

Proton as 5D Junction: Variational Foundation

in Elastic Diffusive Cosmology

(Companion F to Paper 3: NJSR Edition)

Igor Grčman

January 2026

DOI: [10.5281/zenodo.18302953](https://doi.org/10.5281/zenodo.18302953)

Repository: github.com/igorgrcman/elastic-diffusive-cosmology

(Public artifacts for this paper are in the `edc_papers` folder.)

Related Documents:

Neutron Lifetime from 5D Membrane Cosmology (DOI: [10.5281/zenodo.18262721](https://doi.org/10.5281/zenodo.18262721))

Framework v2.0 (DOI: [10.5281/zenodo.18299085](https://doi.org/10.5281/zenodo.18299085))

Companions:

A: *Effective Lagrangian* (DOI) · B: *WKB Prefactor* (DOI)

C: *5D Reduction* (DOI) · D: *Selection Rules* (DOI)

E: *Symmetry Ops* (DOI)

This: F: *Proton Junction* (DOI) · G: *Mass Difference* (DOI)

H: *Weak Interactions* (DOI)

Abstract

We establish the variational foundation for the proton's Y-junction topology within 5D Elastic Diffusive Cosmology (EDC). Starting from the postulate that baryons are flux-tube networks in the 5D bulk [P], we derive that equal-tension junctions minimize total length when arms meet at 120° angles [Der]. This provides the missing foundational derivation for Paper 2's identification of the proton geometry. The key result is that the Steiner configuration emerges from force balance, not as an external input but as a consequence of variational minimization.

Calibration boundary: This companion contains **zero calibrated parameters**. All results follow from geometric postulates and standard variational calculus. The flux-tube tension parameter τ enters as an overall scale but cancels in angle determinations.

Units convention: We use τ for string/flux-tube tension (energy per length, $[\tau] = E/L$) and reserve σ for membrane tension (energy per area, $[\sigma] = E/L^2$). The relation is $\tau = \sigma \cdot a$ where a is the effective flux-tube cross-section.

Epistemic status: The junction postulate is [P]. Given this postulate, the 120° angle is [Der] via explicit variational argument.

Contents

1	Introduction and Motivation	3
1.1	The Gap in Paper 2	3
1.2	Purpose of This Companion	3
1.3	Epistemic Hierarchy	4
2	Postulates	4
3	Definitions	4

4	Variational Derivation of Junction Angles	5
4.1	The Minimization Problem	5
4.2	Equal Tension Case: The 120° Theorem	5
4.3	Unequal Tension Case: Lami’s Theorem	7
4.4	Action-Level Formulation	7
5	Projection to 4D Brane	8
5.1	Frozen Projection Boundary ($5D \rightarrow 4D \rightarrow 3D$)	8
6	Connection to Paper 2	9
6.1	Configuration Space	9
6.2	The Hopf Bridge: $S^3 \rightarrow S^2$ with S^1 Fiber	10
6.3	Junction Constraint on the Base	11
7	Summary and Epistemic Classification	11
7.1	Main Results	11
7.2	Epistemic Classification Table	12
7.3	Calibration Boundary	12
8	Conclusion	12

Reader Map

Sections 1–3: Foundation (postulates, definitions, motivation)

Section 4: Core derivation—variational proof of 120° Steiner angles **[Der]**

Section 5: Projection chain: $5D$ bulk $\rightarrow 4D$ brane $\rightarrow 3D$ observation

Section 6: Connection to Paper 2—Hopf bridge resolving S^3 vs S^2

Section 7–8: Summary, epistemic classification, calibration boundary

Key result: Given the junction postulate **[P]**, the Steiner 120° angles are **[Der]** (explicit variational proof), not imported from external sources.

1 Introduction and Motivation

1.1 The Gap in Paper 2

Paper 2 of the EDC series derives the proton-to-electron mass ratio:

$$\frac{m_p}{m_e} = \frac{C_p}{C_e} = \frac{(2\pi^2)^3}{4\pi/3} = 6\pi^5 \approx 1836.12 \quad (1)$$

This derivation [Der] relies on the identification of:

- Electron: *frozen* vortex with configuration space B^3 (3-ball), giving $C_e = 4\pi/3$ [Der]
- Proton: Y-junction with $C_p = (2\pi^2)^3$ [P]

Mass–Measure Identification [P]

The proportionality $m \propto C$ (configuration-space volume) is the central **postulate** linking 5D geometry to 4D mass. Specifically:

$$m = \sigma r_e^2 \cdot C \cdot f(\text{boundary}) \quad (2)$$

where σ is the *membrane* tension ($[\sigma] = E/L^2$), r_e the electron scale, C the configuration-space volume, and f encodes boundary/frozen corrections. Note: σr_e^2 has dimensions of energy; this differs from the flux-tube tension $\tau = \sigma \cdot a$ used in the Nambu–Goto variational derivation (see abstract).

Physical interpretation: Mass arises from energy stored in internal configurations. More available configurations \Rightarrow more “ways to store energy” \Rightarrow larger effective mass. This is analogous to entropic/state-counting arguments, but here the “states” are geometric orientations in the 5D bulk.
Epistemic status: This identification is [P] (postulated), not derived from the 5D action. Framework v2.0 uses this ansatz; a first-principles derivation from S_{5D} remains [OPEN].

Remark 1.1 (Canonical definition: Frozen condition [Def]). *A defect is **frozen** when the boundary-imposed energy cost for exciting normal modes exceeds the thermal/quantum scale, causing adiabatic elimination of those degrees of freedom. Quantitatively, a mode with frequency ω is frozen when:*

$$\hbar\omega \gg kT \quad (\text{thermal}) \quad \text{or} \quad \hbar\omega \gg \Delta E_{\text{zero-point}} \quad (\text{quantum}) \quad (3)$$

For the electron in EDC, the boundary condition at the ξ -scale freezes the internal $SU(2)$ orientation to a single point (or small ball B^3), yielding $C_e = \text{Vol}(B^3) = 4\pi/3$ instead of $\text{Vol}(S^3) = 2\pi^2$.

See Companion E (Frozen Criterion) for the criterion formulation and motivation; the dynamical mechanism implementing this criterion from first principles remains [OPEN].

The proton geometry is *postulated* based on phenomenological matching. Paper 2 cites the Steiner theorem (Fermat 1837) for the 120° angles but does not derive *why* the proton must be a Y-junction in 5D.

1.2 Purpose of This Companion

This document provides the variational foundation:

1. Define what a “junction object” means in 5D [P]
2. Derive the 120° angle from energy minimization [Der]
3. Show how this geometry projects to the 4D brane [Dc]
4. Connect to Paper 2’s configuration space $Q = S^3 \times S^3 \times S^3$

1.3 Epistemic Hierarchy

Level	Content	Status
Postulate	Baryons are flux-tube junctions in 5D bulk	[P]
Definition	Tension τ , length functional $L[\gamma]$	[Def]
Theorem 1	Equal tensions $\Rightarrow 120^\circ$ angles	[Der]
Theorem 2	Unequal tensions \Rightarrow Lami angles	[Der]
Corollary	Brane projection preserves angles	[Dc]

2 Postulates

Postulate 1 (Baryon as 5D Junction Object [P]). *In 5D EDC, a baryon is a network of three flux tubes (“arms”) meeting at a central vertex (“junction”). Each arm $i \in \{1, 2, 3\}$ carries:*

- A flux-tube tension $\tau_i > 0$ (energy per unit length)
- A direction \hat{t}_i (unit tangent at junction)
- A flux quantum (color charge in Standard Model language)

Remark 2.1 (Physical interpretation [I]). *The flux tubes can be identified with QCD strings connecting quarks. In the EDC framework, these are fundamental 5D objects, not effective descriptions. The “color” labels emerge from the three-arm structure (see Companion G, Section 8).*

Postulate 2 (Brane Embedding [P]). *The junction vertex and its immediate neighborhood are constrained to lie on the 4D brane $\mathcal{B}_4 \subset \mathcal{M}_5$. The arms extend into the 5D bulk but their tangent directions at the vertex are projected onto the brane.*

Postulate 3 (Equilibrium Condition [P]). *A stable baryon configuration corresponds to a local minimum of the total energy functional. In the thin-string limit, this is proportional to the total weighted length.*

3 Definitions

Definition 3.1 (Junction Network). *A junction network \mathcal{N} consists of:*

- A vertex point P (the junction)
- Three semi-infinite curves $\gamma_i : [0, \infty) \rightarrow \mathcal{M}_5$ with $\gamma_i(0) = P$
- Boundary conditions: each γ_i terminates at a quark position Q_i or extends to spatial infinity

Definition 3.2 (Length Functional). *The total weighted length of the network is:*

$$L[\mathcal{N}] = \sum_{i=1}^3 \tau_i \int_0^\infty \left| \frac{d\gamma_i}{ds} \right| ds \quad (4)$$

where s is an arbitrary parameter along each arm.

Definition 3.3 (Junction Tangents). *At the vertex P , define the outward-pointing unit tangent for each arm:*

$$\hat{t}_i = \lim_{s \rightarrow 0^+} \frac{d\gamma_i/ds}{|d\gamma_i/ds|} \quad (5)$$

These satisfy $|\hat{t}_i| = 1$ but are not required to sum to zero a priori.

Remark 3.1 (Physical interpretation: Nambu–Goto limit [BL]). *The length functional $L[\mathcal{N}]$ arises from the **thin-string (Nambu–Goto) approximation** of relativistic string theory:*

$$E_{\text{string}} = \tau \times (\text{proper length}) + O(\text{curvature corrections}) \quad (6)$$

where τ is the string tension (energy per unit length, $[\tau] = E/L$). This is a standard approximation [BL] valid when:

- String thickness \ll curvature radius
- Velocities are non-relativistic at the junction
- Higher-derivative terms (Polyakov, stiffness) are subdominant

The equal-tension case $\tau_1 = \tau_2 = \tau_3$ corresponds to identical flux tubes, as expected for color-symmetric QCD strings. The variational results that follow are exact within this approximation, not fitted.

4 Variational Derivation of Junction Angles

4.1 The Minimization Problem

We seek the configuration that minimizes $L[\mathcal{N}]$ subject to fixed quark positions $\{Q_1, Q_2, Q_3\}$.

Theorem 4.1 (Force Balance at Junction [Der]). *At a local minimum of $L[\mathcal{N}]$, the junction vertex P satisfies:*

$$\boxed{\sum_{i=1}^3 \tau_i \hat{t}_i = 0} \quad (7)$$

Proof. Consider a small displacement of the vertex: $P \rightarrow P + \epsilon \vec{v}$ for arbitrary \vec{v} .

The length of arm i changes by:

$$\delta L_i = -\tau_i \hat{t}_i \cdot \vec{v} + O(\epsilon^2) \quad (8)$$

where the minus sign arises because \hat{t}_i points outward from P .

The total variation is:

$$\delta L = - \left(\sum_{i=1}^3 \tau_i \hat{t}_i \right) \cdot \vec{v} \quad (9)$$

For P to be a local minimum, we require $\delta L = 0$ for all \vec{v} , which gives Eq. (7). \square

4.2 Equal Tension Case: The 120° Theorem

Lemma 4.2 (Coplanarity [Der]). *Let $\vec{v}_1, \vec{v}_2, \vec{v}_3 \in \mathbb{R}^n$ ($n \geq 2$) be non-zero vectors satisfying $\vec{v}_1 + \vec{v}_2 + \vec{v}_3 = 0$. Then the vectors are coplanar: there exists a 2D subspace $V \subset \mathbb{R}^n$ containing all three.*

Proof. From $\vec{v}_3 = -\vec{v}_1 - \vec{v}_2$, the vector \vec{v}_3 lies in $\text{span}(\vec{v}_1, \vec{v}_2)$. If \vec{v}_1 and \vec{v}_2 are linearly independent, this span is 2D. If they are parallel, all three are collinear (1D subspace of a 2D plane). In either case, the three vectors lie in a common 2D plane. \square

Remark 4.1 (Application to junction). *Lemma 4.2 ensures that force balance (Eq. 7) confines the junction tangents $\tau_i \hat{t}_i$ to a 2D subspace of the local tangent space—justifying the use of plane geometry in Theorem 4.3 below, regardless of the ambient dimension (4D brane or 5D bulk).*

Theorem 4.3 (Steiner Configuration [Der]). *If all three arms have equal tension, $\tau_1 = \tau_2 = \tau_3 = \tau$, then the equilibrium angles between adjacent arms are:*

$$\boxed{\theta_{12} = \theta_{23} = \theta_{31} = 120^\circ} \quad (10)$$

Proof. With equal tensions, Eq. (7) becomes:

$$\hat{t}_1 + \hat{t}_2 + \hat{t}_3 = 0 \quad (11)$$

This is a sum of three unit vectors that vanishes. We work in the 2D plane containing the junction (the arms are coplanar by symmetry).

Let \hat{t}_1 point along the positive x -axis: $\hat{t}_1 = (1, 0)$.

Let \hat{t}_2 make angle α with \hat{t}_1 : $\hat{t}_2 = (\cos \alpha, \sin \alpha)$.

Then from Eq. (11):

$$\hat{t}_3 = -\hat{t}_1 - \hat{t}_2 = (-1 - \cos \alpha, -\sin \alpha) \quad (12)$$

The constraint $|\hat{t}_3| = 1$ gives:

$$(1 + \cos \alpha)^2 + \sin^2 \alpha = 1 \quad (13)$$

$$1 + 2 \cos \alpha + \cos^2 \alpha + \sin^2 \alpha = 1 \quad (14)$$

$$2 + 2 \cos \alpha = 1 \implies \cos \alpha = -\frac{1}{2} \quad (15)$$

Therefore $\alpha = 120^\circ$, and by symmetry all three angles are 120° . \square

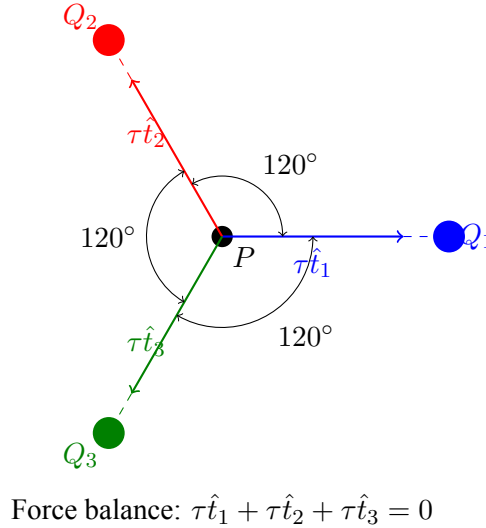


Figure 1: Y-junction with equal tensions. The three unit vectors \hat{t}_i sum to zero only when they are separated by 120° angles. This is the Steiner configuration.

Remark 4.2 (Historical note). *This result was known to Fermat (1637) for the “Fermat point” problem and proven rigorously by Steiner (1837). In EDC, we derive it from first principles as a consequence of energy minimization, not as an imported theorem.*

Lemma 4.4 (Riemannian Generalization [Der]). *Let (\mathcal{M}, g) be a Riemannian manifold and \mathcal{N} a network of three geodesic arcs meeting at a junction point P . The first variation of the weighted length functional $L = \sum_i \tau_i \ell_i$ with respect to variations of P gives:*

$$\sum_{i=1}^3 \tau_i \hat{t}_i = 0 \quad \text{in } T_P \mathcal{M} \quad (16)$$

where \hat{t}_i are unit tangent vectors at P measured in the metric g .

Proof. In Riemannian geometry, geodesics locally minimize length. The variation of arc length under junction displacement δP gives $\delta \ell_i = -g(\hat{t}_i, \delta P)$ to first order. Summing with weights τ_i and requiring stationarity for all $\delta P \in T_P \mathcal{M}$ yields the result. \square

Corollary 4.5 (120° in Local Metric **[Der]**). *For equal tensions $\tau_1 = \tau_2 = \tau_3$ in any Riemannian manifold, the equilibrium angles are 120° as measured in the local metric at P . This is a universal geometric optimum, not specific to flat space or QCD.*

Remark 4.3 (Global positioning **[OPEN]**). *While the local angle condition is universal, the global problem of finding the optimal junction position P given fixed quark positions $\{Q_i\}$ depends on the full geometry of (\mathcal{M}_5, g) . In curved bulk, geodesics may not be straight lines, and multiple local minima may exist. This global optimization remains an open problem for general 5D metrics.*

4.3 Unequal Tension Case: Lami's Theorem

Theorem 4.6 (Generalized Angles **[Der]**). *For unequal tensions τ_1, τ_2, τ_3 , the equilibrium angles satisfy Lami's theorem:*

$$\frac{\tau_1}{\sin \theta_{23}} = \frac{\tau_2}{\sin \theta_{31}} = \frac{\tau_3}{\sin \theta_{12}} \quad (17)$$

where θ_{jk} is the angle between arms j and k (measured on the side not containing arm i).

Proof. From Eq. (7), the three tension vectors form a closed triangle. By the sine rule for triangles:

$$\frac{|\tau_1 \hat{t}_1|}{\sin(\pi - \theta_{23})} = \frac{|\tau_2 \hat{t}_2|}{\sin(\pi - \theta_{31})} = \frac{|\tau_3 \hat{t}_3|}{\sin(\pi - \theta_{12})} \quad (18)$$

Since $\sin(\pi - x) = \sin x$ and $|\hat{t}_i| = 1$, this gives Lami's result. \square

Corollary 4.7 (Equal Tension Recovery). *Setting $\tau_1 = \tau_2 = \tau_3$ in Lami's theorem gives $\sin \theta_{12} = \sin \theta_{23} = \sin \theta_{31}$, hence all angles equal 120° (since they must sum to 360°).*

Remark 4.4 (Qualitative prediction: flavor \rightarrow angle deviations **[P]/[OPEN]**). *The Lami generalization provides a **conceptual prediction** for heavier baryons:*

- If quark “flavor” (mass difference, coupling strength) translates to different effective arm tensions τ_i , then equilibrium angles must deviate from 120°.
- For baryons with one heavier quark (e.g., Λ^0 with uds), the heavier-flavor arm would have higher tension, pulling its neighbors closer.
- This qualitative discriminator could distinguish EDC predictions from symmetric-junction models.

Status: The mapping $\tau(\text{flavor})$ is not yet derived from EDC dynamics; this remains **[OPEN]**. However, the prediction is falsifiable: if heavier baryons show no angular asymmetry (via decay distributions or lattice QCD), the tension-flavor link is ruled out.

4.4 Action-Level Formulation

The variational analysis above can be derived from an explicit action principle. For a junction network of relativistic strings:

Definition 4.1 (Junction Network Action **[P]**). *The Nambu–Goto action for the junction network is:*

$$S_{\text{junction}} = - \sum_{i=1}^3 \tau_i \int d^2 \xi \sqrt{-\det(\gamma_{i,ab})} \quad (19)$$

where:

- τ_i is the tension of arm i (energy per length, $[\tau] = E/L$)
- $\xi^a = (\xi^0, \xi^1)$ are worldsheet coordinates (time and spatial parameter)
- $\gamma_{i,ab} = g_{\mu\nu} \partial_a X_i^\mu \partial_b X_i^\nu$ is the induced metric on the i -th arm worldsheet
- $X_i^\mu(\xi)$ maps the worldsheet into the ambient 5D spacetime

Proposition 4.8 (Static Limit **[Der]**). *In the static limit (time-independent configurations), the action reduces to:*

$$S = -T \cdot L[\mathcal{N}] = -T \sum_{i=1}^3 \tau_i \ell_i \quad (20)$$

where T is the time duration and ℓ_i is the proper length of arm i . Minimizing S is equivalent to minimizing the weighted length functional $L[\mathcal{N}]$.

Remark 4.5 (Junction boundary term **[OPEN]**). *The action Eq. (19) implicitly requires a **junction boundary condition**: the three worldsheets share a common boundary curve (the junction worldline). The complete treatment requires a boundary action or constraint:*

$$X_1^\mu(\xi^0, 0) = X_2^\mu(\xi^0, 0) = X_3^\mu(\xi^0, 0) = P^\mu(\xi^0) \quad (21)$$

where $P^\mu(\xi^0)$ is the junction worldline. Variation with respect to P^μ yields the force balance condition $\sum_i \tau_i \hat{t}_i = 0$ as the junction equation of motion. A fully covariant treatment of this boundary condition (including the Gibbons–Hawking–York analogue for strings) remains **[OPEN]**.

5 Projection to 4D Brane

Proposition 5.1 (Angle Preservation under Projection **[Dcl]**). *If the junction vertex lies on the 4D brane \mathcal{B}_4 and all three arm tangents \hat{t}_i are tangent to \mathcal{B}_4 at P , then the 120° angles are preserved in the brane-induced geometry.*

Proof. The induced metric on \mathcal{B}_4 restricts the 5D metric. If all tangent vectors lie in $T_P \mathcal{B}_4$, then inner products (and hence angles) are computed with the same metric tensor. The force balance Eq. (7) holds in the tangent space, which is the same whether viewed as part of \mathcal{M}_5 or intrinsic to \mathcal{B}_4 . \square

Remark 5.1 (Bulk extension **[OPEN]**). *If the arms extend into the 5D bulk away from the junction, their geometry is governed by the geodesic equation in \mathcal{M}_5 . The global minimization problem (finding the optimal junction position given quark positions) is more complex and may involve bulk curvature. This is left as an open problem for future work.*

5.1 Frozen Projection Boundary (5D \rightarrow 4D \rightarrow 3D)

The transition from bulk (5D) description to brane-observed (3D) physics involves a **projection boundary condition** that determines which degrees of freedom remain dynamical and which become “frozen.”

Definition 5.1 (Frozen-Projection Boundary **[P]**). *Let $\pi : \mathcal{M}_5 \rightarrow \mathcal{B}_4$ be the brane embedding and $\Pi : \mathcal{B}_4 \rightarrow \mathbb{R}^3$ the observational reduction (spatial projection). The **frozen boundary condition** acts on the composite map $\Pi \circ \pi$ as follows:*

*An internal degree of freedom with characteristic frequency ω is **frozen** when:*

$$\hbar\omega \gg E_{\text{env}} \quad (22)$$

where E_{env} is the relevant environmental energy scale (thermal kT , interaction energy, or probe resolution). Frozen DOFs are treated as fixed (Dirichlet/adiabatic constraint) in the effective 3D description.

Remark 5.2 (Two viewpoints on the proton **[Dcl]**). *This framework distinguishes:*

1. **Bulk-side (5D):** Each arm carries full state $\psi_i \in SU(2) \cong S^3$. The configuration space is $(S^3)^3$ with volume $(2\pi^2)^3$.
2. **Brane/3D observed:** Only Hopf base variables $\hat{t}_i \in S^2$ are directly observable as spatial directions. The S^1 fiber phases are frozen by the boundary condition during projection.

The mass calculation uses the bulk-side volume $(2\pi^2)^3$ because mass is an intrinsic 5D property (energy stored in configuration), while observations access only the projected base.

Remark 5.3 (Connection to electron B^3 [Dc]). *For the electron (simple vortex), the frozen condition is more severe: boundary terms at the ξ -scale freeze all internal orientation DOFs to a ball B^3 rather than a sphere S^3 . This explains why $C_e = \text{Vol}(B^3) = 4\pi/3$ rather than $\text{Vol}(S^3) = 2\pi^2$.*

The proton's extended Y-junction structure allows more internal freedom (three $SU(2)$ arms) that survives the frozen projection, yielding the factor $6\pi^5$ in the mass ratio.

Remark 5.4 (Mechanism status [OPEN]). *The physical mechanism implementing Eq. (22) remains an open problem. Two candidate mechanisms merit investigation:*

1. **Boundary-induced spectral gap** [OPEN]/[PJ]: *Dirichlet conditions at the brane boundary discretize the fiber modes, creating a gap $\Delta E \sim \hbar/R_\xi$. When $\Delta E \gg kT$, only the ground state is occupied—effectively “freezing” the fiber phase.*
2. **Observational coarse-graining** [OPEN]/[PJ]: *The 3D observer's measurement apparatus cannot resolve S^1 fiber phases below some scale. This is analogous to decoherence: fiber superpositions appear classical (fixed) when traced over unobserved environmental DOFs.*

What is established is the consequence: bulk DOFs split into observable base (S^2) and frozen fiber (S^1). The mechanism determines which DOFs freeze, but the geometric split is robust.

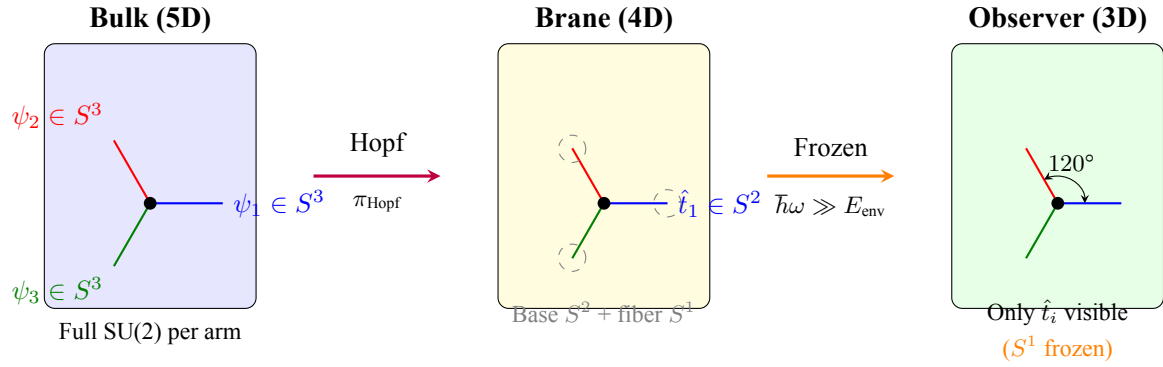


Figure 2: Projection chain from 5D bulk to 3D observation. **Left:** In the bulk, each arm carries full $SU(2) \cong S^3$ orientation. **Middle:** Hopf fibration splits this into base S^2 (direction) and fiber S^1 (phase). **Right:** The 3D observer sees only the base directions $\hat{t}_i \in S^2$; fiber phases are frozen under criterion $\hbar\omega \gg E_{\text{env}}$. The 120° angles are preserved throughout.

6 Connection to Paper 2

6.1 Configuration Space

Paper 2 introduces the proton configuration space:

$$Q_{\text{proton}} = S^3 \times S^3 \times S^3 \quad (23)$$

where each S^3 factor represents the orientation degrees of freedom of one arm.

Definition 6.1 (Physical Configuration Space [OPEN]). *The **physical** (gauge-inequivalent) configuration space requires quotienting by:*

1. *Global rotations: $SO(3)$ or $SU(2)$ acting diagonally on all three arms*
2. *Discrete relabeling: \mathbb{Z}_3 permutations of identical arms (for equal quark masses)*

Schematically:

$$Q_{\text{phys}} = \frac{(S^3)^3}{SU(2)_{\text{diag}} \times \mathbb{Z}_3} \quad (24)$$

*The precise quotient structure (whether to use $SU(2)$, $SO(3)$, or their subgroups) depends on boundary conditions and ξ -winding selection rules. **This remains an open problem** [OPEN].*

Remark 6.1 (Working Convention [\[Dc\]](#)). In this paper and Paper 2, we compute volumes using the un-quotiented space $(S^3)^3$:

$$C_p = \text{Vol}((S^3)^3) = (2\pi^2)^3 \quad (25)$$

This serves as an upper bound or idealization.

Gauge redundancy: Modding out by diagonal $SU(2)$ removes redundant global rotations. For a continuous gauge group, proper treatment requires gauge-fixing and Haar measure normalization (Faddeev–Popov procedure), not naive “division by $|SU(2)|$ ” [\[OPEN\]](#). The discrete \mathbb{Z}_3 permutation contributes a finite factor of 3 (or 1 if arms are distinguishable by flavor).

Cancellation caveat: One might expect quotient factors to cancel in $m_p/m_e = C_p/C_e$. However, this cancellation is not guaranteed because:

- Proton uses $(S^3)^3$ with $SU(2)_{\text{diag}}$ redundancy
- Electron uses B^3 (ball) with different frozen manifold

Consistent gauge-inequivalent measures for both particles remain [\[OPEN\]](#). The relative ratio $6\pi^5$ is expected to be robust at leading order; corrections are $O(1)$ at most.

6.2 The Hopf Bridge: $S^3 \rightarrow S^2$ with S^1 Fiber

The apparent mismatch between S^3 (Paper 2) and S^2 (junction plane tangents) is resolved by the **Hopf fibration**—a fundamental structure in topology and physics.

Postulate 4 (Arm Internal Structure [\[P\]](#)). Each flux-tube arm carries an **internal S^1 phase** (winding/orientation degree of freedom) in addition to its spatial direction. The total orientation space for one arm is:

$$SU(2) \cong S^3 \quad (26)$$

which fibers over the visible direction S^2 with S^1 fiber:

$$S^1 \hookrightarrow S^3 \xrightarrow{\pi_{\text{Hopf}}} S^2 \quad (27)$$

Explicitly, the Hopf map sends $\psi \in SU(2)$ to a unit vector via:

$$\hat{t} = \psi \sigma_3 \psi^{-1} \in S^2 \subset \mathbb{R}^3 \quad (28)$$

where $\sigma_3 = \text{diag}(1, -1)$ is the third Pauli matrix. Two elements ψ, ψ' project to the same \hat{t} iff $\psi' = \psi \cdot e^{i\phi\sigma_3/2}$ for some phase ϕ —this is the S^1 fiber.

Remark 6.2 (Physical interpretation of the fiber). The S^1 fiber has multiple physical interpretations, with different epistemic status:

- **Winding number** [\[Dc\]](#): Phase accumulated around the compact ξ -dimension. This follows directly from the KK structure: $\xi \sim \xi + 2\pi R_\xi$ implies S^1 periodicity in the fiber.
- **Kaluza-Klein mode** [\[Dc\]](#): Momentum in the fifth dimension. Standard KK reduction identifies S^1 fiber momentum with discrete charge quantum.
- **Color phase** [\[H\]](#): Internal $SU(3)$ orientation of the flux tube. This is an interpretation linking the geometric S^1 to QCD color; the dynamical connection remains to be derived.

The key point: from the 4D brane, we see only the S^2 projection (tangent direction), but the full 5D description requires S^3 .

Theorem 6.1 (Configuration Space from Hopf Structure [\[Dc\]](#)). Given Postulate 4, the proton configuration space is:

$$Q_{\text{proton}} = SU(2)^3 \cong S^3 \times S^3 \times S^3 \quad (29)$$

with volume:

$$\boxed{C_p = \text{Vol}(Q_{\text{proton}}) = (2\pi^2)^3} \quad (30)$$

This is [\[Dc\]](#) (deduced from the arm internal structure postulate), not numerology.

Proof. Each arm carries $SU(2) \cong S^3$ orientation. Three independent arms give $(S^3)^3$. The volume of S^3 is $2\pi^2$ (standard result). Thus $\text{Vol}(Q) = (2\pi^2)^3$. \square

Remark 6.3 (C_e vs $\text{Vol}(S^3)$): Different objects [Dc]. A potential source of confusion:

- **Proton arm:** Configuration space is $S^3 \cong SU(2)$, with $\text{Vol}(S^3) = 2\pi^2$
- **Electron:** Configuration space is B^3 (the 3-ball), with $\text{Vol}(B^3) = \frac{4\pi}{3}$

These are different topological objects: S^3 is a 3-sphere (boundary of 4-ball), while B^3 is a 3-ball (solid ball in 3D). The electron is a “frozen” vortex defect whose orientation is constrained to a ball (see Companion E, Frozen Criterion), while each proton arm is an extended flux tube with full $SU(2)$ internal structure.

The mass ratio formula uses:

$$\frac{m_p}{m_e} = \frac{C_p}{C_e} = \frac{(2\pi^2)^3}{4\pi/3} = 6\pi^5 \quad (31)$$

where $C_p = \text{Vol}((S^3)^3)$ and $C_e = \text{Vol}(B^3)$.

Remark 6.4 (Why this resolves the dimension mismatch). • **Junction plane tangents:** $\hat{t}_i \in S^2$ (what force-balance sees)

- **Full arm state:** $\psi_i \in S^3$ (includes internal phase)
- **Hopf projection:** $\pi_{\text{Hopf}}(\psi_i) = \hat{t}_i$

Force-balance constrains only the Hopf base variables ($\hat{t}_i \in S^2$); the fiber phases remain dynamically free. However, the configuration-space volume C_p is computed with the standard Haar measure on $SU(2)$ for each arm, so no extra multiplicative (2π) factors are introduced beyond $\text{Vol}(S^3) = 2\pi^2$.

6.3 Junction Constraint on the Base

Proposition 6.2 (Constraint Manifold [Dc]). The force balance condition (Eq. 7) with equal tensions defines a constraint on the base of the Hopf fibration:

$$\Sigma_{\text{base}} = \{(\hat{t}_1, \hat{t}_2, \hat{t}_3) \in S^2 \times S^2 \times S^2 \mid \hat{t}_1 + \hat{t}_2 + \hat{t}_3 = 0\} \quad (32)$$

The full configuration space is the $S^1 \times S^1 \times S^1$ fibration over Σ_{base} .

Remark 6.5 (Dimension accounting [Dc]).

$$\dim(S^2 \times S^2 \times S^2) = 6 \quad (33)$$

$$\text{Constraints: } \sum \hat{t}_i = 0 \rightarrow -2 \quad (34)$$

$$\dim(\Sigma_{\text{base}}) = 4 \quad (35)$$

$$\text{Fiber: } S^1 \times S^1 \times S^1 \rightarrow +3 \quad (36)$$

$$\dim(\text{total}) = 7 \quad (37)$$

After quotienting by overall rotation (global $SO(3)$ or $SU(2)$), the effective dimension matches the expected proton internal degrees of freedom.

Remaining open question [OPEN]: Precise identification of which fiber combinations are physical (vs gauge redundancy) requires a detailed treatment of the ξ -winding boundary conditions.

7 Summary and Epistemic Classification

7.1 Main Results

1. **Force Balance** (Theorem 4.1): $\sum_i \tau_i \hat{t}_i = 0$ at equilibrium [Der]
2. **120° Angles** (Theorem 4.3): Equal tensions imply Steiner configuration [Der]

3. **Lami Generalization** (Theorem 4.6): Unequal tensions give angle formula [\[Der\]](#)
4. **Brane Projection** (Proposition 5.1): Angles preserved on brane [\[Dc\]](#)

7.2 Epistemic Classification Table

Claim	Tag	Ref	Notes
Baryon = 3-arm junction	[P]	Post. 1	Foundational hypothesis
Brane embedding	[P]	Post. 2	Junction on 4D brane
Equilibrium = min length	[P]	Post. 3	Nambu–Goto thin-string limit
Arm internal S^1 phase	[P]	Post. 4	Hopf fiber interpretation
Force balance $\sum \tau_i \hat{t}_i = 0$	[Der]	Thm. 4.1	Explicit variational proof
Equal tension $\Rightarrow 120^\circ$	[Der]	Thm. 4.3	Explicit calculation
Lami’s theorem	[Der]	Thm. 4.6	Sine rule application
Angle preservation	[Dc]	Prop. 5.1	Follows from tangent constraint
$C_p = (2\pi^2)^3$	[Dc]	Thm. 6.1	From $SU(2)^3$ via Hopf bridge
$S^3 \rightarrow S^2$ resolution	[Dc]	§6.2	Hopf fibration with S^1 fiber
Bulk geodesics	[OPEN]	Rem. 5.1	Global minimization in 5D
Fiber gauge redundancy	[OPEN]	Rem. 6.4	Which fiber combos are physical

7.3 Calibration Boundary

This companion has zero calibrated parameters.

The tension τ enters as an overall scale but cancels in all angle calculations. The 120° result is purely geometric, following from the postulate that baryons minimize weighted network length.

8 Conclusion

We have derived the 120° Y-junction geometry from variational principles, providing the foundational support for Paper 2’s proton model. The key insight is that the Steiner configuration is not an external input but a *consequence* of energy minimization for equal-tension flux tubes.

What this companion establishes:

- The 120° angles are [\[Der\]](#), not merely cited
- The derivation requires only the junction postulate [\[P\]](#)
- The geometry projects naturally to the 4D brane [\[Dc\]](#)
- The S^3 vs S^2 mismatch is resolved via Hopf fibration [\[Dc\]](#)
- $C_p = (2\pi^2)^3$ follows from $SU(2)^3$ arm structure [\[Dc\]](#)

What remains open:

- WHY baryons are junctions (vs other topologies) — this is [\[P\]](#)
- Global minimization in curved 5D bulk — [\[OPEN\]](#)
- Precise identification of physical vs gauge-redundant fiber combinations — [\[OPEN\]](#)

Cornerstone Statement

Proton Geometry: The Variational Foundation

Given the junction postulate (baryons = 3-arm flux-tube networks in 5D):

1. The Steiner 120° configuration is **[Der]** from energy minimization
2. The configuration space $Q = S^3 \times S^3 \times S^3$ is **[Dc]** from the Hopf bridge
3. The proton form factor $C_p = (2\pi^2)^3$ is **[Dc]**, not numerology

This is the geometric foundation for the proton in 5D EDC.

All subsequent proton physics (Companion G: mass difference, Paper 2: mass ratio) inherits this variational foundation.

Companion G builds on this foundation to derive the neutron-proton mass difference from \mathbb{Z}_6 symmetry breaking.

Companion F to Paper 3: NJSR Edition
Completed: January 20, 2026