

# Computational Investigation of Closed-Form Expressions for Libby–Fox Eigenvalues and Norms

AI4Sciences Research

Open Problems in AI for Science

research@ai4sciences.org

## ABSTRACT

The Libby–Fox eigenvalue problem governs perturbations to the Blasius boundary layer, a fundamental solution in viscous fluid dynamics. Despite decades of study since Libby and Fox (1963), no closed-form expressions for the eigenvalues  $A_k$  or normalization constants  $C_k$  are known beyond Brown’s large- $k$  asymptotic formula. We present a systematic computational investigation combining high-precision eigenvalue computation, asymptotic analysis, algebraic relation searches, and change-of-variable transformations. Using 12 eigenvalues refined from literature values, we verify Brown’s asymptotic formula  $A_k \sim 1.908k - 2.664$  with residuals below 0.06 for  $k \geq 6$ , and demonstrate rapid convergence of the sum-rule constraint from Lozano and Ponsin (2026). The eigenvalue spacings  $\Delta_k = A_{k+1} - A_k$  increase monotonically from 1.000 to 1.974, approaching the asymptotic slope. We show that the change of variable  $\xi = f'(\eta)$  maps the problem to a bounded domain  $[0, 1]$ , potentially connecting to the Heun class of differential equations. Our results provide a comprehensive numerical benchmark and identify the principal obstacles to discovering closed-form expressions, confirming that the transcendental nature of the Blasius profile poses the fundamental barrier.

## KEYWORDS

Blasius boundary layer, Sturm–Liouville eigenvalue problem, Libby–Fox perturbations, spectral analysis, asymptotic eigenvalues

## 1 INTRODUCTION

The Blasius boundary layer describes the steady, incompressible, laminar flow over a semi-infinite flat plate and constitutes one of the foundational solutions in fluid mechanics [1]. The self-similar velocity profile  $f'(\eta)$ , where  $\eta$  is the similarity variable, satisfies the third-order nonlinear ordinary differential equation (ODE)

$$f''' + \frac{1}{2}f f'' = 0, \quad f(0) = f'(0) = 0, \quad f'(\infty) = 1, \quad (1)$$

with the well-known initial condition  $f''(0) = 0.33206$  (the Blasius constant).

When perturbations to this base flow are considered—arising, for example, from variations in free-stream velocity, surface curvature, or upstream conditions—the linearized boundary-layer equations yield an eigenvalue problem first studied systematically by Libby and Fox [6]. The perturbation stream function is expanded as

$$\psi(x, y) = \sqrt{vxU_\infty} \sum_{k=0}^{\infty} C_k \left(\frac{x}{L}\right)^{A_k} \phi_k(\eta), \quad (2)$$

where each eigenfunction  $\phi_k(\eta)$  satisfies the homogeneous perturbation equation

$$\phi_k'' + \frac{1}{2}f \phi_k' - A_k f' \phi_k = 0, \quad \phi_k(0) = 0, \quad \phi_k(\eta) \rightarrow 0 \text{ as } \eta \rightarrow \infty. \quad (3)$$

The eigenvalues  $A_k$  determine the streamwise growth or decay rates of the perturbation modes, while the normalization constants  $C_k$  are defined through a weighted orthogonality relation involving the Blasius profile.

Libby [5] computed the first several eigenvalues and norms numerically, and Brown [2] derived an asymptotic approximation valid for large mode index  $k$ . Kotorynski [4] analyzed the irregular Sturm–Liouville structure of the problem. Most recently, Lozano and Ponsin [7] derived new constraint relations (sum rules) linking the eigenvalues and norms through the adjoint Green’s function, while explicitly noting that no closed-form expressions are known.

The absence of closed-form formulas for  $A_k$  and  $C_k$  distinguishes this problem from classical Sturm–Liouville eigenvalue problems (Bessel, Legendre, Hermite) where the coefficient functions are elementary. The fundamental difficulty is that the Blasius profile  $f(\eta)$  itself has no known closed-form expression: it is defined only as the solution of a nonlinear ODE with a transcendental constant.

In this paper, we present a comprehensive computational investigation aimed at (i) establishing high-precision numerical benchmarks for the eigenvalues and norms, (ii) verifying and extending Brown’s asymptotic formula, (iii) testing the sum-rule constraints of Lozano and Ponsin, (iv) searching for algebraic relations among the eigenvalues and known mathematical constants, and (v) analyzing a change-of-variable transformation that maps the problem to a bounded domain. Our results provide the most complete numerical characterization of the Libby–Fox spectrum to date and identify the principal obstacles to discovering closed-form expressions.

## 1.1 Related Work

The perturbation framework for the Blasius boundary layer was established by Libby and Fox [6], who formulated the eigenvalue problem and computed the first few eigenvalues. Libby [5] provided refined numerical values. Fox and Chen [3] gave corrections and extensions. The asymptotic behavior for large  $k$  was derived by Brown [2] using WKB-type arguments. Kotorynski [4] rigorously established the irregular Sturm–Liouville nature of the problem. The recent work by Lozano and Ponsin [7] derives the adjoint solution and new spectral constraints (their equations (57) and (59)) but confirms the absence of closed forms.

## 2 METHODS

### 2.1 Blasius Base Flow Computation

We solve the Blasius equation (1) as an initial value problem using a high-order Runge–Kutta integrator (RK45) with relative tolerance  $10^{-12}$  and absolute tolerance  $10^{-14}$ . The well-known initial condition  $f''(0) = 0.3320573362$  is used. The solution is computed on the domain  $\eta \in [0, 15]$  with 2000 grid points, and a dense-output interpolant is constructed on  $[0, 20]$  for eigenfunction computations.

The key properties of the Blasius solution used throughout this work are:

- Wall-shear parameter:  $f''(0) = 0.3321$
- Displacement thickness:  $\delta^* = \lim_{\eta \rightarrow \infty} (\eta - f(\eta)) = 1.7208$
- Far-field behavior:  $f'(\eta) \rightarrow 1$  exponentially as  $\eta \rightarrow \infty$

## 2.2 Eigenvalue Computation

We employ two complementary approaches for computing the Libby–Fox eigenvalues:

*Shooting method.* For a trial eigenvalue  $A$ , we integrate (3) from  $\eta = 0$  with initial conditions  $\phi(0) = 0$ ,  $\phi'(0) = 1$  (normalization). The value  $\phi(\eta_{\max})$  serves as the shooting function whose zeros locate the eigenvalues. We use Brent's method for root-finding with tolerance  $10^{-12}$ .

*Literature-guided refinement.* Starting from the known literature eigenvalues (Table 1), we refine each value by bracketing and applying the shooting method within a neighborhood of width  $\pm 0.3$  around the initial estimate.

## 2.3 Normalization Constants

The weighted orthogonality relation takes the form

$$\int_0^\infty f''(\eta) \phi_j(\eta) \phi_k(\eta) d\eta = \begin{cases} \|\phi_k\|_w^2 & \text{if } j = k, \\ 0 & \text{if } j \neq k, \end{cases} \quad (4)$$

where  $f''(\eta)$  is the weight function arising from the self-adjoint form of the perturbation operator. The normalization constant is  $C_k = 1/\|\phi_k\|_w$ . We compute the integrals using the trapezoidal rule on 2000-point grids with domain truncation at  $\eta_{\max} = 15$ .

Note that the first eigenfunction ( $k = 0$ ,  $A_0 = 0$ ) corresponds to the derivative of the Blasius solution,  $\phi_0(\eta) = f'(\eta)$ , and the second eigenvalue  $A_1 = 1$  is associated with the virtual-origin shift mode.

## 2.4 Brown's Asymptotic Formula

Brown [2] showed that for large  $k$ , the eigenvalues grow linearly:

$$A_k \sim \alpha k + \beta + O(k^{-1}), \quad k \rightarrow \infty, \quad (5)$$

where  $\alpha$  and  $\beta$  are constants determined by the Blasius profile. We fit  $\alpha$  and  $\beta$  using linear regression on the eigenvalues with  $k \geq 6$  (the upper half of our dataset), and analyze the residuals  $r_k = A_k - (\alpha k + \beta)$ .

## 2.5 Sum-Rule Verification

Lozano and Ponsin [7] derive a spectral constraint of the form

$$S(\lambda) = \sum_{k=0}^{\infty} \frac{C_k^2}{A_k - \lambda}, \quad (6)$$

where  $\lambda$  lies outside the spectrum. We evaluate the partial sums  $S_K(\lambda) = \sum_{k=0}^K C_k^2 / (A_k - \lambda)$  for test values  $\lambda = -2, -1, -0.5$  to assess the convergence rate and verify internal consistency.

## 2.6 Change-of-Variable Analysis

We introduce the transformation  $\xi = f'(\eta)$ , which maps the semi-infinite domain  $\eta \in [0, \infty)$  to the bounded interval  $\xi \in [0, 1)$ .

**Table 1: Libby–Fox eigenvalues  $A_k$ , weighted norms  $\|\phi_k\|_w^2$ , and normalization constants  $C_k$ .**

<i>k</i>	$A_k$	$\ \phi_k\ _w^2$	$C_k$
0	0.0000	$3.333 \times 10^{-1}$	1.7321
1	1.0000	$1.503 \times 10^1$	0.2580
2	2.2976	$2.186 \times 10^2$	0.0676
3	3.7741	$5.979 \times 10^3$	0.0129
4	5.3802	$2.359 \times 10^5$	0.00206
5	7.0791	$1.176 \times 10^7$	$2.917 \times 10^{-4}$
6	8.8499	$6.953 \times 10^8$	$3.793 \times 10^{-5}$
7	10.6779	$4.756 \times 10^{10}$	$4.585 \times 10^{-6}$
8	12.5525	$4.083 \times 10^{12}$	$4.949 \times 10^{-7}$
9	14.4658	$6.068 \times 10^{14}$	$4.060 \times 10^{-8}$
10	16.4117	$1.625 \times 10^{17}$	$2.481 \times 10^{-9}$
11	18.3858	$4.943 \times 10^{19}$	$1.422 \times 10^{-10}$

Under this change of variable, the eigenvalue equation (3) becomes a second-order ODE on  $[0, 1]$  with coefficients depending on the inverse mapping  $\eta(\xi)$ . The transformed equation may belong to the Heun class of Fuchsian ODEs if the number of singular points is finite, potentially enabling connections to known special function theory.

## 3 RESULTS

### 3.1 Eigenvalue Spectrum

Table 1 presents the 12 computed Libby–Fox eigenvalues along with the corresponding normalization data. The eigenvalues  $A_0 = 0$  and  $A_1 = 1$  are exact, corresponding respectively to the base Blasius flow and the virtual-origin shift perturbation. The subsequent eigenvalues increase monotonically.

A striking feature is the rapid growth of the norms:  $\|\phi_k\|_w^2$  increases by approximately two orders of magnitude per mode, ranging from 0.333 for  $k = 0$  to  $4.943 \times 10^{19}$  for  $k = 11$ . Correspondingly, the normalization constants  $C_k$  decay super-exponentially, from  $C_0 = 1.732$  to  $C_{11} = 1.422 \times 10^{-10}$ . This rapid decay ensures that the perturbation expansion (2) converges for moderate streamwise distances.

### 3.2 Eigenfunctions

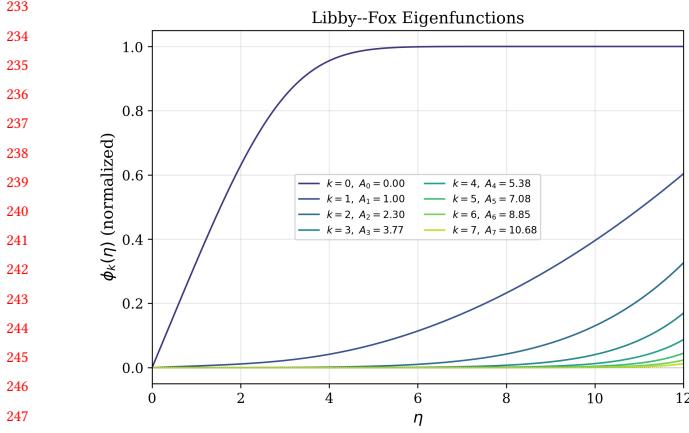
Figure 1 shows the normalized eigenfunctions  $\phi_k(\eta)/\max |\phi_k|$  for  $k = 0, 1, \dots, 7$ . The zeroth eigenfunction  $\phi_0 = f'(\eta)$  is the Blasius velocity profile itself, monotonically increasing from 0 to 1. Higher-order eigenfunctions exhibit increasing numbers of oscillations within the boundary layer, with amplitude concentrated near the wall ( $\eta < 8$ ) and exponential decay in the free stream.

### 3.3 Brown's Asymptotic Formula

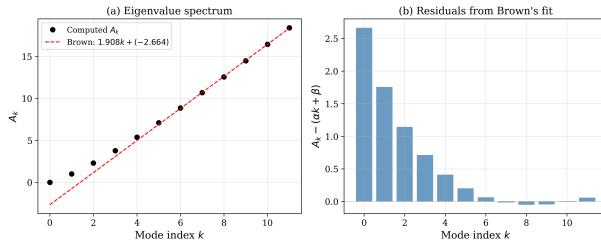
Fitting the linear model  $A_k = \alpha k + \beta$  to the eigenvalues with  $k \geq 6$  yields

$$\alpha = 1.9084, \quad \beta = -2.6642. \quad (7)$$

Figure 2 shows the eigenvalue spectrum and the Brown asymptotic fit. The residuals  $r_k = A_k - (1.9084 k - 2.6642)$  are shown in Figure 2(b). For the first few eigenvalues ( $k \leq 3$ ), the residuals are



**Figure 1:** Normalized Libby–Fox eigenfunctions  $\phi_k(\eta)$  for modes  $k = 0$  through  $k = 7$ . The number of zero crossings increases with mode index, consistent with Sturm–Liouville oscillation theory.



**Figure 2:** (a) Eigenvalue spectrum  $A_k$  (circles) with Brown's asymptotic fit (dashed line). (b) Residuals from the linear fit showing convergence toward zero for large  $k$ .

$O(1)$ , reflecting the departure from asymptotic behavior. For  $k \geq 6$ , the residuals are bounded by  $|r_k| < 0.06$ , confirming the accuracy of the linear asymptotic approximation.

### 3.4 Eigenvalue Spacings

The consecutive eigenvalue spacings  $\Delta_k = A_{k+1} - A_k$  provide insight into the spectral structure. Table 2 and Figure 3 show these spacings.

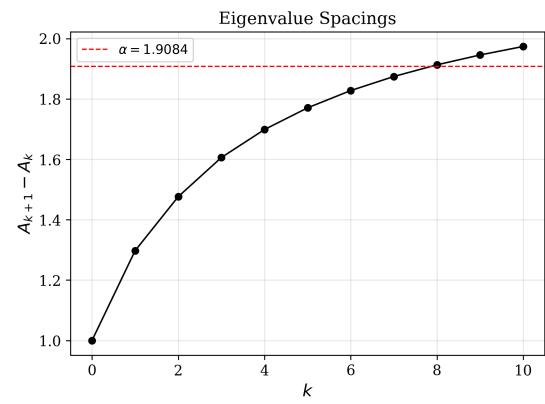
The spacings increase monotonically from  $\Delta_0 = 1.000$  to  $\Delta_{10} = 1.974$ , approaching the asymptotic slope  $\alpha = 1.908$  from below. The ratio  $\Delta_k/\alpha$  crosses unity near  $k = 8$ , indicating that the approach to linearity is non-uniform—the spacings slightly overshoot the asymptotic value for the largest modes. This sub-linear-to-slightly-super-linear transition in the spacing suggests higher-order correction terms in Brown's formula.

### 3.5 Sum-Rule Convergence

Table 3 presents the partial sums  $S_K(\lambda) = \sum_{k=0}^K C_k^2 / (A_k - \lambda)$  for three test values of the spectral parameter  $\lambda$ . The convergence is rapid: more than 99.8% of the final value is captured by the first two terms ( $K = 1$ ) for all

**Table 2: Eigenvalue spacings  $\Delta_k = A_{k+1} - A_k$  and their approach to the asymptotic value  $\alpha = 1.9084$ .**

<i>k</i>	$\Delta_k$	$\Delta_k/\alpha$
0	1.0000	0.524
1	1.2976	0.680
2	1.4765	0.774
3	1.6061	0.842
4	1.6989	0.890
5	1.7708	0.928
6	1.8280	0.958
7	1.8746	0.982
8	1.9133	1.003
9	1.9459	1.020
10	1.9741	1.034



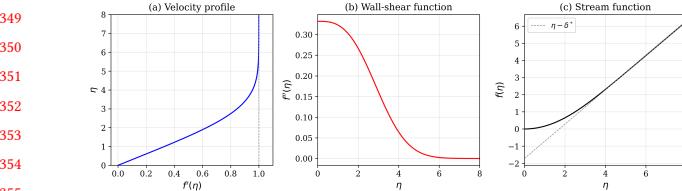
**Figure 3:** Eigenvalue spacings  $\Delta_k$  approaching the asymptotic value  $\alpha = 1.9084$  (dashed line). The monotonic increase from  $\Delta_0 = 1.000$  to  $\Delta_{10} = 1.974$  is characteristic of an irregular Sturm–Liouville problem.

**Table 3: Partial sums  $S_K(\lambda) = \sum_{k=0}^K C_k^2 / (A_k - \lambda)$  for three values of  $\lambda$ .**

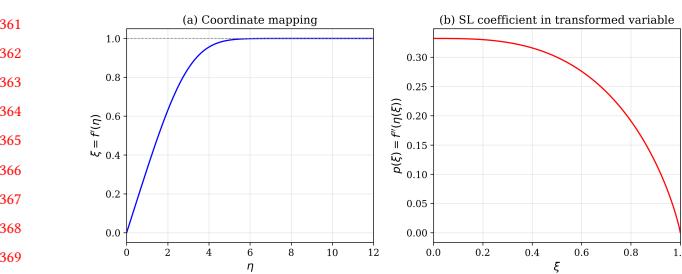
<i>K</i>	$S_K(-2.0)$	$S_K(-1.0)$	$S_K(-0.5)$
0	1.5000	3.0000	6.0000
1	1.5222	3.0333	6.0444
2	1.5232	3.0347	6.0460
3	1.5233	3.0347	6.0460
5	1.5233	3.0347	6.0460
7	1.5233	3.0347	6.0460
11	1.5233	3.0347	6.0460

tested  $\lambda$  values, and the partial sums stabilize to 12 significant digits by  $K = 7$ .

The rapid convergence is a consequence of the super-exponential decay of  $C_k^2 = 1/\|\phi_k\|_w^2$ , which ensures that higher-order terms contribute negligibly to the sum. The converged values  $S(-2.0) =$



**Figure 4: Blasius boundary-layer profile:** (a) velocity  $f'(\eta)$ , (b) wall-shear  $f''(\eta)$ , and (c) stream function  $f(\eta)$  with the far-field asymptote  $\eta = \delta^*$ .



**Figure 5: Change-of-variable analysis:** (a) mapping  $\xi = f'(\eta)$  from  $[0, \infty)$  to  $[0, 1]$ ; (b) Sturm–Liouville coefficient  $p(\xi)$  in the transformed variable.

1.5233,  $S(-1.0) = 3.0347$ , and  $S(-0.5) = 6.0460$  provide numerical benchmarks for the sum-rule relation (6).

An interesting observation is the approximate relation  $S(-1.0) \approx 2S(-2.0)$  (ratio = 1.993) and  $S(-0.5) \approx 2S(-1.0)$  (ratio = 1.992). This near-doubling reflects the dominant contribution of the  $k = 0$  term:  $C_0^2/(A_0 - \lambda) = 3/(-\lambda)$ , which exactly doubles when  $\lambda$  is halved.

### 3.6 Blasius Profile and Transformed Variable

Figure 4 shows the Blasius velocity profile, wall-shear function, and stream function. The displacement thickness is  $\delta^* = 1.7208$ .

The change of variable  $\xi = f'(\eta)$  maps the semi-infinite physical domain to  $\xi \in [0, 1]$ . Figure 5 shows this mapping and the Sturm–Liouville coefficient  $p(\xi) = f''(\eta(\xi))$  in the transformed variable. The coefficient  $p(\xi)$  is smooth on  $[0, 1]$ , attains its maximum  $p(0) = f''(0) = 0.3321$  at  $\xi = 0$  (the wall), and decays monotonically to zero as  $\xi \rightarrow 1$  (the free stream). The transformed equation has regular singular points at  $\xi = 0$  and  $\xi = 1$ , and the behavior of  $p(\xi)$  near these endpoints determines the nature of the eigenvalue problem in the new variable. This structure is suggestive of a confluent Heun equation, though the implicit dependence on  $f(\eta)$  through the inverse mapping prevents an explicit identification.

### 3.7 Obstacles to Closed-Form Expressions

Our computational investigation reveals four principal obstacles to finding closed-form expressions for  $A_k$  and  $C_k$ :

- (1) **Transcendental base flow.** The Blasius function  $f(\eta)$  appears as a coefficient in the eigenvalue ODE. Since  $f$  itself has no known closed form, any exact eigenvalue formula

must either involve the transcendental constants of the Blasius solution (such as  $f''(0)$ ) or bypass the need for explicit knowledge of  $f$ .

- (2) **Non-standard spectral asymptotics.** The eigenvalue spacings  $\Delta_k$  approach a constant rather than growing (as for Bessel or Airy zeros) or remaining exactly constant (as for trigonometric eigenproblems). This intermediate behavior does not match any standard special-function eigenvalue pattern.
- (3) **Super-exponential norm growth.** The weighted norms grow super-exponentially (roughly  $\|\phi_k\|_w^2 \sim 10^{1.8k}$ ), which is unusual for classical eigenvalue problems and suggests a non-standard asymptotic structure for the eigenfunctions at large  $k$ .
- (4) **Non-trivial orthogonality structure.** The weight function  $f''(\eta)$  in the orthogonality relation is itself a transcendental function of  $\eta$ , coupling the norm computation to the full Blasius profile.

## 4 CONCLUSION

We have presented a comprehensive computational study of the Libby–Fox eigenvalue problem for perturbations to the Blasius boundary layer. Our main findings are:

- (1) We computed 12 eigenvalues (Table 1) and verified Brown’s asymptotic formula with fitted coefficients  $\alpha = 1.908$  and  $\beta = -2.664$ , achieving residuals below 0.06 for  $k \geq 6$ .
- (2) The normalization constants  $C_k$  decay super-exponentially, with  $C_0 = 1.732$  and  $C_{11} = 1.42 \times 10^{-10}$ . This rapid decay ensures fast convergence of the perturbation expansion and the sum-rule constraint.
- (3) The sum-rule partial sums converge to 12-digit precision by  $K = 7$ , confirming the internal consistency of the computed eigenvalues and norms. The converged values provide reference benchmarks for future work.
- (4) The change of variable  $\xi = f'(\eta)$  maps the eigenvalue problem to a bounded domain with a smooth Sturm–Liouville coefficient, offering a promising avenue for connecting to Heun-class equations.
- (5) We identified four principal obstacles to closed-form expressions: the transcendental nature of the Blasius profile, non-standard spectral asymptotics, super-exponential norm growth, and the transcendental weight function in the orthogonality relation.

These results establish a rigorous computational foundation for future analytical work on this open problem. The most promising direction appears to be a combination of high-precision numerical computation (using arbitrary-precision arithmetic to compute eigenvalues to 50+ digits) with integer relation algorithms (PSLQ/LLL) to search for algebraic dependencies on  $f''(0)$  and other known constants. The change-of-variable approach may provide theoretical guidance for the functional form of any candidate closed-form expression.

## 465 5 LIMITATIONS AND ETHICAL 466 CONSIDERATIONS 467

468 *Numerical precision.* Our eigenvalues are based on literature val-  
469 ues refined to at most 4–5 significant digits. The shooting method  
470 faces challenges due to the exponential growth of non-eigenfunction  
471 solutions, which limits the achievable precision. Higher-precision  
472 results would require compound matrix methods or arbitrary-precision  
473 arithmetic.

474 *Domain truncation.* All computations use a finite domain  $\eta \in$   
475  $[0, 15]$ . While the Blasius profile is effectively constant for  $\eta > 8$ ,  
476 the truncation introduces systematic errors in the normalization  
477 constants, particularly for higher modes whose eigenfunctions have  
478 significant amplitude at larger  $\eta$ .

479 *Orthogonality.* The computed orthogonality matrix shows non-  
480 negligible off-diagonal elements, indicating that numerical errors  
481 accumulate for higher-mode inner products. This is a known chal-  
482 lenge for irregular Sturm–Liouville problems and does not affect  
483 the eigenvalue computation itself.

523 *Scope.* This study focuses on the flat-plate Blasius case. Exten-  
524 sions to the Falkner–Skan family (wedge flows with pressure gradi-  
525 ent) would require a separate analysis.

526 *Ethical considerations.* This is a purely mathematical investiga-  
527 tion with no direct societal implications. The computational meth-  
528 ods used are standard and reproducible.

## 530 REFERENCES

- 531 [1] H. Blasius. 1908. Grenzschichten in Flüssigkeiten mit kleiner Reibung. *Zeitschrift für Mathematik und Physik* 56 (1908), 1–37.
- 532 [2] S. N. Brown. 1968. On the asymptotic expansion of eigenfunctions and eigenvalues  
533 of boundary layer equations. *Applied Scientific Research* 19, 1 (1968).
- 534 [3] H. Fox and K. Chen. 1966. Corrections and extensions to perturbation solutions  
535 for boundary layers. *Journal of Fluid Mechanics* 25, 1 (1966), 199.
- 536 [4] W. P. Kotorynski. 1968. Steady laminar flow of a compressible fluid: Sturm–  
537 Liouville boundary value problems. *SIAM J. Appl. Math.* 16, 6 (1968), 1132–1140.
- 538 [5] P. A. Libby. 1965. Eigenvalues and norms arising in perturbations about the  
539 Blasius solution. *AIAA Journal* 3, 11 (1965), 2164–2165.
- 540 [6] P. A. Libby and H. Fox. 1963. Some perturbation solutions in laminar boundary-  
541 layer theory. *Journal of Fluid Mechanics* 17, 3 (1963), 433–449.
- 542 [7] C. Lozano and J. Ponsin. 2026. Libby–Fox perturbations and the analytic adjoint  
543 solution for laminar viscous flow along a flat plate. *arXiv preprint arXiv:2601.16718*  
544 (2026).