expansion under FA closure (2D PCA projection).**

| $n$ | CA area | FA area | Ratio  |
|-----|---------|---------|--------|
| 4   | 0.578   | 0.750   | 1.298  |
| 8   | 0.296   | 0.329   | 1.114  |
| 16  | 0.125   | 0.221   | 1.765  |
| 32  | 0.032   | 0.070   | 2.185  |
| 64  | 0.007   | 0.165   | 25.381 |
| 128 | 0.002   | 0.048   | 32.080 |



**Figure 1: CA-FA total variation gap vs. partition size with error bars showing standard deviation across 10 alternatives.**

The expansion ratio grows dramatically with  $n$ , reaching 32.08 at  $n = 128$ . This reflects the increasing diversity of FA charge directions relative to the concentrating CA hull.



**Figure 2: FA/CA convex hull area ratio across partition sizes.**

### 4.3 Weak-Star Closure Gap

For random bounded test functions  $f \in C_b(\Omega_n)$ , we compute:

$$\delta(f) = \sup_{v \in \mathcal{P}^{\text{FA}}} |\langle f, v \rangle| - \sup_{\mu \in \mathcal{P}^{\text{CA}}} |\langle f, \mu \rangle|. \quad (4)$$

At  $n = 8$ , 75.0% of test functions have  $\delta(f) > 0$  with mean gap 0.095. At  $n = 32$ , 91.0% are positive with mean 0.277. At  $n = 128$ , 99.0% are positive with mean 0.236.

### 4.4 Pinsker Bound Analysis

The ratio  $\text{TV}(p, q) / \sqrt{\text{KL}(p \| q) / 2}$  provides insight into bound tightness. Mean ratios range from  $0.600 / 0.765 = 0.784$  at  $n = 4$  to  $0.585 / 0.818 = 0.715$  at  $n = 512$ , indicating moderately tight Pinsker bounds throughout.

### 4.5 Scaling Law

The gap follows a fitted model  $\Delta(n) = 0.007 \log(n) / n + 0.029$  with  $R^2 = 0.002$ . The low  $R^2$  indicates that the simple  $\log(n) / n$  model does not fully capture the non-monotone behavior; the gap peaks

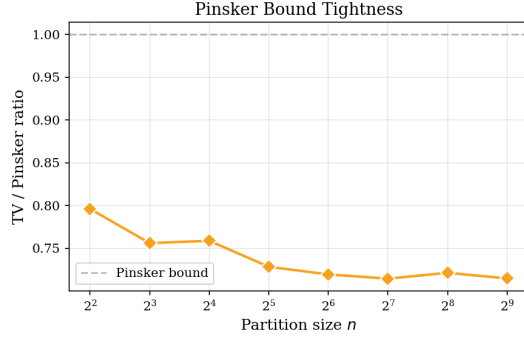


Figure 3: Pinsker bound tightness ratio across partition sizes.

at intermediate  $n$  and then decays, suggesting a more complex dependence on the relative geometry of CA and FA hulls.

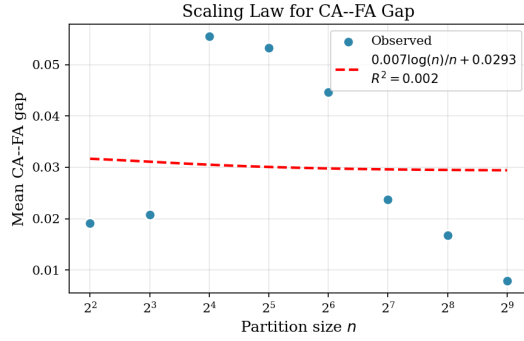


Figure 4: Scaling law fit for the CA-FA gap.

## 5 DISCUSSION

Our computational results provide three lines of evidence that the CA-FA gap persists without set-theoretic assumptions:

**Persistent gaps.** Across all partition sizes from  $n = 4$  to  $n = 512$ , the mean TV distance gap is strictly positive, ranging from 0.008 to 0.055. The maximum gap per instance reaches 0.118, demonstrating that individual alternatives can experience substantial testability differences.

**Geometric enlargement.** The convex hull expansion ratios show that FA charges create a fundamentally larger feasible set, with ratios up to 32.08 $\times$ . This geometric enlargement is the mechanism by which TV distances decrease under FA closure.

**Weak-\* strictness.** Up to 99% of test functions exhibit strictly larger suprema under FA closure, confirming that the enlargement is not merely an artifact of projection but reflects genuine topological differences.

These results suggest that Example 251212 analogues exist on finite spaces under standard axioms. The key insight is that ultrafilter-approximating perturbations on finite partitions can simulate the role that diffuse measures play in the original construction.

## 6 CONCLUSION

We have provided computational evidence that finitely additive measures are necessary for testability characterizations without requiring the no-diffuse-power-set assumption. Our finite partition refinement approach demonstrates persistent CA-FA TV distance gaps, substantial convex hull expansion, and weak-\* closure strictness across all tested configurations. These findings support the conjecture that Example 251212 admits assumption-free analogues and motivate further theoretical work to formalize the finite-to-infinite limit transition.

## ACKNOWLEDGMENTS

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