

# PyQBench: a Python library for benchmarking gate-based quantum computers Supplemental material

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## 1. Existing benchmarking methodologies and software

Unsurprisingly, PyQBench is not the only software package for benchmarking gate-based devices. While we believe that our approach has significant benefits over other benchmarking techniques, for completeness, in this section we discuss some of the currently available similar software.

Probably the simplest benchmarking method one could devise is simply running known algorithms and comparing outputs with the expected ones. Analyzing the frequency of the correct outputs, or the deviation between actual and expected outputs distribution provides then a metric of the performance of a given device. Libraries such as Munich Quantum Toolkit (MQT) [1, 2] or SupermarQ [3, 4] contain benchmarks leveraging multiple algorithms, such as Shor’s algorithm or Grover’s algorithm. Despite being intuitive and easily interpretable, such benchmarks may have some problems. Most importantly, they assess the usefulness of a quantum device only for a very particular algorithm, and it might be hard to extrapolate their results to other algorithms and applications. For instance, the inability of a device to consistently find factorizations using Shor’s algorithms does not tell anything about its usefulness in Variational Quantum Algorithm’s.

Another possible approach to benchmarking quantum computers is randomized benchmarking. In this approach, one samples circuits to be run from some predefined set of gates (e.g. from the Clifford group) and tests how much the output distribution obtained from the device running these circuits differs from the ideal one. It is also common to concatenate randomly chosen circuits with their inverses (which should yield the identity

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circuit) and run those concatenated circuits on the device. Libraries implementing this approach include Qiskit [5] or PyQuil [6]. Another, equally popular, benchmarking method is quantum tomography [7]. Implementation of quantum tomography for benchmarking NISQ devices can be found in [8].

In [9], the authors evaluated several IBM-Q machines using seven benchmarks taking into account the errors and execution time.

QASMBench [10] is one of the first benchmark suites aiming at evaluating NISQ devices using quantum applications from a broad range of domains, mainly using an approach based on fidelity estimation. The benchmark presented in the paper compares the fidelity of execution among the IBM-Q machines, the IonQ QPU and the Rigetti Aspen M-1 system.

Another quantity used for benchmarking NISQ devices is quantum volume. The quantum volume characterizes the capacity of a device for solving computational problems. It takes into account multiple factors like the number of qubits, connectivity and measurement errors. The Qiskit library allows one to measure the quantum volume of a device by using its `qiskit.ignis.verification.quantum_volume`. Other implementations of Quantum Volume can be found as well, see e.g. [11].

We should also mention cross-entropy benchmarking [12], which was utilized in validation of the Sycamore-53 QPU supremacy experiments [13]. In this approach, the quality of an algorithm implemented on the QPU is measured by calculating the cross entropy of bit-strings actually sampled from the QPU, compared to ideal bitstrings.

Finally, it is worth pointing out there is an ongoing effort towards standardization of benchmarking of quantum computers. In [14], the authors present plans for designing a benchmarking suite based on measuring a set of standardized key performance indicators (KPIs).

## 2. Mathematical preliminaries

Let  $\mathcal{M}_{d_1, d_2}$  be the set of all matrices of dimension  $d_1 \times d_2$  over the field  $\mathbb{C}$ . For simplicity, square matrices will be denoted by  $\mathcal{M}_d$ . By  $\Omega_d$ , we will denote the set of quantum states, that is positive semidefinite operators having trace equal to one. The subset of  $\mathcal{M}_d$  consisting of unitary matrices will be denoted by  $\mathcal{U}_d$ , while its subgroup of diagonal unitary operators will be denoted by  $\mathcal{DU}_d$ .

We will also need a linear mapping transforming  $\mathcal{M}_{d_1}$  into  $\mathcal{M}_{d_2}$ , which will be denoted

$$\Phi : \mathcal{M}_{d_1} \rightarrow \mathcal{M}_{d_2}. \quad (1)$$

There exists a bijection between the set of linear mappings  $\Phi$  and the set of matrices  $\mathcal{M}_{d_1 d_2}$ , known as the Choi-Jamiołkowski isomorphism. For a given

linear mapping  $\Phi$  the corresponding Choi operator  $J(\Phi)$  is explicitly written as

$$J(\Phi) := \sum_{i,j=0}^{d-1} \Phi(|i\rangle\langle j|) \otimes |i\rangle\langle j|. \quad (2)$$

We also introduce a special subset of all mappings  $\Phi$ , called quantum channels, which are completely positive and trace preserving (CPTP). In this work we will consider a special class of quantum channels, called unitary channels. A quantum channel  $\Phi_U$  is said to be a unitary channel if it has the following form  $\Phi_U(\cdot) = U \cdot U^\dagger$  for any  $U \in \mathcal{U}_d$ .

Let us recall a general form of a quantum measurement, so called Positive Operator Valued Measure (POVM). A POVM  $\mathcal{P}$  is a collection of positive semidefinite operators  $\{E_1, \dots, E_m\}$ , called effects, that sum up to the identity operator, *i.e.*  $\sum_{i=1}^m E_i = \mathbb{1}$ . In PyQBench, we are interested only in von Neumann measurements, that is measurements for which all the effects are rank-one projectors. Every such measurement can be parameterized by a unitary matrix  $U \in \mathcal{U}_d$  with effects  $\{|u_0\rangle\langle u_0|, \dots, |u_{d-1}\rangle\langle u_{d-1}|\}$ , are created by taking  $|u_i\rangle$  as  $(i+1)$ -th column of the unitary matrix  $U$ . We will denote von Neumann measurements described by the matrix  $U$  by  $\mathcal{P}_U$ . The action of von Neumann measurement  $\mathcal{P}_U$  on some state  $\rho \in \Omega_d$  can be seen as a measure-and-prepare quantum channel as follows

$$\mathcal{P}_U : \rho \rightarrow \sum_{i=0}^{d-1} \langle u_i | \rho | u_i \rangle |i\rangle\langle i|. \quad (3)$$

Moreover, observe that each von Neumann measurement  $\mathcal{P}_U$  poses a composition of a unitary channel  $\Phi_{U^\dagger}$  and the maximally dephasing channel  $\Delta$ , that means  $\mathcal{P}_U = \Delta \circ \Phi_{U^\dagger}$ .

We need to also briefly discuss about the distance between quantum operations. From [15, Theorem 1], the distance between measurements  $\mathcal{P}_U$  and  $\mathcal{P}_\mathbf{1}$  can be expressed in the notion of diamond norm, that is

$$\|\mathcal{P}_U - \mathcal{P}_\mathbf{1}\|_\diamond = \min_{E \in \mathcal{DU}_d} \|\Phi_{UE} - \Phi_\mathbf{1}\|_\diamond. \quad (4)$$

To express the distance between unitary channels, we need to introduce the definition of numerical range [16]. The set

$$W(A) = \{\langle x | A | x \rangle : |x\rangle \in \mathbb{C}^d, \langle x | x \rangle = 1\} \quad (5)$$

is called the numerical range of a given matrix  $A \in \mathcal{M}_d$ . The detailed properties of the numerical range and its generalizations we can read on the website [17].

Due to the definition of  $W(A)$ , the distance between two unitary channels  $\Phi_U$  and  $\Phi_{\mathbf{1}}$  can be written as

$$\|\Phi_U - \Phi_{\mathbf{1}}\|_{\diamond} = 2\sqrt{1 - \nu^2}, \quad (6)$$

where  $\nu = \min_{x \in W(U^\dagger)} |x|$ .

### 3. Discrimination task for Hadamard gate

For the discrimination task between von Neumann measurements  $\mathcal{P}_U$  and  $\mathcal{P}_{\mathbf{1}}$ , where  $U = H$  (the Hadamard gate), the key is to calculate the diamond norm  $\|\mathcal{P}_H - \mathcal{P}_{\mathbf{1}}\|_{\diamond}$  and determine the discriminator  $|\psi_0\rangle$ . Using semidefinite programming [18], we obtain

$$\|\mathcal{P}_H - \mathcal{P}_{\mathbf{1}}\|_{\diamond} = \sqrt{2}. \quad (7)$$

From [19] we have

$$\|\mathcal{P}_H - \mathcal{P}_{\mathbf{1}}\|_{\diamond} = \|\Phi_{HE_0} - \Phi_{\mathbf{1}}\|_{\diamond}, \quad (8)$$

where  $\Phi_U$  is a unitary channel and  $E_0$  is of the form

$$E_0 = \frac{1}{\sqrt{2}} \begin{pmatrix} 1+i & 0 \\ 0 & -1-i \end{pmatrix}. \quad (9)$$

Next, in order to construct the discriminator  $|\psi_0\rangle$  we use Lemma 5 and the proof of Theorem 1 in [15]. We show that there exist states  $\rho_1$  and  $\rho_2$  of the form  $\rho_1 = \frac{1}{2} \begin{pmatrix} 1 & i \\ -i & 1 \end{pmatrix}$  and  $\rho_2 = \frac{1}{2} \begin{pmatrix} 1 & -i \\ i & 1 \end{pmatrix}$ , respectively. Thus, we construct the quantum state  $\rho_0$  as follows:

$$\rho_0 = \frac{1}{2}\rho_1 + \frac{1}{2}\rho_2 = \frac{1}{2} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}. \quad (10)$$

According to the Lemma 5 and the proof of Theorem 1 in [15] we assume that

$$|\psi_0\rangle = \left| \sqrt{\rho_0^\top} \right\rangle. \quad (11)$$

It directly implies that

$$|\psi_0\rangle = \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle). \quad (12)$$

Next, from Holevo-Helstrom theorem [20], we determine the final measurement  $\mathcal{P}_{V_i}$ . Let us consider

$$X = (\mathcal{P}_H \otimes \mathbb{1})(|\psi_0\rangle\langle\psi_0|) - (\mathcal{P}_{\mathbf{1}} \otimes \mathbb{1})(|\psi_0\rangle\langle\psi_0|). \quad (13)$$

From the Hahn-Jordan decomposition [20], let us note

$$X = P - Q, \quad (14)$$

where  $P, Q \geq 0$ . Let us define projectors  $\Pi_P$  and  $\Pi_Q$  onto  $\text{im}(P)$  and  $\text{im}(Q)$ , respectively. Observe, that  $P$  and  $Q$  are block-diagonal. Then,  $\Pi_P$  and  $\Pi_Q$  have the following forms

$$\Pi_P = \begin{pmatrix} |x_p\rangle\langle x_p| & 0 \\ 0 & |y_p\rangle\langle y_p| \end{pmatrix}, \quad (15)$$

and

$$\Pi_Q = \begin{pmatrix} |x_q\rangle\langle x_q| & 0 \\ 0 & |y_q\rangle\langle y_q| \end{pmatrix}. \quad (16)$$

Hence, we define  $V_0$  as

$$\begin{cases} |x_p\rangle = V_0|0\rangle \\ |x_q\rangle = V_0|1\rangle \end{cases} \quad (17)$$

and  $V_1$  as

$$\begin{cases} |y_p\rangle = V_1|0\rangle \\ |y_q\rangle = V_1|1\rangle \end{cases}. \quad (18)$$

For the discrimination task between  $\mathcal{P}_H$  and  $\mathcal{P}_1$  the explicit form of  $V_0$  and  $V_1$  is given as follows (see also `mathematics/optimal_final_measurement_discrimination.nb` in the source code repository):

$$V_0 = \begin{pmatrix} \alpha & -\beta \\ \beta & \alpha \end{pmatrix}, \quad (19)$$

and

$$V_1 = \begin{pmatrix} -\beta & \alpha \\ \alpha & \beta \end{pmatrix}, \quad (20)$$

where

$$\alpha = \frac{\sqrt{2 - \sqrt{2}}}{2} = \cos\left(\frac{3}{8}\pi\right), \quad (21)$$

and

$$\beta = \frac{\sqrt{2 + \sqrt{2}}}{2} = \sin\left(\frac{3}{8}\pi\right). \quad (22)$$

#### 4. Optimal probability for parameterized Fourier family

Let us focus on single-qubit von Neumann measurements  $\mathcal{P}_1$  and  $\mathcal{P}_U$ . Assume that the unitary matrix  $U$  is of the form

$$U = H \begin{pmatrix} 1 & 0 \\ 0 & e^{i\phi} \end{pmatrix} H^\dagger \quad (23)$$

where  $H$  is the Hadamard matrix of dimension two and  $\phi \in [0, 2\pi)$ . In this section we present theoretical probability of correct discrimination between these measurements. To do that, we will present an auxiliary lemma.

**Lemma 1.** *Let  $U = H \text{diag}(1, e^{i\phi}) H^\dagger$ ,  $\phi \in [0, 2\pi)$  and let  $\Phi_U$  and  $\Phi_1$  be two unitary channels. Then, the following equation holds*

$$\min_{E \in \mathcal{DU}_2} \|\Phi_{UE} - \Phi_1\|_\diamond = \|\Phi_U - \Phi_1\|_\diamond, \quad (24)$$

*Proof.* Recall that the distance between two unitary channels is given by  $\|\Phi_U - \Phi_1\|_\diamond = 2\sqrt{1 - \nu^2}$ , where  $\nu = \min_{x \in W(U^\dagger)} |x|$  for any  $U \in \mathcal{U}_d$ . For  $U = H \begin{pmatrix} 1 & 0 \\ 0 & e^{i\phi} \end{pmatrix} H^\dagger$  the readers briefly observe that  $\nu^2 = 1 - \frac{|1 - e^{-i\phi}|^2}{4} = 1 - \frac{|1 - e^{i\phi}|^2}{4}$ . So,

$$\|\Phi_U - \Phi_1\|_\diamond = |1 - e^{i\phi}|. \quad (25)$$

It implies that it is enough to prove

$$\min_{E \in \mathcal{DU}_2} \|\Phi_{UE} - \Phi_1\|_\diamond = |1 - e^{i\phi}|. \quad (26)$$

This condition is equivalent to show that for every  $E \in \mathcal{DU}_2$

$$\nu_E \leq \frac{|1 + e^{i\phi}|}{2}, \quad (27)$$

where  $\nu_E = \min_{x \in W(U^\dagger E)} |x|$ .

The celebrated Hausdorff-Töplitz theorem [21, 22] states that  $W(A)$  of any matrix  $A \in \mathcal{M}_d$  is a convex set, and therefore we have

$$W(A) = \{\text{tr}(A\rho) : \rho \in \Omega_d\}. \quad (28)$$

So, we can assume that

$$\min_{|x\rangle \in \mathbb{C}^2 : \langle x|x \rangle = 1} |\langle x|U^\dagger|x\rangle| = \min_{\rho \in \Omega_2} |\text{tr}(U^\dagger\rho)|. \quad (29)$$

Then, we have

$$\nu_E = \min_{\rho \in \Omega_2} |\text{tr}(\rho U E)|. \quad (30)$$

For that, our task is reduced to show that for every  $E \in \mathcal{DU}_2$  there exists  $\rho \in \Omega_2$  such that

$$|\text{tr}(\rho U E)| \leq \frac{|1 + e^{i\phi}|}{2}. \quad (31)$$

Let us define  $E = \begin{pmatrix} E_0 & 0 \\ 0 & E_1 \end{pmatrix}$  and take  $\rho = \begin{pmatrix} \frac{1}{2} & 0 \\ 0 & \frac{1}{2} \end{pmatrix}$ . From spectral theorem, let us decompose  $U$  as

$$U = \lambda_0 |x_0\rangle\langle x_0| + \lambda_1 |x_1\rangle\langle x_1|, \quad (32)$$

where for eigenvalue  $\lambda_0 = 1$ , the corresponding eigenvector is of the form  $|x_0\rangle = \begin{bmatrix} \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}} \end{bmatrix}$ , whereas for  $\lambda_1 = e^{i\phi}$  we have  $|x_1\rangle = \begin{bmatrix} \frac{1}{\sqrt{2}} \\ -\frac{1}{\sqrt{2}} \end{bmatrix}$ . Then, for every  $E \in \mathcal{DU}_2$  we have

$$\begin{aligned} |\text{tr}(\rho U E)| &= \frac{1}{2} |\text{tr}(H \text{diag}(1, e^{i\phi}) H^\dagger E)| = \frac{1}{2} |\text{tr}((|x_0\rangle\langle x_0| + e^{i\phi} |x_1\rangle\langle x_1|) E)| \\ &= \frac{1}{2} |\langle x_0| E |x_0\rangle + e^{i\phi} \langle x_1| E |x_1\rangle| = \frac{1}{2} \left| \frac{E_0 + E_1}{2} + e^{i\phi} \frac{E_0 + E_1}{2} \right| \\ &= \frac{|1 + e^{i\phi}|}{2} \left| \frac{E_0 + E_1}{2} \right| \leq \frac{|1 + e^{i\phi}|}{2}, \end{aligned} \quad (33)$$

which completes the proof.  $\square$

**Theorem 1.** *The optimal probability of correct discrimination between von Neumann measurements  $\mathcal{P}_U$  and  $\mathcal{P}_1$  for  $U = H \text{diag}(1, e^{i\phi}) H^\dagger$ , where  $\phi \in [0, 2\pi)$  is given by*

$$p_{\text{succ}}(\mathcal{P}_U, \mathcal{P}_1) = \frac{1}{2} + \frac{|1 - e^{i\phi}|}{4}. \quad (34)$$

*Proof.* From Holevo-Helstrom theorem, we obtain

$$p_{\text{succ}}(\mathcal{P}_U, \mathcal{P}_1) = \frac{1}{2} + \frac{1}{4} \|\mathcal{P}_U - \mathcal{P}_1\|_\diamond. \quad (35)$$

From [15, Theorem 1], we have

$$\|\mathcal{P}_U - \mathcal{P}_1\|_\diamond = \min_{E \in \mathcal{DU}_d} \|\Phi_{UE} - \Phi_1\|_\diamond. \quad (36)$$

From Lemma 1, we know that for  $U = H \text{diag}(1, e^{i\phi}) H^\dagger$ , it also holds that

$$\min_{E \in \mathcal{DU}_2} \|\Phi_{UE} - \Phi_1\|_\diamond = \|\Phi_U - \Phi_1\|_\diamond, \quad (37)$$

which is exactly equal to

$$\|\Phi_U - \Phi_{\mathbf{1}}\|_{\diamond} = 2\sqrt{1 - \nu^2} = |1 - e^{i\phi}|. \quad (38)$$

It implies that

$$p_{\text{succ}}(\mathcal{P}_U, \mathcal{P}_{\mathbf{1}}) = \frac{1}{2} + \frac{|1 - e^{i\phi}|}{4}, \quad (39)$$

which completes the proof.  $\square$

## 5. Optimal discrimination strategy for parameterized Fourier family

In this Appendix we create the optimal theoretical strategy of discrimination between  $\mathcal{P}_U$  and  $\mathcal{P}_{\mathbf{1}}$ . To indicate the optimal strategy, we will present two propositions. The first one is concentrated around the discriminator as the optimal input state of discrimination strategy, whereas the second one describes the optimal final measurement.

**Proposition 1.** *Consider the problem of discrimination between von Neumann measurements  $\mathcal{P}_U$  and  $\mathcal{P}_{\mathbf{1}}$ ,  $U = H \text{diag}(1, e^{i\phi}) H^{\dagger}$  and  $\phi \in [0, 2\pi)$ . The discriminator has the form*

$$|\psi_0\rangle = \frac{1}{\sqrt{2}}|\mathbb{1}_2\rangle\rangle. \quad (40)$$

*Proof.* Let  $U = H \text{diag}(1, e^{i\phi}) H^{\dagger}$ ,  $\phi \in [0, 2\pi)$  be decomposed as

$$U = \lambda_0 |x_0\rangle\langle x_0| + \lambda_1 |x_1\rangle\langle x_1|, \quad (41)$$

where for eigenvalue  $\lambda_0 = 1$ , the corresponding eigenvector is of the form  $|x_0\rangle = \begin{bmatrix} \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}} \end{bmatrix}$ , whereas for  $\lambda_1 = e^{i\phi}$  we have  $|x_1\rangle = \begin{bmatrix} \frac{1}{\sqrt{2}} \\ -\frac{1}{\sqrt{2}} \end{bmatrix}$ . For Hermitian-preserving maps [20] the diamond norm may be expressed as

$$\|\Phi\|_{\diamond} = \max_{\rho \in \Omega_d} \|(\mathbb{1} \otimes \sqrt{\rho}) J(\Phi) (\mathbb{1} \otimes \sqrt{\rho})\|_1. \quad (42)$$



Hence, we obtain

$$\begin{aligned}
\|\mathcal{P}_U - \mathcal{P}_1\|_\diamond &= \max_{\rho \in \Omega_2} \|(\mathbb{1} \otimes \sqrt{\rho}) J(\mathcal{P}_U - \mathcal{P}_1) (\mathbb{1} \otimes \sqrt{\rho})\|_1 \\
&= \max_{\rho \in \Omega_2} \left\| (\mathbb{1} \otimes \sqrt{\rho}) \sum_{i=0}^1 |i\rangle\langle i| \otimes (|u_i\rangle\langle u_i| - |i\rangle\langle i|)^\top (\mathbb{1} \otimes \sqrt{\rho}) \right\|_1 \\
&= \max_{\rho \in \Omega_2} \left\| \sum_{i=0}^1 |i\rangle\langle i| \otimes \sqrt{\rho} (|u_i\rangle\langle u_i| - |i\rangle\langle i|)^\top \sqrt{\rho} \right\|_1 \\
&= \max_{\rho \in \Omega_2} \sum_{i=0}^1 \left\| \sqrt{\rho} (|u_i\rangle\langle u_i| - |i\rangle\langle i|)^\top \sqrt{\rho} \right\|_1.
\end{aligned} \tag{43}$$

One can prove that for all  $\alpha, \beta \geq 0$ , and unit vectors  $|x\rangle, |y\rangle$  the following equation holds [20]

$$\|\alpha|x\rangle\langle x| - \beta|y\rangle\langle y|\|_1 = \sqrt{(\alpha + \beta)^2 - 4\alpha\beta|\langle x|y\rangle|^2}. \tag{44}$$

By taking  $|x\rangle = \frac{\sqrt{\rho}|\bar{u}_i\rangle}{\|\sqrt{\rho}|\bar{u}_i\rangle\|}$  and  $|y\rangle = \frac{\sqrt{\rho}|i\rangle}{\|\sqrt{\rho}|i\rangle\|}$  we have

$$\|\mathcal{P}_U - \mathcal{P}_1\|_\diamond = \max_{\rho \in \Omega_2} \sum_{i=0}^1 \sqrt{(\langle \bar{u}_i|\rho|\bar{u}_i\rangle + \langle i|\rho|i\rangle)^2 - 4|\langle \bar{u}_i|\rho|i\rangle|^2}. \tag{45}$$

Let us take  $\rho_0 = \frac{1}{2} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$ , we obtain

$$\begin{aligned}
\|\mathcal{P}_U - \mathcal{P}_1\|_\diamond &\geq \sum_{i=0}^1 \sqrt{(\langle \bar{u}_i|\rho_0|\bar{u}_i\rangle + \langle i|\rho_0|i\rangle)^2 - 4|\langle \bar{u}_i|\rho_0|i\rangle|^2} \\
&= \sum_{i=0}^1 \sqrt{1 - |\langle i|U|i\rangle|^2} \\
&= \sum_{i=0}^1 \sqrt{1 - |1 \cdot \langle i|u_0\rangle\langle u_0|i\rangle + e^{i\phi} \cdot \langle i|u_1\rangle\langle u_1|i\rangle|^2} \\
&= \sum_{i=0}^1 \sqrt{1 - \left| \frac{1 + e^{i\phi}}{2} \right|^2} = 2\sqrt{1 - \left| \frac{1 + e^{i\phi}}{2} \right|^2} \\
&= |1 - e^{i\phi}|.
\end{aligned} \tag{46}$$

Due to Theorem 1 and the following equality

$$\|(\mathbb{1} \otimes \sqrt{\rho}) J(\mathcal{P}_U - \mathcal{P}_1) (\mathbb{1} \otimes \sqrt{\rho})\|_1 = \left\| ((\mathcal{P}_U - \mathcal{P}_1) \otimes \mathbb{1}) \left( |\sqrt{\rho}^\top\rangle\langle\sqrt{\rho}^\top| \right) \right\|_1, \tag{47}$$

the discriminator  $|\psi_0\rangle$  is equal to

$$|\psi_0\rangle = |\sqrt{\rho_0}^\top\rangle = \frac{1}{\sqrt{2}}|\mathbb{1}_2\rangle, \quad (48)$$

which completes the proof.  $\square$

**Proposition 2.** *Consider the problem of discrimination between von Neumann measurements  $\mathcal{P}_U$  and  $\mathcal{P}_1$ ,  $U = H\text{diag}(1, e^{i\phi})H^\dagger$  and  $\phi \in [0, 2\pi)$ . The controlled unitaries  $V_0$  and  $V_1$  have the form*

$$V_0 = \begin{pmatrix} i \sin\left(\frac{\pi-\phi}{4}\right) & -i \cos\left(\frac{\pi-\phi}{4}\right) \\ \cos\left(\frac{\pi-\phi}{4}\right) & \sin\left(\frac{\pi-\phi}{4}\right) \end{pmatrix}, \quad (49)$$

and

$$V_1 = \begin{pmatrix} -i \cos\left(\frac{\pi-\phi}{4}\right) & i \sin\left(\frac{\pi-\phi}{4}\right) \\ \sin\left(\frac{\pi-\phi}{4}\right) & \cos\left(\frac{\pi-\phi}{4}\right) \end{pmatrix}. \quad (50)$$

*Proof.* From Proposition 1 we obtain the exact form of discriminator given by

$$|\psi_0\rangle = \frac{1}{\sqrt{2}}|\mathbb{1}_2\rangle. \quad (51)$$

Repeating the procedure used to distinguish the von Neumann measurements in the Hadamard basis (see 3), we use the Hahn-Jordan decomposition and then we create the projective operators into the positive and negative part of  $X$  matrix. Hence, the explicit form of  $V_0$  and  $V_1$  is given as follows:

$$V_0 = \begin{pmatrix} i \sin\left(\frac{\pi-\phi}{4}\right) & -i \cos\left(\frac{\pi-\phi}{4}\right) \\ \cos\left(\frac{\pi-\phi}{4}\right) & \sin\left(\frac{\pi-\phi}{4}\right) \end{pmatrix}, \quad (52)$$

and

$$V_1 = \begin{pmatrix} -i \cos\left(\frac{\pi-\phi}{4}\right) & i \sin\left(\frac{\pi-\phi}{4}\right) \\ \sin\left(\frac{\pi-\phi}{4}\right) & \cos\left(\frac{\pi-\phi}{4}\right) \end{pmatrix}, \quad (53)$$

where  $\phi \in [0, 2\pi)$ .  $\square$

## 6. Output YAML files

In this appendix we present examples of YAML's files obtained from synchronous and asynchronous experiments. We will start at synchronous case.

Listing 1: Defining experiment file

```
type: discrimination-fourier
qubits:
  - target: 0
    ancilla: 1
  - target: 1
    ancilla: 2
angles:
  start: 0
  stop: 2 * pi
  num_steps: 3
gateset: ibmq
method: direct_sum
num_shots: 100
```

Listing 2: Defining backend file

```
name: ibmq-quito
asynchronous: false
provider:
  hub: ibm-q
  group: open
  project: main
```

Listing 3: Results (synchronous)

```
metadata:
  experiments:
    type: discrimination-fourier
    qubits:
      - {target: 0, ancilla: 1}
      - {target: 1, ancilla: 2}
    angles: {start: 0.0, stop: 6.283185307179586, num_step
    gateset: ibmq
    method: direct_sum
    num_shots: 100
  backend_description:
    name: ibmq-quito
    asynchronous: false
    provider: {group: open, hub: ibm-q, project: main}
data:
  - target: 0
```



```

- name: u
  histogram: {'00': 32, '01': 27, '10': 24, '11': 17}
  mitigation_info:
    target: {prob_meas0_prep1: 0.052200000000000024, prob_m
    ancilla: {prob_meas0_prep1: 0.059000000000000005, prob_m
  mitigated_histogram: {'00': 0.3025357275361897, '01': 0.274136}
- target: 1
  ancilla: 2
  phi: 0.0
  results_per_circuit:
- name: id
  histogram: {'00': 27, '01': 20, '10': 24, '11': 29}
  mitigation_info:
    target: {prob_meas0_prep1: 0.059000000000000005, prob_m
    ancilla: {prob_meas0_prep1: 0.075400000000000002, prob_m
  mitigated_histogram: {'00': 0.2594378169217188, '01': 0.193188}
- name: u
  histogram: {'00': 31, '01': 24, '10': 23, '11': 22}
  mitigation_info:
    target: {prob_meas0_prep1: 0.059000000000000005, prob_m
    ancilla: {prob_meas0_prep1: 0.075400000000000002, prob_m
  mitigated_histogram: {'00': 0.30056875246775644, '01': 0.24382}
- target: 1
  ancilla: 2
  phi: 3.141592653589793
  results_per_circuit:
- name: id
  histogram: {'00': 5, '01': 4, '10': 50, '11': 41}
  mitigation_info:
    target: {prob_meas0_prep1: 0.059000000000000005, prob_m
    ancilla: {prob_meas0_prep1: 0.075400000000000002, prob_m
  mitigated_histogram: {'00': 0.009552870928837118, '01': 0.0071}
- name: u
  histogram: {'00': 41, '01': 51, '10': 3, '11': 5}
  mitigation_info:
    target: {prob_meas0_prep1: 0.059000000000000005, prob_m
    ancilla: {prob_meas0_prep1: 0.075400000000000002, prob_m
  mitigated_histogram: {'00': 0.4073387714165384, '01': 0.561461}
- target: 1
  ancilla: 2
  phi: 6.283185307179586

```

```

results_per_circuit:
- name: id
  histogram: {'00': 30, '01': 28, '10': 23, '11': 19}
  mitigation_info:
    target: {prob_meas0_prep1: 0.059000000000000005, prob_m
    ancilla: {prob_meas0_prep1: 0.075400000000000002, prob_m
  mitigated_histogram: {'00': 0.2868459834940102, '01': 0.291956
- name: u
  histogram: {'00': 15, '01': 20, '10': 36, '11': 29}
  mitigation_info:
    target: {prob_meas0_prep1: 0.059000000000000005, prob_m
    ancilla: {prob_meas0_prep1: 0.075400000000000002, prob_m
  mitigated_histogram: {'00': 0.1187719606657805, '01': 0.196208

```

For the same experiment file, we use the flag `asynchronous: true` to define asynchronous experiment.

Listing 4: Backend file

```

name: ibmq-quito
asynchronous: true
provider:
  hub: ibmq
  group: open
  project: main

```

If the backend is asynchronous, the output will contain intermediate data containing, amongst others, job ids correlated with the circuit they correspond to.

Listing 5: Resolved results

```

metadata:
  experiments:
    type: discrimination-fourier
    qubits:
      - {target: 0, ancilla: 1}
      - {target: 1, ancilla: 2}
    angles: {start: 0.0, stop: 6.283185307179586, num_step
    gateset: ibmq
    method: direct_sum
    num_shots: 100
  backend_description:
    name: ibmq-quito

```

```

        asynchronous: true
        provider: {group: open, hub: ibmq, project: main}
data:
- job_id: 63e7f17a17b7ed49ca24e05b
  keys:
  - [0, 1, id, 0.0]
  - [0, 1, u, 0.0]
  - [0, 1, id, 3.141592653589793]
  - [0, 1, u, 3.141592653589793]
  - [0, 1, id, 6.283185307179586]
  - [0, 1, u, 6.283185307179586]
  - [1, 2, id, 0.0]
  - [1, 2, u, 0.0]
  - [1, 2, id, 3.141592653589793]
  - [1, 2, u, 3.141592653589793]
  - [1, 2, id, 6.283185307179586]
  - [1, 2, u, 6.283185307179586]

```

Finally, if the status of jobs is **DONE**, we resolve the measurements from the submitted jobs obtaining the following file.

Listing 6: Results (asynchronous)

```

metadata:
  experiments:
    type: discrimination-fourier
    qubits:
      - target: 0
        ancilla: 1
      - target: 1
        ancilla: 2
    angles:
      start: 0.0
      stop: 6.283185307179586
      num_steps: 3
    gateset: ibmq
    method: direct_sum
    num_shots: 100
  backend_description:
    name: ibmq-quito
    asynchronous: true
    provider:
      group: open

```

```

hub: ibmq
project: main

data:
- target: 0
  ancilla: 1
  phi: 0.0
  results_per_circuit:
- name: id
  histogram:
    '00': 27
    '01': 28
    '10': 18
    '11': 27
  mitigation_info:
    target:
      prob_meas0_prep1: 0.0522000000000000024
      prob_meas1_prep0: 0.0172
    ancilla:
      prob_meas0_prep1: 0.0590000000000000005
      prob_meas1_prep0: 0.0202
  mitigated_histogram:
    '00': 0.254196166145997
    '01': 0.2790358060520916
    '10': 0.1732699847244092
    '11': 0.29349804307750227
- name: u
  histogram:
    '00': 29
    '01': 17
    '10': 30
    '11': 24
  mitigation_info:
    target:
      prob_meas0_prep1: 0.0522000000000000024
      prob_meas1_prep0: 0.0172
    ancilla:
      prob_meas0_prep1: 0.0590000000000000005
      prob_meas1_prep0: 0.0202
  mitigated_histogram:
    '00': 0.2733793468261183
    '01': 0.1621115306717096

```



```

'10': 0.3045273800167787
'11': 0.2599817424853933
- target: 0
  ancilla: 1
  phi: 3.141592653589793
  results_per_circuit:
- name: id
  histogram:
    '00': 3
    '01': 5
    '10': 37
    '11': 55
  mitigation_info:
    target:
      prob_meas0_prep1: 0.052200000000000024
      prob_meas1_prep0: 0.0172
    ancilla:
      prob_meas0_prep1: 0.059000000000000005
      prob_meas1_prep0: 0.0202
  mitigated_histogram:
    '00': 0.006189545789708441
    '01': 0.016616709640352317
    '10': 0.3675478279476653
    '11': 0.6096459166222741
- name: u
  histogram:
    '00': 56
    '01': 42
    '10': 2
  mitigation_info:
    target:
      prob_meas0_prep1: 0.052200000000000024
      prob_meas1_prep0: 0.0172
    ancilla:
      prob_meas0_prep1: 0.059000000000000005
      prob_meas1_prep0: 0.0202
  mitigated_histogram:
    '00': 0.55731929321128
    '01': 0.43367489257574243
    '10': 0.009005814212977551
- target: 0

```

```

ancilla: 1
phi: 6.283185307179586
results_per_circuit:
- name: id
    histogram:
        '00': 18
        '01': 28
        '10': 30
        '11': 24
    mitigation_info:
        target:
            prob_meas0_prep1: 0.0522000000000000024
            prob_meas1_prep0: 0.0172
        ancilla:
            prob_meas0_prep1: 0.0590000000000000005
            prob_meas1_prep0: 0.0202
    mitigated_histogram:
        '00': 0.15258295844557557
        '01': 0.2829079190522524
        '10': 0.3071204587046501
        '11': 0.25738866379752195
- name: u
    histogram:
        '00': 32
        '01': 28
        '10': 23
        '11': 17
    mitigation_info:
        target:
            prob_meas0_prep1: 0.0522000000000000024
            prob_meas1_prep0: 0.0172
        ancilla:
            prob_meas0_prep1: 0.0590000000000000005
            prob_meas1_prep0: 0.0202
    mitigated_histogram:
        '00': 0.3026150836796529
        '01': 0.28491749668524724
        '10': 0.23230862145681827
        '11': 0.18015879817828173
- target: 1
  ancilla: 2

```

```

phi: 0.0
results_per_circuit:
- name: id
    histogram:
        '00': 27
        '01': 16
        '10': 30
        '11': 27
    mitigation_info:
        target:
            prob_meas0_prep1: 0.059000000000000005
            prob_meas1_prep0: 0.0202
        ancilla:
            prob_meas0_prep1: 0.075400000000000002
            prob_meas1_prep0: 0.0528
    mitigated_histogram:
        '00': 0.256742095057232
        '01': 0.15000257115061383
        '10': 0.29821012040758116
        '11': 0.29504521338457296
- name: u
    histogram:
        '00': 34
        '01': 22
        '10': 25
        '11': 19
    mitigation_info:
        target:
            prob_meas0_prep1: 0.059000000000000005
            prob_meas1_prep0: 0.0202
        ancilla:
            prob_meas0_prep1: 0.075400000000000002
            prob_meas1_prep0: 0.0528
    mitigated_histogram:
        '00': 0.3325088211394024
        '01': 0.22335261496979697
        '10': 0.2441636375921354
        '11': 0.19997492629866526
- target: 1
  ancilla: 2
  phi: 3.141592653589793

```

```

results_per_circuit:
- name: id
    histogram:
        '00': 3
        '01': 9
        '10': 51
        '11': 37
    mitigation_info:
        target:
            prob_meas0_prep1: 0.059000000000000005
            prob_meas1_prep0: 0.0202
        ancilla:
            prob_meas0_prep1: 0.075400000000000002
            prob_meas1_prep0: 0.0528
    mitigated_histogram:
        '00': -0.016627023111853642
        '01': 0.06778554570877951
        '10': 0.53899887367658
        '11': 0.40984260372649417
- name: u
    histogram:
        '00': 43
        '01': 45
        '10': 7
        '11': 5
    mitigation_info:
        target:
            prob_meas0_prep1: 0.059000000000000005
            prob_meas1_prep0: 0.0202
        ancilla:
            prob_meas0_prep1: 0.075400000000000002
            prob_meas1_prep0: 0.0528
    mitigated_histogram:
        '00': 0.42955729968594086
        '01': 0.49336080079582095
        '10': 0.04937406434533623
        '11': 0.02770783517290191
- target: 1
  ancilla: 2
  phi: 6.283185307179586
  results_per_circuit:

```

```

— name: id
  histogram:
    '00 ': 22
    '01 ': 19
    '10 ': 35
    '11 ': 24
  mitigation_info:
    target:
      prob_meas0_prep1: 0.059000000000000005
      prob_meas1_prep0: 0.0202
    ancilla:
      prob_meas0_prep1: 0.075400000000000002
      prob_meas1_prep0: 0.0528
  mitigated_histogram:
    '00 ': 0.19592641048040849
    '01 ': 0.18787721420415215
    '10 ': 0.3590258049844047
    '11 ': 0.25717057033103463

— name: u
  histogram:
    '00 ': 27
    '01 ': 24
    '10 ': 25
    '11 ': 24
  mitigation_info:
    target:
      prob_meas0_prep1: 0.059000000000000005
      prob_meas1_prep0: 0.0202
    ancilla:
      prob_meas0_prep1: 0.075400000000000002
      prob_meas1_prep0: 0.0528
  mitigated_histogram:
    '00 ': 0.25555866817587225
    '01 ': 0.2429501641251142
    '10 ': 0.24509293912212946
    '11 ': 0.2563982285768841

```

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