Harmonic Analysis

Ikhan Choi

June 4, 2023

Contents

Ι	Fourier analysis	3
1	Fourier series 1.1 Fourier series in L^p spaces	4 4 4 6
2	Fourier transform2.1 Fourier transform in L^p space	7 7 7
3	Hilbert transform3.1 Harmonic conjugate3.2 Kernel representation3.3 Fourier series in L^p space	9 9 9
II	Singular integral operators	10
4	Calderón-Zygmund theory 4.1 Convolution type operators 4.2 Truncated integrals 4.3 A_p weights 4.4 Bounded mean oscillation	11 11 12 13 13
5 6	Littlewood-Paley theory 5.1 Littlewood-Paley decomposition 5.2 Multiplier theorems Almost orthogonality 6.1 Coltar lemma	14 14 14 15 15
II	I Oscillatory integral operators	16
7	Oscillatory integrals	17
8	Foureir restriction	19
9		20

IV	Pseudo-differential operators	21
	Pseudo-differential calculus	22
	10.2	
	Semiclassical analysis 11.1 Heisenberg group 11.2 Phase space transforms	
12	Microlocal analysis	27

Part I Fourier analysis

Fourier series

1.1 Fourier series in L^p spaces

1.1.

$$\|\widehat{f}\|_{\ell^1(\mathbb{Z})} \lesssim \|f\|_{W^{1,1+\varepsilon}(\mathbb{T})}.$$

Inversion theorem is an approximation problem given by $\mathcal{F}^*\mathcal{F}=\lim_{n\to\infty}\mathcal{F}_n^*\mathcal{F}$. The condition $\widehat{f}\in \ell^1(\mathbb{Z})$ is a condition just for defining $\mathcal{F}^*\widehat{f}$ without using distribution theory, and it does not affect the inversion phenomena. The approximation, in other words, can be seen as an extension method for $\mathcal{F}^*:\ell^1(\mathbb{Z})\to C(\mathbb{T})$ on $c_0(\mathbb{Z})$. Note that \mathcal{F}_n^* on $c_0(\mathbb{Z})$ cannot be bounded directly without distribution theory, but $\mathcal{F}_n^*\mathcal{F}$ on $L^p(\mathbb{T})$ can be bounded well.

1.2 Summability methods

- If \mathcal{F}_n^* is the standard partial sum, then $\mathcal{F}_n^*\mathcal{F}$ is the Dirichlet kernel.
- If \mathcal{F}_n^* is the Cesàro mean, then $\mathcal{F}_n^*\mathcal{F}$ is the Fejér kernel.
- If \mathcal{F}_r^* is the Abel sum, then $\mathcal{F}_r^*\mathcal{F}$ is the Poisson kernel.
- In Fourier transform, we often use the Gauss-Weierstrass kernel.

The injectivity of $\mathcal F$ is not an easy problem, which comes from the inversion theorem.

1.2 (Dirichlet kernel). The *Dirichlet kernel* is a function $D_n: \mathbf{T} \to \mathbb{R}$ defined by

$$D_n = \widehat{\mathbf{1}_{|k| \le n}}$$
, or equivalently, $\widehat{D_n} = \mathbf{1}_{|k| \le n}$.

This is because they are invariant under inverse, in other words, they are even.

(a) $D_n(x) = \frac{\sin \frac{2n+1}{2} x}{\sin \frac{1}{2} x}.$

(b) If $f \in \text{Lip}(\mathbf{T})$, then $D_n * f \to f$ pointwisely as $n \to \infty$.

(c) $||D_n||_{L^1(\mathbf{T})} \gtrsim \log n.$

Proof.

$$D_n(x) = \sum_{k=-n}^{n} e^{ikx}$$

$$= \frac{e^{i\frac{2n+1}{2}x} - e^{-i\frac{2n+1}{2}x}}{e^{i\frac{1}{2}x} - e^{-i\frac{1}{2}x}}$$

$$= \frac{\sin\frac{2n+1}{2}x}{\sin\frac{1}{2}x}.$$

(c) By (2) $\sin x \le x$ for $x \in [0, \pi/2]$, (3) change of variable,

$$||D_n||_{L^1(\mathbf{T})} = \frac{1}{2\pi} \int_{-\pi}^{\pi} |\frac{\sin\frac{2n+1}{2}x}{\sin\frac{1}{2}x}| dx$$

$$\geq \frac{2}{\pi} \int_{0}^{\pi} \frac{|\sin\frac{2n+1}{2}x|}{x} dx$$

$$= \frac{2}{\pi} \int_{0}^{\frac{2n+1}{2}\pi} \frac{|\sin x|}{x} dx$$

$$= \frac{2}{\pi} \sum_{k=0}^{2n} \int_{\frac{k}{2}\pi}^{\frac{k+1}{2}\pi} \frac{|\sin x|}{x} dx$$

$$\geq \frac{2}{\pi} \sum_{k=0}^{2n} \int_{0}^{\frac{1}{2}\pi} \frac{\sin x}{\frac{k+1}{2}\pi} dx$$

$$\geq \frac{4}{\pi^2} \sum_{k=0}^{2n} \frac{1}{1+k}$$

$$\geq \frac{4}{\pi^2} \log(2n+2).$$

..?

1.3 (Fejér kernel). The Fejér kernel is

(a)

$$K_n(x) = \frac{1}{n+1} \frac{\sin^2 \frac{n+1}{2} x}{\sin^2 \frac{1}{2} x}.$$

Proof. Since

$$\begin{split} D_n(x) &= \frac{e^{i\frac{2n+1}{2}x} - e^{-i\frac{2n+1}{2}x}}{e^{i\frac{1}{2}x} - e^{-i\frac{1}{2}x}} \\ &= \frac{\left[e^{i\frac{2n+1}{2}x} - e^{-i\frac{2n+1}{2}x}\right] \left[e^{i\frac{1}{2}x} - e^{-i\frac{1}{2}x}\right]}{\left[e^{i\frac{1}{2}x} - e^{-i\frac{1}{2}x}\right]^2} \\ &= \frac{\left[e^{i(n+1)x} + e^{-i(n+1)x}\right] - \left[e^{inx} + e^{-inx}\right]}{\left[e^{i\frac{1}{2}x} - e^{-i\frac{1}{2}x}\right]^2}, \end{split}$$

by telescoping, we get

$$\begin{split} \sum_{k=0}^{n} D_k(x) &= \frac{\left[e^{i(n+1)x} + e^{-i(n+1)x}\right] - \left[e^{i0x} + e^{-i0x}\right]}{\left[e^{i\frac{1}{2}x} - e^{-i\frac{1}{2}x}\right]^2} \\ &= \frac{\left[e^{i\frac{n+1}{2}x} - e^{-i\frac{n+1}{2}x}\right]^2}{\left[e^{i\frac{1}{2}x} - e^{-i\frac{1}{2}x}\right]^2} \\ &= \frac{\sin^2\frac{n+1}{2}x}{\sin^2\frac{1}{2}x}. \end{split}$$

Two important results from Fejér kernel:

- 1. If f(x-), f(x+) exist and $S_n f(x)$ converges, then $S_n f(x) \to \frac{1}{2} (f(x-) + f(x+))$.
- 2. (If $f \in L^1(\mathbf{T})$, then $\sigma_n f \to f$ a.e.)
- 3. If $f \in L^1(\mathbf{T})$, then $S_n f \to f$ in L^1 and L^2 .
- 4. If f is continuous and $\hat{f} \in L^1(\mathbb{Z})$, then $S_n f \to f$ uniformly.
- 5. Since $\sigma_n f$ is a trigonometric polynomial, the set of trigonometric polynomials are dense in $L^1(\mathbf{T})$ and $L^2(\mathbf{T})$.

1.3 Pointwise convergence of Fourier series

BV function: Dini, Jordan's criterion

1.4 (Riemann localization principle).

Exercises

1.5 (Gibbs phenomenon).

1.6 (Du Bois-Reymond function).

Fourier transform

2.1 Fourier transform in L^p space

2.1 (Riemann-Lebesgue lemma).

Lp extension

Gaussian function computation: differential equation method, contour integral method inversion theorem

2.2 (Plancherel theorem).

2.2 Distributions

2.3 (Cauchy principal value). indented contour, imaginary shift, Feynman's trick

Exercises

2.4 (Sampling theorem).

$$\mathcal{F}\mathbf{1}_{[-\frac{1}{2},\frac{1}{2}]}(\xi) = \operatorname{sinc}(\xi/2)$$

 $\operatorname{sinc} \in L^{1+\varepsilon}(\mathbb{R}).$

2.5 (Poisson summation formula).

2.6 (Uncertainty principle).

2.7 (Multipole expansion). Let ρ be a compactly supported distribution on \mathbb{R}^d . We want to investigate the limit behavior of $\rho(\varepsilon^{-1}x)$ as $\varepsilon \to 0$. More precisely, we want to compute an integer $k \ge d$ such that $\lim_{\varepsilon \to 0+} \varepsilon^{-k} \rho(\varepsilon^{-1}x)$ defines a distribution supported at $\{0\}$, and the coefficients of derivatives of Dirac measures.

We need to introduce quantities called monopole, dipole, quadrapole, octupole, etc.

(a) A distribution supported on {0} is a linear combination of the Dirac measure and its derivatives.

(b)

Problems

1. Find all
$$\alpha > 0$$
 such that

$$\lim_{x \to \infty} x^{-\alpha} \int_0^x f(y) \, dy = 0$$

for all
$$f \in L^3([0,\infty))$$
.

Hilbert transform

- 3.1 Harmonic conjugate
- 3.2 Kernel representation
- **3.3** Fourier series in L^p space

Part II Singular integral operators

Calderón-Zygmund theory

4.1 Convolution type operators

4.1 (Calderón-Zygmund decomposition of sets). Let $f \in L^1(\mathbb{R}^d)$. Let $E_n f$ be the conditional expectation with repect to the σ -algebra generated by dyadic cubes with side length 2^{-n} . Let $Mf := \sup_n E_n |f|$ be the maximal function, and let $\Omega := \{x : Mf(x) > \lambda\}$ for fixed $\lambda > 0$. For $x \in \Omega$ let Q_x be the maximal dyadic cube such that $x \in Q_x$ and

$$\frac{1}{|Q_x|} \int_{Q_x} |f| > \lambda.$$

- (a) $\{Q_x : x \in \Omega\}$ is a countable partition of Ω .
- (b) We have an weak type estimate $|\Omega| \leq \frac{1}{\lambda} ||f||_{L^1}$.
- (c) $||f||_{L^{\infty}(\mathbb{R}^d\setminus\Omega)} \leq \lambda$.
- (d) For $x \in \Omega$

$$\frac{1}{|Q_x|} \int_{Q_x} |f| \le 2^d \lambda.$$

4.2 (Calderón-Zygmund decomposition of functions). Let

$$g(x) := \begin{cases} |f(x)| & , x \notin \Omega \\ \frac{1}{|Q_x|} \int_{Q_x} |f| & , x \in \Omega \end{cases}$$

and $b_i := (|f| - g)\chi_{Q_i}$ so that |f| = g + b where $b = \sum_i b_i$.

- (a) $||g||_{L^1} = ||f||_{L^1}$ and $||g||_{L^{\infty}} \lesssim_d \lambda$.
- (b) $||b||_{L^1} \le 2||f||_{L^1}$ and $\int b_i = 0$.

Proof.

- **4.3** (L^p boundedness of Calderón-Zygmund operators). Let $T: C_c^{\infty}(\mathbb{R}^d) \to \mathcal{D}'(\mathbb{R}^d)$ be a *singular integral operator of convolution type* in the sense that there is a function $K \in L^1_{loc}(\mathbb{R}^d \setminus \{0\}) \cap \mathcal{D}'(\mathbb{R}^d)$ such that Tf(x) = K * f(x) for all $f \in \mathcal{D}(\mathbb{R}^d)$, whenever $x \notin \text{supp } f$. We say T is called a *Calderón-Zygmund* operator if
 - (i) T is L^2 -bounded: we have

$$||Tf||_{L^2} \lesssim ||f||_{L^2},$$

(ii) T satisfies the Hörmander condition: we have

$$\int_{|x|>2|y|} |K(x-y)-K(x)| \, dx \lesssim 1$$

for every y > 0.

Let $f=g+b=g+\sum_i b_i$ be the Calderón-Zygmund decomposition, and let $\Omega^*:=\bigcup_i Q_i^*$ where Q_i^* is the cube with the same center as Q_i and whose sides are $2\sqrt{d}$ times longer.

(a) The L^2 -boundedness implies

$$|\{x: |Tg(x)| > \frac{\lambda}{2}\}| \lesssim_d \frac{1}{\lambda} ||f||_{L^1}.$$

(b) The Hörmander condition implies

$$|\{x: |Tb(x)| > \frac{\lambda}{2}\} \setminus \Omega^*| \lesssim_d \frac{1}{\lambda} ||f||_{L^1}.$$

(c)

Proof. (a) Using the Chebyshev inequality and the Hölder inequality,

$$|\{x: |Tg(x)| > \frac{\lambda}{2}\}| \le \frac{4}{\lambda^2} ||Tg||_{L^2(\Omega)}^2 \le \frac{4C}{\lambda^2} ||g||_{L^2(\Omega)}^2 \le \frac{4C}{\lambda^2} ||g||_{L^1(\Omega)} ||g||_{L^\infty(\Omega)}.$$

(b) Write

$$|\{x: |Tb(x)| > \frac{\lambda}{2}\} \setminus \Omega^*| \le \frac{2}{\lambda} \int_{\mathbb{R}^d \setminus \Omega^*} |Tb(x)| \, dx \le \frac{2}{\lambda} \sum_i \int_{\mathbb{R}^d \setminus \Omega^*} |Tb_i(x)| \, dx.$$

Since $x \in \mathbb{R}^d \setminus Q_i^*$ does not belong to supp $b_i \subset Q_i$ and $\int b_i = 0$, we have

$$Tb_{i}(x) = \int_{Q_{i}} K(x - y)b_{i}(y) dy = \int_{Q_{i}} [K(x - y) - K(x)]b_{i}(y) dy,$$

and

$$\int_{\mathbb{R}^d \setminus Q_i^*} |Tb_i(x)| \, dx = \int_{Q_i} |b_i(y)| \int_{\mathbb{R}^d \setminus Q_i^*} |K(x-y) - K(x)| \, dx \, dy \lesssim \|b_i\|_{L^1}.$$

(We need to show it is valid even though b_i is not smooth)

(c)

4.4 (Hölder boundedness of Calderón-Zygmund operators).

4.2 Truncated integrals

Homogeneous kernels

4.3 A_p weights

4.4 Bounded mean oscillation

Exercises

4.5 (Size and cancellation condition). Let $K \in L^1_{loc}(\mathbb{R}^d \setminus \{0\}) \cap \mathcal{D}'(\mathbb{R}^d)$. We say the condition $|K(x)| \lesssim |x|^{-d}$ for $x \neq 0$ as the *size condition*, and say the condition $\int_{r < |x| < R} K(x) \, dx = 0$ for all $0 < r < R < \infty$ as the *cancellation condition*. If K satisfies the size, cancellation, and Hörmander condition, then it is L^2 bounded, hence Calderón-Zygmund.

4.6 (Gradient size condition). Let $|\nabla K(x)| \lesssim |x|^{-d-1}$ for $x \neq 0$. Then, convolution with K is a Calderón-Zygmund operator.

4.7 (Riesz potential).

Littlewood-Paley theory

- 5.1 Littlewood-Paley decomposition
- 5.2 Multiplier theorems

Almost orthogonality

Carleson measures, paraproducts

- 6.1 Coltar lemma
- **6.2** T(1) theorem

Part III Oscillatory integral operators

Oscillatory integrals

7.1 (Justification of oscillatory integral). For ϕ , we define a linear functional $O_{\phi}: A^m_{\delta}(\mathbb{R}^d) \to \mathbb{C}$ such that

$$O_{\phi}(a) := \int_{\mathbb{R}^d} e^{i\phi(x)} a(x) \, dx$$

for all $a \in A^m_{\delta}(\mathbb{R}^d) \cap L^1(\mathbb{R}^d)$. As a notation, we will use the above integral to denote the value of O_{ϕ} even for $a \in A^m_{\delta}(\mathbb{R}^d) \setminus L^1(\mathbb{R}^d)$.

- (a) O_{ϕ} is well-defined and continuous.
- (b) The change of variables is justified as follows:
- (c) The integral by parts is justified as follows: for $\phi \in$ and $a \in$, we have

$$\int_{\mathbb{R}^d} e^{i\phi(y)} i\partial \phi(y) a(x+y) dy = -\int_{\mathbb{R}^d} e^{i\phi(y)} \partial a(x+y) dy.$$

- (d) The Fubini theorem is justified as follows:
- (e) The Fourier inversion is justified as follows:

$$a(x) = (2\pi)^{-d} \int_{\mathbb{R}^{2d}} e^{i(x-y)\xi} a(y) \, dy \, d\xi.$$

Proof. (a) The most difficult part is the construction and the computation of L and its transpose.

- (e) Note that the function $(y, \xi) \mapsto a(y)$ belongs to $A_{\delta}^{m'}(\mathbb{R}^{2d})$ since
- **7.2** (Point evaluation of multiplier). We want to show the following point evaluation holds with previously justified oscillatory integral: for each x at which the left-hand side is continuous, we have

$$\Phi(D)a(x) = (2\pi)^{-d} \int_{\mathbb{R}^d} e^{i\phi(y)} a(x+y) \, dy$$

for all $a \in A^m_{\delta}$, where $\Phi := \mathcal{F}^* e^{i\phi}$.

7.3 (Stationary phase approximation).

Proof.
$$\Box$$

7.4 (Van der Corput lemma).

Dispersive equations and strichartz estimates

Exercises

7.5 (Fresnel phase). We compute L with a specific example

Proof.

$$(1 + xQ^{-1}D)e^{\frac{i}{2}xQx} = \langle x \rangle^2 e^{\frac{i}{2}xQx}.$$

The transpose of $\langle x \rangle^{-2} (1 + xQ^{-1}D)$ is $\langle x \rangle^{-2} (1 + di - 2ix^2 - xD)$ for Q = I.

Note that $\langle x \rangle^{-2n} \langle D \rangle^{2n}$ is self-adjoint.

Let Q be a non-degenerate symmetric bilinear form on \mathbb{R}^d . Consider a multiplier operator $e^{\frac{i}{2}DQD}$: $\mathcal{S} \to \mathcal{S}$ such that

$$e^{\frac{i}{2}DQD}a(x) := \mathcal{F}^*e^{\frac{i}{2}\xi Q\xi}\mathcal{F}a(x).$$

(a) If $a \in A_{\delta}^{m}(\mathbb{R}^{d})$, then the pointwise evaluation is given by the oscillatory integral.

$$e^{\frac{i}{2}DQD}a(x) = (2\pi)^{-d} \frac{e^{\frac{i\pi \operatorname{sgn} Q}{4}}}{|\det Q|^{\frac{1}{2}}} \int_{\mathbb{R}^d} e^{-\frac{i}{2}yQ^{-1}y} a(x+y) \, dy.$$

(b)
$$e^{\frac{i}{2}DQD}a(x) = \sum_{k=0}^{n} \frac{i^{k}}{2^{k}k!} (DQD)^{k} a(x) + r_{n}(x)$$

Foureir restriction

Kakeya Bochner-Riesz Geometric measure theory

Part IV Pseudo-differential operators

Pseudo-differential calculus

10.1

10.1 (Hörmander symbol classes). Let $m, \rho, \delta \in \mathbb{R}$. The Hörmander class $S_{\rho, \delta}^m(\mathbb{R}^{2d})$ of symbols is the set of smooth functions $a \in C^{\infty}(\mathbb{R}^d_x \times \mathbb{R}^d_{\varepsilon})$ such that

$$|\partial_x^{\alpha}\partial_\xi^{\beta}a(x,\xi)| \lesssim_{\alpha,\beta} \langle \xi \rangle^{m+\delta|\alpha|-\rho|\beta|}$$

for each $\alpha, \beta \in \mathbb{Z}_{\geq 0}^d$.

(a) Fréchet space

10.2 (Asymptotic expansion). Let $\rho, \delta \in \mathbb{R}$ and $(m_k)_{k=0}^{\infty} \subset \mathbb{R}$ be a sequence with m_0 and $m_k \downarrow -\infty$. Given $a_k \in S_{\rho,\delta}^{m_k}(\mathbb{R}^{2d})$, we want to construct $a \in S_{\rho,\delta}^{m_0}(\mathbb{R}^{2d})$ such that

$$a - \sum_{k=0}^{n} a_k \in S_{\rho,\delta}^{m_{n+1}}(\mathbb{R}^{2d}). \tag{\dagger}$$

The symbol a_0 is called the *principal symbol* of a, or the operator $Op^t(a)$.

Let $\chi \in C_c^{\infty}(\mathbb{R}^d_{\xi}, [0, 1])$ be a cutoff function such that

$$\chi(\xi) = \begin{cases} 1, & \text{if } |\xi| \le 1 \\ 0, & \text{if } |\xi| \ge 2 \end{cases}.$$

- (a) If $a \in S^m_{\rho,\delta}$, then $\chi(\varepsilon\xi)a(x,\xi)$ is uniformly bounded in $S^m_{\rho,\delta}$ for $\varepsilon \in (0,1)$ if $\rho \le 1$.
- (b) There is $a \in S_{\rho,\delta}^{m_0}$ such that (†) if $\rho \leq 1$.

Proof. (a) On the support of $\xi \mapsto \chi(\varepsilon \xi)$ holds $\langle \xi \rangle < 2|\xi| \le 4\varepsilon^{-1}$ because $1 < \varepsilon^{-1}$, so for each $\alpha, \beta \in \mathbb{Z}_{\ge 0}^d$ we have

$$\begin{split} |\partial_{x}^{\alpha}\partial_{\xi}^{\beta}(\chi(\varepsilon\xi)a(x,\xi))| &= |\sum_{\tau}\binom{\beta}{\tau}\partial_{\xi}^{\beta-\tau}(\chi(\varepsilon\xi))\partial_{x}^{\alpha}\partial_{\xi}^{\tau}a(x,\xi)| \\ &= |\sum_{\tau}\binom{\beta}{\tau}\varepsilon^{|\beta|-|\tau|}\partial_{\xi}^{\beta-\tau}\chi(\varepsilon\xi)\partial_{x}^{\alpha}\partial_{\xi}^{\tau}a(x,\xi)| \\ &(\because \langle \xi \rangle \leq 4\varepsilon^{-1}) \quad \leq \sum_{\tau}\binom{\beta}{\tau}(4\langle \xi \rangle^{-1})^{|\beta|-|\tau|}|\partial_{\xi}^{\beta-\tau}\chi(\varepsilon\xi)||\partial_{x}^{\alpha}\partial_{\xi}^{\tau}a(x,\xi)| \\ &\lesssim \sum_{\tau}\binom{\beta}{\tau}\langle \xi \rangle^{-(|\beta|-|\tau|)}\langle \xi \rangle^{m+\delta|\alpha|-\rho|\tau|} \\ &(\because \rho \leq 1) \quad \leq \langle \xi \rangle^{m+\delta|\alpha|-\rho|\beta|}. \end{split}$$

(b) Because we have $\varepsilon^{-1} \leq \langle \xi \rangle$ on the support of $1 - \chi(\varepsilon \xi)$, for each k we can take a sequence ε_k small enough such that

$$\max_{\substack{\alpha,\beta\in\mathbb{Z}_0^{\perp}\\|\alpha|+|\beta|\leq k}} |\partial_x^{\alpha}\partial_{\xi}^{\beta}((1-\chi(\varepsilon_k\xi))a_k(x,\xi))| \leq 2^{-k}\langle\xi\rangle^{m_k+1+\delta|\alpha|-\rho|\beta|}.$$

We may assume $\varepsilon_k \downarrow 0$ so that the following sum is locally finite:

$$a(x,\xi) := \sum_{k=0}^{\infty} (1 - \chi(\varepsilon_k \xi)) a_k(x,\xi).$$

If we choose *n* such that $m_0 \ge m_{n+1} + 1$, then in the expansion

$$a(x,\xi) = \sum_{k=0}^{n} (1 - \chi(\varepsilon_k \xi)) a_k(x,\xi) + \sum_{k=n+1}^{\infty} (1 - \chi(\varepsilon_k \xi)) a_k(x,\xi)$$

the first sum clearly belongs to $S_{\rho,\delta}^{m_0}$ and so is the second sum because

$$\begin{split} |\partial_x^\alpha \partial_\xi^\beta \sum_{k=n+1}^\infty (1-\chi(\varepsilon_k \xi)) a_k(x,\xi)| &\leq \sum_{k=n+1}^\infty 2^{-k} \langle \xi \rangle^{m_{k+1}+1+\delta|\alpha|-\rho|\beta|} \\ &\leq \langle \xi \rangle^{m_{n+1}+1+\delta|\alpha|-\rho|\beta|} \\ &\leq \langle \xi \rangle^{m_0+\delta|\alpha|-\rho|\beta|} \end{split}$$

for every $\alpha, \beta \in \mathbb{Z}_{\geq 0}^d$. Therefore, $a \in S_{\rho, \delta}^{m_0}$.

Write

$$(a-\sum_{k=0}^{n}a_{k})(x,\xi)=\sum_{k=0}^{n}\chi(\varepsilon_{k}\xi)a_{k}(x,\xi)+\sum_{k=n+1}^{\infty}(1-\chi(\varepsilon_{k}\xi))a_{k}(x,\xi).$$

The first sum belongs to $S^{-\infty}$ because it is compactly supported, and we can also show that the second sum belongs to $S^{m_{n+1}}_{\rho,\delta}$ by decomposing with n' such that $m_{n+1} \geq m'_n + 1$ and by considering the multiplication with a cutoff remains in the same symbol class.

10.3 (Quantization). The *t*-quantization of a symbol *a* is the pseudo-differential operator $\operatorname{Op}^t(a)$ on $\mathcal{S}(\mathbb{R}^d_*)$ defined by

$$\operatorname{Op}^{t}(a)f(x) := (2\pi)^{-d} \int_{\mathbb{R}^{2d}} e^{i(x-y)\xi} a((1-t)x + ty, \xi)f(y) \, dy \, d\xi$$

for all $f \in \mathcal{S}(\mathbb{R}^d)$. Kohn-Nirenberg calculus for t = 0, Weyl calculus for $t = \frac{1}{2}$.

- (a) $\operatorname{Op}^0(a): \mathcal{S}(\mathbb{R}^d) \to \mathcal{S}'(\mathbb{R}^d)$ is continuous for $a \in \mathcal{S}'(\mathbb{R}^d)$.
- (b) $\operatorname{Op}^0(a): \mathcal{S}(\mathbb{R}^d) \to \mathcal{S}(\mathbb{R}^d)$ is continuous for $a \in S^m_{o,\delta}(\mathbb{R}^{2d})$ if $\delta \leq 1$.

Proof. (b) Since $(D_{\gamma})^2$ is a self-adjoint partial differential operator, for any $n \in \mathbb{Z}_{\geq 0}$ we have

$$\operatorname{Op^{0}}(a)f(x) = (2\pi)^{-d} \int_{\mathbb{R}^{2d}} e^{i(x-y)\xi} a(x,\xi) f(y) \, dy \, d\xi$$

$$(\because D_{y}e^{i(x-y)\xi} = \xi e^{i(x-y)\xi}) = (2\pi)^{-d} \int_{\mathbb{R}^{2d}} \langle \xi \rangle^{-2n} \langle D_{y} \rangle^{2n} e^{i(x-y)\xi} a(x,\xi) f(y) \, dy \, d\xi$$

$$(\because \operatorname{IBP}) = (2\pi)^{-d} \int_{\mathbb{R}^{2d}} e^{i(x-y)\xi} \langle \xi \rangle^{-2n} a(x,\xi) \langle D_{y} \rangle^{2n} f(y) \, dy \, d\xi.$$

The derivatives of the integrand is integrable with respect to ξ for a sufficiently large n with $m + |\beta| - 2n < -d$ because

$$\begin{split} |\partial_x^\beta (e^{i(x-y)\xi} \langle \xi \rangle^{-2n} a(x,\xi) \langle D_y \rangle^{2n} f(y))| \\ &= |\sum_\tau \binom{\beta}{\tau} (i\xi)^{\beta-\tau} e^{i(x-y)\xi} \langle \xi \rangle^{-2n} \partial_x^\tau a(x,\xi) \langle D_y \rangle^{2n} f(y)| \\ &\leq \sum_\tau \binom{\beta}{\tau} \langle \xi \rangle^{|\beta|-|\tau|} \langle \xi \rangle^{-2n} |\partial_x^\tau a(x,\xi)| |\langle D_y \rangle^{2n} f(y)| \\ (\because a \in S_{\rho,\delta}^m) &\lesssim \sum_\tau \binom{\beta}{\tau} \langle \xi \rangle^{|\beta|-|\tau|} \langle \xi \rangle^{-2n} \langle \xi \rangle^{m+\delta|\tau|} |\langle D_y \rangle^{2n} f(y)| \\ (\because \delta \leq 1) &\lesssim \langle \xi \rangle^{m+|\beta|-2n} |\langle D_y \rangle^{2n} f(y)|, \end{split}$$

so the partial derivative ∂_x commutes with the integral. Since

$$x^{\alpha}e^{i(x-y)\xi} = (y+D_{\xi})^{\alpha}e^{i(x-y)\xi} = \sum_{\sigma} {\alpha \choose \sigma} y^{\alpha-\sigma}D_{\xi}^{\sigma}e^{i(x-y)\xi},$$

we have an expansion

$$\begin{split} x^{\alpha}\partial_{x}^{\beta}\operatorname{Op^{0}}(a)f(x) &= x^{\alpha}\partial_{x}^{\beta}\int_{\mathbb{R}^{2d}}e^{i(x-y)\xi}\langle\xi\rangle^{-2n}a(x,\xi)\langle D_{y}\rangle^{2n}f(y))\,dy\,d\xi\\ &= \int_{\mathbb{R}^{2d}}x^{\alpha}\partial_{x}^{\beta}(e^{i(x-y)\xi}\langle\xi\rangle^{-2n}a(x,\xi)\langle D_{y}\rangle^{2n}f(y))\,dy\,d\xi\\ &= \int_{\mathbb{R}^{2d}}\sum_{\sigma,\tau}\binom{\alpha}{\sigma}\binom{\beta}{\tau}y^{\alpha-\sigma}D_{\xi}^{\sigma}e^{i(x-y)\xi}(i\xi)^{\beta-\tau}\langle\xi\rangle^{-2n}\partial_{x}^{\tau}a(x,\xi)\langle D_{y}\rangle^{2n}f(y)\,dy\,d\xi\\ &= \int_{\mathbb{R}^{2d}}\sum_{\sigma,\tau}\binom{\alpha}{\sigma}\binom{\beta}{\tau}e^{i(x-y)\xi}(-D_{\xi})^{\sigma}[(i\xi)^{\beta-\tau}\langle\xi\rangle^{-2n}\partial_{x}^{\tau}a(x,\xi)]y^{\alpha-\sigma}\langle D_{y}\rangle^{2n}f(y)\,dy\,d\xi. \end{split}$$

Here

$$\sup_{x \in \mathbb{R}^d} |(-D_{\xi})^{\sigma} [(i\xi)^{\beta-\tau} \langle \xi \rangle^{-2n} \partial_x^{\tau} a(x,\xi)]|$$

is integrable with respect to ξ for sufficiently large n, so with this n we have

$$\sup_{x \in \mathbb{R}^d} |x^{\alpha} \partial_x^{\beta} \operatorname{Op}^0(a) f(x)| \lesssim \sum_{\alpha \leq \alpha} \sup_{y \in \mathbb{R}^d} |y^{\alpha - \sigma} \langle D_y \rangle^{2n} f(y)|$$

for each $\alpha, \beta \in \mathbb{Z}^d_{\geq 0}$ and all $f \in \mathcal{S}(\mathbb{R}^d)$, which implies $\operatorname{Op}^0(a) f \in \mathcal{S}(\mathbb{R}^d)$.

10.4 (Change of quantization). Let $m \in \mathbb{R}$, .

- (a) $Op^{t}(a) = Op^{0}(e^{itD_{x}D_{\xi}}a)$.
- (b) $a \in S^m_{\rho,\delta}(\mathbb{R}^{2d})$ if and only if $e^{itD_xD_\xi}a \in S^m_{\rho,\delta}(\mathbb{R}^{2d})$, if $0 \le \delta \le \rho \le 1$ and $\delta < 1$.
- (c) We have the formal adjoint

$$\operatorname{Op}^{t}(a)^{*} = \operatorname{Op}^{1-t}(\overline{a}).$$

In particular, we have $\operatorname{Op}^{t}(a): \mathcal{S}' \to \mathcal{S}'$.

Proof. (a) Note that

$$\operatorname{Op}^{t}(a)f(x) = (2\pi)^{-d} \int_{\mathbb{R}^{2d}} e^{i(x-y)\xi} a((1-t)x + ty, \xi) f(y) \, dy \, d\xi$$

$$(\because \text{Inversion on } \mathbb{R}^{2d}) = (2\pi)^{-3d} \int_{\mathbb{R}^{4d}} e^{i(x-y)\xi} e^{i((1-t)x + ty)x^* + i\xi\xi^*} \hat{a}(x^*, \xi^*) f(y) \, dx^* \, d\xi^* \, dy \, d\xi$$

$$= (2\pi)^{-3d} \int_{\mathbb{R}^{4d}} e^{i(x-y + \xi^*)\xi} \hat{a}(x^*, \xi^*) e^{i((1-t)x + ty)x^*} f(y) \, dx^* \, d\xi^* \, dy \, d\xi$$

$$(\because \text{Inversion on } \mathbb{R}^d) = -(2\pi)^{-2d} \int_{\mathbb{R}^{2d}} \hat{a}(x^*, y - x) e^{i((1-t)x + ty)x^*} f(y) \, dx^* \, dy$$

$$(\because [\xi^*/y - x]) = -(2\pi)^{-2d} \int_{\mathbb{R}^{2d}} e^{i(x + t\xi^*)x^*} \hat{a}(x^*, \xi^*) f(x + \xi^*) \, dx^* \, d\xi^*.$$

(b) We have the oscillatory integral

$$e^{itD_xD_\xi}a(x,\xi) = (2\pi)^{-d}|t|^{-d} \int_{\mathbb{R}^{2d}} e^{-it^{-1}y\eta}a(x+y,\xi+\eta)\,dy\,d\eta.$$

Enough to show

$$\left| \int_{\mathbb{R}^{2d}} e^{-it^{-1}y\eta} a(x+y,\xi+\eta) \, dy \, d\eta \right| \lesssim \langle \xi \rangle^m.$$

Fix ξ and $\delta \leq \rho$

10.5 (Moyal product). Let $a \in S^m_{\rho,\delta}(\mathbb{R}^{2d})$ and $b \in S^l_{\rho,\delta}(\mathbb{R}^{2d})$.

(a) there exists a unique function $a\#^t b \in S^{m+l}_{\rho,\delta}(\mathbb{R}^{2d})$ such that

$$a^{t}(x,D)b^{t}(x,D) = (a\#^{t})^{t}(x,D).$$

(b) It is concretely described by

$$(a\#^t b)(x,\xi) = (2\pi)^{-2} \int_{\mathbb{R}^{4d}} e^{-i(y\eta - z\zeta)} a(x+tz,\xi+\eta) b((1-t)y+x,\xi+\zeta) \, dy \, d\eta \, dz \, d\zeta.$$

(c) If $\delta < \rho$, then

$$a^{\#t}b(x,\xi) \sim \sum_{k \in \mathbb{Z}_{\geq 0}} \frac{1}{i^k k!} (\partial_y \partial_\eta - \partial_z \partial_\zeta)^k a((1-t)x + tz, \eta) b(tx + (1-t)y, \zeta) \Big|_{\substack{y=z=x,\\ \eta=\zeta=\xi}}.$$

10.6 (Parametirx and elliptic operators).

10.2

10.7 (Calderón-Vaillancourt theorem).

Semiclassical analysis

For parameters $0 \le \lambda \le 1$ and h > 0, let

$$\widehat{a}\psi(x) := \frac{1}{(2\pi h)^d} \int \int e^{\frac{i}{h}\langle x-y,\xi\rangle} a((1-\lambda)x + \lambda y,\xi)\psi(y) \, dy \, d\xi.$$

For example, regardless of h and λ ,

$$\hat{\xi}\psi(x) = \frac{h}{i}\psi'(x)$$

and

$$\hat{H}\psi(x) = -h^2\Delta\psi(x) + V(x)\psi(x),$$

where $V: \mathbb{R}^d_x \times \mathbb{R}^d_\xi \to \mathbb{R}$ and $H: \mathbb{R}^d_x \times \mathbb{R}^d_\xi \to \mathbb{R}$ such that

$$H(x,\xi) := |\xi|^2 + V(x).$$

$$\begin{split} \frac{d}{dt}a(t) &= \{a(t), H\} = X_H a(t) \\ \frac{d}{dt}\hat{a}(t) &= \frac{d}{dt}e^{\frac{i}{\hbar}t\hat{H}}\hat{a}e^{-\frac{i}{\hbar}t\hat{H}} = -\frac{i}{\hbar}[\hat{a}(t), \hat{H}] \end{split}$$

11.1 Heisenberg group

11.2 Phase space transforms

Microlocal analysis